

# Measuring the Jeans Scale of the IGM with Close Quasar Pairs

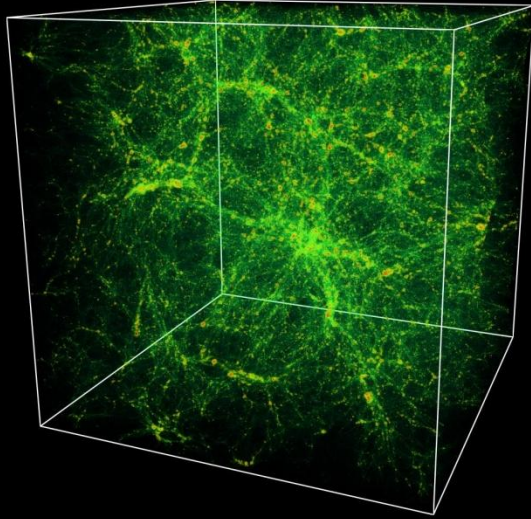
Alberto Rorai, Joseph Hennawi, Martin White



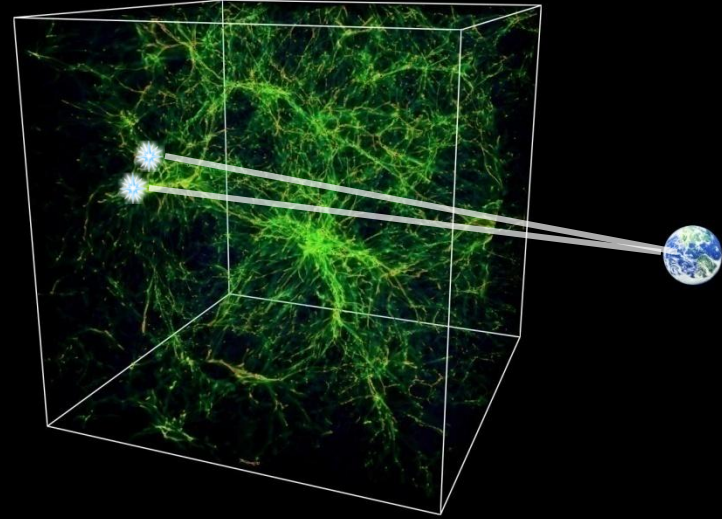
arXiv:1305.0210

# Baryons in the IGM and the Lyman- $\alpha$ Forest

## Dark Matter

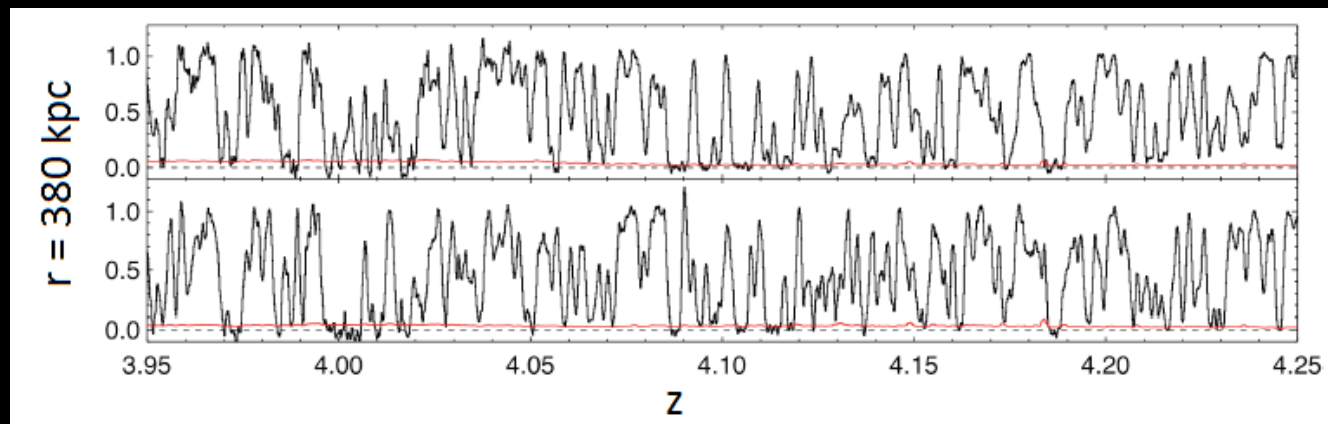


## Baryons



Credit: R. Cen

## Observations



# The Jeans Scale as a Thermal Record

$$\lambda_J = \frac{c_s}{1+z} \sqrt{\frac{\pi}{G\rho}} \approx 1.1 \text{ Mpc} \left( \frac{T}{10^4 \text{K}} \right)^{1/2} \left( \frac{1+z}{1+3} \right)^{-1/2}$$

T increases with reionization events → Jeans scale increases

- Small-scale-structure growth and Galaxy formation suppressed

**Filtering mass:**

$$M_F = \frac{4}{3} \pi \bar{\rho} \lambda_J^3$$

- **Clumpiness of the IGM** decreased by truncating power spectrum

$$C = 1 + \sigma_{\text{IGM}}^2 \equiv 1 + \int d \ln k \frac{k^3 P_{\text{IGM}}(k)}{2\pi^2}$$

- The Jeans filtering scale depends on the *whole thermal history* (Gnedin & Hui, 1998), because the sound-crossing time  $\lambda_J/c_s$  is comparable to the Hubble time

$$\lambda_J^2(t) = \int_0^t f(T(t')) dt'$$

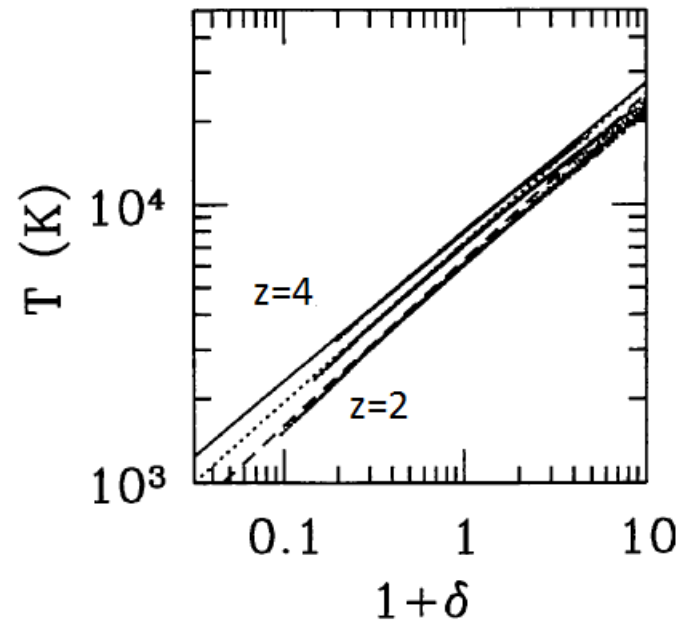
- It is a source of uncertainty in IGM studies

# The Thermal State of the IGM

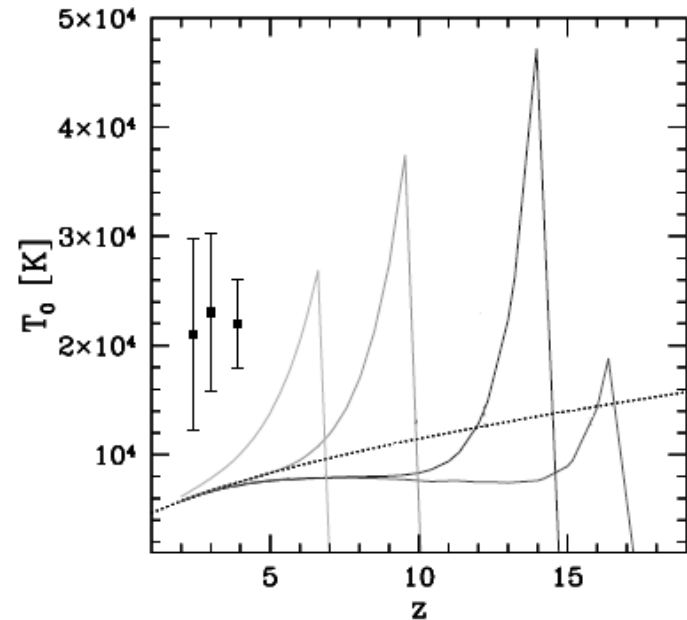
At photoionization equilibrium :

$$T(\rho) = T_0 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$$

Hui and Gnedin, 1997



Reionization at  
 $z = 10$



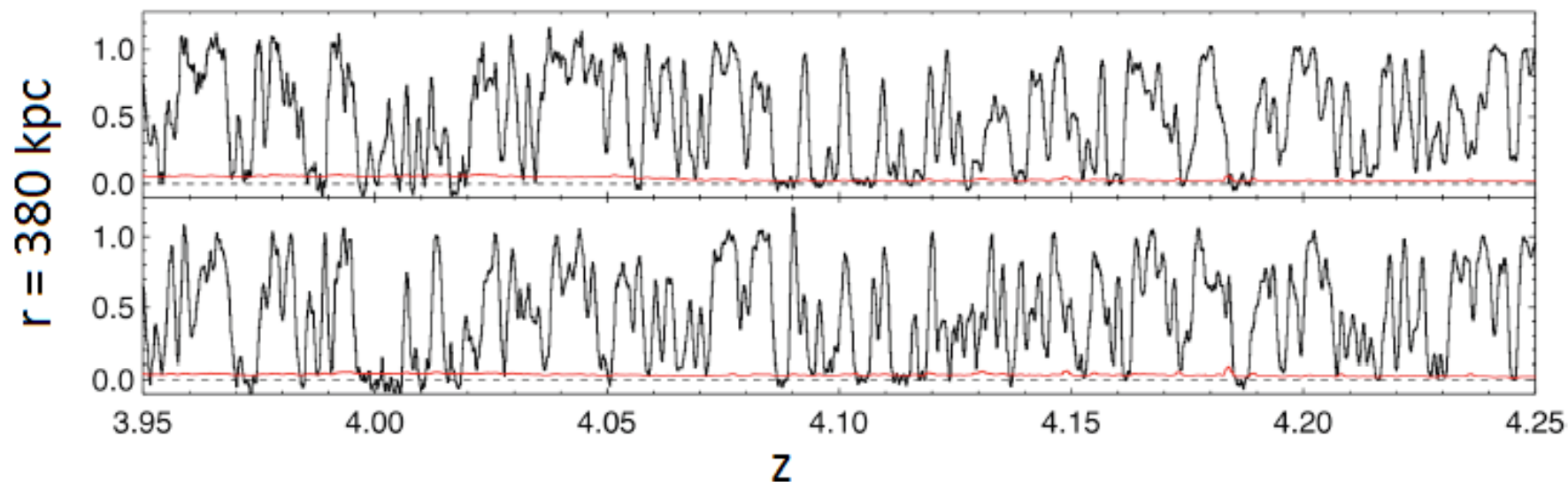
Hui and Haiman 2003, measurements  
from Zaldarriaga et al. 2001

The temperature of the IGM provides little information of early reionization events

# Probing the IGM with Quasar Pairs

We have a sample of  $\sim 300$  pairs at  $1.6 < z < 4.3$

(Hennawi et al. 2004,2006,2009)



- Close separation (380 comoving kpc)



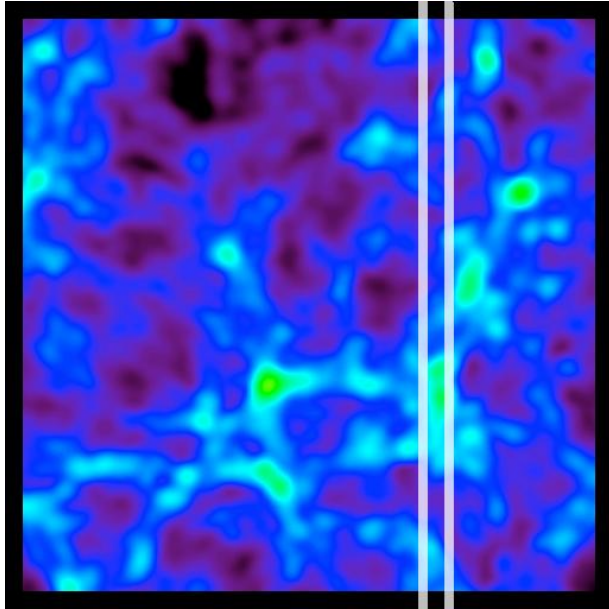
We are observing the  
Jeans scale at  $z=4$

- High degree of coherence

The effect of temperature on the Lyman- $\alpha$  forest is highly uncertain:

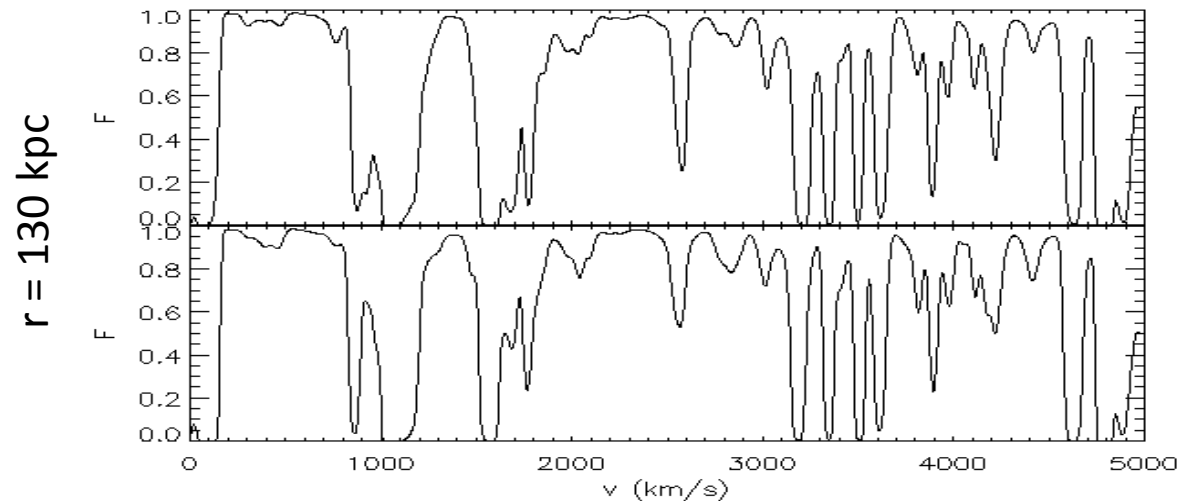
Resolving the Jeans scale in real space can give independent constraints

# Methodology: Simulations and Parameter Study



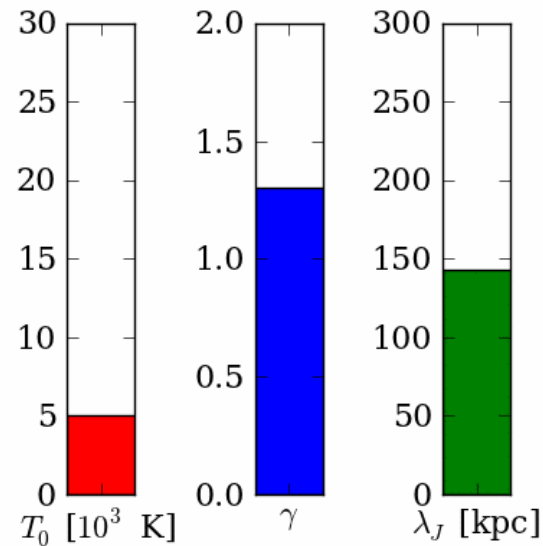
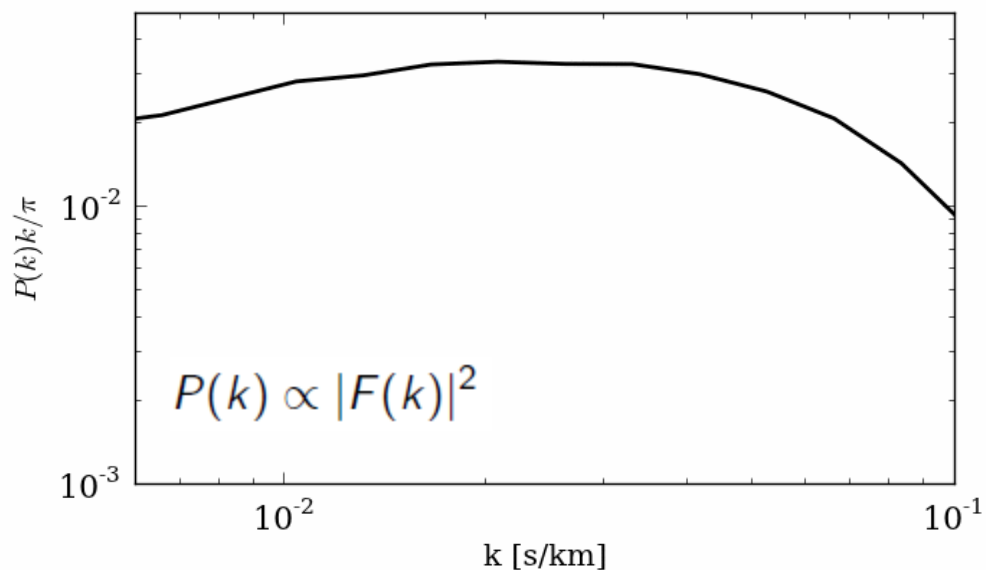
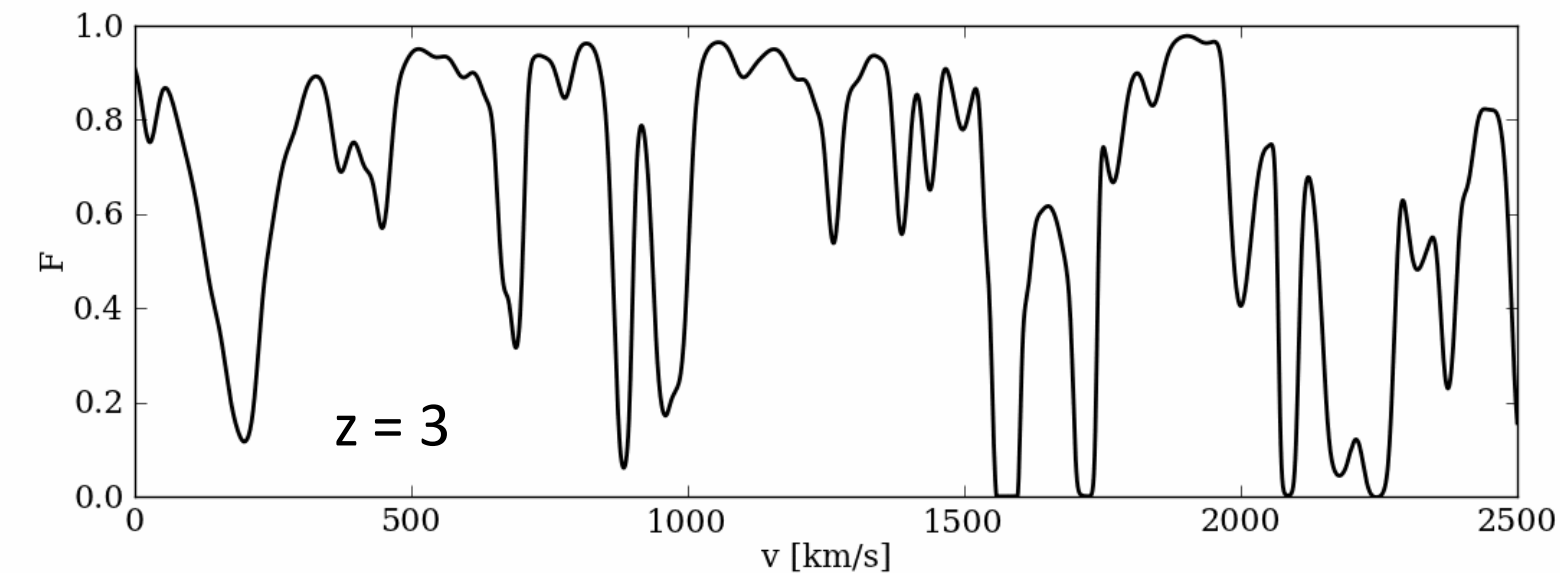
1. Dark matter only simulation, snapshot at  $z=3$   
( $Box=50 Mpc/h, n=1500^3$ )
2. Mimic pressure support by a smoothing density with a kernel of radius  $\lambda_J$
3. Definition of the equation of state

$$T(\rho) = T_0 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$$



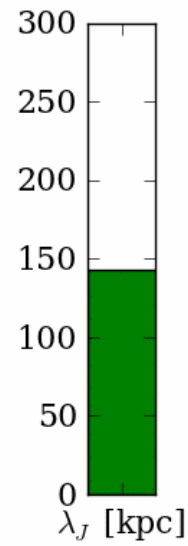
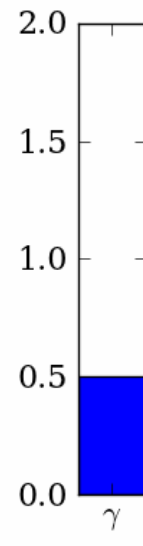
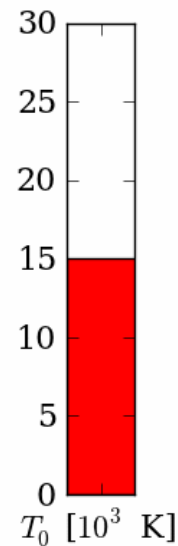
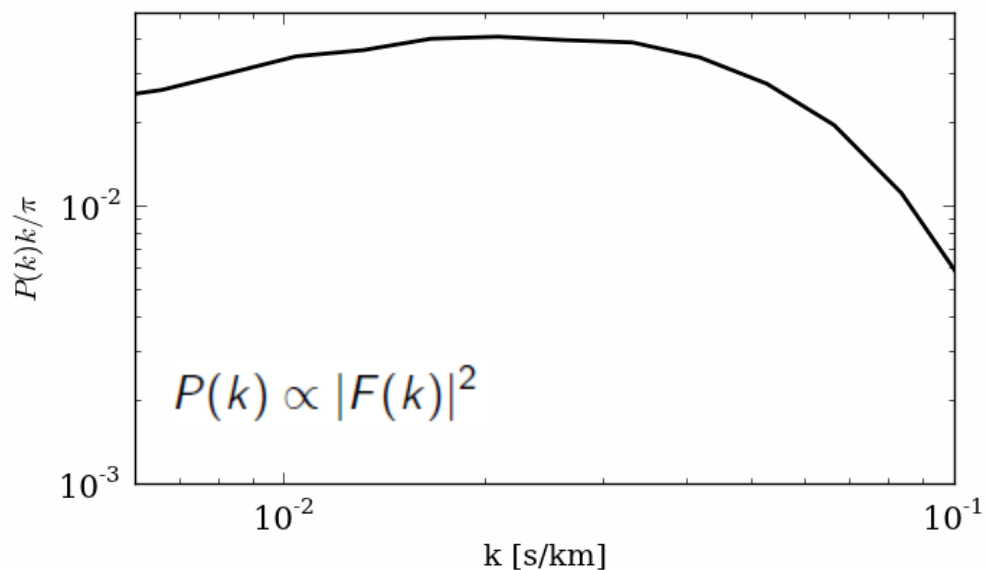
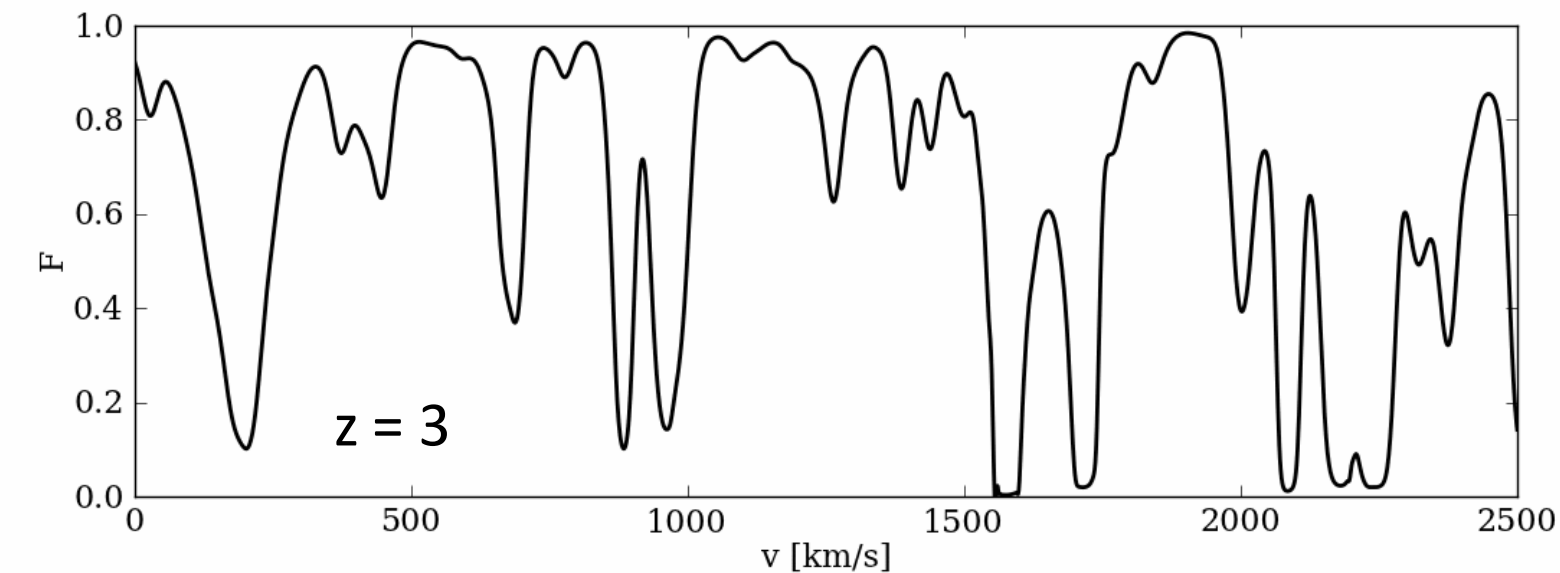
We calculate 500 different thermal models, defined by  $T_0, \gamma$  and  $\lambda_J$

# Sensitivity of the LOS Power Spectrum to 1D smoothing



$$T(\rho) = T_0 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$$

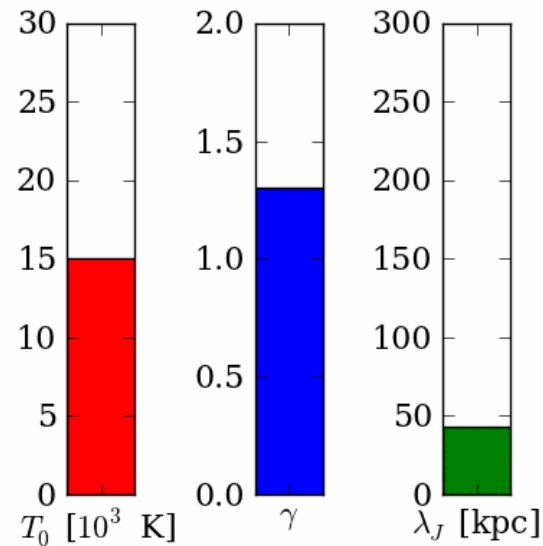
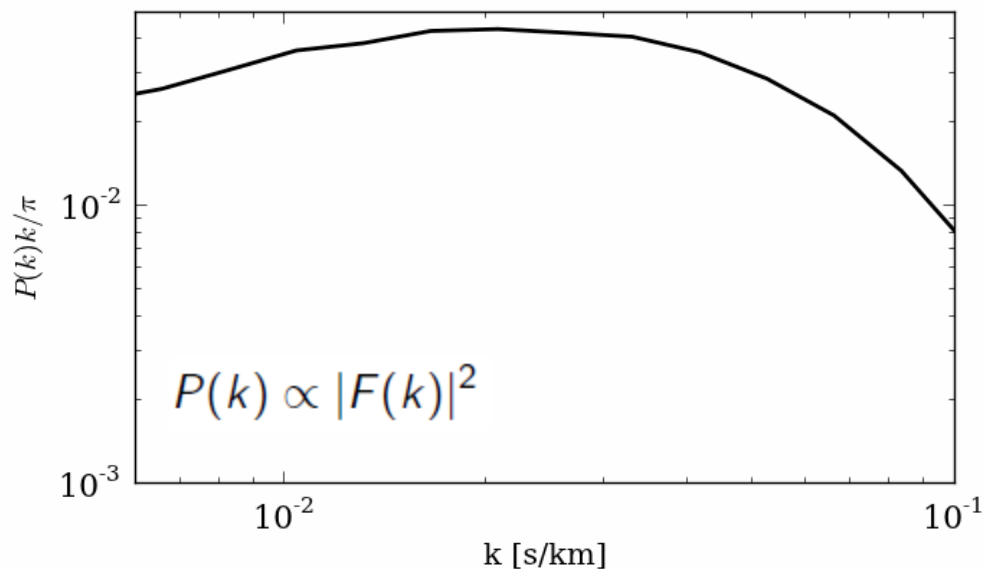
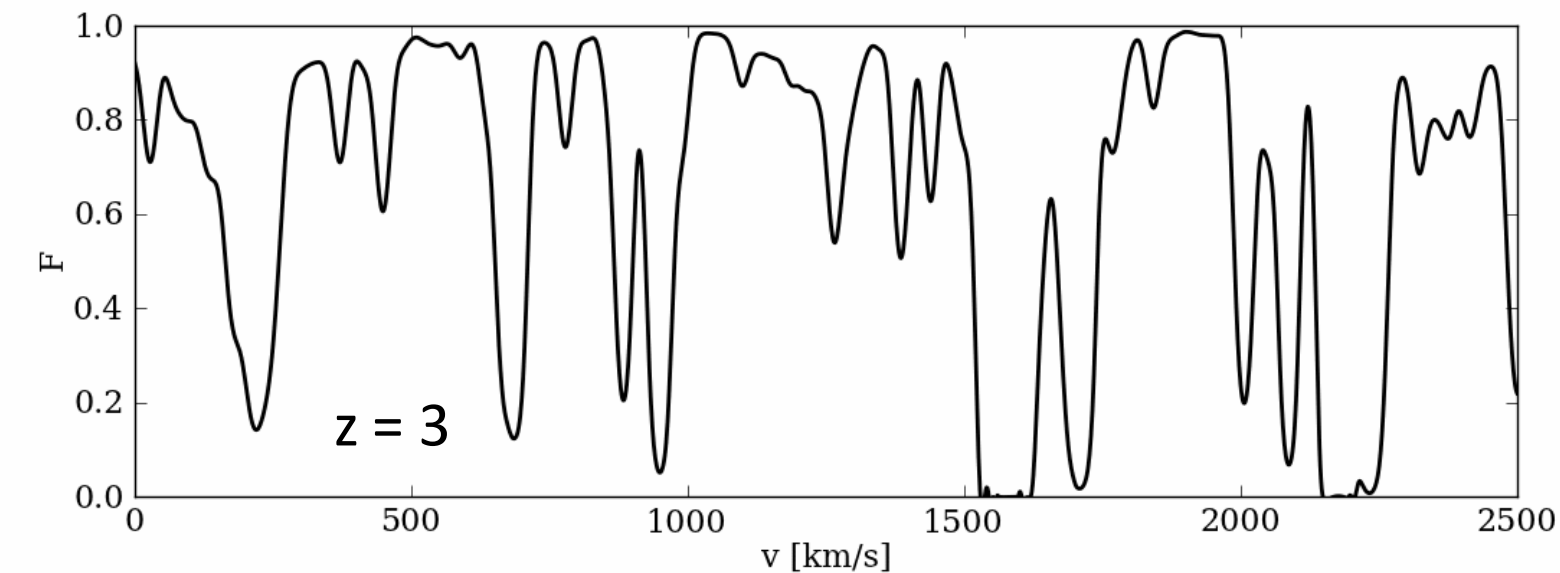
# Sensitivity of the LOS Power Spectrum to 1D smoothing



$$T(\rho) = T_0 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$$



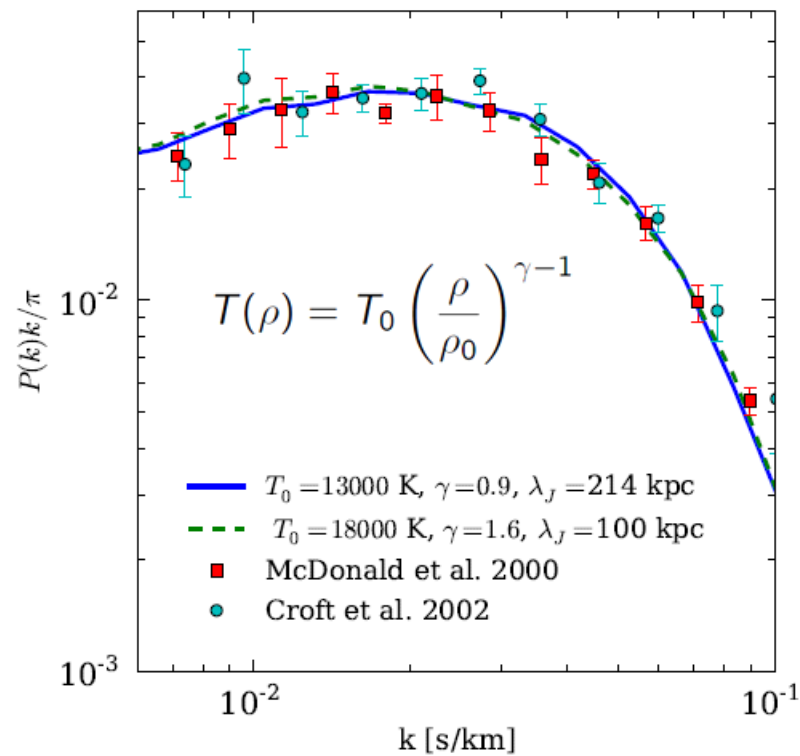
# Sensitivity of the LOS Power Spectrum to 3D smoothing



$$T(\rho) = T_0 \left( \frac{\rho}{\rho_0} \right)^{\gamma-1}$$

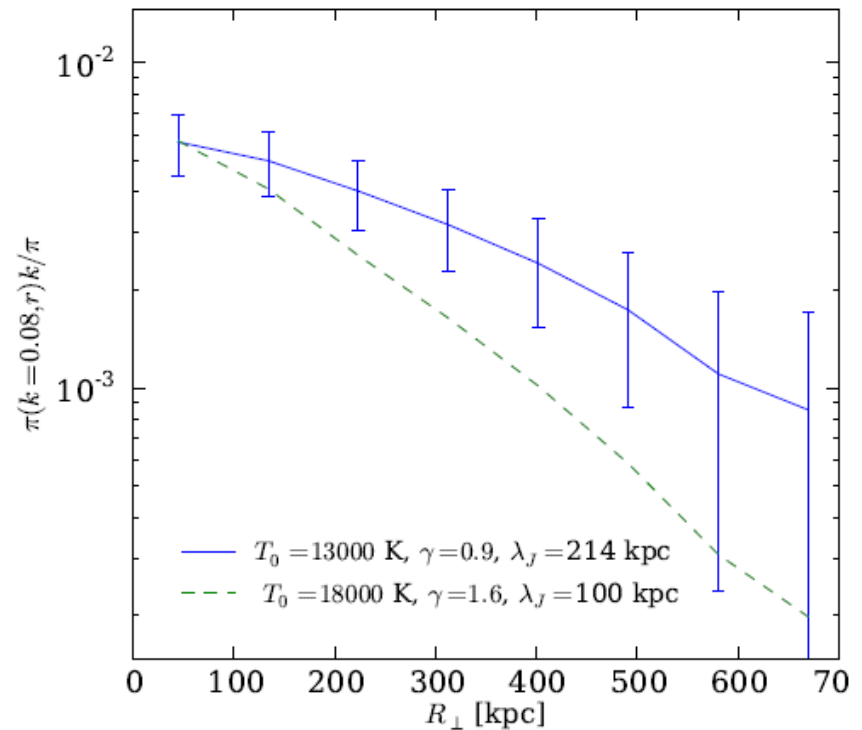
# 1D Degeneracies Broken by Pairs

Different thermal models compatible with the measured line-of-sight power



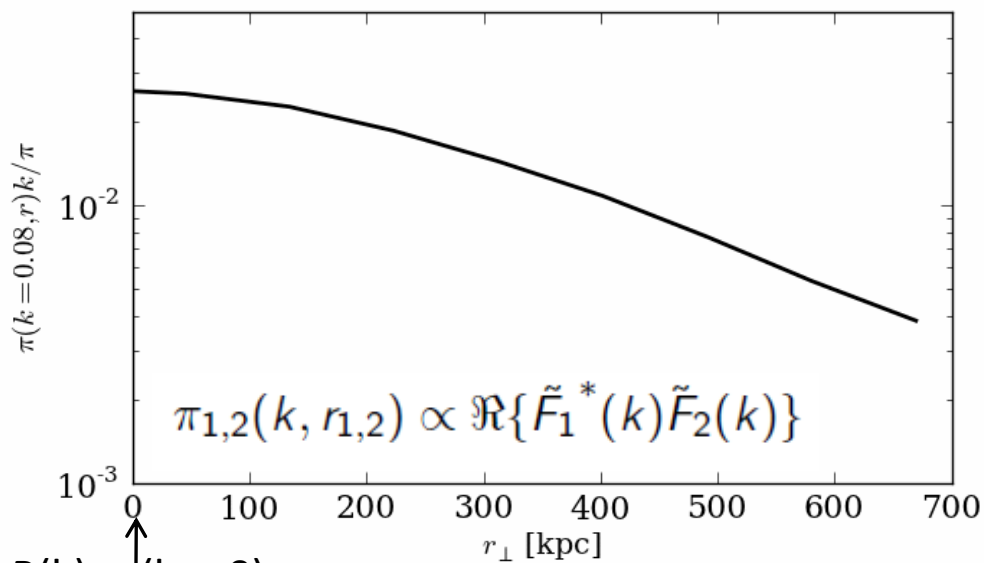
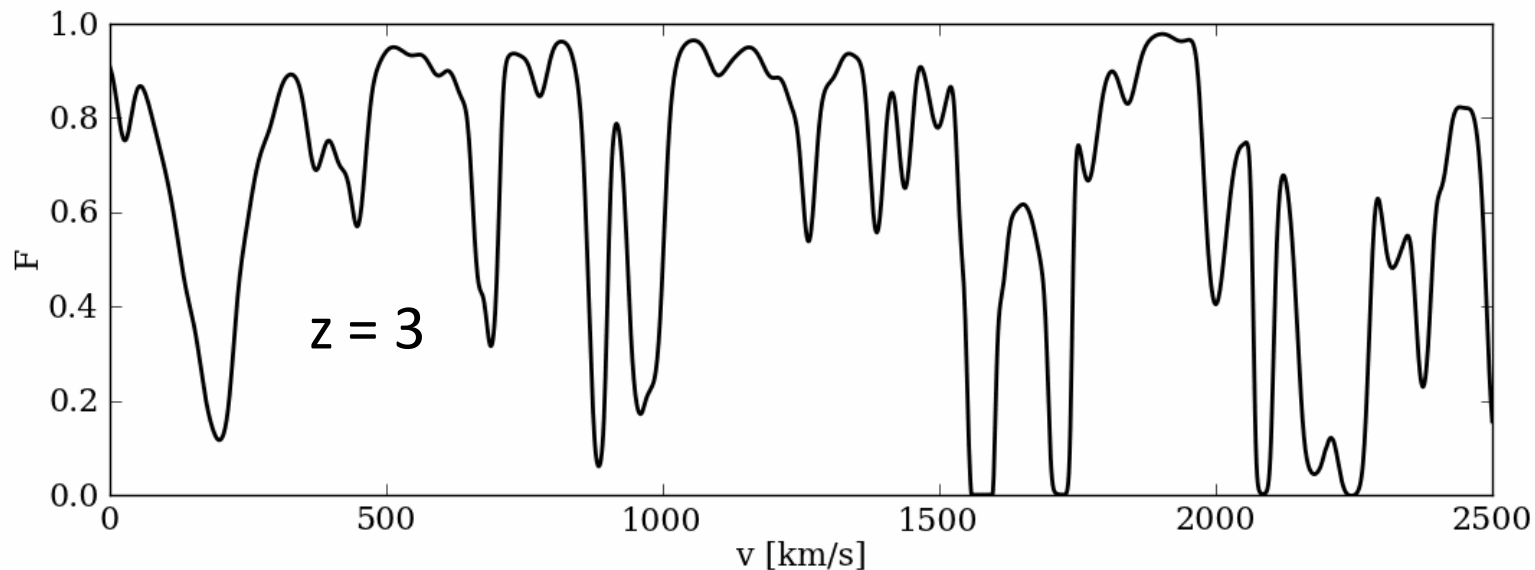
$$P(k) \propto |F(k)|^2$$

The Jeans scale dominates the transverse coherence

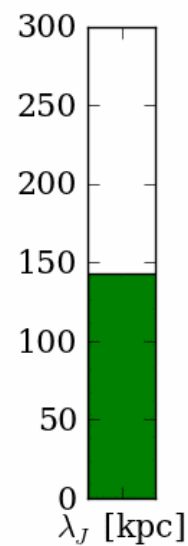
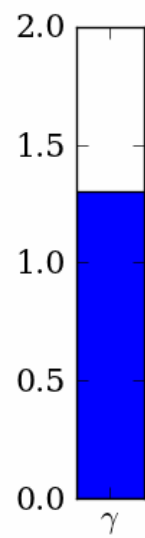
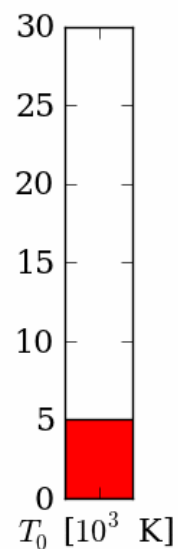


$$\pi_{1,2}(k, r_{1,2}) \propto \Re\{\tilde{F}_1^*(k)\tilde{F}_2(k)\}$$

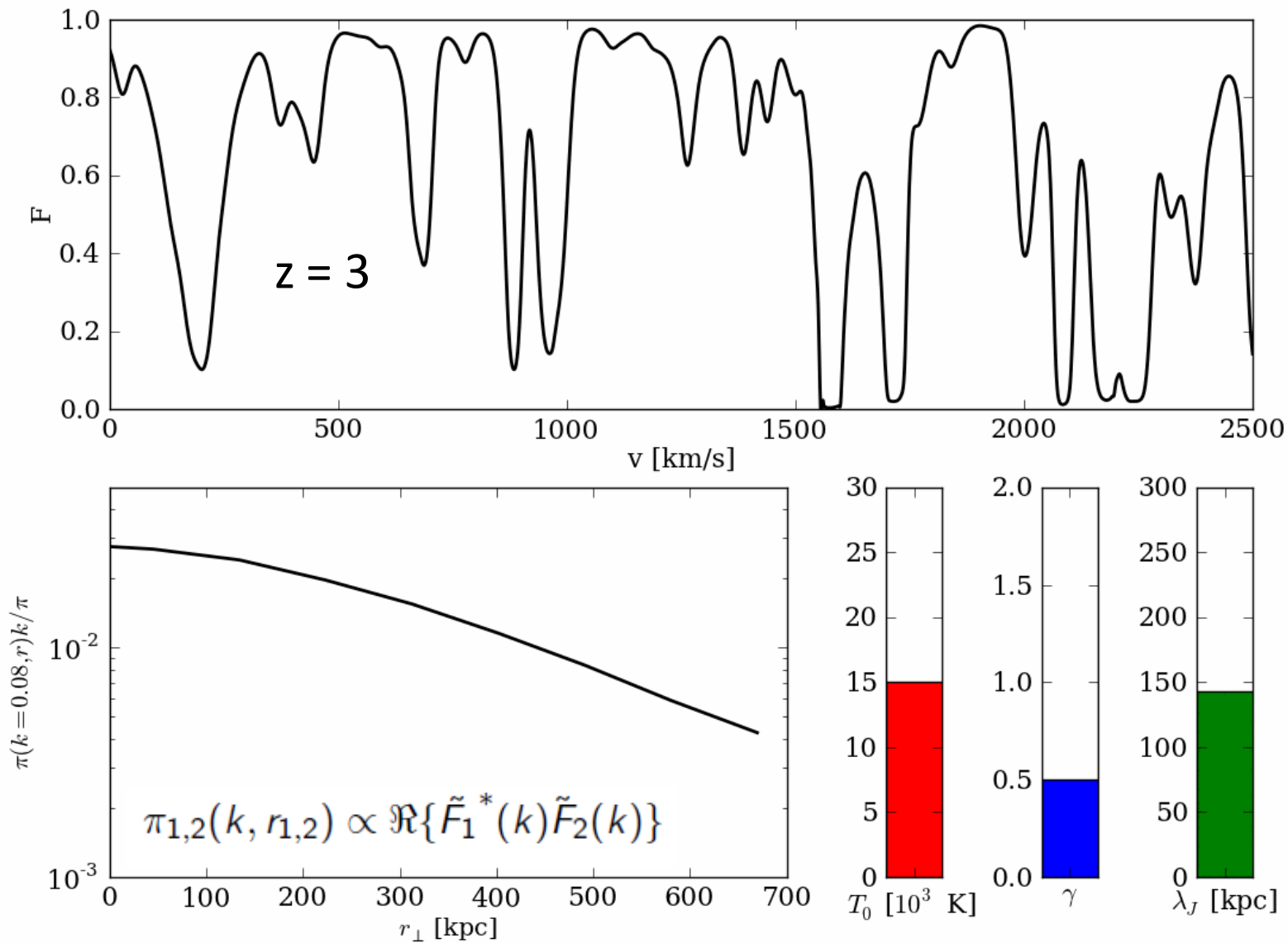
# Sensitivity of the Cross Power Spectrum to 1D smoothing



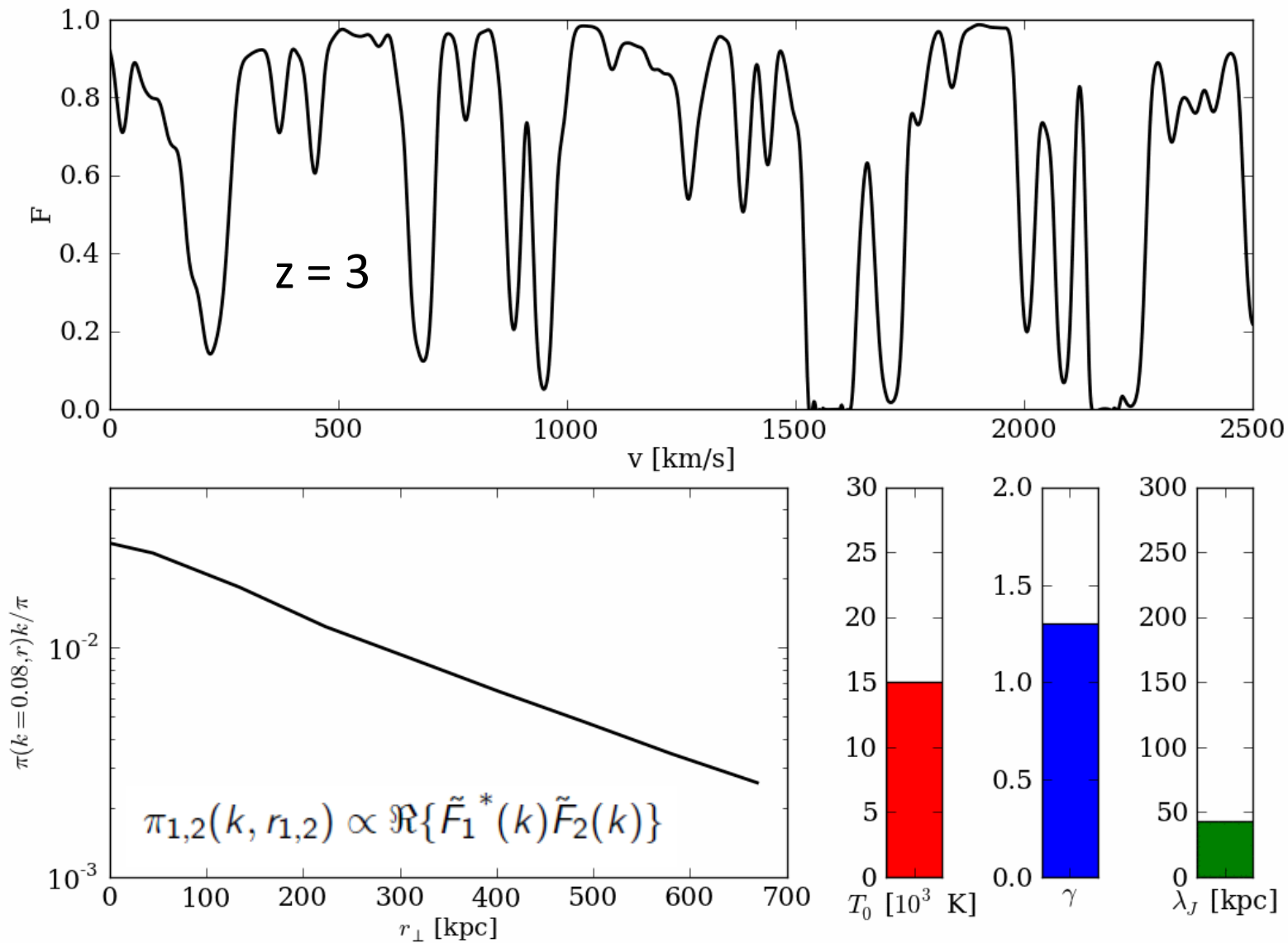
$$\pi_{1,2}(k, r_{1,2}) \propto \Re\{\tilde{F}_1^*(k)\tilde{F}_2(k)\}$$



# Sensitivity of the Cross Power Spectrum to 1D smoothing

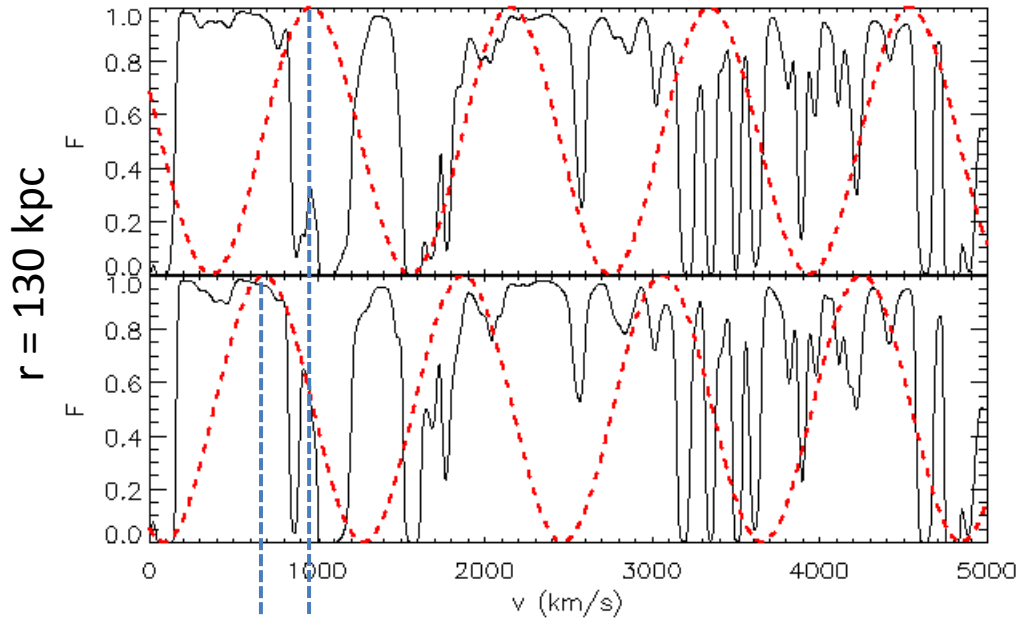


# Sensitivity of the Cross Power Spectrum to 3D smoothing



# Is Cross Power the Optimal Statistic?

Cross power defined in Fourier space as:  $\pi_{1,2}(k, r_{1,2}) \propto \Re\{\tilde{F}_1^*(k)\tilde{F}_2(k)\}$



In terms of moduli and phases:

$$\tilde{F}(k) = \rho(k)e^{i\theta(k)}$$

$$\pi(k, r_{1,2}) \propto \underbrace{\rho_1(k)\rho_2(k)}_{\text{amplitudes}} \underbrace{\cos\theta_{12}(k)}_{\text{angular part}}$$

The amplitudes of the modes depend on the LOS power:

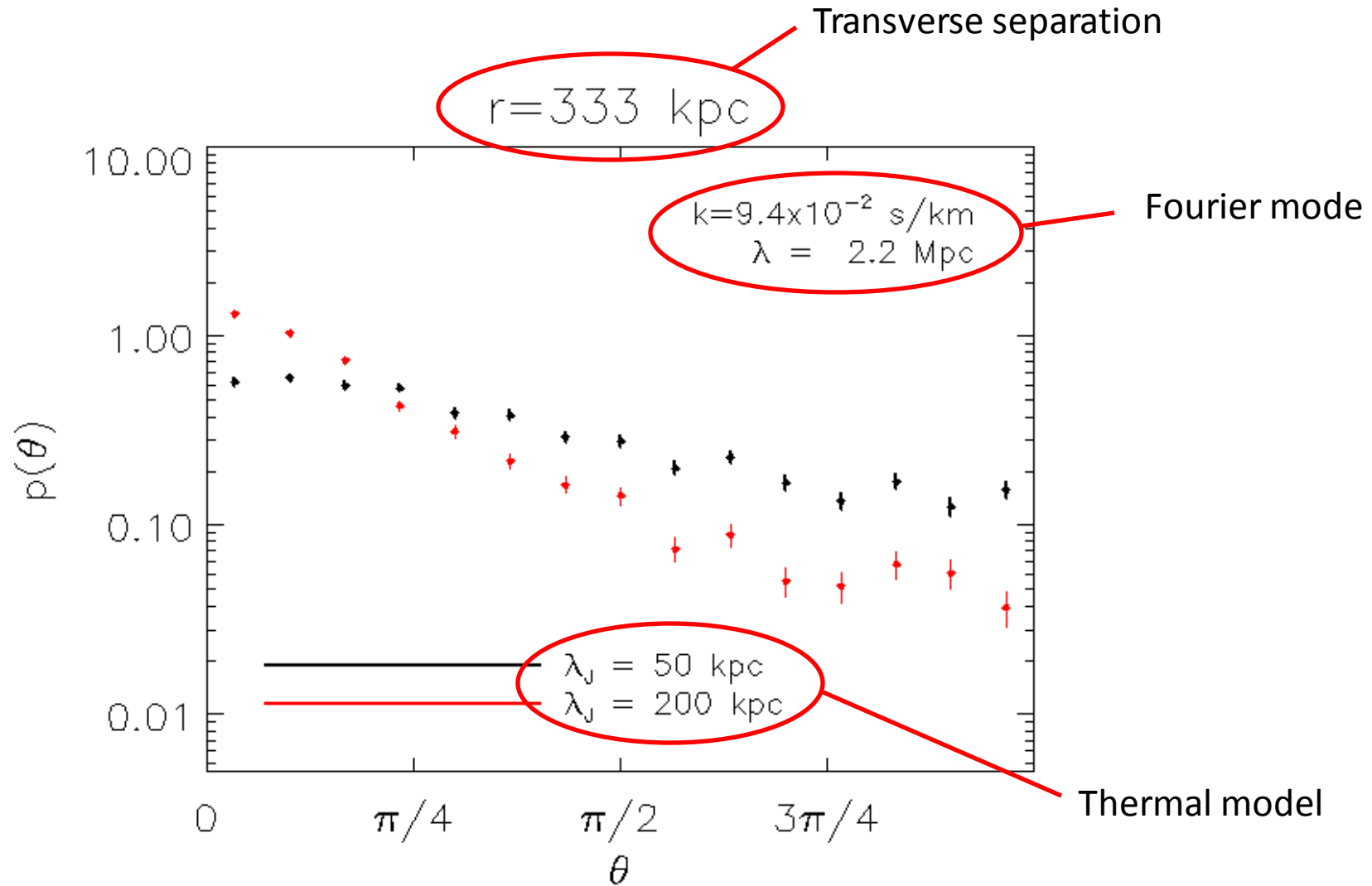
$$P(k) \sim \langle \rho_1\rho_2 \rangle$$

Angular part retains the new 3D information of pairs

$$\theta_{12}(k) = \arccos \left( \frac{\Re[\tilde{F}_1^*(k)\tilde{F}_2(k)]}{\sqrt{|\tilde{F}_1(k)|^2|\tilde{F}_2(k)|^2}} \right)$$

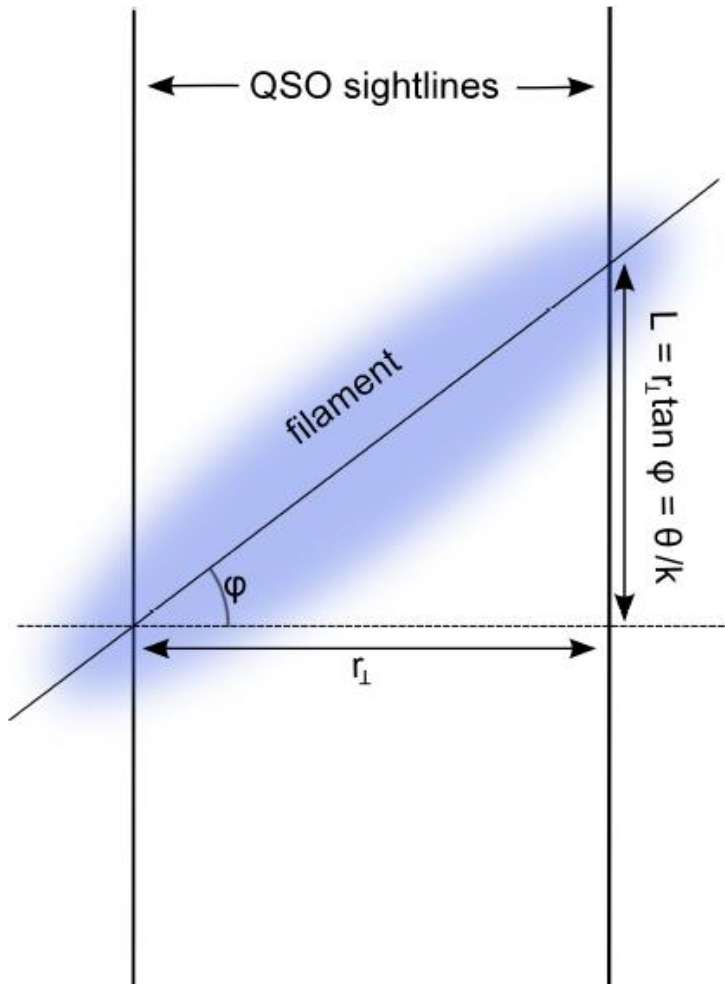
Problem: what is the best estimator for the angular component?

# Probability Distribution of Phases



We need a model for the probability distribution

# Statistical Distribution of the Phases: Ansatz



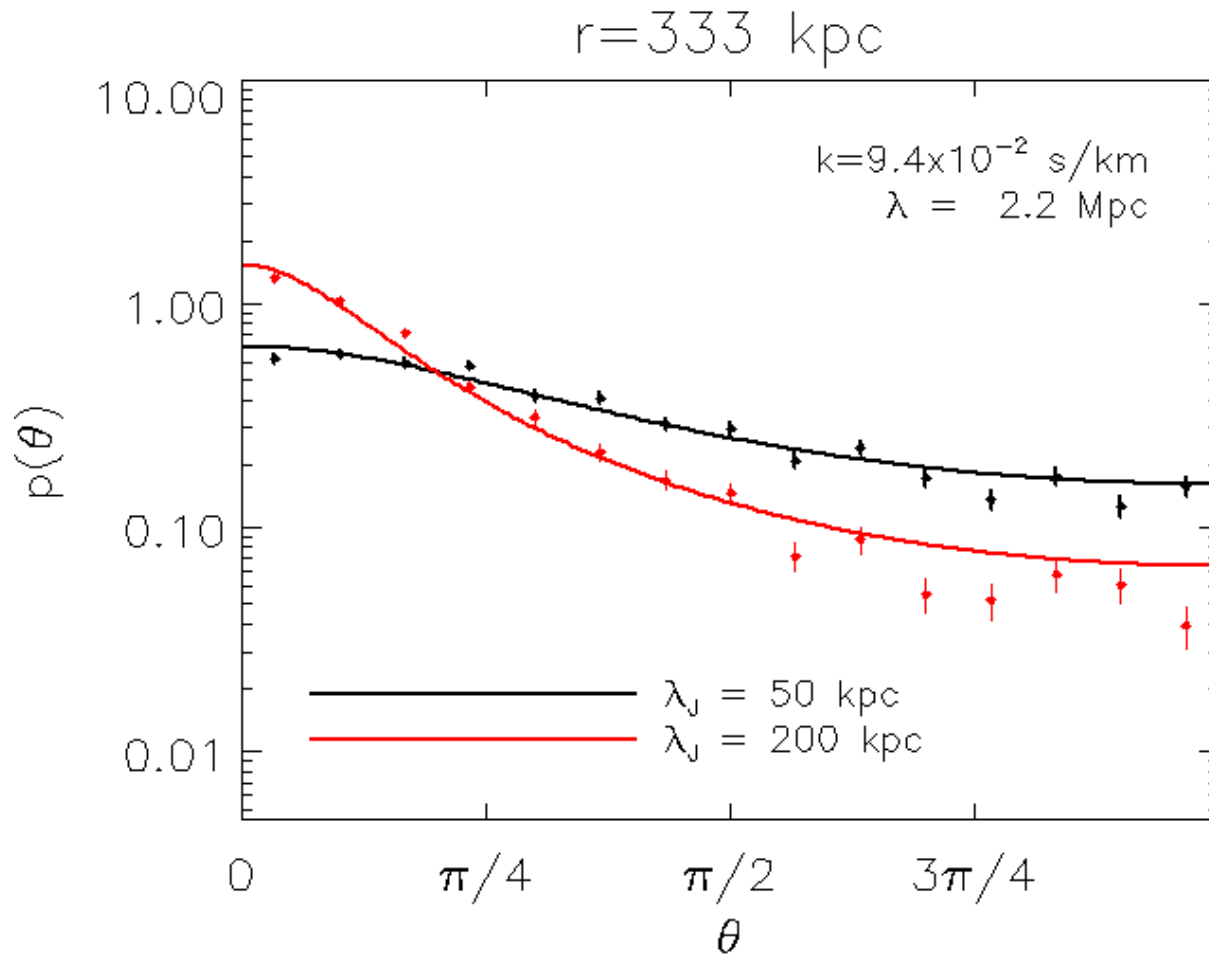
- A filament produce absorption features in two close spectra
- Phase differences are driven by the orientation of these filaments
- If the orientation  $\varphi$  is uniformly distributed, then the phase  $\Theta$  follows a *wrapped Cauchy distribution*

$$P_{WC}(\theta; \zeta) = \frac{1}{2\pi} \frac{1 - \zeta^2}{1 + \zeta^2 - 2\zeta \cos(\theta)}$$

Defines a 1-parameter family of functions

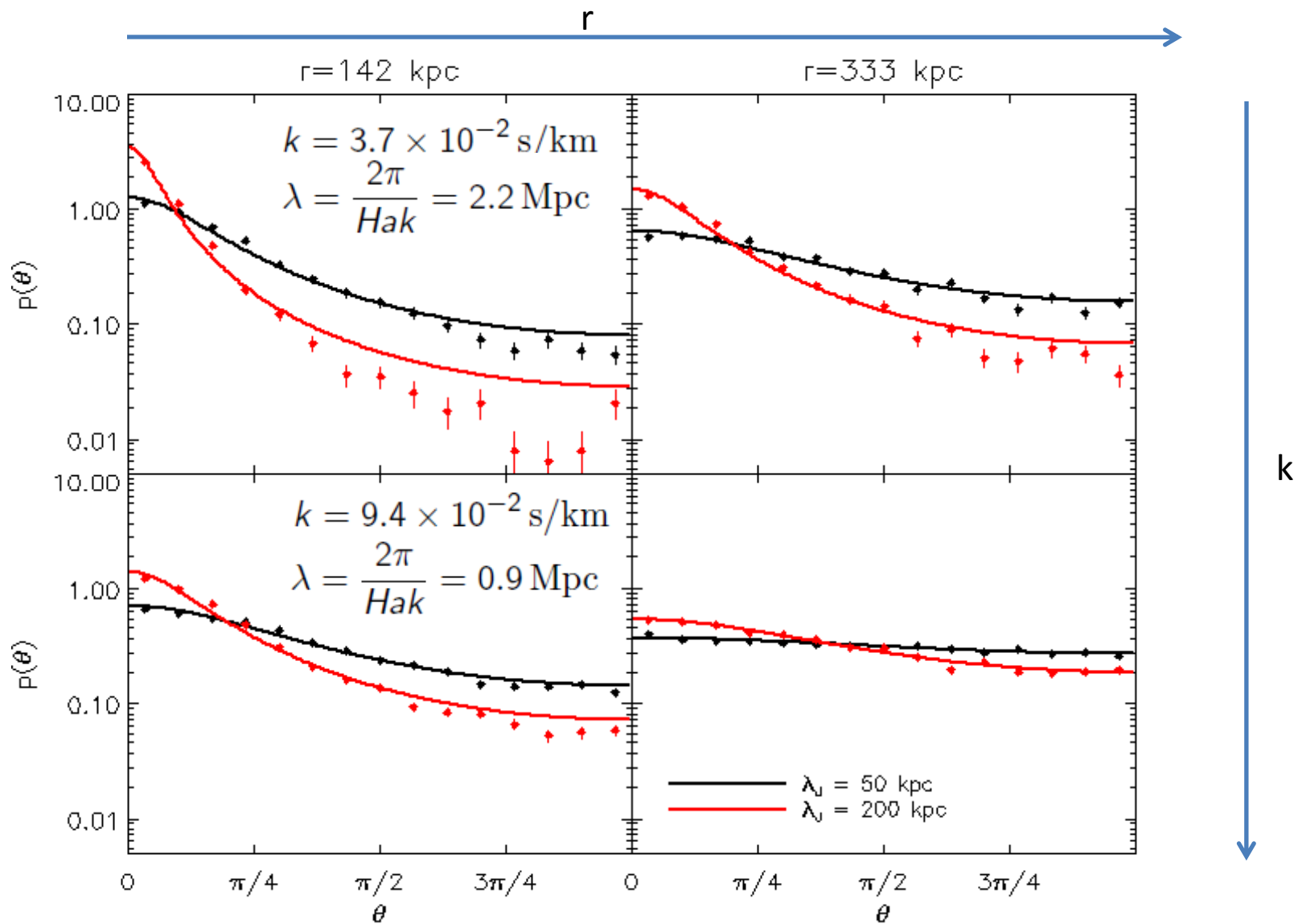


# Probability Distribution of Phases



Wrapped-Cauchy fits well our simulations

# Dependencies of the Phase Distributions



# A Bayesian Likelihood for the Phase PDF

1 - A quasar pair data set provides a set of phase differences  $\{\theta(k_i, r_j)\}$

2 - We predict the probability of measuring those phases for each of the 500 thermal models (*Likelihood function*)

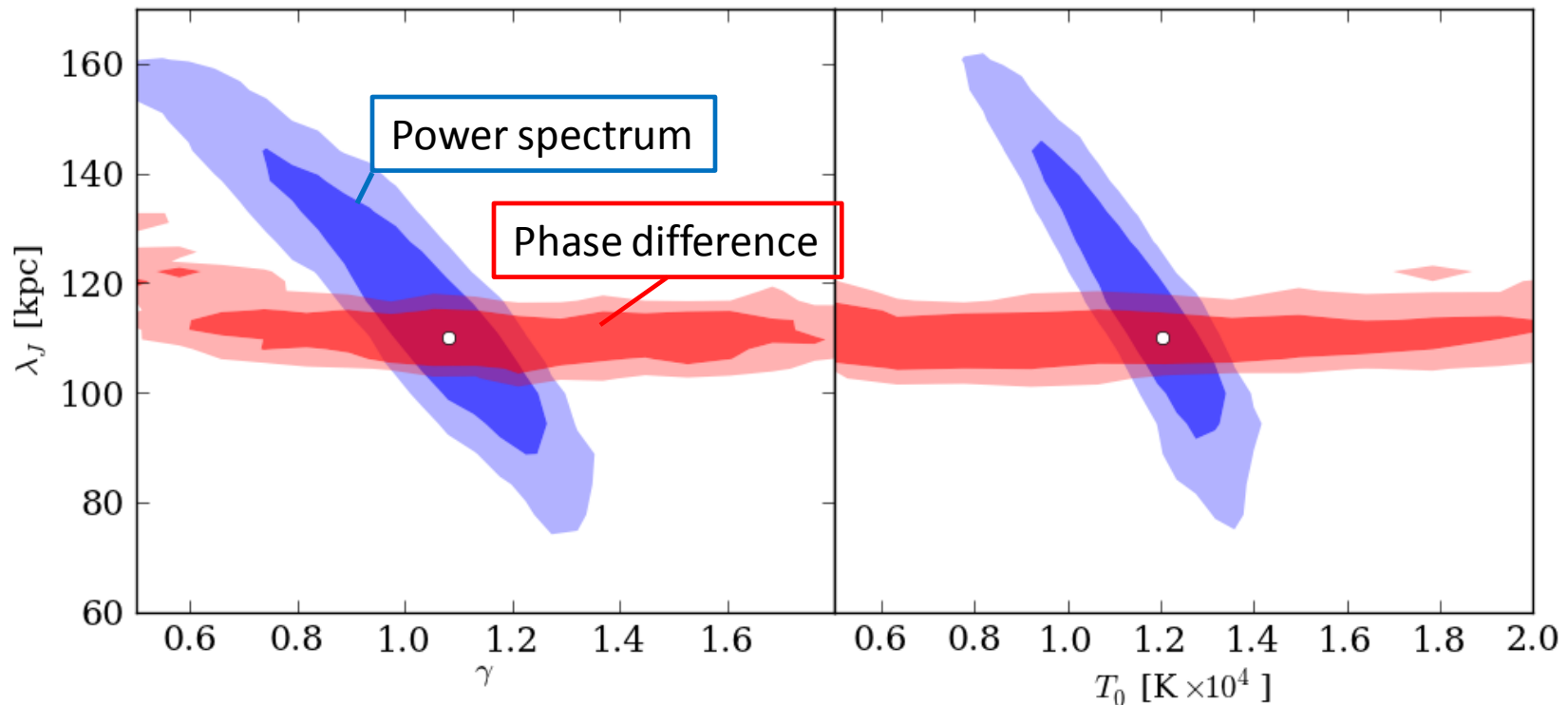
$$L(\{\theta\} | T_0, \gamma, \lambda_J) = \prod_{i,j} P(\theta(k_i, r_j) | T_0, \gamma, \lambda_J)$$

3 - Via MCMC techniques, we do a Bayesian analysis of parameter space, estimating:

- **Degeneracies** between parameters
- **Accuracy** of a measurement

# Estimated Constraints on Thermal Parameters

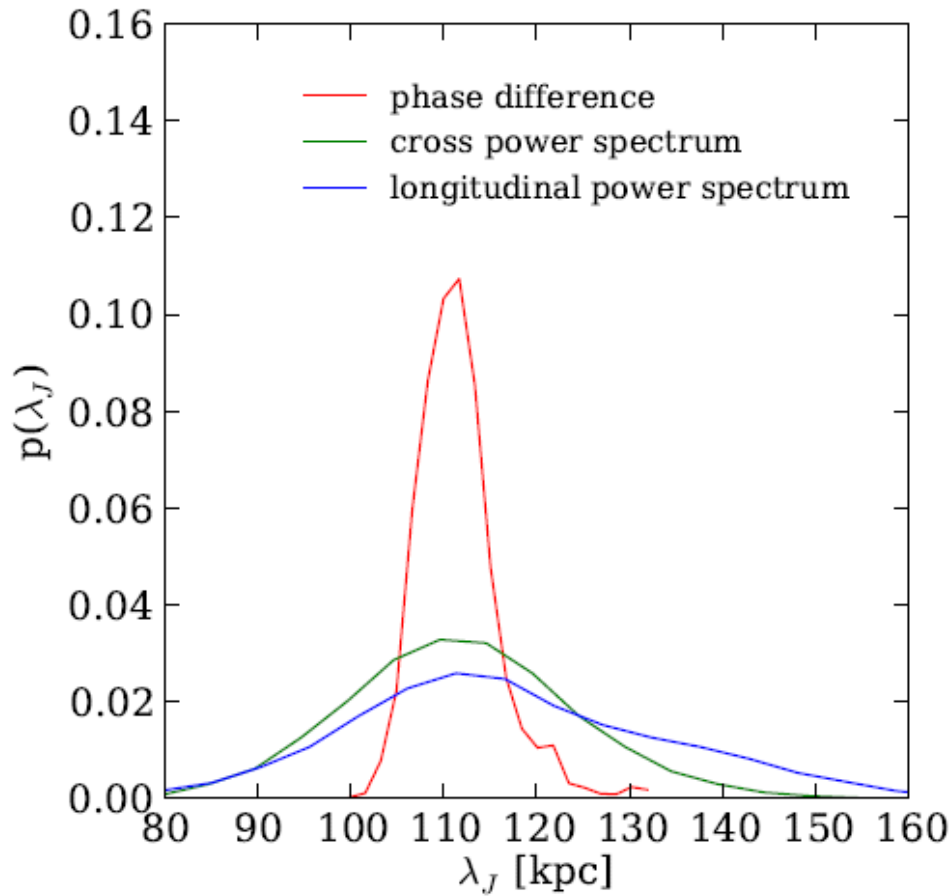
We assume a sample of 20 spectra (LOS power)  
20 pairs (phase difference)



Rorai et al. 2013

Jeans scale measurements are independent on the equation of state

# Accuracy of the Jeans Scale Measurement



We can achieve a precision up to 5% with only 20 pairs

Phase difference statistic is by far more efficient than LOS power and cross-power

Typical precision of IGM temperature measurements is  $\sim 30\%$

Noise decreases the precision by few %

# Nice Properties of Phase Differences

- Invariant under convolution with symmetric Kernel  $W$

$$W(k) = \int W(x) \cos(kx) + i \int \underbrace{W(x)}_{\text{even}} \underbrace{\sin(kx)}_{\text{odd}} = \text{real} \longrightarrow \text{No change in phases}$$

- Thermal Broadening has very small effect
- Resolution does not need to be precisely modeled

- Sensitive to the Jeans scale also at low  $k$

- High-resolution spectra are not required if we resolve the Jeans scale in the transverse dimension

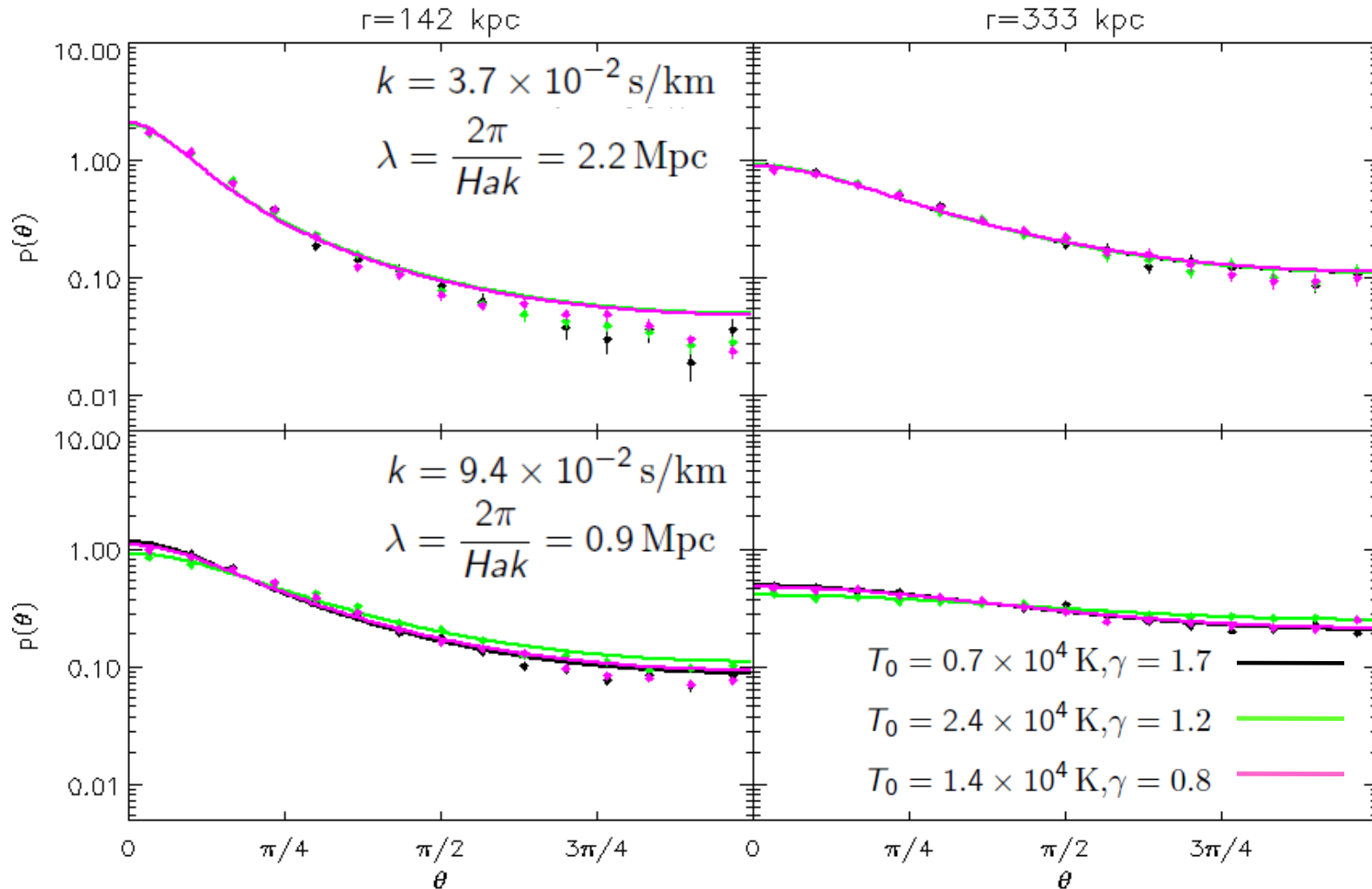
$$P(k) = \int_k^\infty 2\pi k' P_{3D}(k') dk'$$

- Invariant under rescaling of the *normalized* flux

$$\theta_{12}(k) = \arccos \left( \frac{\Re[\tilde{F}_1^*(k) \tilde{F}_2(k)]}{\sqrt{|\tilde{F}_1(k)|^2 |\tilde{F}_2(k)|^2}} \right)$$

- Robustness to uncertainties in continuum fitting

# Sensitivity to the Equation-of-State Parameters



Phases are almost independent of the equation of state

# Summary

- The Jeans scale is a key quantity for galaxy formation, reionization and thermal history
- Close QSO pairs will provide a first accurate measurement of the Jeans scale
- The phase analysis is insensitive to the equation of state parameters  $T_0$  and  $\Upsilon$
- Combining this with others Ly- $\alpha$  forest statistic can break degeneracies and give tight constraints on the thermal history of the IGM