



Automating BSM for Neutrino Generators

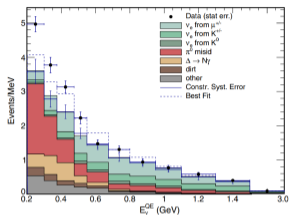
Joshua Isaacson

In Collaboration with with Stefan H"ochel, Diego Lopez Gutierrez, and Noemi Rocco

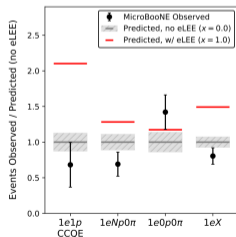
Based on arxiv:2110.15319

15 December 2021

Motivation: MiniBooNE and MicroBooNE

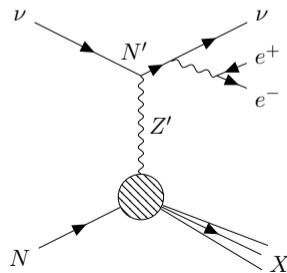


[PRL 121, 221801]



[arXiv:2110.14054]

- MiniBooNE sees excess of events
- MicroBooNE does not see excess of single electron events
- Excess can be from multiple lepton final states
- Event generators can not simulate these processes



Motivation: Theory

Summary of Workshop on Common Neutrino Event Generator Tools

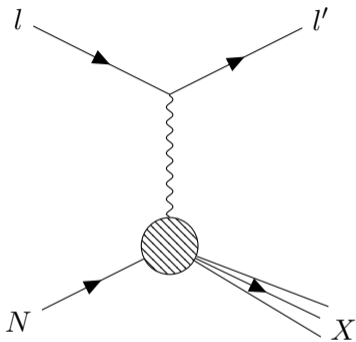
2.5.2 Factors to consider for interface design

While our goal should be to standardize the interface as much as possible, the workshop raised several issues that should be considered when evaluating possible approaches. These issues are summarized below:

4. **Human factors** Our interface should be designed for ease of use, and should consider the skills and limitations of the theorists likely to use it. It was pointed out that many theorists are PhD students or postdoctoral researchers working on limited-term contracts. To fit in with this way of working, **it should be possible to develop, implement, and test a model against data on timescales of the order of a year**. We must also bear in mind that many theorists are not primarily programmers, and **that models may be developed using tools, such as Mathematica, that are not natively compatible with the languages used in generator software**. **If we restricted ourselves to an interface in a particular programming language, we could severely limit the accessibility of the interface to new models.**

[<https://arxiv.org/abs/2008.06566>]

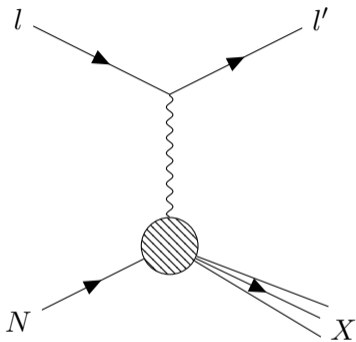
Leptonic and Hadronic Tensor



Observations:

- Nuclear physics calculations are hard
- Calculating arbitrary perturbative diagrams is a solved problem
- BSM effects of interest only weakly couple to quarks and gluons \Rightarrow no corrections to the nuclear physics

Leptonic and Hadronic Tensor



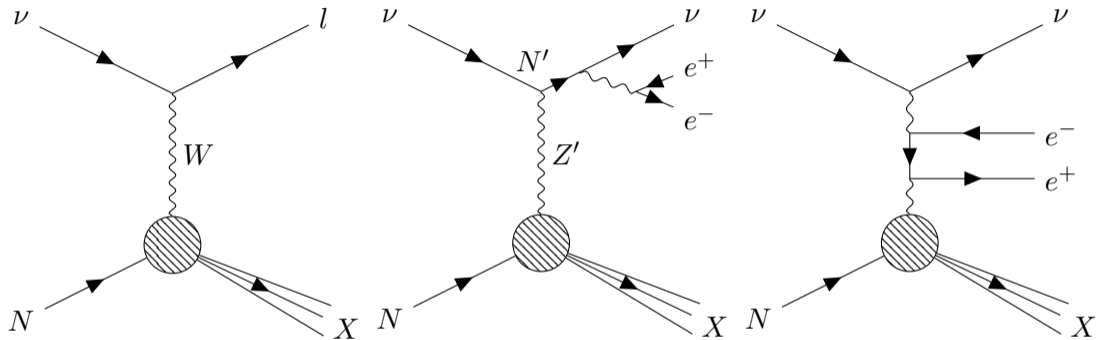
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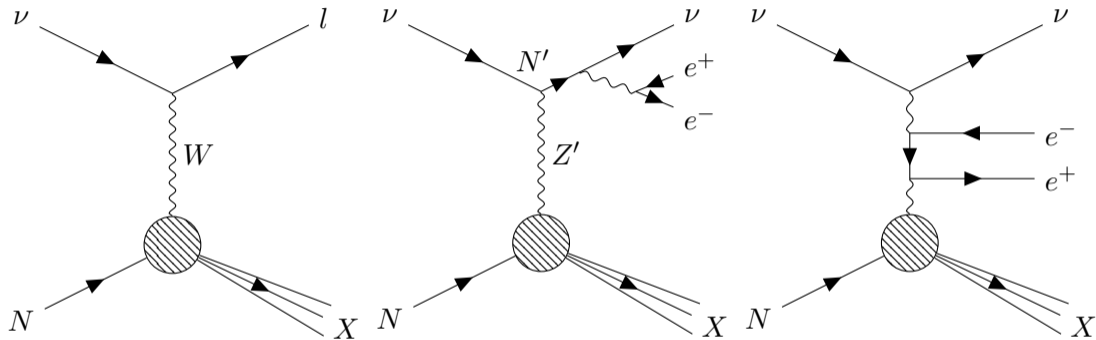
Question

Can we use these observations to automate Beyond the Standard Model Physics?

Leptonic and Hadronic Tensor



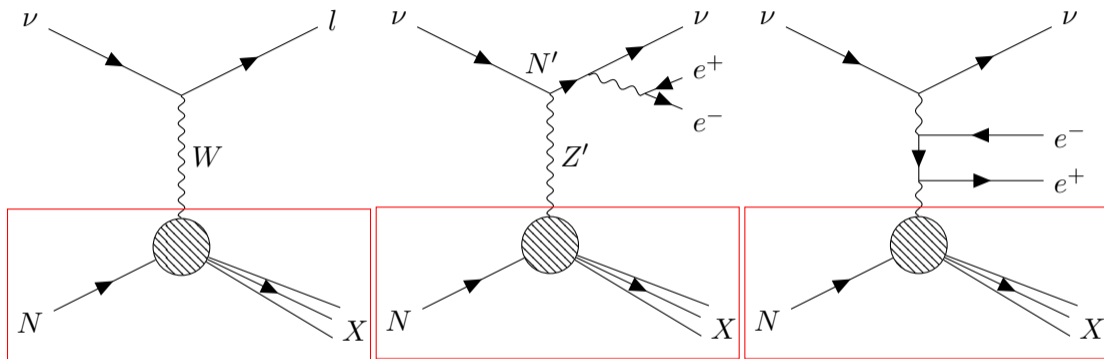
Leptonic and Hadronic Tensor



Question

What do all the diagrams above have in common?

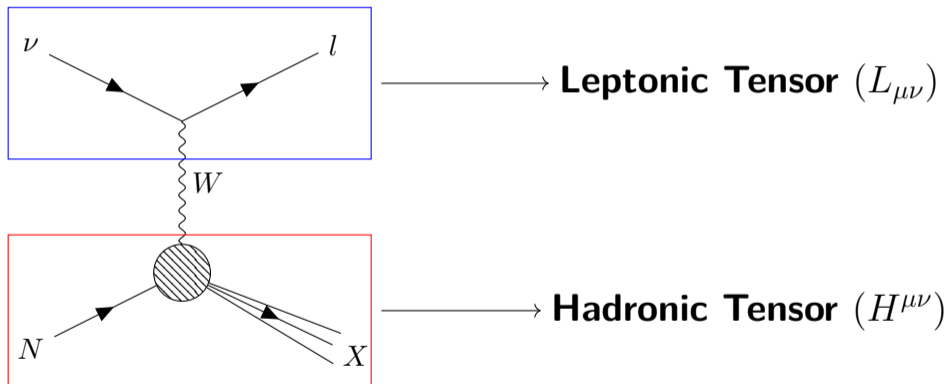
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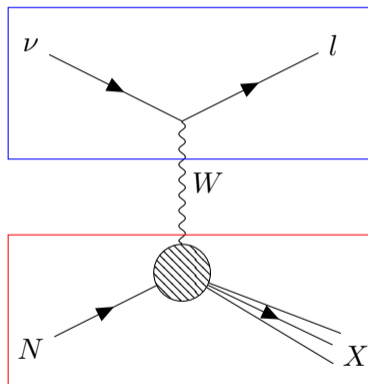
Question

What do all the diagrams above have in common?

Leptonic and Hadronic Tensor



Leptonic and Hadronic Tensor



Notes:

- Leptonic tensor only contains perturbative physics.
- Can use LHC tools to calculate Leptonic tensor
- Hadronic tensor is difficult, but event generators have these calculations implemented already.

Using Currents

Using tensors:

$$\frac{d\sigma}{d\Omega} = \sum_{i,j} L_{\mu\nu}^{(ij)} W^{(ij)\mu\nu} = L_{\mu\nu}^{(\gamma\gamma)} W^{(\gamma\gamma)\mu\nu} + L_{\mu\nu}^{(\gamma Z)} W^{(\gamma Z)\mu\nu} + L_{\mu\nu}^{(Z\gamma)} W^{(Z\gamma)\mu\nu} + L_{\mu\nu}^{(ZZ)} W^{(ZZ)\mu\nu} + \dots$$

Using Currents:

$$\frac{d\sigma}{d\Omega} = \left| \sum_i L_\mu^{(i)} W^{(i)\mu} \right|^2$$

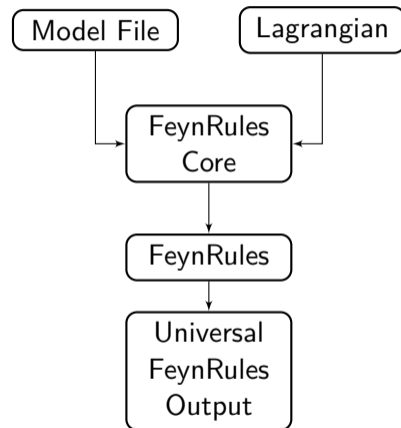
Interferences handled automatically using currents

Interface to tensors provided for nuclear calculations that **must** be expressed using tensors.

FeynRules

- *Mathematica* Program
- Takes model file and Lagrangian as input
- Calculates the Feynman rules
- Outputs in Universal FeynRules Output (UFO) format

[[arXiv:0806.4194](https://arxiv.org/abs/0806.4194), [arXiv:1310.1921](https://arxiv.org/abs/1310.1921)]



Universal FeynRules Output (UFO)

- Python output files
- Contains model-independent files and model-dependent files
- Contains all information to calculate any tree level matrix element
- Has parameter file to adjust model parameters to scan allowed regions

Example QED ($e^+e^-\gamma$ Vertex):

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi} (iD^\mu\gamma_\mu - m) \psi$$

$$V_{e^+e^-\gamma} = ie\gamma^\mu = \gamma \text{ ~~~~~ } \begin{array}{l} \nearrow \\ \searrow \end{array}$$

[arXiv:1108.2040]

Handling Form Factors

Nuclear one-body current operators:

$$\mathcal{J}^\mu = (\mathcal{J}_V^\mu + \mathcal{J}_A^\mu)$$

$$\mathcal{J}_V^\mu = \gamma^\mu \mathcal{F}_1^a + i\sigma^{\mu\nu} q_\nu \frac{\mathcal{F}_2^a}{2M}$$

$$\mathcal{J}_A^\mu = -\gamma^\mu \gamma_5 \mathcal{F}_A^a - q^\mu \gamma_5 \frac{\mathcal{F}_P^a}{M}$$

Standard Model Form Factors:

$$\mathcal{F}_i^{\gamma(p,n)} = F_i^{p,n}, \quad \mathcal{F}_A^\gamma = 0$$

$$\mathcal{F}_i^{W(p,n)} = F_i^p - F_i^n, \quad \mathcal{F}_A^W = F_A$$

$$\mathcal{F}_i^{Z(p)} = \left(\frac{1}{2} - 2 \sin^2 \theta_W \right) F_i^p - \frac{1}{2} F_i^n,$$

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Straight-forward to extend to BSM if CVC is valid

Recursive Matrix Element Generation

$$\mathcal{J}_\alpha(\pi) = P_\alpha(\pi) \sum_{\nu_\alpha^{\alpha_1, \alpha_2}} \sum_{\mathcal{P}_2(\pi)} \mathcal{S}(\pi_1, \pi_2) V_\alpha^{\alpha_1, \alpha_2}(\pi_1, \pi_2) \mathcal{J}_{\alpha_1}(\pi_1) \mathcal{J}_{\alpha_2}(\pi_2)$$

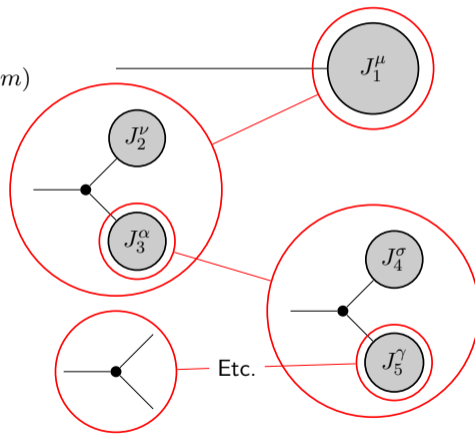
$$L_\mu^{(i)}(1, \dots, m) = \mathcal{J}_\mu^{(i)}(1, \dots, m)$$

$$L_{\mu\nu}^{(i,j)}(1, \dots, m) = \mathcal{J}_\mu^{(i)}(1, \dots, m) \mathcal{J}_\nu^{(j)\dagger}(1, \dots, m)$$

Berends-Giele Recursion

- Reuse parts of calculation
- Most efficient for high multiplicity
- Reduces computation from $\mathcal{O}(n!)$ to $\mathcal{O}(3^n)$

[Nucl. Phys. B306(1988), 759]



Phase Space Generation

$$d\Phi_n(a, b; 1, \dots, n) = \delta^{(4)}\left(p_a + p_b - \sum_{i=1}^n p_i\right) \left[\prod_{i=1}^n \frac{d^4 p_i}{(2\pi)^3} \delta(p_i^2 - m_i^2) \Theta(p_{i0}) \right]$$

The above phase space definition does not contain the handling of initial states.

Algorithms for n -body phase space generation

- RAMBO [[Comput. Phys. Commun. 40\(1986\) 359](#)]
- Multi-channel techniques [[hep-ph/9405257](#)]
- Recursive Phase Space [[arXiv:0808.3674](#)]

Results

Processes Considered:

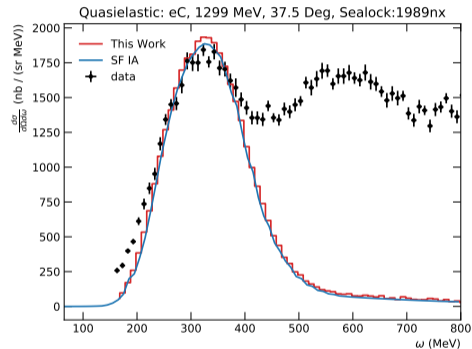
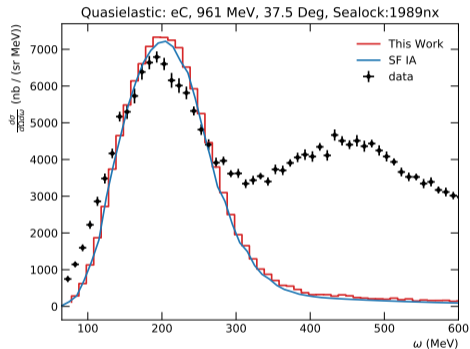
- Electron-Carbon Scattering
- Neutrino-Carbon Scattering
- Neutrino Tridents

NOTE: All processes are fully differential

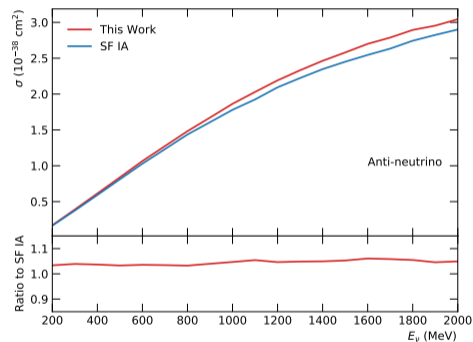
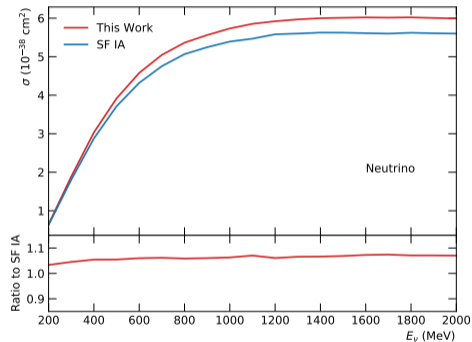
Parameters:

- Only quasielastic scattering is included and no FSI
- EM Form Factors: Kelly parameterization [[PRC 70, 068202 \(2004\)](#)]
- Axial Form Factor:
 - Dipole
 - $M_A = 1.0$ GeV
 - $g_A = 1.2694$
- $\alpha = 1/137$
- $G_F = 1.16637 \times 10^{-5}$
- $M_Z = 91.1876$ GeV

Electron Scattering

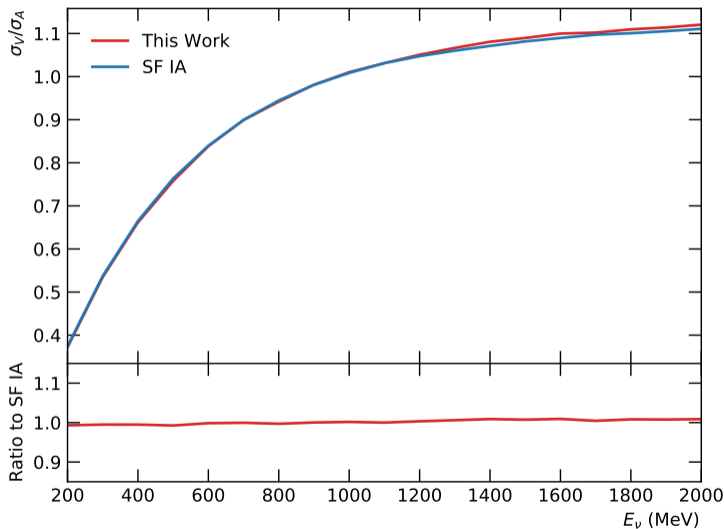


Neutrino Total Cross Section

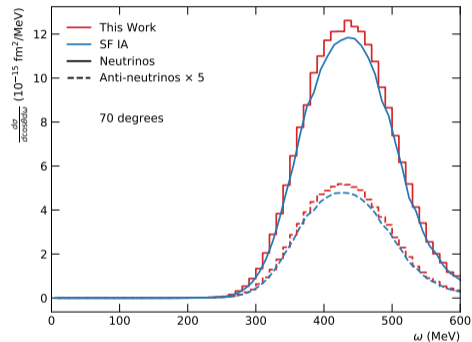
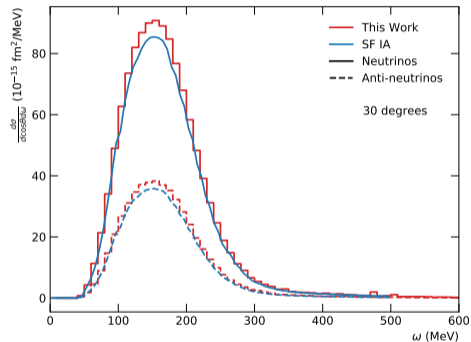


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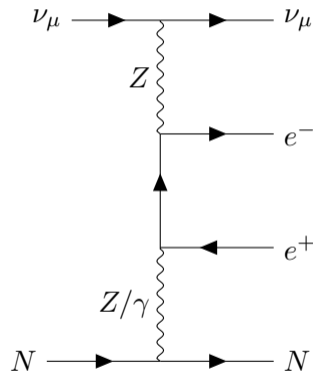
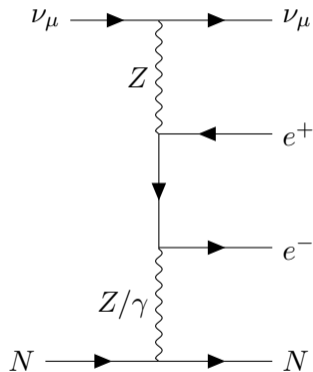
- Difference due to how couplings are handled
- Ratio of $\frac{\sigma_V}{\sigma_A}$ consistent



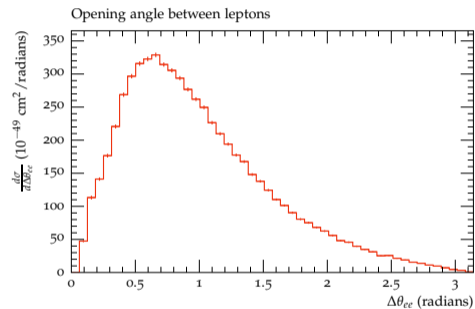
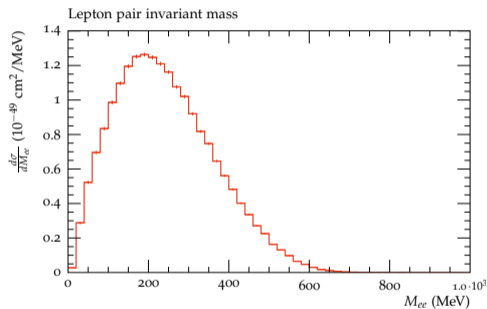
Neutrino Differential Cross Section



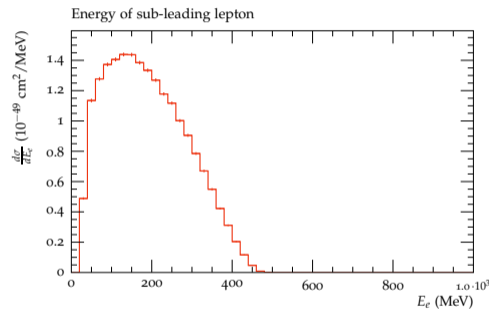
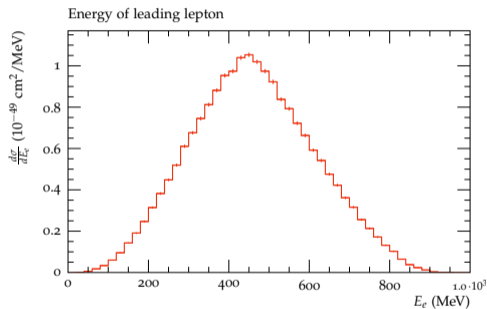
Neutrino Tridents



Neutrino Tridents



Neutrino Tridents



Conclusions

- BSM important for the current and next generation neutrino experiments
- Robust BSM program requires automating theory calculations
- Developed method for handling arbitrary form factors
- Proof of principle for SM processes

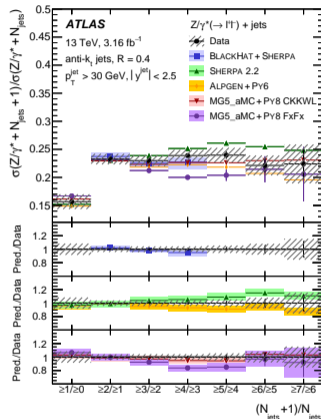
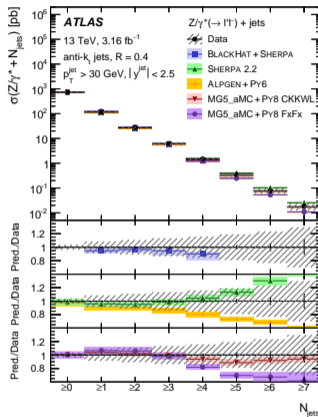
Universal FeynRules Output (UFO)

Example for photon-electron vertex

```
e__minus__ = Particle(pdg_code=11, name='e-', antiname='e+',
                    spin=2, color=1, mass=Param.ZERO,
                    width=Param.ZERO, texname='e-',
                    antitexname='e+', charge=-1,
                    GhostNumber=0, LeptonNumber=1,
                    Y=0)
V_77 = Vertex(name='V_77',
             particles=[ P.e__plus__, P.e__minus__, P.a ],
             color=[ '1' ], lorentz=[ L.FFV1 ],
             couplings={(0,0):C.GC_3})
FFV1 = Lorentz(name='FFV1', spins=[ 2, 2, 3 ],
              structure='Gamma(3,2,1)')
GC_3 = Coupling(name='GC_3', value='-(ee*complex(0,1))',
               order={'QED':1})
```

Tree Level Matrix Element Generators

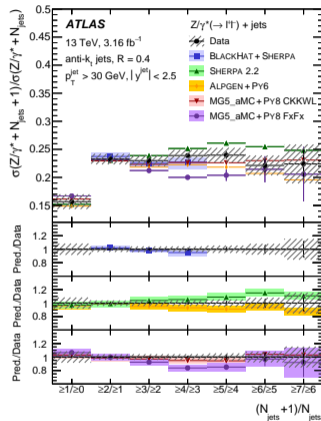
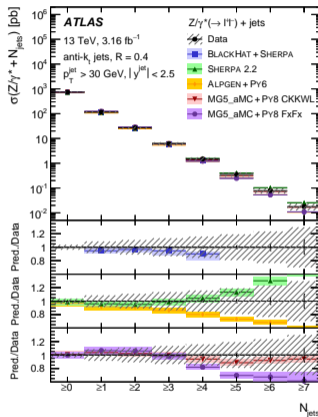
- ALPGEN [arXiv:hep-ph/0206293]
- AMEGIC [arXiv:hep-ph/0109036]
- COMIX [arXiv:0808.3674]
- CALCHEP [arXiv:1207.6082]
- HERWIG [arXiv:0803.0883]
- MADGRAPH [arXiv:1405.0301]
- WHIZARD [arXiv:0708.4233]
- etc.



[arXiv:1702.05725]

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2 \rightarrow 2 Phase Space Example

Consider $l + {}^{12}\text{C} \rightarrow l' + N + X$ in the quasielastic regime.

$$d\sigma \propto d\Phi_2(a, b; 1, 2) d^4p_a d^3p_b$$

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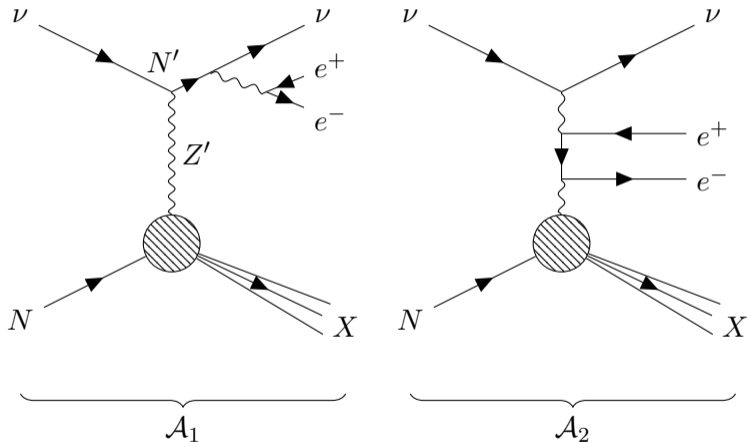
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Quasielastic Delta Function: $\delta(E_b - E_1 - E_r + m - E_2)$

Phase Space Delta Function: $\delta(E_a + E_b - E_1 - E_2)$

Define initial nucleon energy as $E_a = m - E_r$. Allows use of phase space tools developed at LHC.

Multi-channel Integration



- Both diagrams contribute to cross section
- They have different pole structures
- Need method to sample these structures efficiently (i.e. $|\mathcal{A}_1 + \mathcal{A}_2|^2$)

Multi-channel Integration and VEGAS

Multi-channel Integration

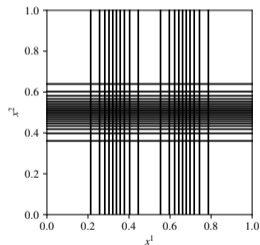
- Generate PS efficiently for $|\mathcal{A}_1|^2$ or $|\mathcal{A}_2|^2$
- Do not know how to efficiently sample $2\text{Re}(\mathcal{A}_1\mathcal{A}_2^\dagger)$
- Define channels: C_1 and C_2
- Generate events according to distributions g_i for channel i

$$\int d\vec{x} f(\vec{x}) = \sum_i \alpha_i \int d\vec{x} g_i(\vec{x}) \frac{f(\vec{x})}{g_i(\vec{x})}$$

- Optimize α_i to minimize variance

VEGAS

- Adaptive importance sampling
- Use this to get interference terms more accurately



VEGAS grid for $\int_0^1 d^4x \left(e^{-100(\vec{x}-\vec{r}_1)^2} + e^{-100(\vec{x}-\vec{r}_2)^2} \right)$

[J.Comput.Phys. 27 (1978) 291, 2009.05112]

Recursive Phase Space Decomposition

Phase space can be decomposed as:

$$d\Phi_n(a, b; 1, \dots, n) = d\Phi_{n-m+1}(a, b; m+1, \dots, n) \frac{ds_\pi}{2\pi} d\Phi_m(\pi; 1, \dots, m)$$

Iterate until only $1 \rightarrow 2$ phase spaces remain.

Basic building blocks:

$$S_\pi^{\rho, \pi \setminus \rho} = \frac{\lambda(s_\pi, s_\rho, s_{\pi \setminus \rho})}{16\pi^2 2s_\pi} d\cos\theta_\rho d\phi_\rho$$
$$T_{\alpha, b}^{\pi, \overline{\alpha b \pi}} = \frac{\lambda(s_{\alpha b}, s_\pi, s_{\overline{\alpha b \pi}})}{16\pi^2 2s_{\alpha b}} d\cos\theta_\pi d\phi_\pi$$

Momentum conservation: $(2\pi)^4 d^4 p_{\overline{\alpha b}} \delta^{(4)}(p_\alpha + p_b - p_{\overline{\alpha b}})$

Proposed Interface

