

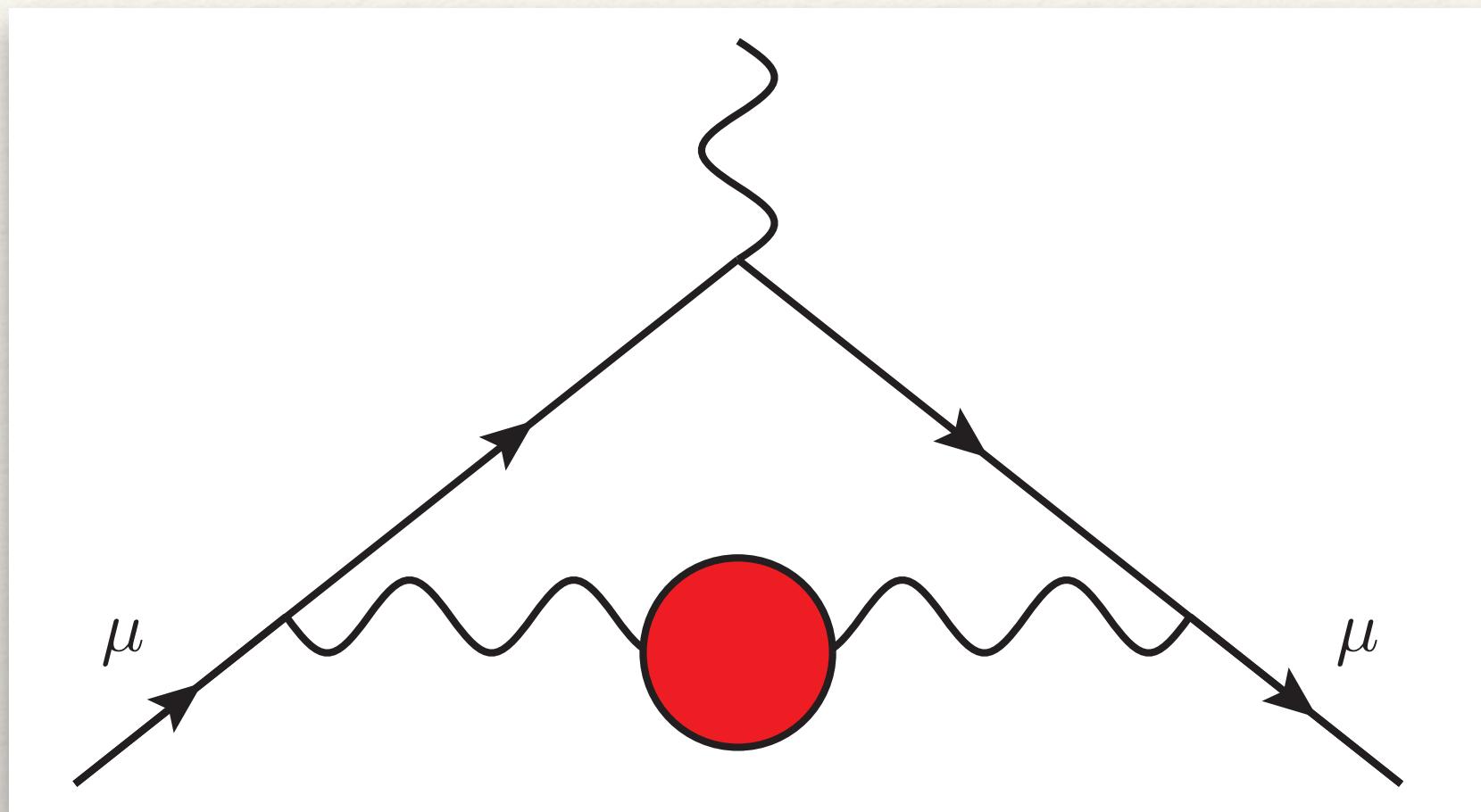
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NLO and NNLO HVP contributions to the muon $g - 2$ in the space-like region

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Work in Collaboration with S. Laporta and M. Passera, arXiv: 2112.05704

The hadronic LO contribution: time-like method



$$a_\mu^{HVP}(LO) = \frac{\alpha}{\pi^2} \int_{s_0^2}^{\infty} \frac{ds}{s} K^{(2)}(s/m^2) \text{Im}\Pi_h(s)$$

$$\text{Im}\Pi_h(s) = \frac{\alpha}{3} R(s) \quad R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{4\pi\alpha^2/(3s)}$$

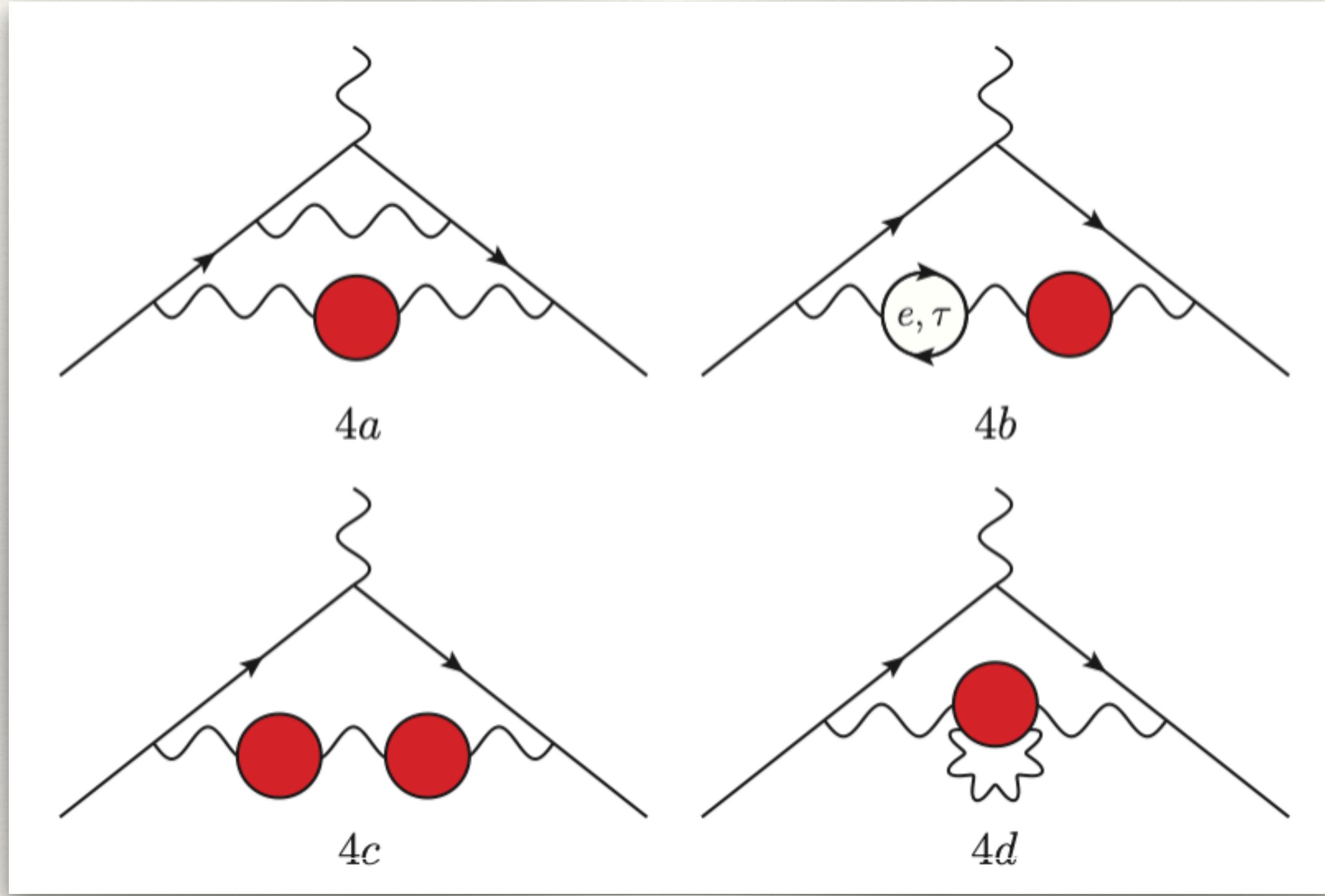
$$K^{(2)}(s/m^2) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)}$$

Defining $z = \frac{q^2}{m^2}$ & $y(z) = \frac{z - \sqrt{z(z-4)}}{z + \sqrt{z(z-4)}}$:

$$K^{(2)}(z) = \frac{1}{2} - z + \left(\frac{z^2}{2} - z \right) \ln z + \frac{\ln y(z)}{\sqrt{z(z-4)}} \left(z - 2z^2 + \frac{z^3}{2} \right)$$

The hadronic NLO VP contribution

- $O(\alpha^3)$ contributions containing HVP insertion

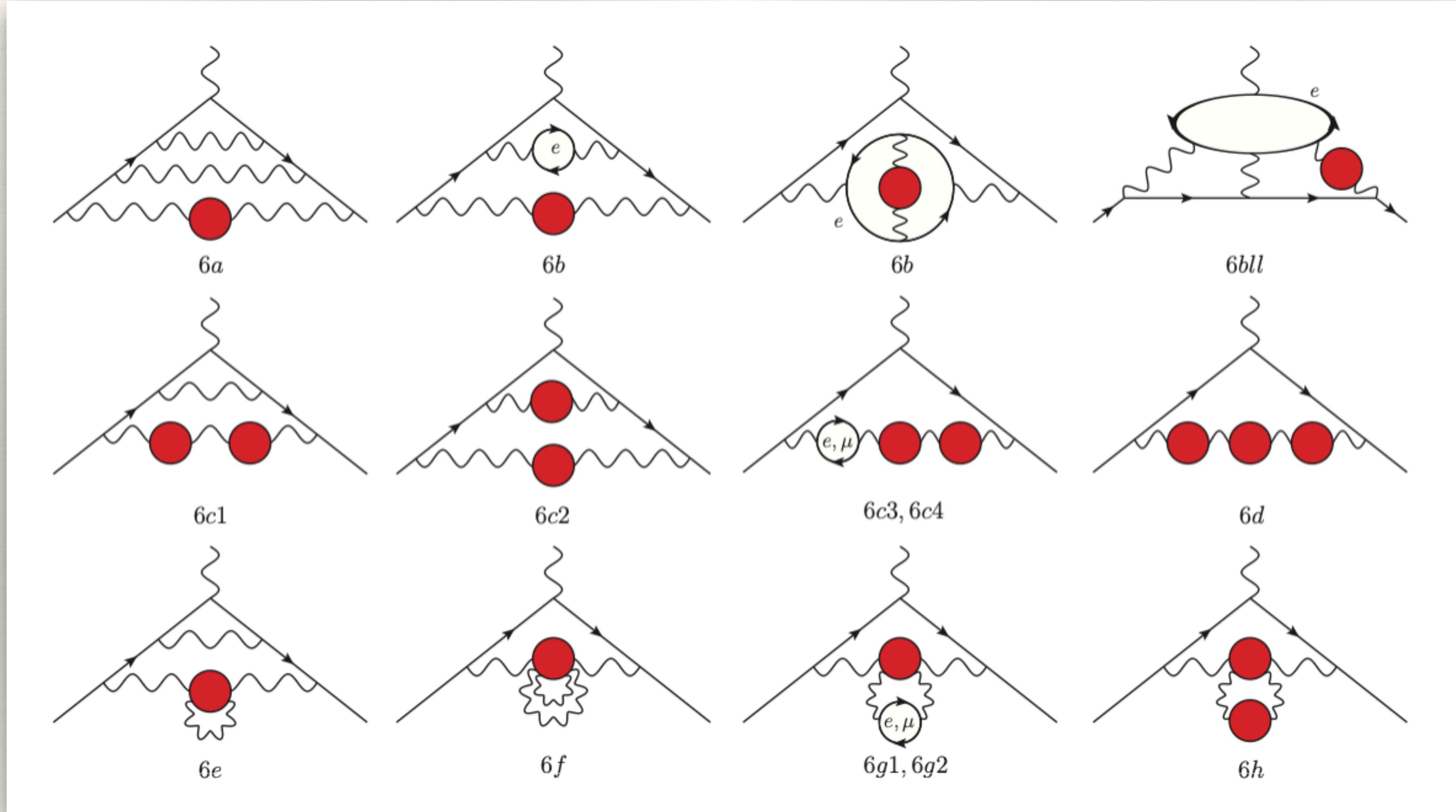


$$a_\mu^{HVP}(NLO) = -98.3(7) \times 10^{-11}$$

Krause '96; Keshevarzi, Nomura, Teubner 2019; WP20

The hadronic NNLO VP contribution

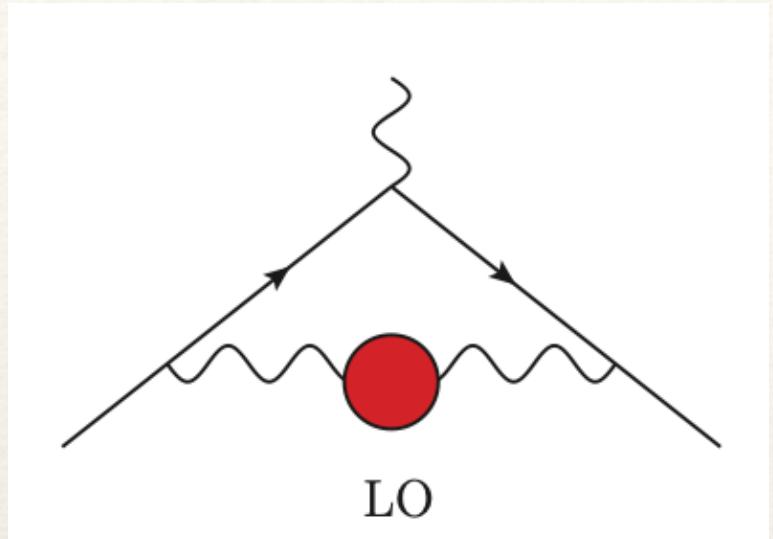
- $O(\alpha^4)$ contributions containing HVP insertion



$$a_\mu^{HVP}(NNLO) = 12.4(1) \times 10^{-11}$$

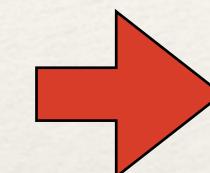
Kurz, Liu, Marquard, Steinhauser 2014

Space-like method: LO HVP contribution



$$K^{(2)}(z) = \frac{1}{\pi} \int_{-\infty}^0 dz' \frac{\text{Im}K^{(2)}(z')}{z' - z}, \quad z > 0$$

$$\frac{\Pi(q^2)}{q^2} = \frac{1}{\pi} \int \frac{ds}{s} \frac{\text{Im}\Pi(s)}{s - q^2}, \quad q^2 < 0$$



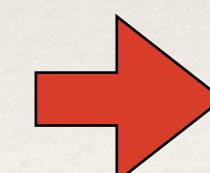
$$a_\mu^{HVP}(LO) = -\frac{\alpha}{\pi^2} \int_{-\infty}^0 \frac{dt}{t} \Pi_h(t) \text{Im}K^{(2)}(t/m^2)$$

With the imaginary part:

$$\text{Im}K^{(2)}(z + i\epsilon) = \pi\theta(-z) \left[\frac{z^2}{2} - z + \frac{z - 2z^2 + z^3/2}{\sqrt{z(z-4)}} \right] = \pi\theta(-z) F^{(2)}(1/y(z)), \quad F^{(2)}(u) = \frac{u+1}{u-1} u^2$$

Changing the variable in the dispersive integral $t \rightarrow y \rightarrow x$: $t(x) = \frac{m^2 x^2}{1 - x^2}$, $x = 1 + y$

$$\Delta\alpha_h(t) = -\Pi_h(t)$$

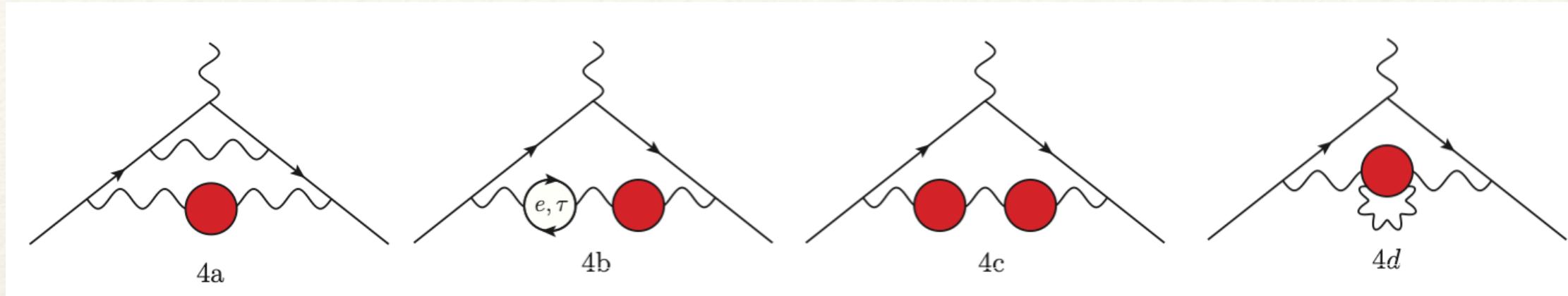


$$a_\mu^{HVP}(LO) = \frac{\alpha}{\pi} \int_0^1 dx \kappa^{(2)}(x) \Delta\alpha_h(t(x))$$

$$\kappa^{(2)}(x) = 1 - x$$

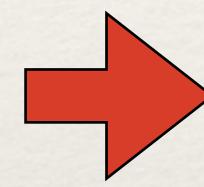
Lautrup, Peterman, de Rafael 1972

Space-like method: NLO HVP contribution



$$a_\mu^{(4a)} = \frac{\alpha^2}{\pi^3} \int_{s_0}^\infty \frac{ds}{s} 2K^{(4)}(s/m^2) \text{Im}\Pi_h(s)$$

$K^{(4)}$ from Barbieri Remiddi 1975



$$a_\mu^{(4a)} = -\frac{\alpha^2}{\pi^3} \int_{-\infty}^0 \frac{dt}{t} \Pi_h(t) \text{Im}K^{(4)}(t/m^2)$$

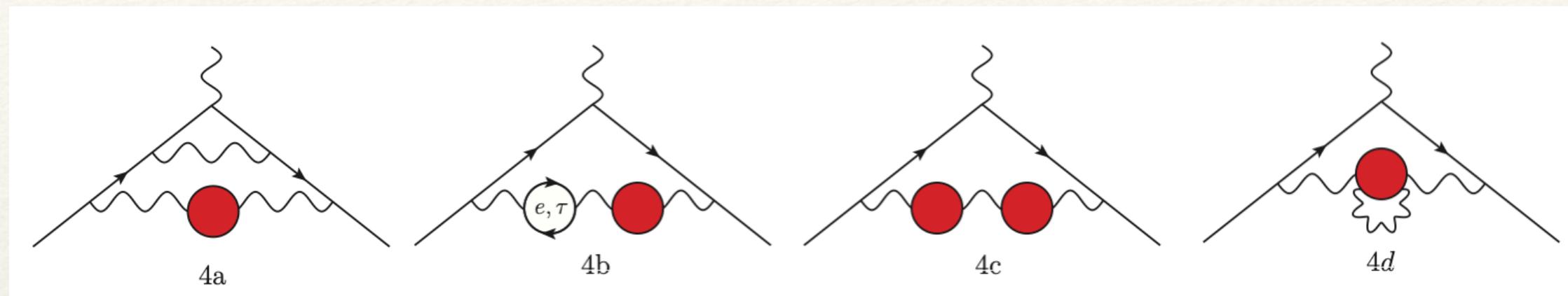
$$\begin{aligned} F^{(4)}(u) &= R_1(u) + R_2(u) \ln(-u) \\ &\quad + R_3(u) \ln(1+u) + R_4(u) \ln(1-u) \\ &\quad + R_5(u) [4 \operatorname{Li}_2(u) + 2 \operatorname{Li}_2(-u) \\ &\quad + \ln(-u) \ln((1-u)^2(1+u))] \end{aligned}$$

$$\text{Im}K^{(4)}(z + i\epsilon) = \pi \theta(-z) F^{(4)}(1/y(z))$$

$$\begin{aligned} R_1 &= \frac{23u^6 - 37u^5 + 124u^4 - 86u^3 - 57u^2 + 99u + 78}{72(u-1)^2u(u+1)}, \\ R_2 &= \frac{12u^8 - 11u^7 - 78u^6 + 21u^5 + 4u^4 - 15u^3 + 13u + 6}{12(u-1)^3u(u+1)^2}, \\ R_3 &= \frac{(u+1)(-u^3 + 7u^2 + 8u + 6)}{12u^2}, \\ R_4 &= \frac{-7u^4 - 8u^3 + 8u + 7}{12u^2}, \\ R_5 &= -\frac{3u^4 + 5u^3 + 7u^2 + 5u + 3}{6u^2}. \end{aligned}$$

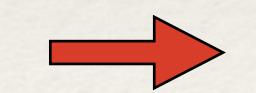
Obtained independently by Nesterenko, arXiv: 2112.05009

Space-like method: NLO hadronic vacuum polarization contribution



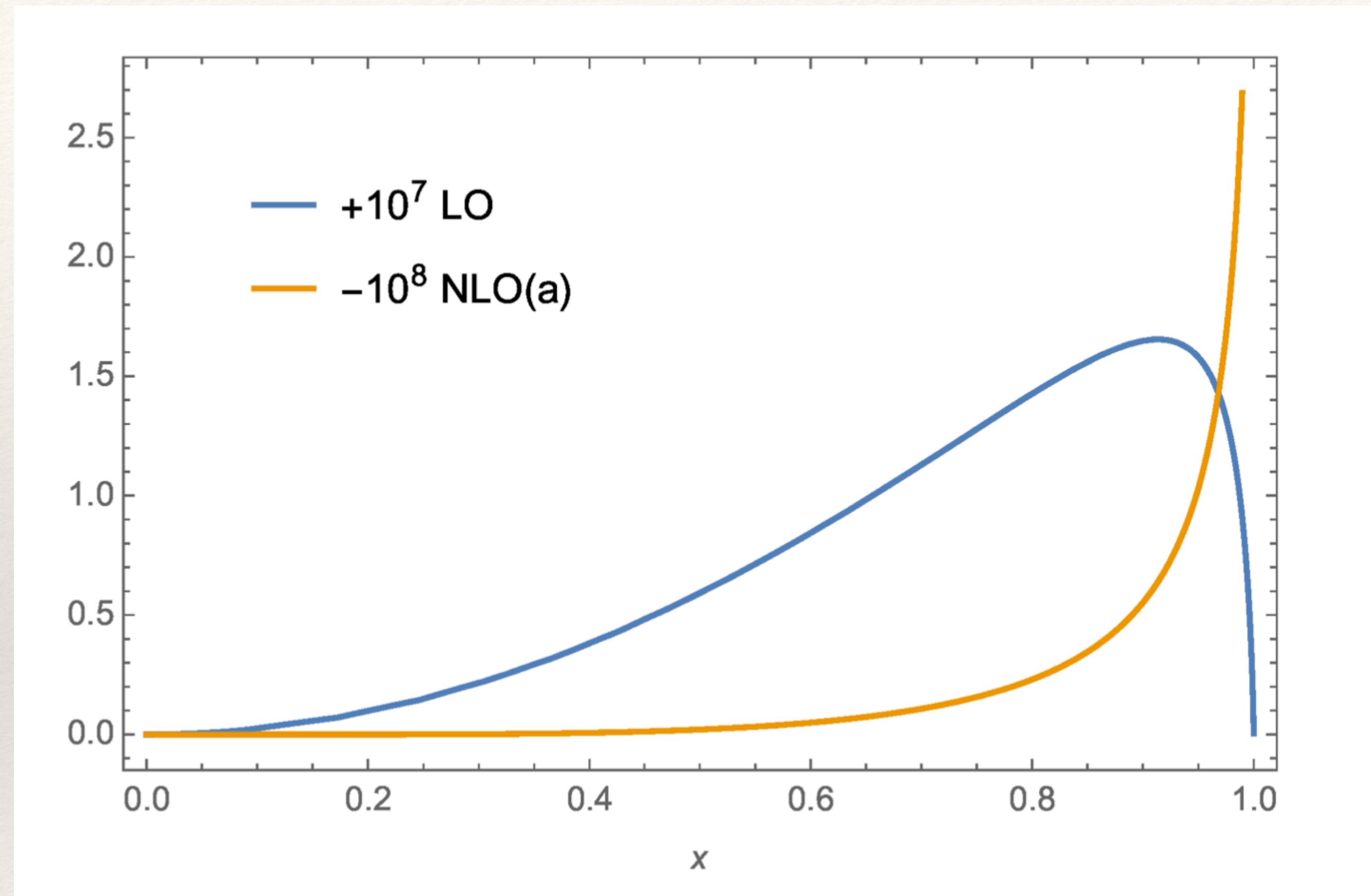
$$\begin{aligned} \rightarrow a_{\mu}^{(4a)} &= \left(\frac{\alpha}{\pi}\right)^2 \int_0^1 dx \kappa^{(4)}(x) \Delta \alpha_h(t(x)), & \kappa^{(4)} &= \frac{2(2-x)}{x(x-1)} F^{(4)}(x-1) \\ \rightarrow a_{\mu}^{(4b)} &= \left(\frac{\alpha}{\pi}\right) \int_0^1 dx \kappa^{(2)}(x) \Delta \alpha_h(t(x)) \times 2 [\Delta \alpha_e^{(2)}(t(x)) + \Delta \alpha_{\tau}^{(2)}(t(x))] \\ \rightarrow a_{\mu}^{(4c)} &= \left(\frac{\alpha}{\pi}\right) \int_0^1 dx \kappa^{(2)}(x) [\Delta \alpha_h(t(x))]^2 \end{aligned}$$

Chakraborty et al. (arXiv: 1806.08190) provided an approximated expression for the space-like kernel function. They added a $O(10\%)$ uncertainty to their final result.



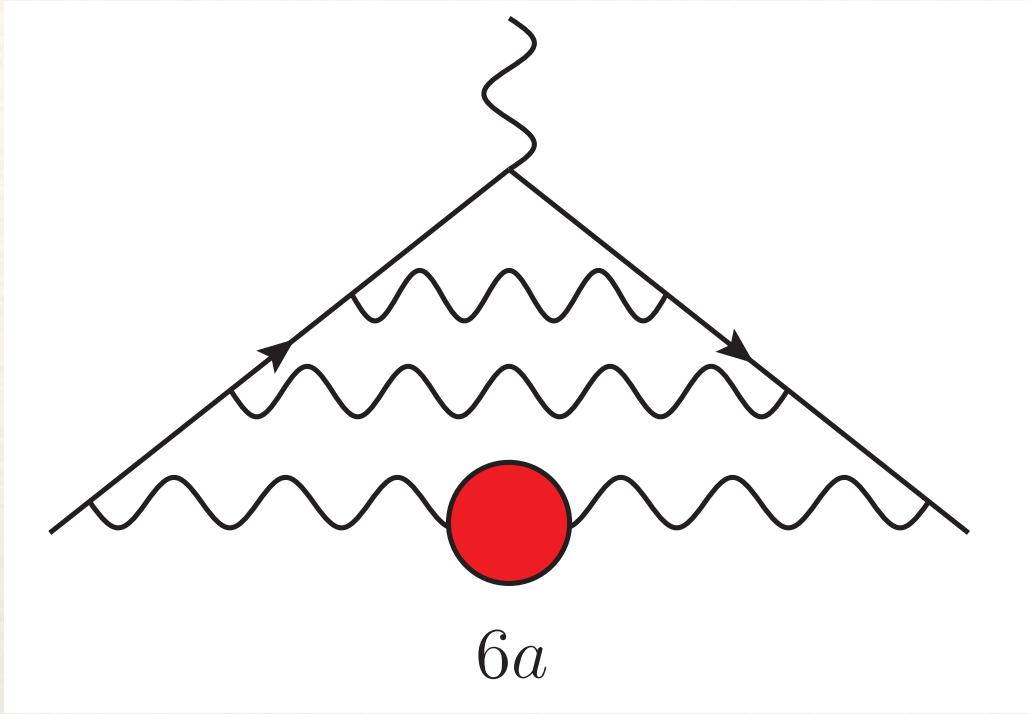
This uncertainty can be eliminated using our exact formula $\kappa^{(4)}(x)$.

Space-like method: NLO hadronic vacuum polarization contribution



- The LO integrand has a peak at $x \sim 0.914$
- The NLO integrand has an integrable logarithmic singularity at $x \rightarrow 1$

Space-like method: NNLO HVP contribution



$$a_\mu^{(6a)} = \frac{\alpha^3}{\pi^4} \int_{s_0}^\infty \frac{ds}{s} K^{(6a)}(s/m^2) \text{Im}\Pi_h(s)$$

Approximated series expansion of $K^{(6a)}(z)$ in the parameter $r = m^2/s$ contains powers r^n of degree $n = 1, 2, 3, 4$, multiplied by constants, $\ln r$, $\ln^2 r$, $\ln^3 r$ terms.

$$K^{(6a)}(s/m^2) = \int_0^1 d\xi \left[\frac{P^{(6a)}(\xi)}{1 + r\xi} + \frac{L^{(6a)}(\xi)}{\xi + r} \right]$$

$$L^{(6a)}(\xi) = G^{(6a)}(\xi) + H^{(6a)}(\xi) \ln \xi + j^{(6a)}(\xi) \ln^2 \xi$$

Exploiting generating integral representation to fit all the term.
Groote, Körner, Pivovarov, arXiv:0111206v2

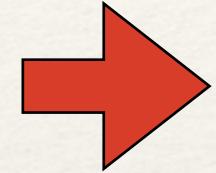
$$P^{(6a)}(\xi) = p_0^{(6a)} + p_1^{(6a)}\xi + p_2^{(6a)}\xi^2 + p_3^{(6a)}\xi^3$$

$$G^{(6a)}(\xi) = g_0^{(6a)} + g_1^{(6a)}\xi + g_2^{(6a)}\xi^2 + g_3^{(6a)}\xi^3$$

$$H^{(6a)}(\xi) = h_0^{(6a)} + h_1^{(6a)}\xi + h_2^{(6a)}\xi^2 + h_3^{(6a)}\xi^3$$

$$J^{(6a)}(\xi) = j_0^{(6a)} + j_1^{(6a)}\xi + j_2^{(6a)}\xi^2 + j_3^{(6a)}\xi^3$$

Space-like method: NNLO HVP contribution



$$a_\mu^{(6a)} = \left(\frac{\alpha}{\pi}\right)^3 \int_0^1 dx \bar{k}^{(6a)}(x) \Delta\alpha_h(t(x))$$

For $0 < x < x_\mu = (\sqrt{5} - 1)/2 = 0.618\dots$

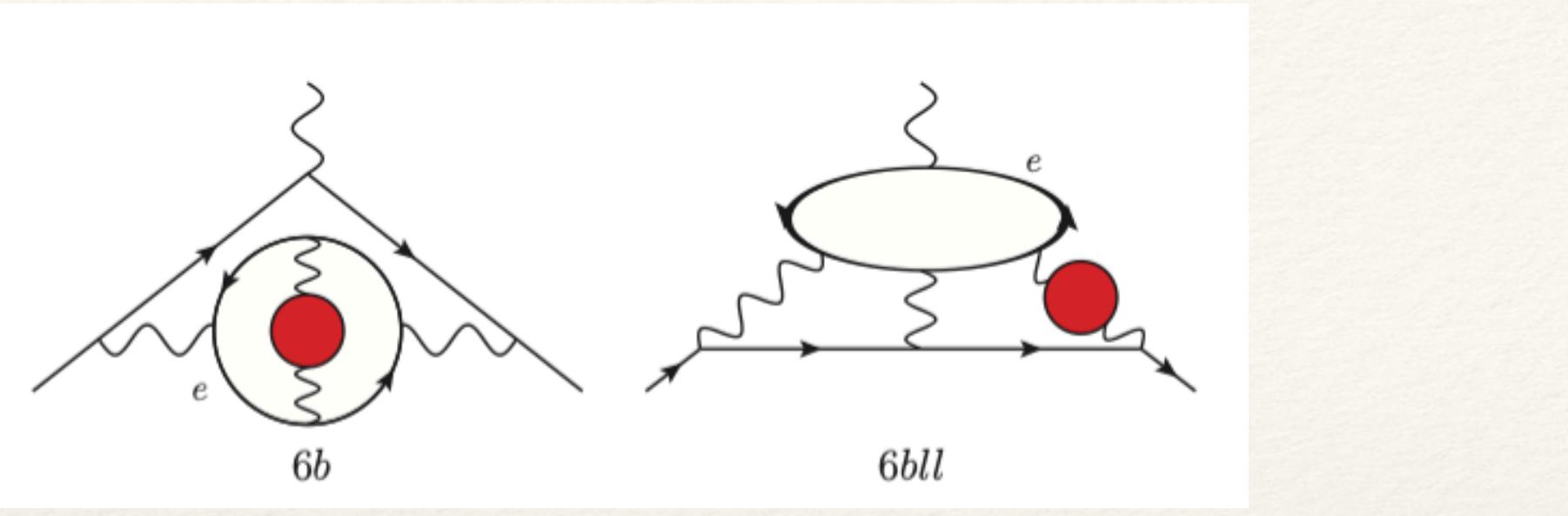
$$\bar{k}^{(6a)} = \frac{2-x}{x(1-x)} P^{(6a)} \left(\frac{x^2}{1-x} \right)$$

For $x_\mu < x < 1$

$$\bar{k}^{(6a)} = \frac{2-x}{x^3} L^{(6a)} \left(\frac{1-x}{x^2} \right)$$

(6a)	
$j_0 = 0;$	$h_0 = -\frac{359}{36};$
$j_1 = -\frac{3793}{864};$	$h_1 = \frac{122293}{5184};$
$j_2 = \frac{35087}{21600};$	$h_2 = -\frac{43879427}{648000};$
$j_3 = \frac{1592093}{43200};$	$h_3 = \frac{14388407}{48000};$
$g_0 = \frac{1301}{144} - \frac{19\pi^2}{9};$	
$g_1 = \frac{441277}{10368} + \pi^2 \left(-\frac{355}{648} + \ln 4 \right) + \frac{25}{2} \zeta(3);$	
$g_2 = -\frac{5051645167}{38880000} + \pi^2 \left(\frac{221411}{32400} - 18 \ln 2 \right) - \frac{3919}{60} \zeta(3);$	
$g_3 = \frac{14588342017}{38880000} + \pi^2 \left(-\frac{2479681}{64800} + 112 \ln 2 \right) + \frac{3113}{10} \zeta(3);$	
$p_0 = -\frac{1808080780513}{14580000} + \frac{41851\pi^4}{15} + \frac{8432 \ln^4 2}{3} + 67456 a_4 + \frac{2085448}{15} \zeta(3) +$	
$+ \pi^2 \left(-\frac{11944163099}{194400} + \frac{272}{3} (180 - 31 \ln 2) \ln 2 + \frac{115072}{3} \zeta(3) \right) - \frac{575360}{3} \zeta(5);$	
$p_1 = \frac{134017456919}{96000} - \frac{4481182\pi^4}{135} - \frac{98420 \ln^4 2}{3} - 787360 a_4 + 2255200 \zeta(5) +$	
$+ \pi^2 \left(\frac{23549054249}{32400} - 201122 \ln 2 + \frac{98420 \ln^2 2}{3} - 451040 \zeta(3) \right) - \frac{57189259}{36} \zeta(3);$	
$p_2 = -\frac{13069081405453}{3888000} + \frac{330073\pi^4}{4} + 80790 \ln^4 2 + 1938960 a_4 + \frac{77371609}{20} \zeta(3) +$	
$+ \pi^2 \left(-\frac{729995599}{405} + 6(85313 - 13465 \ln 2) \ln 2 + 1114360 \zeta(3) \right) - 5571800 \zeta(5);$	
$p_3 = \frac{1274611832039}{583200} - \frac{986377\pi^4}{18} - 53340 \ln^4 2 - 1280160 a_4 + \frac{11057200}{3} \zeta(5) +$	
$+ \pi^2 \left(\frac{5809659289}{4860} + 420 \ln 2 (-823 + 127 \ln 2) - \frac{2211440}{3} \zeta(3) \right) - \frac{22833188}{9} \zeta(3);$	

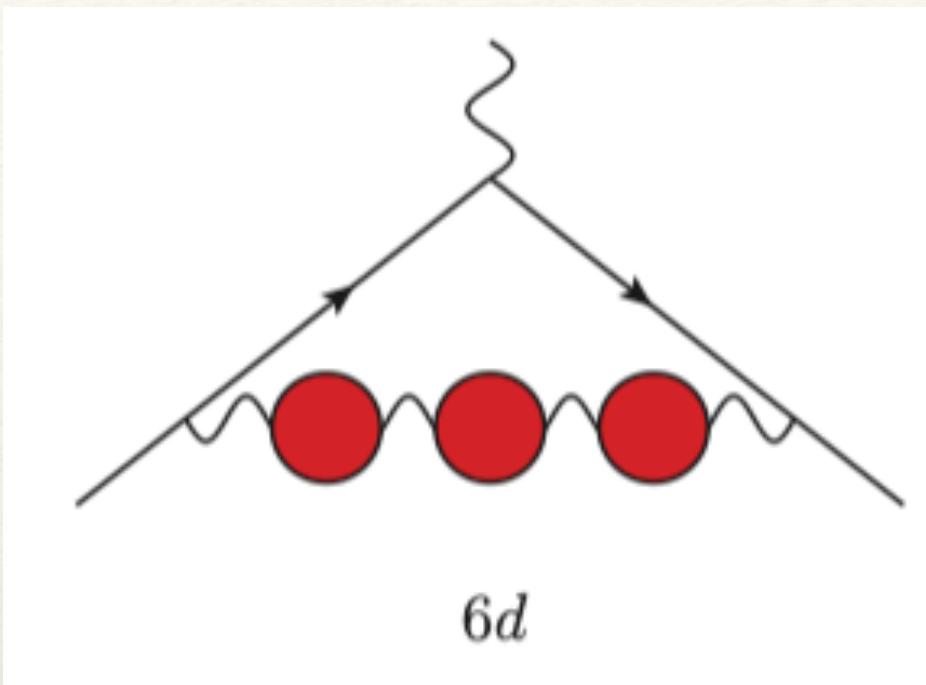
Space-like method: NNLO HVP contribution



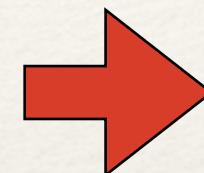
(6b)	
$j_0 = 0;$	$h_0 = \frac{65}{54};$
$j_1 = \frac{11}{27};$	$h_1 = -\frac{3559}{1296} + \rho^2 + \frac{5}{18} \ln \rho;$
$j_2 = \frac{41}{120};$	$h_2 = \frac{3917}{432} - \frac{82\rho^2}{3} + \frac{61}{10} \ln \rho;$
$j_3 = -\frac{507}{40};$	$h_3 = -\frac{4109}{80} + \frac{2211\rho^2}{10} - \frac{1763}{30} \ln \rho;$
$g_0 = \frac{1}{108} (259 - 72\rho^2 + 276 \ln \rho);$	
$g_1 = -\frac{9215}{1296} + \frac{65\pi^2}{162} - \frac{3\pi^2\rho}{4} + \frac{49\rho^2}{36} + (-\frac{301}{54} + 8\rho^2) \ln \rho + \frac{4}{3} \ln^2 \rho + 2 \zeta(3);$	
$g_2 = \frac{501971}{40500} - \frac{113\pi^2}{36} + \frac{270\pi^2\rho}{36} - \frac{8417\rho^2}{180} + (\frac{3479}{900} - 44\rho^2) \ln \rho - 8 \ln^2 \rho - 12 \zeta(3);$	
$g_3 = -\frac{2523823}{324000} + \frac{625\pi^2}{36} - 49\pi^2\rho + \frac{84946\rho^2}{225} + (\frac{987}{50} + 200\rho^2) \ln \rho + \frac{112}{3} \ln^2 \rho + 56 \zeta(3);$	
$p_0 = -\frac{95519053063}{486000} - 7275\pi^2\rho + \left(-\frac{587150693}{5400} + \frac{75272\rho^2}{3} + \frac{120800\pi^2}{9}\right) \ln \rho + \left(\frac{1135508}{9} + 96\rho^2\right) \zeta(3) +$	
$+ 4720 \ln^2 \rho + \frac{1067115409\rho^2}{5400} + \pi^2 \left(\frac{24382331}{810} - \frac{285184}{9} \ln 2\right) - 32\pi^2\rho^2 (687 + \ln 4);$	
$p_1 = \frac{279489728279}{121500} + \frac{179283\pi^2\rho}{2} + \left(\frac{2280933773}{1800} - 309540\rho^2 - \frac{1419328\pi^2}{9}\right) \ln \rho - \frac{10}{3} (446023 + 216\rho^2) \zeta(3) +$	
$- \frac{174712}{3} \ln^2 \rho - \frac{174350167\rho^2}{75} + \pi^2 \left(-\frac{143574463}{405} + \frac{3352256 \ln 2}{9}\right) + \frac{16}{3} \pi^2\rho^2 (48481 + 90 \ln 2);$	
$p_2 = -\frac{229560199193}{40500} - \frac{912495\pi^2\rho}{4} + \left(-\frac{1867939691}{600} + 788488\rho^2 + \frac{1168336\pi^2}{3}\right) \ln \rho + \left(\frac{11034553}{3} + 1440\rho^2\right) \zeta(3) +$	
$+ 148348 \ln^2 \rho + \frac{258653648\rho^2}{45} + \frac{4}{135} \pi^2 (29597029 - 31048560 \ln 2) - \frac{320}{3} \pi^2\rho^2 (5989 + \ln 512);$	
$p_3 = \frac{72762177677}{19440} + 154035\pi^2\rho - \frac{7}{108} (-31650719 + 3973440\pi^2 + 8220240\rho^2) \ln \rho - \frac{280}{9} (78283 + 27\rho^2) \zeta(3) +$	
$- 100240 \ln^2 \rho - \frac{513692207\rho^2}{135} + \frac{35}{162} \pi^2 (-2687659 + 2816064 \ln 2) + \frac{140}{3} \pi^2\rho^2 (9055 + \ln 4096);$	

(6bll)
$j_0 = 0;$
$j_1 = \frac{4}{27} - \frac{9\rho^2}{2};$
$j_2 = -\frac{41}{48} + \frac{2201\rho^2}{216};$
$j_3 = \frac{3037}{900} - \frac{5909\rho^2}{216};$
$h_0 = -\frac{9}{2};$
$h_1 = \frac{59}{9} - \frac{275\rho^2}{36} - 18\rho^2 \ln \rho;$
$h_2 = -\frac{485}{32} + \frac{1351\rho^2}{48} + \frac{659\rho^2}{18} \ln \rho;$
$h_3 = \frac{282617}{6750} - \frac{10481\rho^2}{108} - \frac{851\rho^2}{9} \ln \rho;$
$g_0 = \frac{43}{8} - 4\pi^2\rho + 15\rho^2 + \pi^2\rho^2 - 18\rho^2 \ln \rho + 6\rho^2 \ln^2 \rho;$
$g_1 = -\frac{73}{81} + \frac{8\pi^2}{81} + \frac{40\pi^2\rho}{9} + \frac{2437\rho^2}{108} + \frac{17\pi^2\rho^2}{9} + \frac{607\rho^2}{18} \ln \rho - \frac{20\rho^2}{3} \ln^2 \rho + \frac{2}{3} \zeta(3) + 2\rho^2 \zeta(3);$
$g_2 = -\frac{385}{162} - \frac{41\pi^2}{72} - \frac{28\pi^2\rho}{3} - \frac{89873\rho^2}{5184} - \frac{997\pi^2\rho^2}{324} - \frac{1961\rho^2}{72} \ln \rho + 14\rho^2 \ln^2 \rho - \frac{5}{2} \zeta(3) - \frac{16\rho^2}{3} \zeta(3);$
$g_3 = \frac{2691761}{202500} + \frac{3037\pi^2}{1350} + 24\pi^2\rho + \frac{655429\rho^2}{97200} + \frac{2359\pi^2\rho^2}{324} + \frac{6943\rho^2}{360} \ln \rho - 36\rho^2 \ln^2 \rho + \frac{42}{5} \zeta(3) + 15\rho^2 \zeta(3);$
$p_0 = -\frac{343277101}{45000} - \frac{33156604927\rho^2}{583200} + \pi^2 \left(-\frac{615427}{4050} + \frac{6776\rho}{3} + \frac{763121\rho^2}{972}\right) - \frac{4\pi^4}{135} (7817 + 3212\rho^2) +$
$+ \left(-\frac{7290521}{3240} + \frac{49622\pi^2}{27} - \frac{128\pi^4}{9}\right) \rho^2 \ln \rho + \left(-3388 - \frac{80\pi^2}{3}\right) \rho^2 \ln^2 \rho +$
$+ \left(25642 + \frac{1515724\rho^2}{27} - 128\pi^2\rho^2 - 160\rho^2 \ln \rho\right) \zeta(3) - \frac{1280}{3} \rho^2 \zeta(5);$
$p_1 = \frac{89280434843}{972000} + \frac{248834878697\rho^2}{388800} - \frac{1}{324} \pi^2 (-533001 + 9110736\rho + 3110417\rho^2) + \frac{2}{135} \pi^4 (180247 + 73530\rho^2) +$
$+ \left(\frac{11101973}{1080} - \frac{193400\pi^2}{9} + \frac{320\pi^4}{3}\right) \rho^2 \ln \rho + \frac{2}{3} (63269 + 300\pi^2) \rho^2 \ln^2 \rho +$
$+ \frac{1}{45} (-13410977 + 100 (-292301 + 432\pi^2) \rho^2 + 54000\rho^2 \ln \rho) \zeta(3) + 3200\rho^2 \zeta(5);$
$p_2 = -\frac{6209532853}{27000} - \frac{29997466847\rho^2}{19440} + \pi^2 \left(-\frac{114521}{30} + 71840\rho + \frac{1970140\rho^2}{81}\right) - \frac{4}{9} \pi^4 (14685 + 6032\rho^2) +$
$- \frac{1}{54} (190613 - 2847360\pi^2 + 11520\pi^4) \rho^2 \ln \rho - 80 (1347 + 5\pi^2) \rho^2 \ln^2 \rho +$
$- \frac{10}{9} (-658509 + (-1431463 + 1728\pi^2) \rho^2 + 2160\rho^2 \ln \rho) \zeta(3) - 6400\rho^2 \zeta(5);$
$p_3 = \frac{49726331179}{324000} + \frac{7324831423\rho^2}{7290} + \pi^2 \left(\frac{3897971}{1620} - \frac{145880\rho}{3} - \frac{3977785\rho^2}{243}\right) + \frac{14}{27} \pi^4 (8269 + 3419\rho^2) +$
$+ \frac{7}{81} (-81551 - 401520\pi^2 + 1440\pi^4) \rho^2 \ln \rho + \frac{140}{3} (1563 + 5\pi^2) \rho^2 \ln^2 \rho +$
$+ \frac{35}{27} (-371889 + 16 (-50437 + 54\pi^2) \rho^2 + 1080\rho^2 \ln \rho) \zeta(3) + \frac{11200}{3} \rho^2 \zeta(5);$

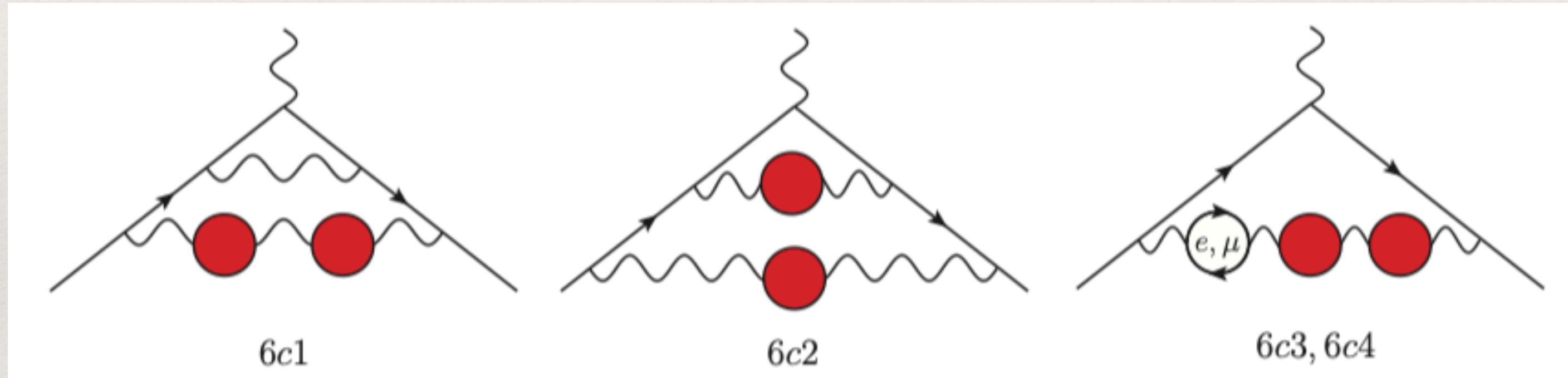
Space-like method: NNLO HVP contribution



$$a_\mu^{(6d)} = \frac{\alpha}{\pi^4} \int_{s_0}^{\infty} \frac{ds}{s} \frac{ds'}{s'} \frac{ds''}{s''} K^{(6d)}(s, s', s'') \times \text{Im}\Pi_h(s) \text{Im}\Pi_h(s') \text{Im}\Pi_h(s'')$$



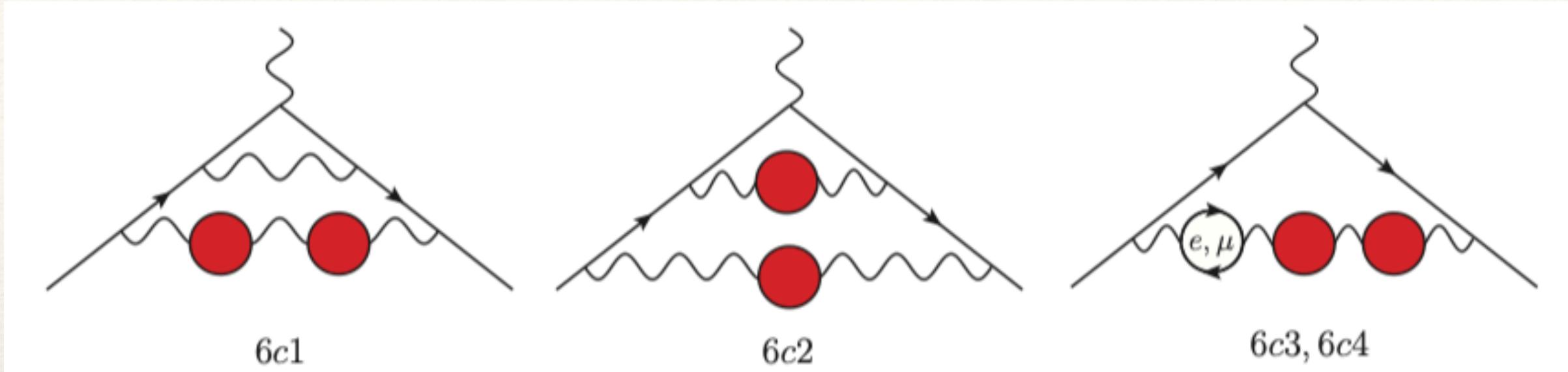
$$a_\mu^{(6d)} = \left(\frac{\alpha}{\pi} \right) \int_0^1 dx \kappa^{(2)}(x) [\Delta\alpha_h(t(x))]^3$$



$$a_\mu^{(6c)} = \frac{\alpha^2}{\pi^4} \int_{s_0}^{\infty} \frac{ds}{s} \frac{ds'}{s'} K^{(6c)}(s/m^2, s'/m^2) \times \text{Im}\Pi_h(s) \text{Im}\Pi_h(s')$$

(6c) diagrams contain two HVP insertions \rightarrow $a_\mu^{(6c)}$ requires two dispersive integrations

Space-like method: NNLO hadronic vacuum polarization contribution



$$a_\mu^{(6c1)} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^1 dx \lambda^{(4)}(x) [\Delta\alpha_h(t(x))]^2$$

$$\lambda^{(4)}(x) = \kappa^{(4)}(x) - \frac{2\pi}{\alpha} \kappa^{(2)}(x) \Delta\alpha_\mu^{(2)}(t(x))$$

We start from the time-like asymptotic expansions:

- $K^{(6c)}(s/m^2, s'/m^2)$ in $s' \approx s \gg m^2$
- $K^{(6c)}(s/m^2, s'/m^2)$ in $s' \gg s \gg m^2$

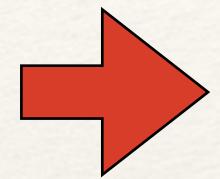
$$a_\mu^{(6c3)} = \frac{\alpha}{\pi} \int_0^1 dx \kappa^{(2)}(x) [3\Delta\alpha_h(t(x))^2 \Delta\alpha_e^{(2)}(t(x))]$$

$$a_\mu^{(6c4)} = \frac{\alpha}{\pi} \int_0^1 dx \kappa^{(2)}(x) [3\Delta\alpha_h(t(x))^2 \Delta\alpha_\mu^{(2)}(t(x))]$$

Two steps:

1. Computed the approximated time-like kernel $K^{(6c2)}(s/m^2, s'/m^2)$
2. Matched the LO of the approximated time-like kernels $K^{(6c2)}(s/m^2, s'/m^2)$, with those of a series expansion of a two dimensional generating integral representation

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$$a_\mu^{(6c2)} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^1 dx \int_0^1 dx' \bar{\kappa}^{(6c2)} \times \Delta\alpha_h(t(x)) \Delta\alpha_h(t(x'))$$

For $x_\mu < \{x, x'\} < 1$

$$\bar{\kappa}^{(6c2)}(x, x') = \frac{2-x}{x^3} \frac{2-x'}{x'^3} G^{(6c2)}\left(\frac{1-x}{x^2}, \frac{1-x'}{x'^2}\right)$$

$$G^{(6c2)}(\xi, \xi') = \frac{1}{4(32\pi^2 - 315)} \times \\ \times \left[(1855 - 188\pi^2) \frac{\min(\xi, \xi')}{\max(\xi, \xi')^2} + (988\pi^2 - 9765) \frac{\min(\xi, \xi')^2}{\max(\xi, \xi')^3} + 24(435 - 44\pi^2) \frac{\min(\xi, \xi')^3}{\max(\xi, \xi')^4} \right]$$

Conclusions

- ❖ We provide simple analytic expressions to calculate the HVP contributions to the muon $g - 2$ in the space-like region up to NNLO.
- ❖ These results will allow to extend MUonE's original proposal to determine the LO contribution to $a_\mu^{(HVP)}$ to NNLO. They can also be employed in lattice QCD computations.
- ❖ An existing approximation for the NLO kernel induced large uncertainties. These can be eliminated using our exact expression $\kappa^{(4)}(x)$.
- ❖ For NNLO a combination of exact and approximated one- and two-dimensional kernels is provided. The uncertainty due to the kernel approximations is estimated to be less than $O(10^{-12})$.
- ❖ Calculation of higher-order HVP corrections require a precise treatment of the QED radiative corrections to the HVP function. MUonE will naturally include their leading effects in the space-like approach.
- ❖ These results allow to compare time-like and space-like calculations of $a_\mu^{(HVP)}$ at NNLO accuracy.

Thank you for your attention!