New tools towards fast evaluation of Feynman integrals

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Outline

I. Introduction

- II. Block-triangular system and Blade
- **III.** Auxiliary mass flow and AMFlow
- **IV.** Applications
- V. Summary and outlook

Feynman integrals

> Feynman integrals in dimensional regularization

$$I(\vec{\nu}) = \int \prod_{i=1}^{L} \frac{\mathrm{d}^{D}\ell_{i}}{\mathrm{i}\pi^{D/2}} \frac{\mathcal{D}_{K+1}^{-\nu_{K+1}} \cdots \mathcal{D}_{N}^{-\nu_{N}}}{(\mathcal{D}_{1} + \mathrm{i}0)^{\nu_{1}} \cdots (\mathcal{D}_{K} + \mathrm{i}0)^{\nu_{K}}}$$

- interesting in physics
 - key objects for fixed order computation
 - Standard Model & New Physics
- challenging in mathematics
 - algebraic complexity (integrals reduction)
 - function space (master integrals computation)
- state of the art

- massive two-loop five-point integrals
- massive three-loop four-point integrals

Feynman integrals

Methods and tools

- Reduction:
 - integration-by-parts reduction [Chetyrkin and Tkachov, Nucl. Phys. B, 1981] [Laporta, Int. J. Mod. Phys. A, 2000]
 - AIR [Anastasiou and Lazopoulos, JHEP, 2004] FIRE [Smirnov, JHEP, 2008] [Smirnov, Smirnov, Comput. Phys. Commun., 2013] [Smirnov, Comput. Phys. Commun., 2015] [Smirnov and Chuharev, Comput. Phys. Commun., 2020]
 Reduze [Studerys, Comput. Phys. Commun., 2010] [Manteuffel and Studerus, e-Print: 1201.4330] Kira
 [Maierhofer, Usovitsch and Uwer, Comput. Phys. Commun., 2018] [Klappert, Lange, Maierhofer and Usovitsch, Comput. Phys. Commun., 2021]
 - Syzygy equations [Gluza, Kajda and Kosower, Phys. Rev. D, 2011] [Larsen and Zhang, Phys. Rev. D, 2016]
 - NeatIBP [Wu, Boehm, Ma et al, e-Print: 2305.08783]
 - finite field reconstruction [Manteuffel and Schabinger, Phys. Lett. B, 2015] [Peraro, JHEP, 2016]
 - FiniteFlow [Peraro, JHEP, 2019] FireFly [Klappert and Lange, Comput. Phys. Commun., 2020] RATRACER [Magerya, e-Print: 2211.03572]
 - block-triangular system and Blade

Feynman integrals

- Computation:
 - canonical differential equations [Kotikov, Phys. Lett. B, 1991] [Henn, Phys. Rev. Lett., 2013]
 - CANONICA [Meyer, 2018] epsilon [Prausa, 2017] Fuchsia [Gituliar and Magerya, 2017] Libra [Lee, 2020] DLogBasis [Henn, Mistlberger, Smirnov and Wasser, JHEP 2020]
 - numerical differential equations [Caffo, Czyz, Gunia and Remiddi, Comput. Phys. Commun., 2009] [Czakon, Phys. Lett. B, 2008]
 - DiffExp [Hidding, Comput.Phys.Commun., 2021] SeaSyde [Armadillo, Bonciani, Devoto, Rana and Vicini, Comput.Phys.Commun., 2023]
 - sector decomposition [Binoth and Heinrich, Nucl. Phys. B, 2000]
 - FIESTA [Smirnov and Tentyukov, 2009] [Smirnov, Smirnov and Tentyukov, 2011] [Smirnov, 2014] [Smirnov, 2015]
 SecDec [Carter and Heinrich, 2011] [Borowka, Carter and Heinrich, 2013] [Borowka, Heinrich, Jones, et al, 2015]
 pySecDec [Borowka, Heinrich, Jahn, et al, 2018]
 - tropical geometry [Borinsky, 2008.12310] [Borinsky, Munch and Tellander, 2302.08955]
 - feyntrop [Borinsky, Munch and Tellander, 2302.08955]
 - auxiliary mass flow and AMFlow

> Other good methods

- Intersection theory [Mastralia and Mizera, JHEP, 2019] [Frellesvig, Gasparotto, Laporta, et al, JHEP 2019]
- Partial fraction decomposition [Boehm, Wittmann, Wu, et al, JHEP, 2020]
- Method of region [Beneke and Smirnov, Nucl. Phys. B, 1998] [Smirnov, Phys. Lett. B, 1999]
- Mellin-Barnes representation [Boos and Davydychev, Theor.Math.Phys., 1991][Smirnov, Phys. Lett. B, 1999]
- Dimensional recurrence and Analyticity [Tarasov, Phys. Rev. D, 1996] [Lee, Nucl. Phys. B, 2010]
- Positivity constraints [Zeng, 2303.15624]
- ...



I. Introduction

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Block-triangular system

Motivation

- IBP system
 - large: many equations and integrals, typically $10^5 \sim$
 - coupled: the systems usually become denser during the solutions
 - time- and resource-consuming
- How to improve?
 - small: as few equations as possible
 - decoupled or semi-decoupled: the system should never become denser
- A possible solution: block-triangular system
- A simple example:

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- 2 equations reducing I_6 , I_7 to I_8
- 2 equations reducing I_4 , I_5 to I_8
- 3 equations reducing I_1, I_2, I_3 to I_8

$$\begin{split} 2I_1 + 3sI_2 + tI_3 + 5I_4 + DI_5 - I_6 + DI_7 + sI_8 &= 0\\ 2sI_1 + tI_2 + DI_3 + stI_4 + I_5 - 6I_6 + sI_7 + DI_8 &= 0\\ 2tI_1 + DI_2 - I_3 - 2sI_4 - 8I_5 + tI_6 - I_7 + sI_8 &= 0\\ I_4 + I_5 - sI_6 + tI_7 - DI_8 &= 0\\ I_4 - I_5 - DI_6 + stI_7 + D^2I_8 &= 0\\ sI_6 + tI_7 - I_8 &= 0\\ stI_6 - I_7 + I_8 &= 0 \end{split}$$

Relations among Feynman integrals

> Feynman integrals are vectors



- $\vec{s} = \{s_{12}, s_{23}, s_{34}, s_{45}, s_{15}\}$
- 108 master integrals: $\{M_1, ..., M_{108}\}$
- each integral can be written as a linear combination
 - $I = c_1 M_1 + \dots + c_{108} M_{108} \iff (c_1 \dots c_{108})^{\mathrm{T}}$
 - c_1, \ldots, c_{108} are numerically computable from IBP system over finite fields

Relations among Feynman integrals

- > Relations among $G \equiv \{I_1, I_2, ..., I_n\}$ $Q_1I_1 + Q_2I_2 + \cdots + Q_nI_n = 0$
 - Q_i : polynomials of \vec{s} and D

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- ansatz: $Q_1 = a + bD$, $Q_2 = cs_{23} + ds_{34}$, $Q_3 = es_{45} + fs_{15}$, ...
- *a*, *b*, *c*, *d*, *e*, *f*, ... : unknown rational numbers to be determined (The total number of unknowns is finite!)
- For a given ansatz (search algorithm)
 - compute the numerical reduction coefficients of $I_1, ..., I_n$ at $D = D_0$ and $\vec{s} = \vec{s}_0$
 - this can provide at most 108 linear constraints for the unknowns

$$a\begin{pmatrix} c_{1,1}\\ \vdots\\ c_{1,108} \end{pmatrix} + bD_0 \begin{pmatrix} c_{1,1}\\ \vdots\\ c_{1,108} \end{pmatrix} + cs_{23,0} \begin{pmatrix} c_{2,1}\\ \vdots\\ c_{2,108} \end{pmatrix} + \dots = 0$$

• repeat the above process until there are enough constraints

Reduction

≻ Reducing $G_1 = \{I_1, ..., I_m\}$ to $G_2 = \{I_{m+1}, ..., I_n\}$

- 1. Set $G = G_1 \cup G_2 = \{I_1, \dots, I_n\};$
- for a given ansatz of Q₁, ..., Q_n, use the search algorithm to determine all the relations;
- if the obtained relations are sufficient to reduce G₁, stop; else make another ansatz and go back to step 2.

Making ansatz

- from simple to complicated, for instance:
 - degree-0: $Q_1 = a_1$, $Q_2 = a_2$, $Q_3 = a_3$, ...
 - degree-1: $Q_1 = a_1 + b_1 D + c_1 s_{23} + d_1 s_{34} + e_1 s_{45} + f_1 s_{15}, Q_2 = a_2 + b_2 D + c_2 s_{23} + d_2 s_{34} + e_2 s_{45} + f_2 s_{15}, Q_3 = \cdots, \cdots$
 - degree-2: $Q_1 = a_1 + b_1 D + c_1 s_{23} + d_1 s_{34} + e_1 s_{45} + f_1 s_{15} + g_1 D^2 + h_1 D s_{23} + i_1 D s_{34} + \cdots$

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Block-triangular system

\succ Constructing G_1 and G_2 for each block

- Basic construction
 - *G*₁: target integrals from some sector
 - *G*₂: integrals in its subsectors
 - if G_1 does not contain dotted integrals, then G_2 as well
- Example (top-sector):

- *G*₁: top-sector integrals with up to degree-5 numerators
 - {I(1,1,1,1,1,1,1,1,0,0,0), I(1,1,1,1,1,1,1,1,1,1,0,0), I(1,1,1,1,1,1,1,1,1,0,0,-1,0), ..., I(1,1,1,1,1,1,1,1,0,-5,0), I(1,1,1,1,1,1,1,0,0,-5)}
- *G*₂: integrals in subsectors
 - {I(0,1,1,1,1,1,1,1,0,0,0), ..., I(0,1,1,1,1,1,1,1,0,0,-4), I(1,0,1,1,1,1,1,1,0,0,0), ..., I(1,0,1,1,1,1,1,1,0,0,-4), ..., I(0,0,1,1,1,1,1,1,0,0,0), ...}

Results

- IBP system
 - $\sim 3 \times 10^5$ equations
 - numerical sampling over finite fields: 7s per phase-space point
- Block-triangular system
 - 3806 equations reducing 3914 integrals to 108 master integrals
 - 0.17s per phase-space point
 - improvement: $7/0.17 \approx 40$



Blade (in progress)



- BLock-triAngular form improved Feynman integral DEcomposition
- Available at <u>https://gitlab.com/multiloop-pku/blade</u>
- Installation
 - \$ chmod +x auto_install
 - \$./auto_install

Blade (in progress)

Basic usage

```
family = dbox;
dimension = 4-2*eps;
loop = {l1,l2};
leg = {p1,p2,p3,p4};
conservation = {p4->-p1-p2-p3};
replacement = {p1^2 -> 0, p2^2 -> 0, p3^2 -> msq, p1 p2 -> s/2, p1 p3 -> t/2, p2 p3 -> (-s-t)/2};
propagator = { (l1)^2, (l1+p1)^2, (l1+p1+p2)^2, l2^2-msq, (l1+l2)^2-msq, (l2-p1-p2)^2-msq, (l2-p3)^2, (l1-p3)^2, (l2+p1)^2);
topsector = {1,1,1,1,1,1,1,0,0};
numeric = {s -> 1};
```

BLFamilyDefine[family,dimension, propagator,loop,leg,conservation, replacement,topsector,numeric]

```
target={BL[dbox,{1,1,1,1,1,1,0,-3}],BL[dbox,{1,1,1,1,1,1,1,1,-1,-2}],
BL[dbox,{1,1,1,1,1,1,1,-2,-1}],BL[dbox,{1,1,1,1,1,1,1,-3,0}],
BL[dbox,{2,1,1,1,1,1,0,0,-2}], BL[dbox,{2,1,1,1,1,1,0,-2,0}],
BL[dbox,{0,1,2,2,1,1,0,0,-2}], BL[dbox,{0,1,2,2,1,1,0,-2,0}]);
```

res = BLReduce[target];

de = BLDifferentialEquation[{t}]

- Define the integral family using "BLFamilyDefine"
- Reduce target integrals using "BLReduce"
- Construct the differential equations using "BLDifferentialEquation"

```
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```



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Motivation

- Traditional methods
 - Analytic: less systematic, case-by-case analysis
 - Numerical: less efficient, hard to obtain high precision results
- A possible solution: semi-analytic (semi-numerical) method
 - reserve analytic structure as far as possible
 - perform quickly-converged numerical computation
- A good candidate: differential equations + power series expansion
 - + systematic boundary conditions

Dimensionally regulated Feynman integrals

• integrals with auxiliary mass parameter η

$$I_{\text{aux}}(\vec{\nu};\eta) = \int \prod_{i=1}^{L} \frac{\mathrm{d}^{D}\ell_{i}}{\mathrm{i}\pi^{D/2}} \frac{\mathcal{D}_{K+1}^{-\nu_{K+1}} \cdots \mathcal{D}_{N}^{-\nu_{N}}}{(\mathcal{D}_{1}-\eta)^{\nu_{1}} \cdots (\mathcal{D}_{K}-\eta)^{\nu_{K}}}$$

• obtain physical integrals through

$$I(\vec{\nu}) = \lim_{\eta \to i0^{-}} I_{\text{aux}}(\vec{\nu};\eta)$$

- method of region [Beneke and Smirnov, Nucl. Phys. B, 1998]
 - the only contributing region: $\ell_i^{\mu} \sim \sqrt{\eta}$

$$\frac{1}{((\ell+p)^2 - m^2 - \eta)^{\nu}} = \frac{1}{(\ell^2 - \eta)^{\nu}} \sum_{i=0}^{\infty} \frac{(\nu)_i}{i!} \left(-\frac{2\ell \cdot p + p^2 - m^2}{\ell^2 - \eta} \right)^i$$

fully massive vacuum integrals [Davydychev and Tausk, Nucl, Phys. B, 1993] [Broadhurst, Eur. Phys. J.
 C, 1999][Schroder and Vuorinen, JHEP, 2005] [Kniehl, Pikelner and Veretin, JHEP, 2017][Luthe, phdthesis, 2015]
 [Luthe, Maier, Marquard et al, JHEP, 2017]

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> Analytic continuation

• differential equations

 $\frac{\partial}{\partial \eta} \vec{\mathcal{I}}_{\mathrm{aux}}(\eta) = A(\eta) \vec{\mathcal{I}}_{\mathrm{aux}}(\eta)$

- boundary conditions at $\eta = \infty$
- define a path: $\{\eta_0, \eta_1, \dots, \eta_l\}$
- expand at $\eta = \infty$ to estimate $I(\eta_0)$
- expand at $\eta = \eta_i$ to estimate $I(\eta_{i+1})$
- expand formally at $\eta = 0$ and match at

 $\eta=\eta_l$

- η_0 : outside the larger circle
- η_l : inside the smaller circle
- $|\eta_{i+1} \eta_i| < r_i$

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AMFlow Free Project ID: 32748265



-0- 73 Commits 🖇 1 Branch 🛷 2 Tags 🗔 594 KB Project Storage 🕑 1 Release

A proof-of-concept implementation of auxiliary mass flow method.

- available at https://gitlab.com/multiloop-pku/amflow
- basic features
 - systematic: works for arbitrary integrals in principle
 - one-loop six-point[Henn, Matijasic and Miczajka, JHEP 2023], two-loop five-point[Badger, Becchetti, Chaubey and Marzucca, JHEP 2023], three-loop three-point [Fael, Lange, Schonwald and Steinhauser, 2302.00693] ...
 - efficient: easy to reach high precision
 - O(100) digits [Badger, Becchetti, Chaubey and Marzucca, JHEP 2023] ..
 - O(1000) digits [Abreu, Becchetti, Duhr and Ozcelik, JHEP 2022]
 - user-friendly: press the button & wait for the results

AMFlow

examples/automatic_vs_manual

AMFlowInfo["Conservation"] = (p4 -> -p1-p2-p3);

```
(*load the package*)
current = If[$FrontEnd===Null,$InputFileName,NotebookFileName[]]//DirectoryName;
Get[FileNameJoin[{current, "..", "..", "AMFlow.m"}]];
(*set ibp reducer, could be "FiniteFlow+LiteRed", "Kira" or "FIRE+LiteRed"*)
SetReductionOptions["IBPReducer" -> "Kira"];
(*configuration of the integral family*)
AMFlowInfo["Family"] = tt;
AMFlowInfo["Loop"] = {l1, l2};
AMFlowInfo["Leg"] = (p1, p2, p3, p4);
```

AMFlowInfo["Propagator"] = (l1^2, (l1+p1)^2, (l1+p1+p2)^2, l2^2, -msq+(l2+p3)^2, (l2+p3+p4)^2, (l1+l2)^2, (l1-p3)^2, (l2+p1)^2);

- SetReductionOptions["IBPReducer" -> "reducer"];
 - "reducer": "Blade", "FIRE+LiteRed", "FiniteFlow+LiteRed", "Kira"

AMFlowInfo["Replacement"] = {p1^2 -> 0, p2^2 -> 0, p3^2 -> msq, p4^2 -> msq, (p1+p2)^2 -> s, (p1+p3)^2 -> t};

SolveIntegrals[targets, precision, epsorder];

•

AMFlow

- other features
 - compute integrals at any space-time dimension, $D = 6 2\epsilon$, $3 2\epsilon$, $11/2 2\epsilon$
 - see examples/spacetime_dimension
 - compute integrals with complex masses, $m^2 = 100 5i$
 - see examples/complex_kinematics
 - compute phase-space integrals based on reverse unitarity [Anastasiou and Melnikov, Nucl. Phys. B, 2002]
 - $\delta(p_i^2 m_i^2) \rightarrow \frac{1}{[p_i^2 m_i^2]_c}$
 - see examples/automatic_phase and examples/feynman_prescription



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Applications

 $\geq e^+e^- \rightarrow q \overline{q} @ \text{NNNLO QCD} [Xiang Chen, Xin Guan, Chuan-Qi He, XL and Yan-Qing Ma, e-Print: 2209.14259] }$

- Compute the total cross section as a piecewise function of $x \coloneqq 4m^2/s$
 - reduce to master integrals using Blade

$$A = \sum c_i M_i$$

• construct differential equations of master integrals with respect to *x* using Blade

$$\frac{\partial}{\partial x}M_i = \sum A_{ij}M_j$$

• compute the master integrals at a regular point x = 4/23 using AMFlow, which serves as boundary conditions

$$M_i(4/23) = \cdots$$

- solve the differential equations numerically to obtain piecewise functions of master integrals, using the differential equation solver in AMFlow
- substitute the results of master integrals into the cross section



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Summary and Outlook

- Integrals reduction: block-triangular system, Blade
- Master integrals computation: auxiliary mass flow, AMFlow
- Solution of univariate problems? Completely automated.
- Solution of multivariate problems?
 - From point to point? OK.
 - More efficient methods? Not known yet.

Thank you!

$\succ I(\eta)$ as an analytic function of η

- there should be a maximal threshold $\eta = \eta_{th}$ on the real axis
 - $I(\eta)$ is real-valued for $\eta > \eta_{\text{th}}$ and complex-valued for $\eta < \eta_{\text{th}}$
- branch cut can be chosen as the straight line connecting $\eta = -\infty$ and $\eta = \eta_{\text{th}}$ along the real axis, such that $I(\eta^*) = I^*(\eta)$



> A simple example: one-loop massless bubble

$$I(\nu_1, \nu_2) = \int \frac{\mathrm{d}^D \ell}{\mathrm{i}\pi^{D/2}} \frac{1}{(\ell^2)^{\nu_1} ((\ell+p)^2)^{\nu_2}}$$

• master integral: I(1, 1)

$$I(1,1) = (-p^2 - i0)^{D/2 - 2} \times \frac{\Gamma(2 - D/2)\Gamma(D/2 - 1)^2}{\Gamma(D - 2)}$$
$$I(1,1)|_{p^2 = 1, D = 4 - 2\epsilon} = \frac{1}{\epsilon} + (2 - \gamma + i\pi) + O(\epsilon^1)$$
$$= \frac{1}{\epsilon} + (1.42278 + 3.14159i) + O(\epsilon^1)$$

• insert auxiliary mass

$$- \int \frac{\mathrm{d}^D \ell}{\mathrm{i}\pi^{D/2}} \frac{1}{(\ell^2 - \eta)^{\nu_1} ((\ell + p)^2 - \eta)^{\nu_2}}$$

- master integrals: $\vec{I}_{aux}(\eta) = \{I_{aux}(1,0;\eta), I_{aux}(1,1;\eta)\}$
- construct differential equations using IBP reduction

$$\frac{\partial}{\partial \eta} \vec{\mathcal{I}}_{aux}(\eta) = \begin{pmatrix} \frac{1-\epsilon}{\eta} & 0\\ \frac{2(\epsilon-1)}{\eta(4\eta-1)} & -\frac{2(2\epsilon-1)}{4\eta-1} \end{pmatrix} \vec{\mathcal{I}}_{aux}(\eta)$$

- $\eta_{\rm th} = 1/4$
- boundary conditions

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$$I_{\text{aux}}(1,0;\eta) = \eta^{1-\epsilon} \times (-\Gamma(\epsilon-1))$$
$$I_{\text{aux}}(1,1;\eta) \sim \eta^{-\epsilon} \times (\Gamma(\epsilon) + \mathcal{O}(\eta^{-1}))$$



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- expand near $\eta = \eta_0 = -i/2$
 - $I_{\text{aux}}(1,1;\eta) = \sum_{n=0}^{\infty} a_n(\epsilon)(\eta \eta_0)^n$
 - $a_0(\epsilon) = \epsilon^{-1} + 0.0548501 + 1.88709i$
 - $a_1(\epsilon) = 0.5714 1.77538i$
 - ...
 - $a_{100}(\epsilon) = -1.29958 \times 10^{24} + 1.28029 \times 10^{26}$
- estimate at $\eta = \eta_1 = -i/4$ to obtain
 - $I_{\text{aux}}(1,1;-i/4) = \epsilon^{-1} + 0.609168 + 2.13174i$

- expand near $\eta = \eta_1 = -i/4$ and estimate at $\eta = -i/8$ to obtain
 - $I_{\text{aux}}(1,1;-i/8) = \epsilon^{-1} + 0.994236 + 2.42639i$

- expand near $\eta = 0$
 - $I_{\text{aux}}(1,1;\eta) = \sum_{n=0}^{\infty} a_n(\epsilon)\eta^n + \eta^{1-\epsilon} \sum_{n=0}^{\infty} b_n(\epsilon)\eta^n$
 - $b_n(\epsilon)$ can be totally determined by sub-topology
 - $b_0(\epsilon) = -2\Gamma(\epsilon 1)$
 - $b_1(\epsilon) = 4\Gamma(\epsilon 1)/(\epsilon 2)$
 - ...
 - $a_n(\epsilon)$ cannot be totally determined but can be reduced to $a_0(\epsilon)$
 - $a_1(\epsilon) = 2(2\epsilon 1)a_0(\epsilon)$
 - $a_2(\epsilon) = 2(2\epsilon 1)(2\epsilon + 1)a_0(\epsilon)$

- match at $\eta = \eta_2 = -i/8$ to obtain $a_0(\epsilon) = \epsilon^{-1} + 1.42278 + 3.14159i$
- take the limit $\eta \rightarrow i0^-$
 - $\lim_{\eta \to i0^{-}} I_{aux}(1,1;\eta) = a_0(\epsilon) = \epsilon^{-1} + 1.42278 + 3.14159i$

^{• ...}