



# Unitarity Triangle analysis in the SM and beyond from UTfit



LAPHROAIG®

Marcella Bona



Heavy Flavour 2016: Quo Vadis?

Thursday July 14<sup>th</sup>

Islay, Scotland

*still UK.. still EU..*

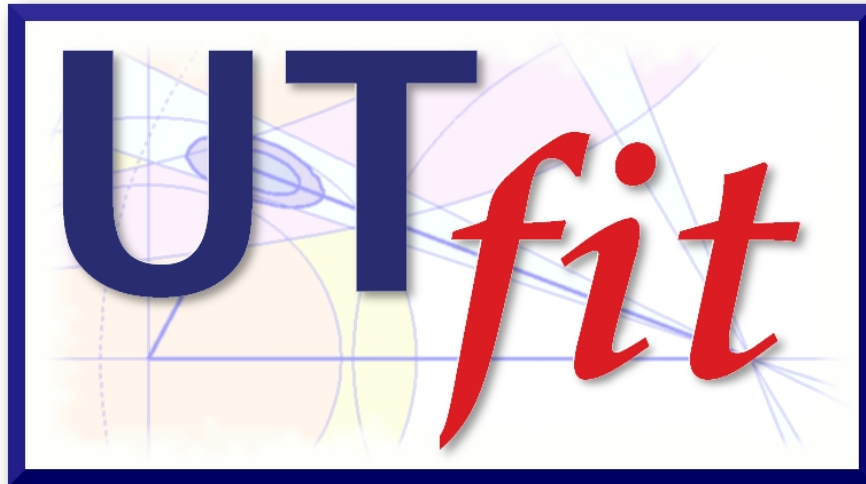


# Unitarity Triangle analysis in the SM

- SM UT analysis:
  - provide the best determination of CKM parameters
  - test the consistency of the SM (“*direct*” vs “*indirect*” determinations)
  - provide predictions (from data..) for SM observables

## .. and beyond

- NP UT analysis:
  - model-independent analysis
  - provides limit on the allowed deviations from the SM
  - obtain the NP scale



[www.utfit.org](http://www.utfit.org)

A. Bevan, M.B., M. Ciuchini, D. Derkach,  
E. Franco, V. Lubicz, G. Martinelli, F. Parodi,  
M. Pierini, C. Schiavi, L. Silvestrini, A. Stocchi,  
V. Sordini, C. Tarantino and V. Vagnoni

Other UT analyses exist, by:

CKMfitter (<http://ckmfitter.in2p3.fr/>),

Laiho&Lunghi&Van de Water (<http://latticeaverages.org/>)

Lunghi&Soni (1010.6069)

the method and the inputs:

$$f(\bar{\rho}, \bar{\eta}, X | c_1, \dots, c_m) \sim \prod_{j=1, m} f_j(C | \bar{\rho}, \bar{\eta}, X) * \prod_{i=1, N} f_i(x_i) f_0(\bar{\rho}, \bar{\eta})$$

Bayes Theorem

$$X \equiv x_1, \dots, x_n = m_t, B_K, F_B, \dots$$

$$C \equiv c_1, \dots, c_m = \epsilon, \Delta m_d / \Delta m_s, A_{CP}(J/\psi K_S), \dots$$

$$(b \rightarrow u) / (b \rightarrow c)$$

$$\bar{\rho}^2 + \bar{\eta}^2$$

$$\bar{\Lambda}, \lambda_1, F(1), \dots$$

$$\epsilon_K$$

$$\bar{\eta}[(1 - \bar{\rho}) + P]$$

$$B_K$$

$$\Delta m_d$$

$$(1 - \bar{\rho})^2 + \bar{\eta}^2$$

$$f_B^2 B_B$$

$$\Delta m_d / \Delta m_s$$

$$(1 - \bar{\rho})^2 + \bar{\eta}^2$$

$$\xi$$

$$A_{CP}(J/\psi K_S)$$

$$\sin 2\beta$$

Standard Model +  
OPE/HQET/  
Lattice QCD  
to go  
from quarks  
to hadrons

M. Bona *et al.* (UTfit Collaboration)  
JHEP 0507:028,2005 hep-ph/0501199

M. Bona *et al.* (UTfit Collaboration)  
JHEP 0603:080,2006 hep-ph/0509219

# $V_{cb}$ and $V_{ub}$

$$V_{cb} (excl) = (39.19 \pm 0.70) 10^{-3}$$

$$V_{cb} (incl) = (42.20 \pm 0.70) 10^{-3}$$

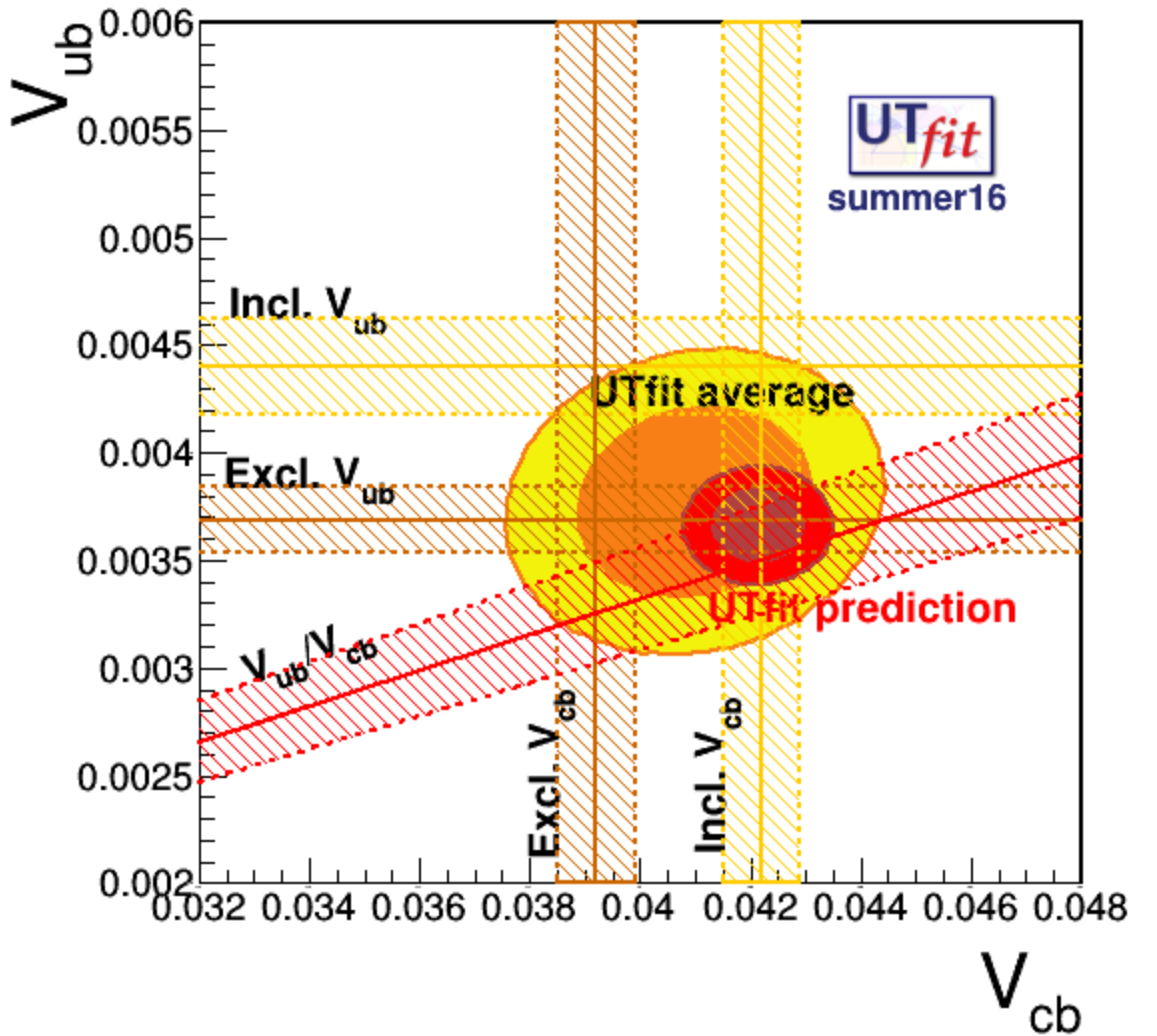
$\sim 3.0\sigma$  discrepancy

$$V_{ub} (excl) = (3.69 \pm 0.14) 10^{-3}$$

$$V_{ub} (incl) = (4.40 \pm 0.22) 10^{-3}$$

$\sim 2.7\sigma$  discrepancy

$$V_{ub} / V_{cb} (LHCb) = (8.3 \pm 0.6) 10^{-2}$$



# $V_{cb}$ and $V_{ub}$

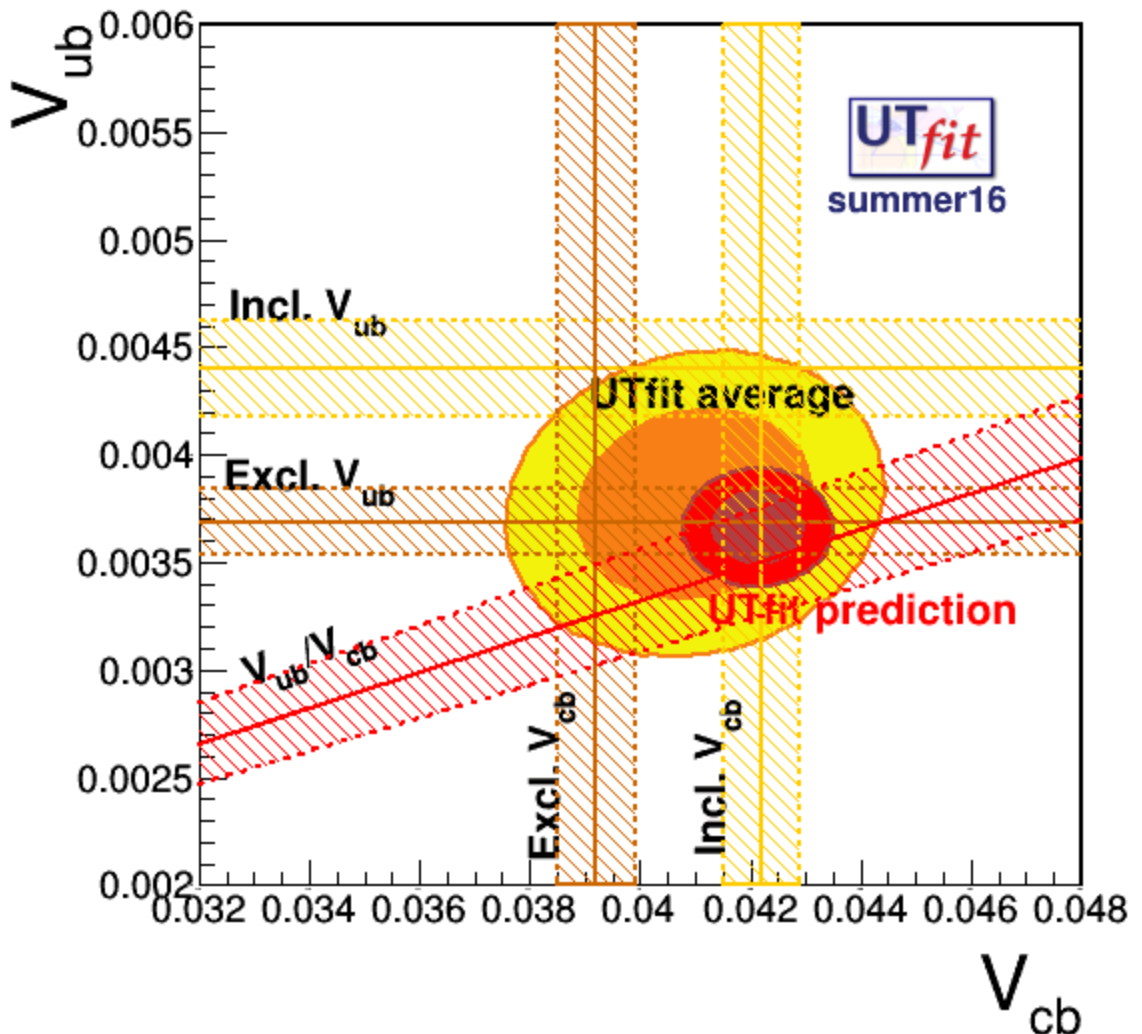
Average inspired by D'Agostini skeptical procedure (hep-ex/9910036). However very similar results are obtained from a 2D a la PDG procedure.

$$V_{cb} = (41.0 \pm 1.4) 10^{-3}$$

uncertainty  $\sim 3.4\%$

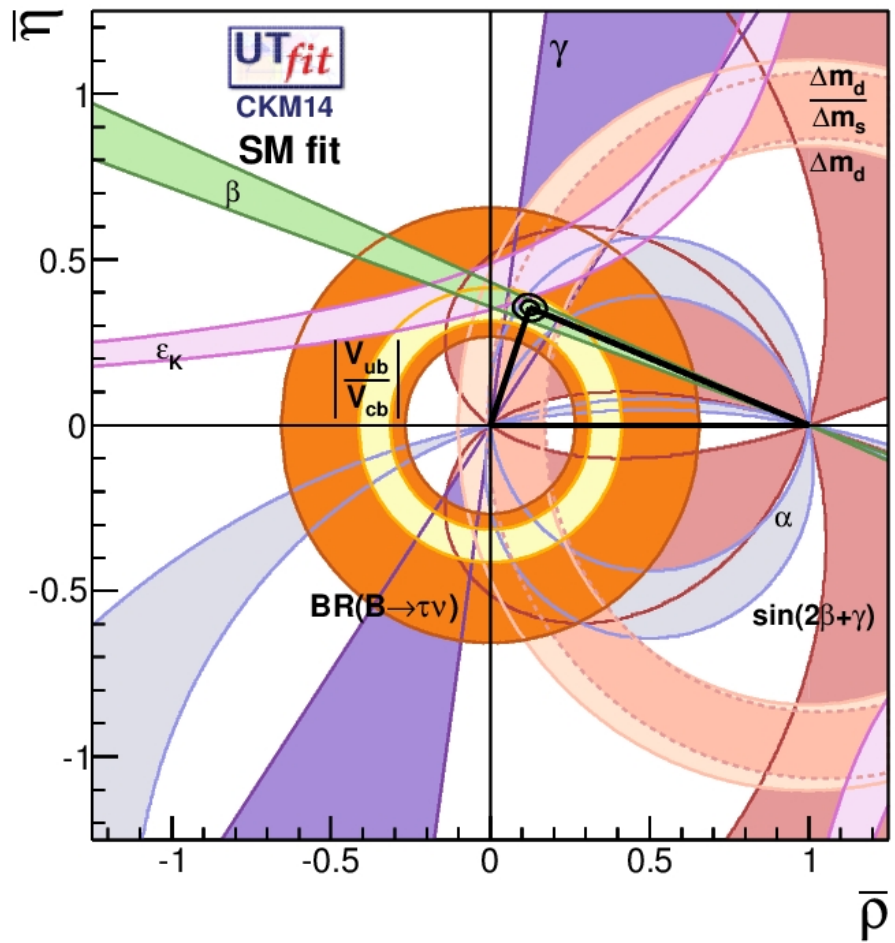
$$V_{ub} = (3.77 \pm 0.29) 10^{-3}$$

uncertainty  $\sim 7.7\%$

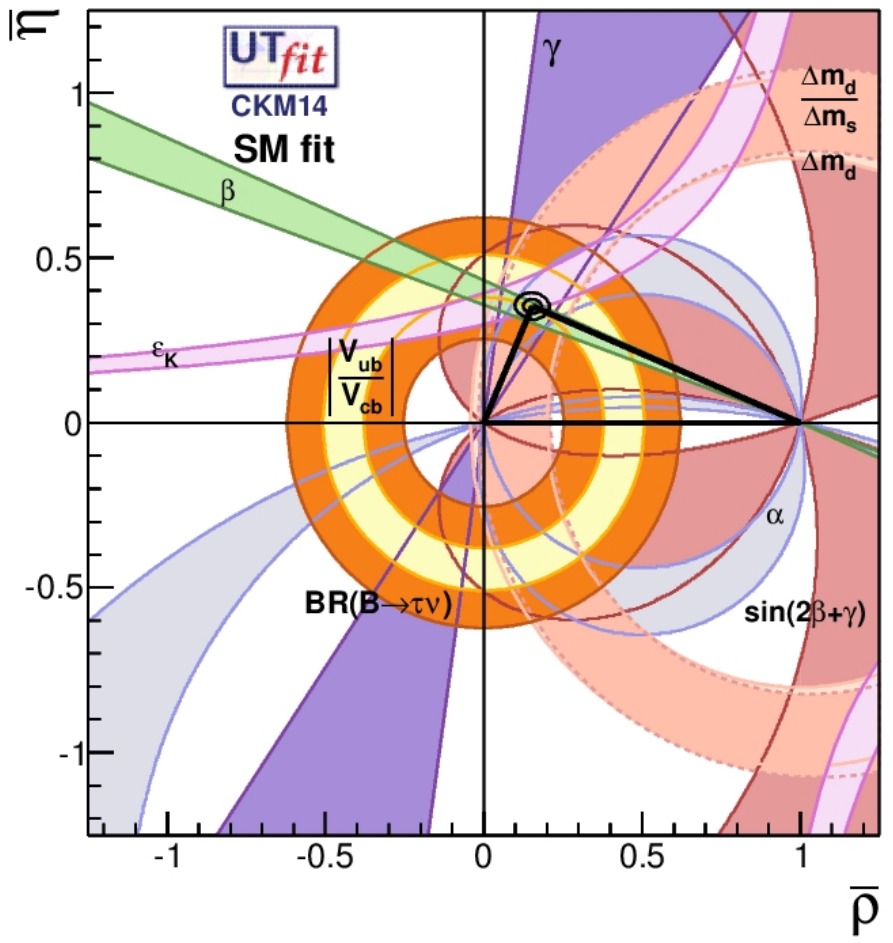


# exclusives vs inclusives

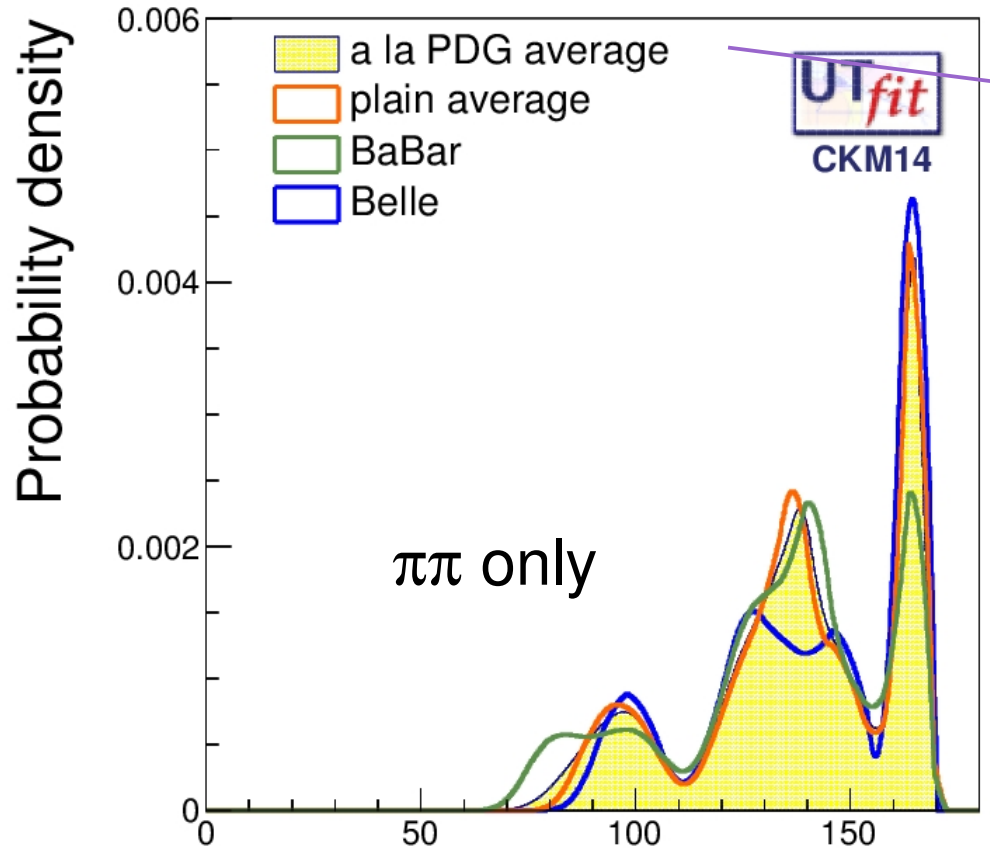
only exclusive values



only inclusive values



# sin2α (ϕ<sub>2</sub>) from charmless B decays: ρρ, (ρρ, πρ)

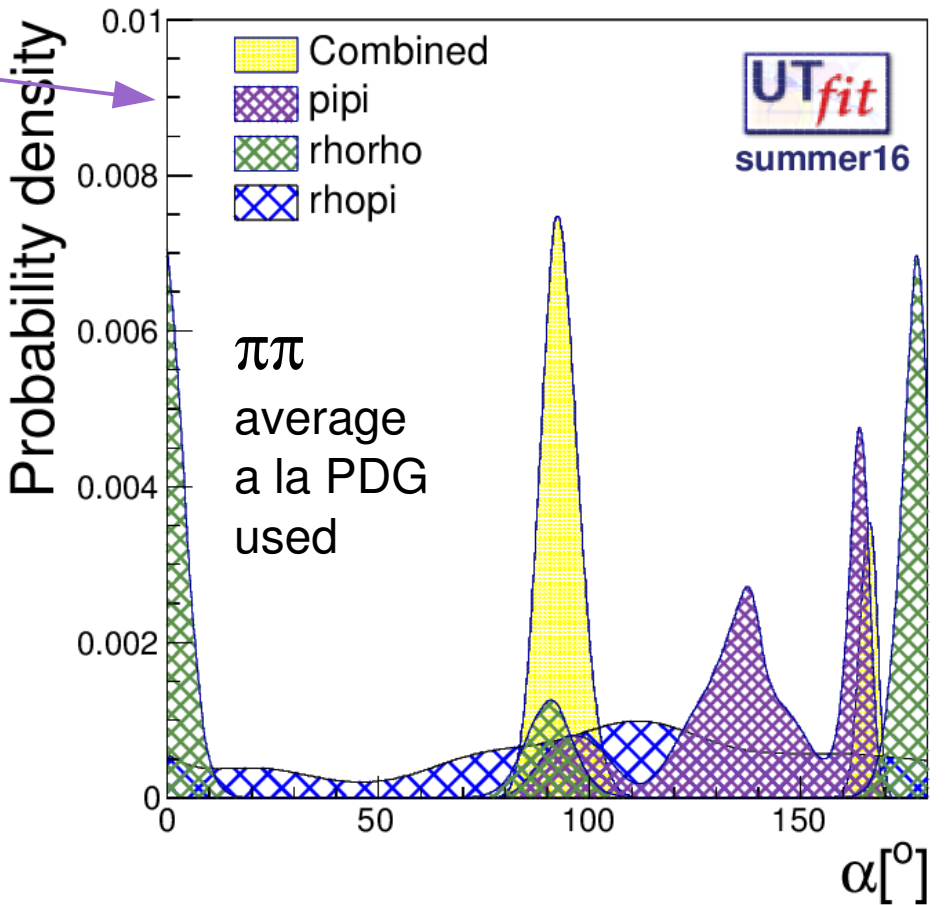


newest  $\pi^0\pi^0$  from Belle  $\alpha[^\circ]$

à la PDG average inflates the error

$$\text{BR}(\pi^0\pi^0) = (1.15 \pm 0.41) 10^{-6}$$

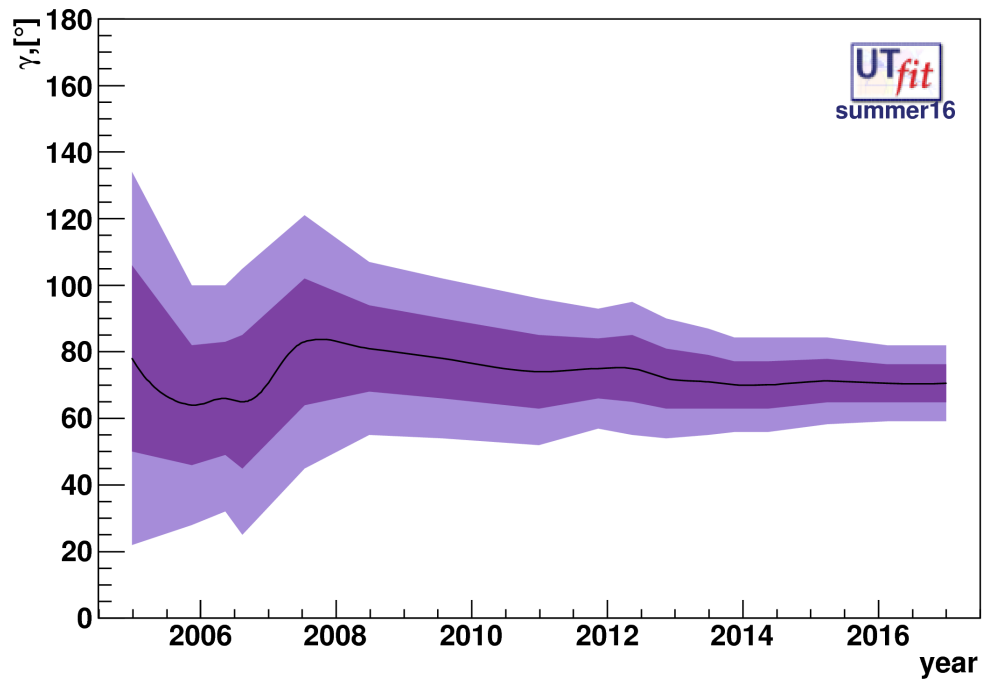
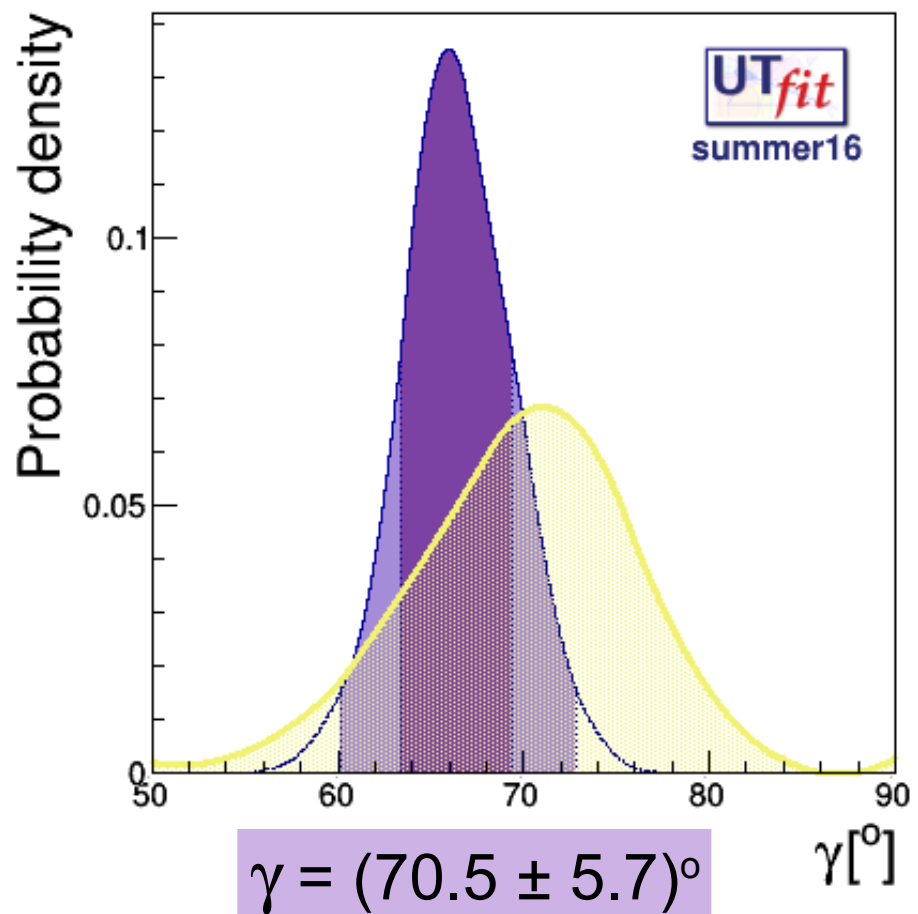
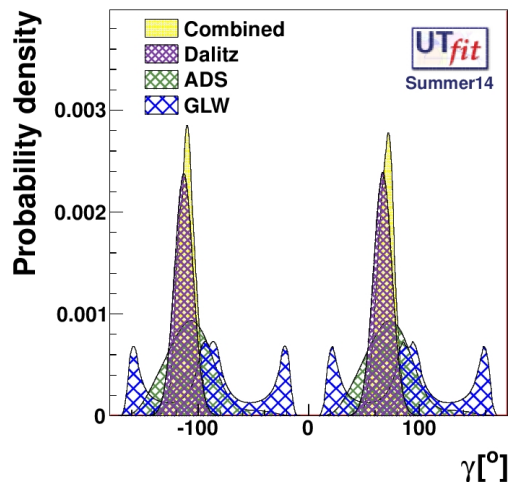
$$\text{wrt plain average: } (1.15 \pm 0.13) 10^{-6}$$



$\alpha$  from  $\pi\pi$ ,  $\rho\rho$ ,  $\pi\rho$  decays:  
combined:  $(92.5 \pm 5.5)^\circ$

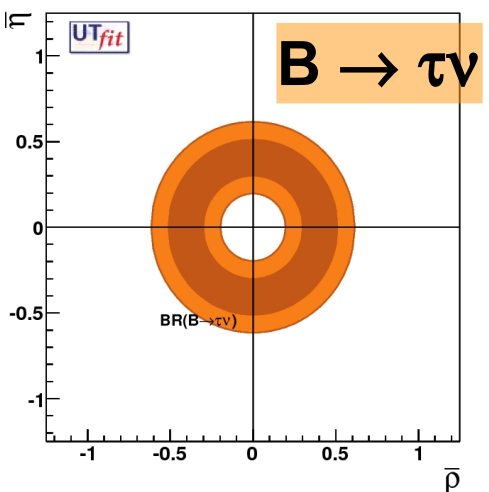
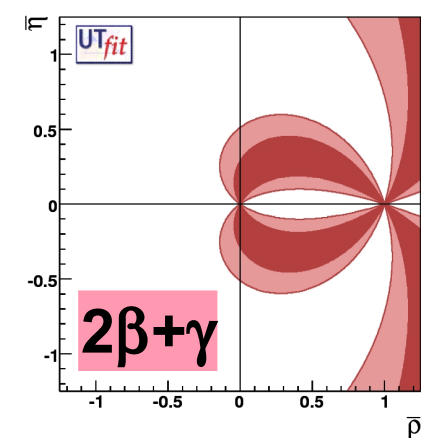
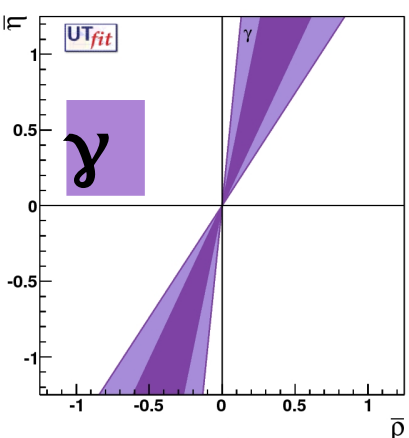
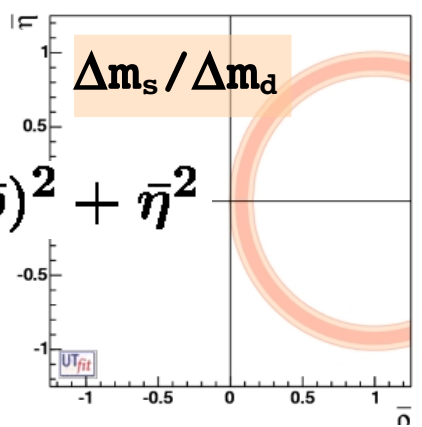
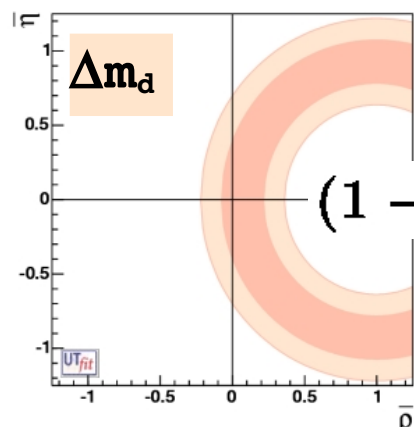
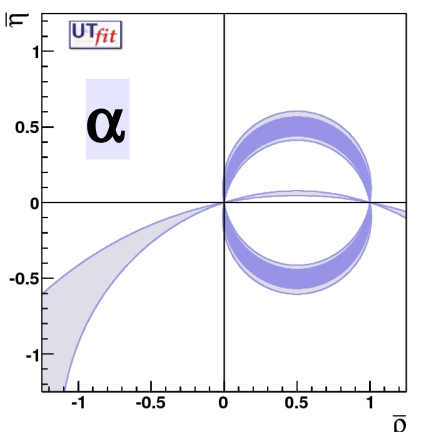
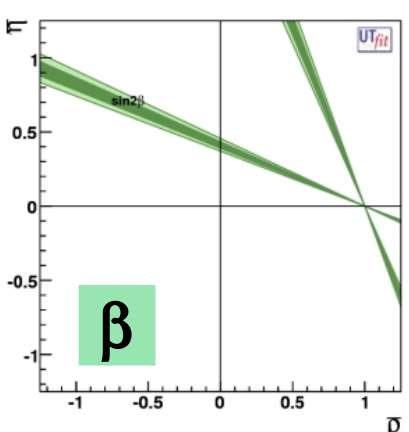
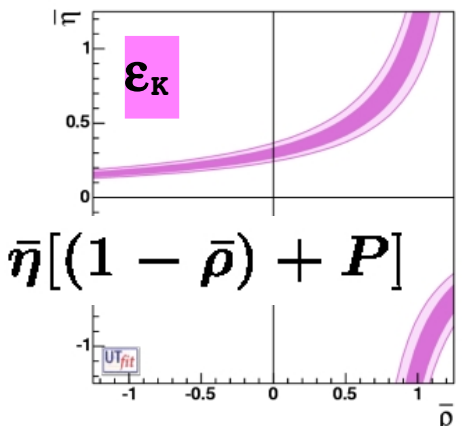
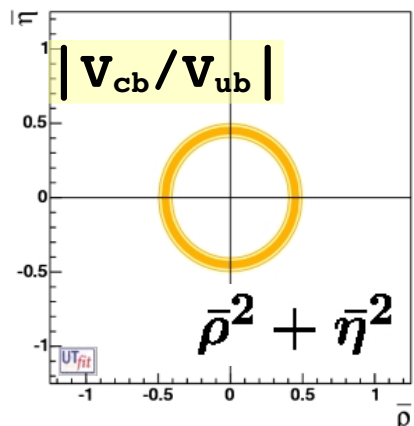


# $\gamma$ and DK trees

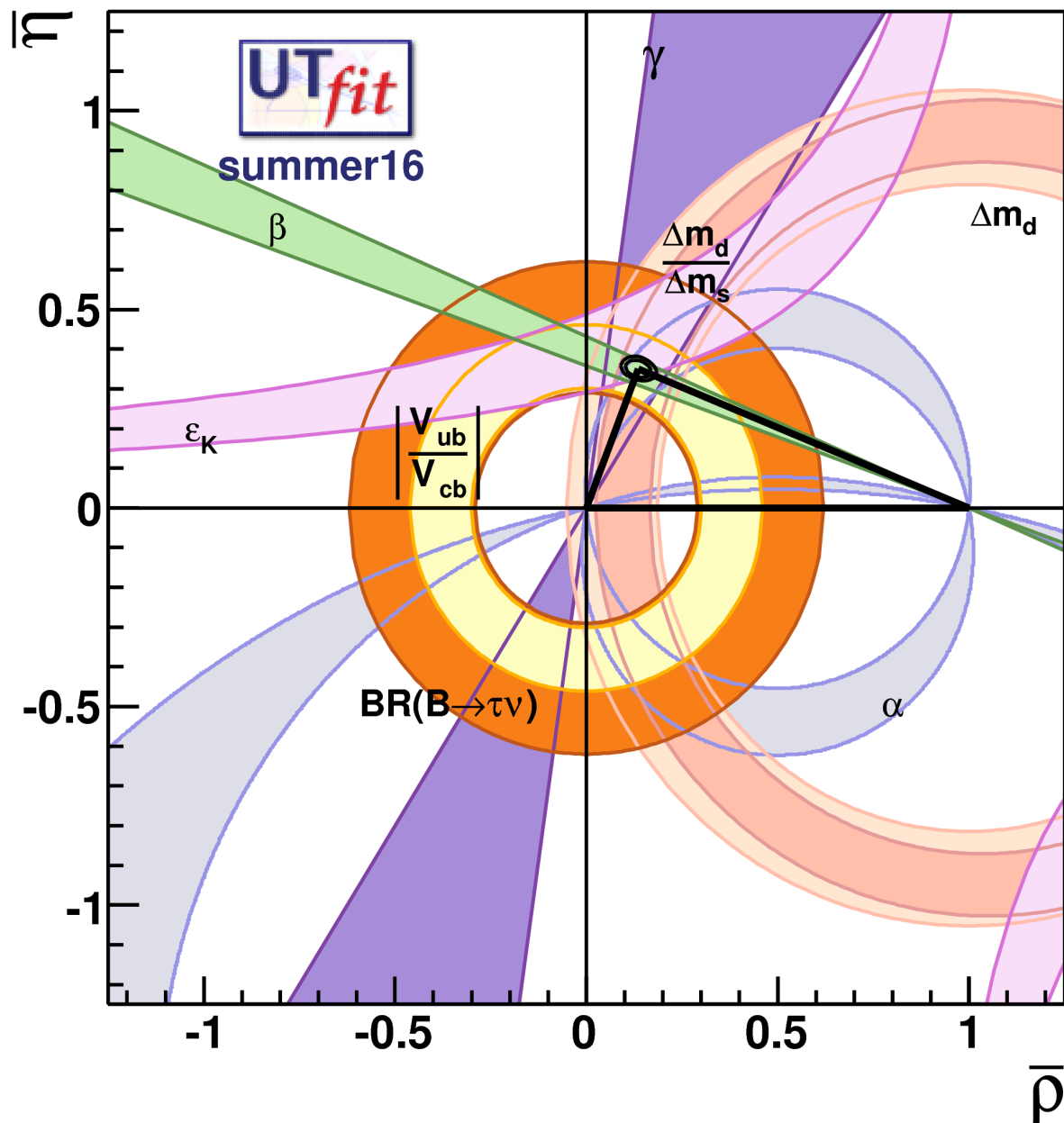


After a decade of analyses and almost 50 papers published, the world average uncertainty has decreased by a factor 3

# Unitarity Triangle analysis in the SM:



# Unitarity Triangle analysis in the SM:

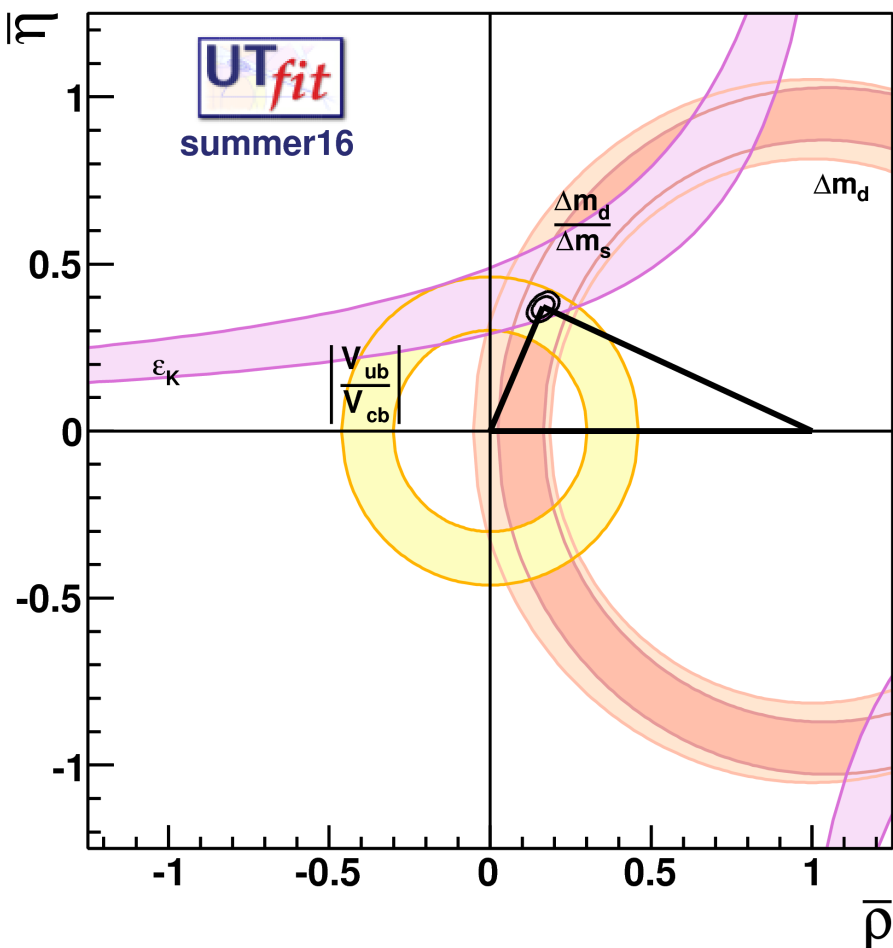
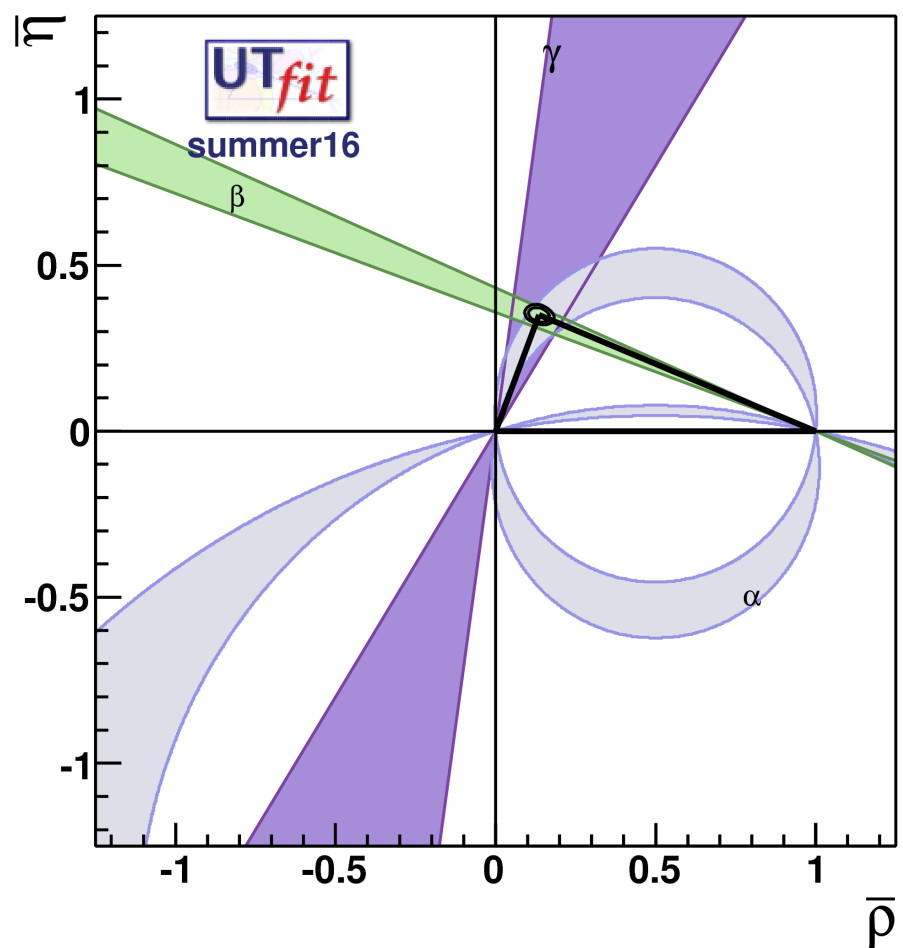


levels @  
95% Prob

$$\begin{aligned} \bar{\rho} &= 0.144 \pm 0.018 \\ \bar{\eta} &= 0.346 \pm 0.012 \end{aligned}$$

# CP violating vs CP conserving

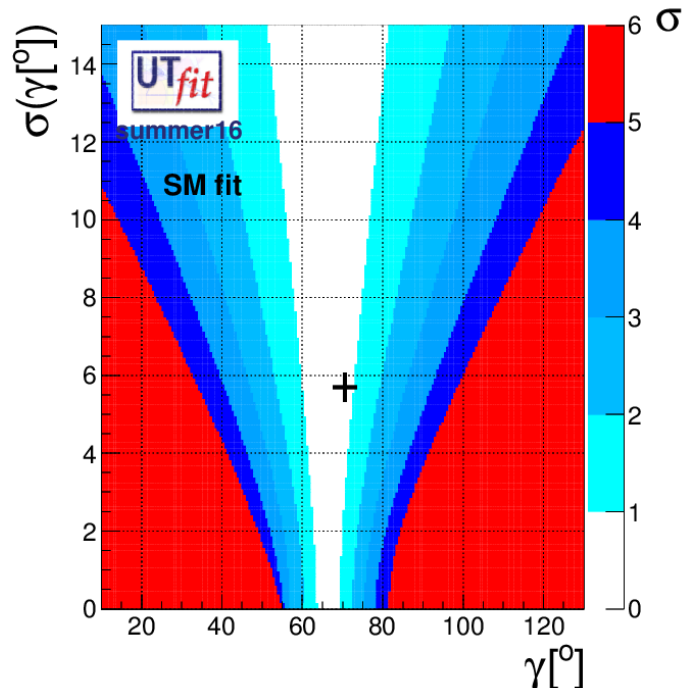
levels @  
95% Prob



# compatibility plots

A way to “measure” the agreement of a single measurement with the indirect determination from the fit using all the other inputs: test for the SM description of the flavour physics

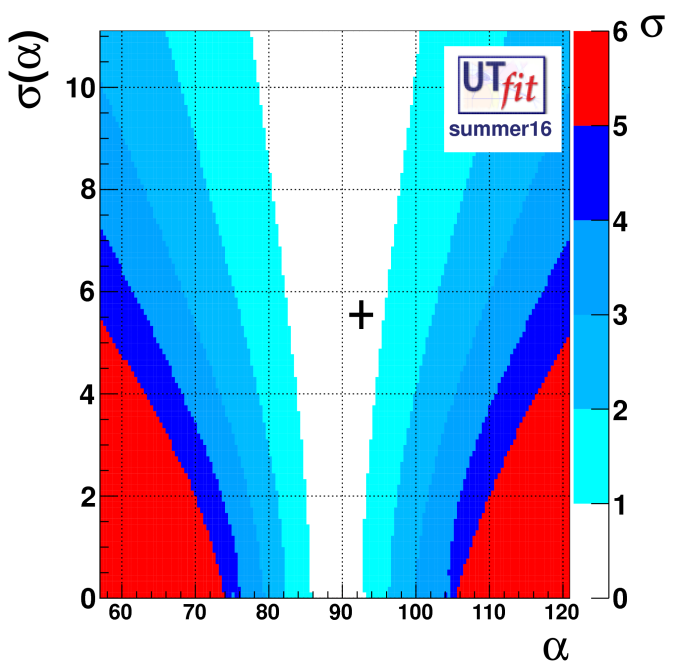
Color code: agreement between the predicted values and the measurements at better than 1, 2, ... $n\sigma$



$$\gamma_{\text{exp}} = (70.5 \pm 5.7)^\circ$$

$$\gamma_{\text{UTfit}} = (66.3 \pm 3.0)^\circ$$

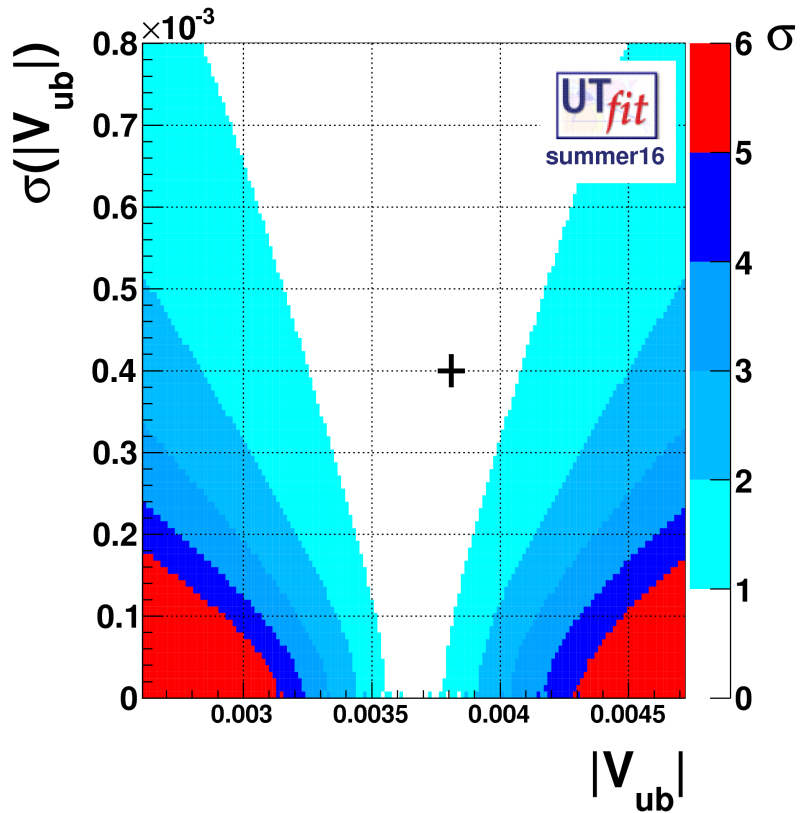
The cross has the coordinates  $(x,y)=(\text{central value, error})$  of the direct measurement



$$\alpha_{\text{exp}} = (92.5 \pm 5.5)^\circ$$

$$\alpha_{\text{UTfit}} = (88.9 \pm 3.5)^\circ$$

tensions? not any more..



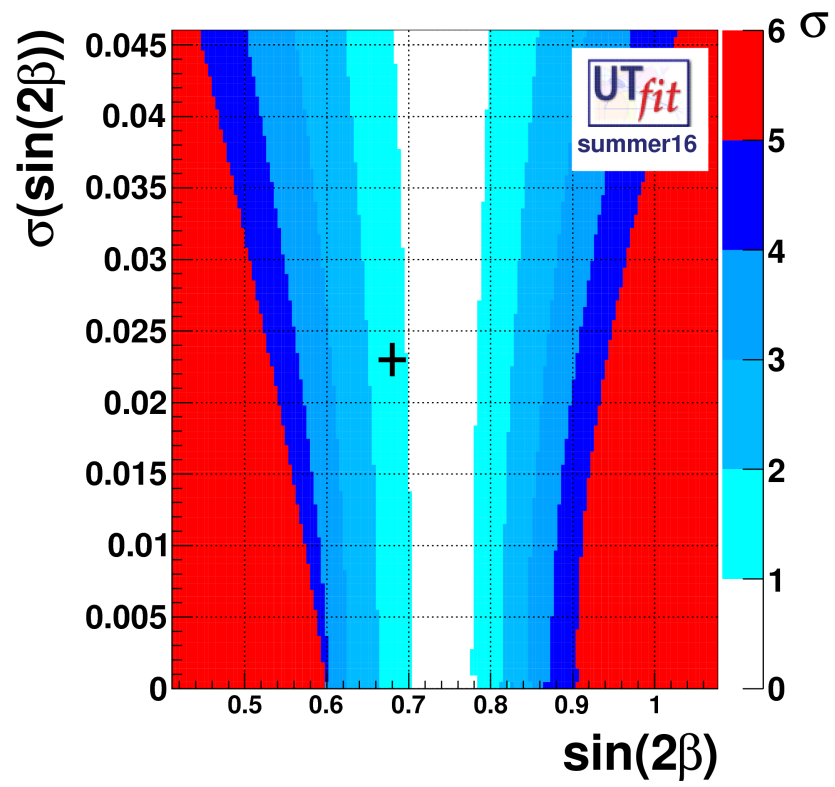
$$V_{ub_{\text{exp}}} = (3.77 \pm 0.29) \cdot 10^{-3}$$

$$V_{ub_{\text{UTfit}}} = (3.66 \pm 0.11) \cdot 10^{-3}$$

$\sim 1.4\sigma$

$$\sin 2\beta_{\text{exp}} = 0.680 \pm 0.023$$

$$\sin 2\beta_{\text{UTfit}} = 0.740 \pm 0.037$$



# Unitarity Triangle analysis in the SM:

obtained excluding  
the given constraint  
from the fit

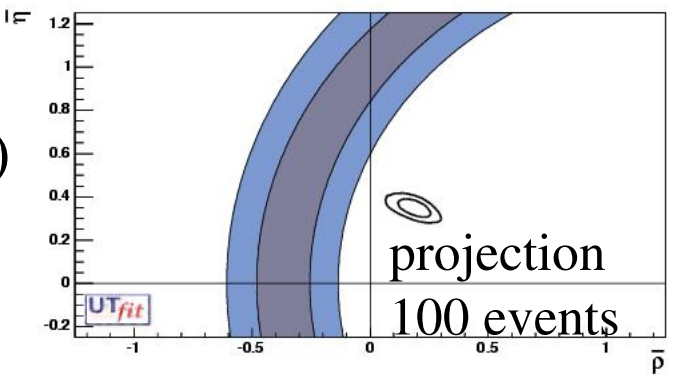
Observables	Measurement	Prediction	Pull ( $\# \sigma$ )
$\sin 2\beta$	$0.680 \pm 0.023$	$0.740 \pm 0.037$	$\sim 1.4$
$\gamma$	$70.5 \pm 5.7$	$66.3 \pm 3.0$	$< 1$
$\alpha$	$92.2 \pm 6.2$	$87.3 \pm 3.9$	$< 1$
$ V_{ub}  \cdot 10^3$	$3.77 \pm 0.29$	$3.66 \pm 0.11$	$< 1$
$ V_{ub}  \cdot 10^3$ (incl)	$4.40 \pm 0.22$	—	$\sim 3.0$
$ V_{ub}  \cdot 10^3$ (excl)	$3.69 \pm 0.14$	—	$< 1$
$ V_{cb}  \cdot 10^3$	$41.0 \pm 1.4$	$42.3 \pm 0.6$	$< 1$
$\beta_s$	$0.97 \pm 0.94$	$1.05 \pm 0.04$	$< 1$
$B_K$	$0.766 \pm 0.010$	$0.832 \pm 0.081$	$< 1$
$BR(B \rightarrow \tau \nu)[10^{-4}]$	$1.06 \pm 0.20$	$0.81 \pm 0.06$	$\sim 1.3$
$A_{SL}^d \cdot 10^3$	$-1.5 \pm 1.7$	$-0.283 \pm 0.024$	$< 1$
$A_{SL}^s \cdot 10^3$	$-7.5 \pm 4.1$	$0.013 \pm 0.001$	$\sim 1.8$

# some old plots coming back to fashion:

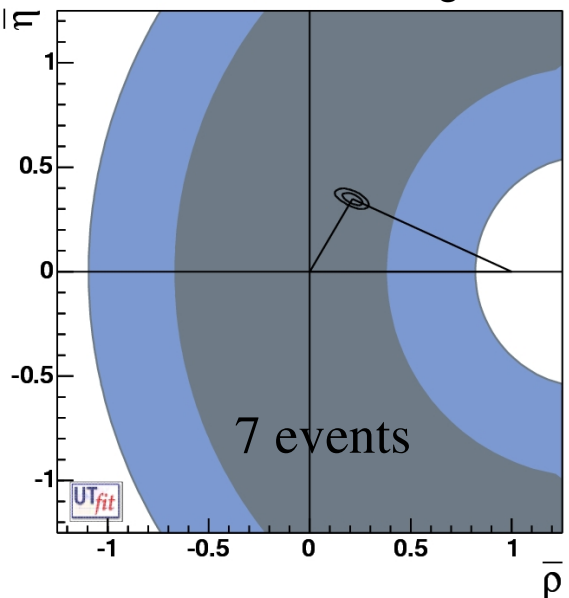
As NA62 and KOTO are approaching data taking:

$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$$

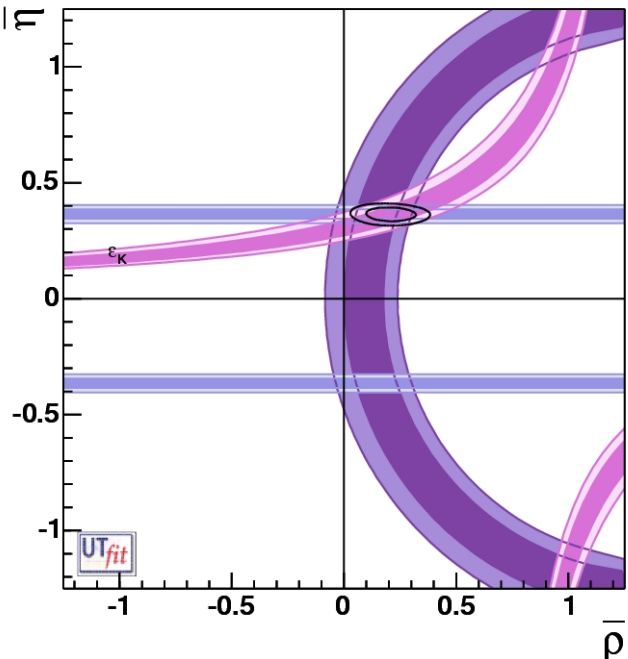
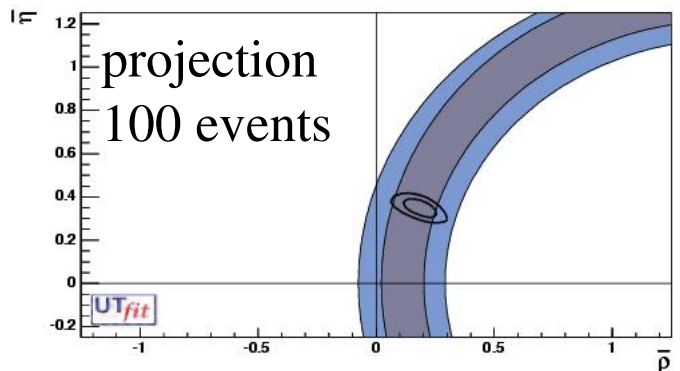
E949 central value



2007 global fit area



SM central value



including  
 $\text{BR}(K^0 \rightarrow \pi^0 \nu \bar{\nu})$   
 SM central value



# UT analysis including new physics

fit simultaneously for the CKM and the NP parameters (generalized UT fit)

- ▶ add most general loop NP to all sectors
- ▶ use all available experimental info
- ▶ find out NP contributions to  $\Delta F=2$  transitions

$B_d$  and  $B_s$  mixing amplitudes (2+2 real parameters):

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left( 1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

$$\Delta m_{q/K} = C_{B_q/\Delta m_K} (\Delta m_{q/K})^{SM}$$

$$A_{CP}^{B_d \rightarrow J/\psi K_s} = \sin 2(\beta + \phi_{B_d})$$

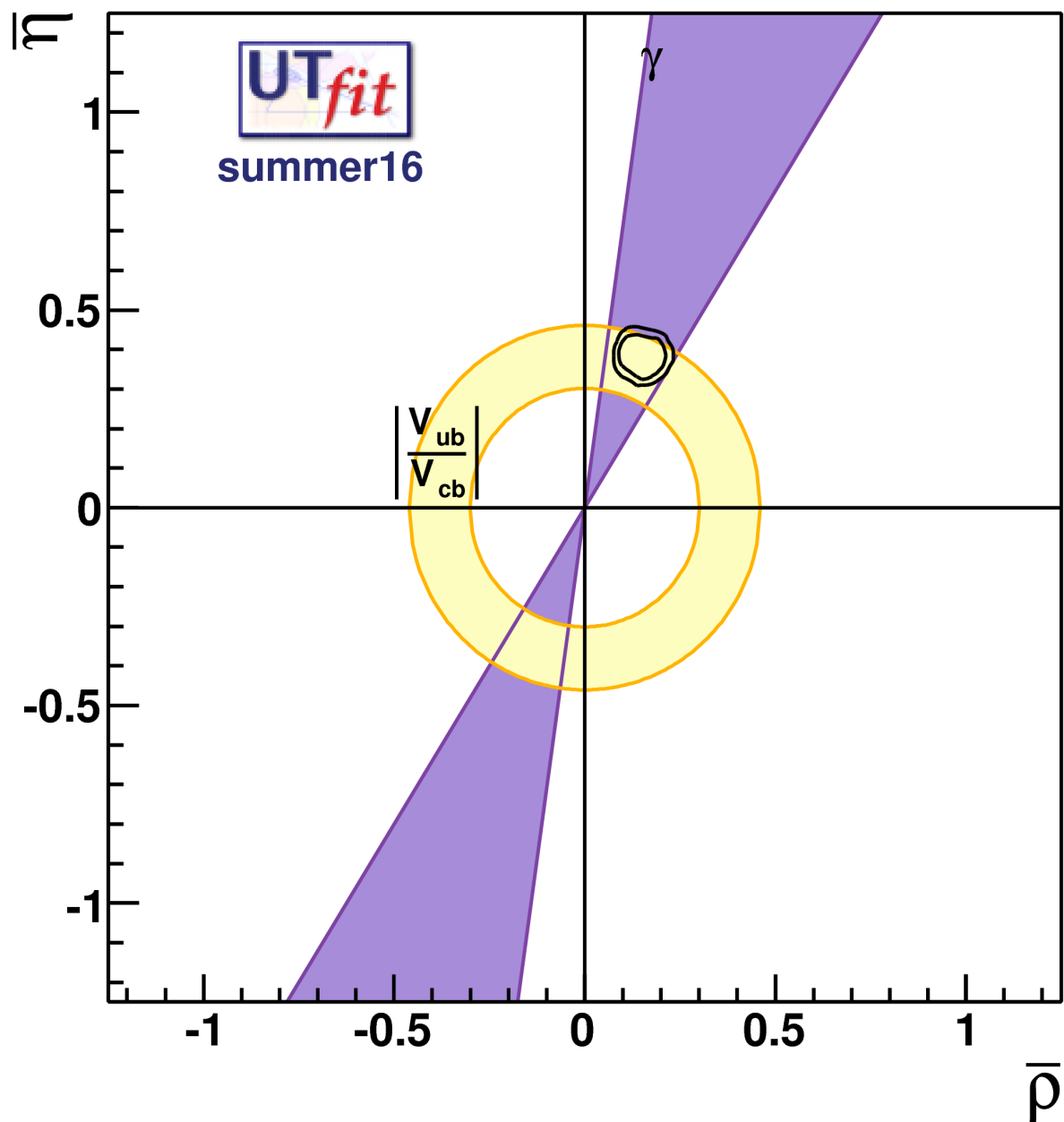
$$A_{SL}^q = \text{Im}(\Gamma_{12}^q / A_q)$$

$$\varepsilon_K = C_\varepsilon \varepsilon_K^{SM}$$

$$A_{CP}^{B_s \rightarrow J/\psi \phi} \sim \sin 2(-\beta_s + \phi_{B_s})$$

$$\Delta \Gamma^q / \Delta m_q = \text{Re}(\Gamma_{12}^q / A_q)$$

# NP analysis results



$$\begin{aligned} \bar{\rho} &= 0.151 \pm 0.040 \\ \bar{\eta} &= 0.384 \pm 0.037 \end{aligned}$$

**SM is**

$$\begin{aligned} \bar{\rho} &= 0.144 \pm 0.018 \\ \bar{\eta} &= 0.346 \pm 0.012 \end{aligned}$$

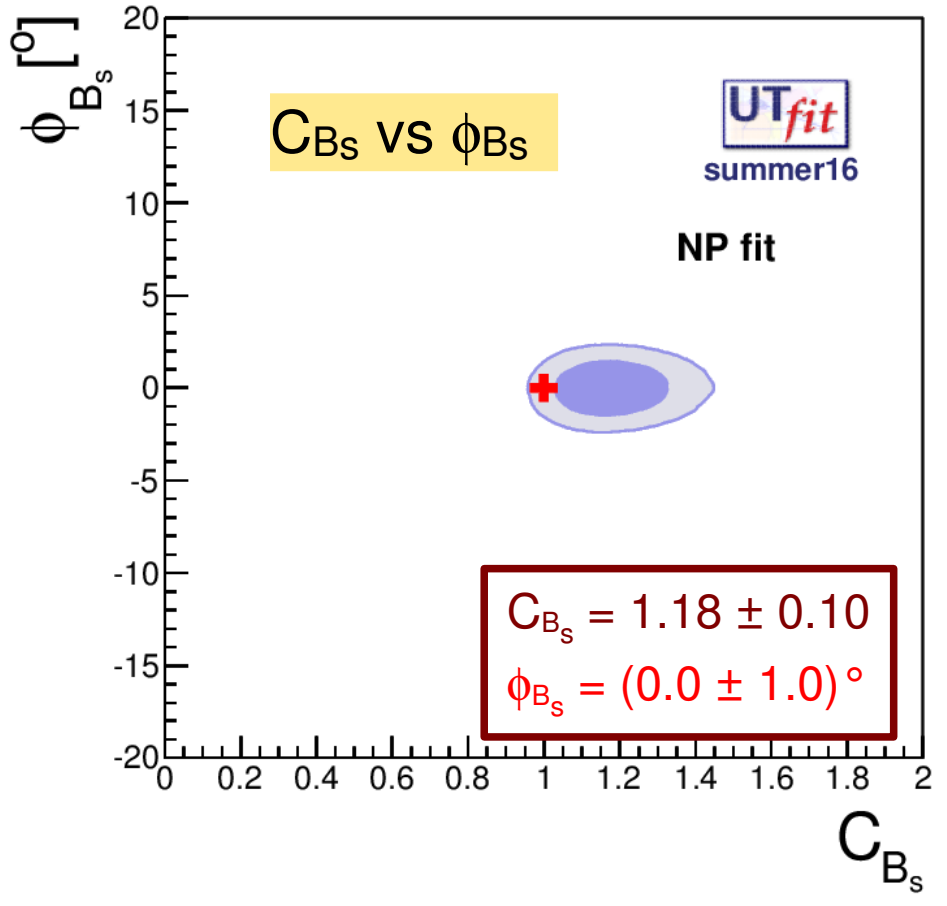
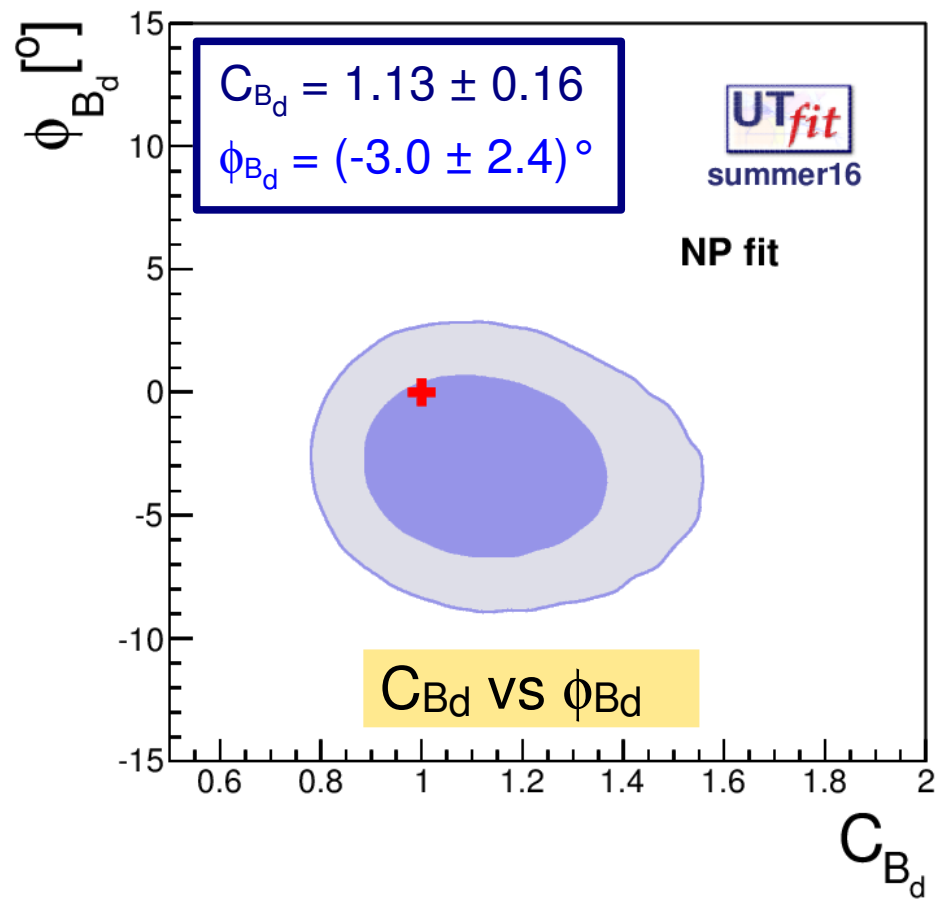
# NP parameter results

dark: 68%  
light: 95%  
SM: red cross

K system

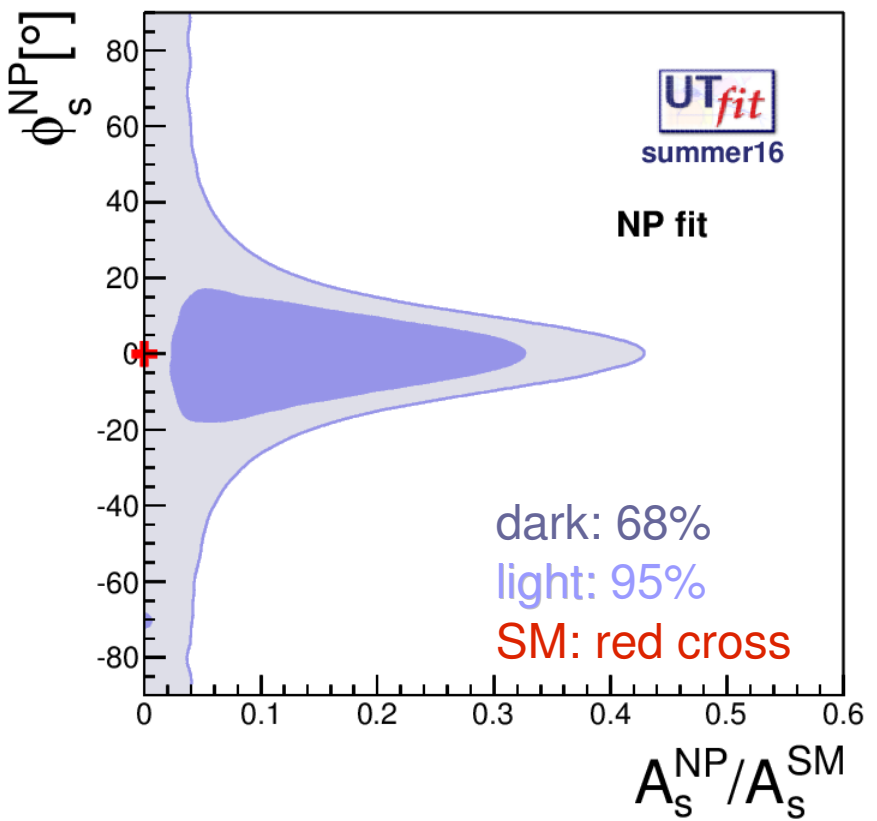
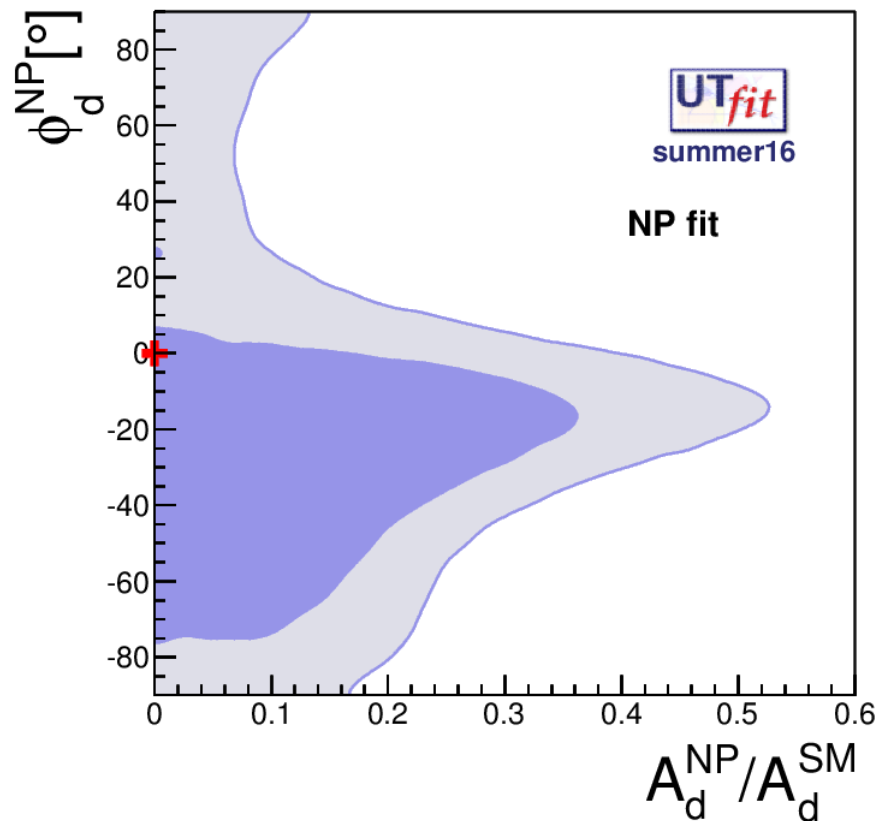
$$C_{\epsilon_K} = 1.14 \pm 0.16$$

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}}$$



# NP parameter results

$$A_q = \left( 1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\varphi_q^{NP} - \varphi_q^{SM})} \right) A_q^{SM} e^{2i\varphi_q^{SM}}$$



The ratio of NP/SM amplitudes is:

< 35% @68% prob. (50% @95%) in  $B_d$  mixing

< 30% @68% prob. (40% @95%) in  $B_s$  mixing

see also Lunghi & Soni, Buras et al., Ligeti et al.

## testing the new-physics scale

R  
G  
E

### At the high scale

new physics enters according to its specific features

### At the low scale

use OPE to write the most general effective Hamiltonian.

the operators have different chiralities than the SM

NP effects are in the Wilson Coefficients  $C$

### NP effects are enhanced

- up to a factor 10 by the values of the matrix elements especially for transitions among quarks of different chiralities
- up to a factor 8 by RGE

$$\mathcal{H}_{\text{eff}}^{\Delta B=2} = \sum_{i=1}^5 C_i Q_i^{bq} + \sum_{i=1}^3 \tilde{C}_i \tilde{Q}_i^{bq}$$

$$Q_1^{q_i q_j} = \bar{q}_{jL}^{\alpha} \gamma_{\mu} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} \gamma^{\mu} q_{iL}^{\beta},$$

$$Q_2^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jR}^{\beta} q_{iL}^{\beta},$$

$$Q_3^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jR}^{\beta} q_{iL}^{\alpha},$$

$$Q_4^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} q_{iR}^{\beta},$$

$$Q_5^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jL}^{\beta} q_{iR}^{\alpha}.$$

M. Bona *et al.* (UTfit)  
JHEP 0803:049,2008  
arXiv:0707.0636

## effective BSM Hamiltonian for $\Delta F=2$ transitions

The Wilson coefficients  $C_i$  have in general the form

$$C_i(\Lambda) = \frac{F_i L_i}{\Lambda^2}$$

Putting bounds on the Wilson coefficients give insights into the NP scale in different NP scenarios that enter through  $F_i$  and  $L_i$

- $F_i$ : function of the NP flavour couplings
- $L_i$ : loop factor (in NP models with no tree-level FCNC)
- $\Lambda$ : NP scale (typical mass of new particles mediating  $\Delta F=2$  transitions)

## testing the TeV scale

The dependence of  $C$  on  $\Lambda$  changes on flavor structure.

We can consider different flavour scenarios:

- **Generic:**  $C(\Lambda) = \alpha/\Lambda^2$   $F_i \sim 1$ , arbitrary phase
- **NMFV:**  $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$   $F_i \sim |F_{SM}|$ , arbitrary phase
- **MFV:**  $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$   $F_1 \sim |F_{SM}|$ ,  $F_{i \neq 1} \sim 0$ , SM phase

$\alpha (L_i)$  is the coupling among NP and SM

- ⊙  $\alpha \sim 1$  for strongly coupled NP
- ⊙  $\alpha \sim \alpha_w (\alpha_s)$  in case of loop coupling through **weak** (**strong**) interactions

If no NP effect is seen  
lower bound on NP scale  $\Lambda$   
if NP is seen  
upper bound on NP scale  $\Lambda$

$F$  is the flavour coupling and so

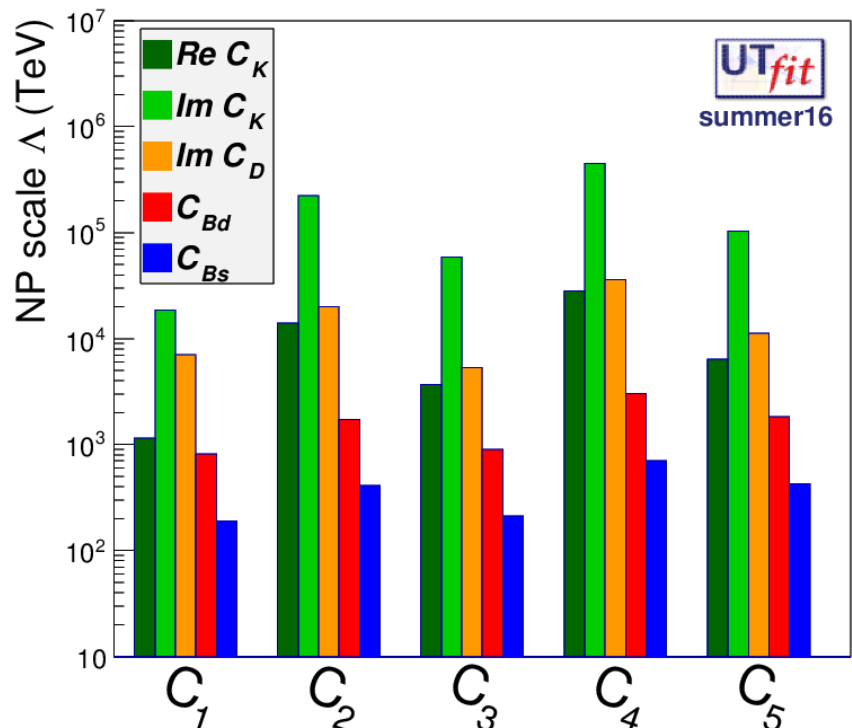
$F_{SM}$  is the combination of CKM factors for the considered process

$$C_i(\Lambda) = \frac{L_i}{\Lambda^2}$$

# results from the Wilson coefficients

**Generic:**  $C(\Lambda) = \alpha/\Lambda^2$ ,  $F_i \sim 1$ , arbitrary phase

$\alpha \sim 1$  for strongly coupled NP



Lower bounds on NP scale  
(in TeV at 95% prob.)

**Non-perturbative NP**  
 $\Lambda > 4.5 \cdot 10^5 \text{ TeV}$

To obtain the lower bound for loop-mediated contributions, one simply multiplies the bounds by  $\alpha_s$  ( $\sim 0.1$ ) or by  $\alpha_w$  ( $\sim 0.03$ ).

$\alpha \sim \alpha_w$  in case of loop coupling through **weak** interactions

NP in  $\alpha_w$  loops  
 $\Lambda > 1.4 \cdot 10^4 \text{ TeV}$

Best bound from  $\epsilon_K$   
dominated by CKM error  
CPV in charm mixing follows,  
exp error dominant

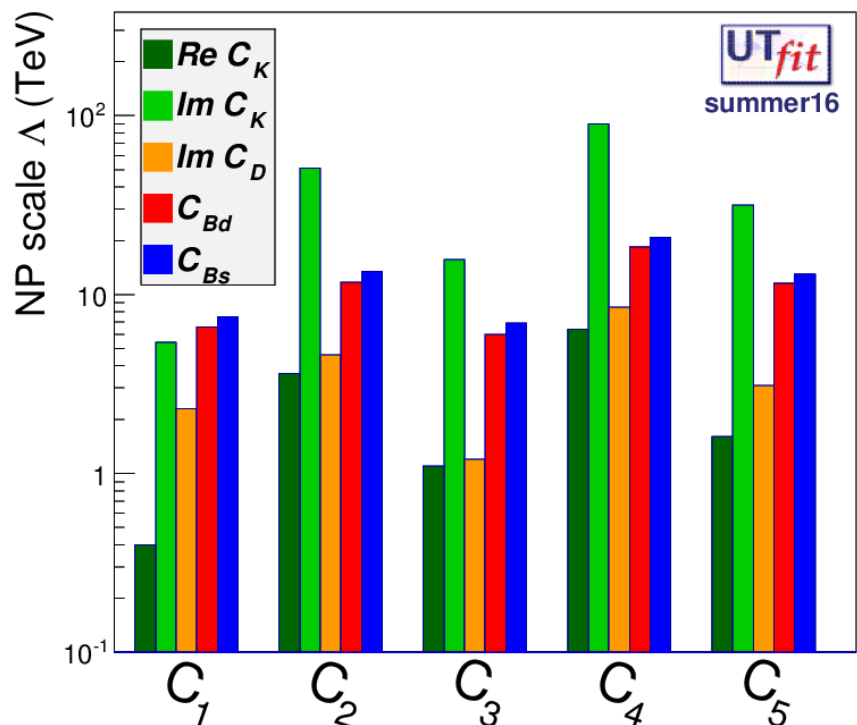
Best CP conserving from  $\Delta m_K$ ,  
dominated by long distance  
 $B_d$  and  $B_s$  behind,  
errors from both CKM  
and B-parameters



# results from the Wilson coefficients

**NMFV:**  $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ ,  $F_i \sim |F_{SM}|$ , arbitrary phase

$\alpha \sim 1$  for strongly coupled NP



Lower bounds on NP scale  
(in TeV at 95% prob.)

**Non-perturbative NP**  
 $\Lambda > 90$  TeV

To obtain the lower bound for loop-mediated contributions, one simply multiplies the bounds by  $\alpha_s$  ( $\sim 0.1$ ) or by  $\alpha_w$  ( $\sim 0.03$ ).

$\alpha \sim \alpha_w$  in case of loop coupling through **weak** interactions

NP in  $\alpha_w$  loops  
 $\Lambda > 2.7$  TeV

If new chiral structures present,  
 $\epsilon_K$  still leading  
 $B_{(s)}$  mixing provides very stringent constraints, especially if no new chiral structures are present  
Constraining power of the various sectors depends on unknown NP flavour structure.

## Look at the future

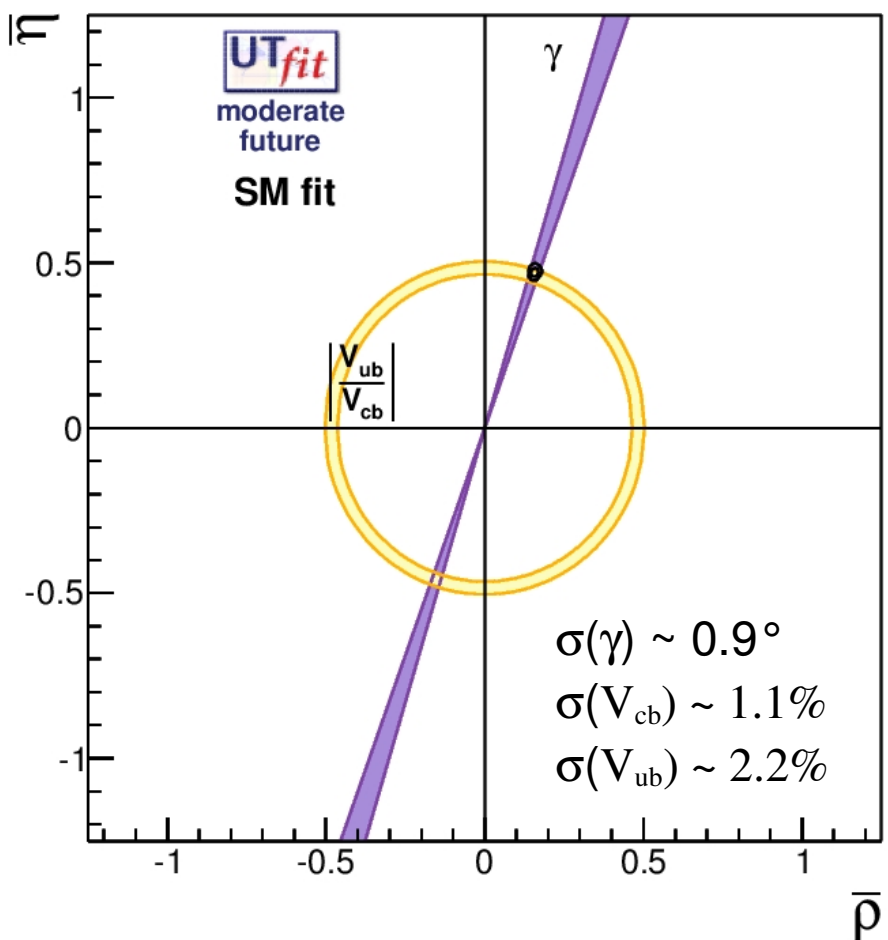
In the next decade, Belle-II and LHCb upgrade will push down the exp. error on  $\sin 2\beta(s)$  to less than 0.01

Theory error can be kept below 0.01 using control channels as  $S(B \rightarrow J/\psi\pi)$

B-parameters will go below the % level, new ideas to attack long-distance in K and D

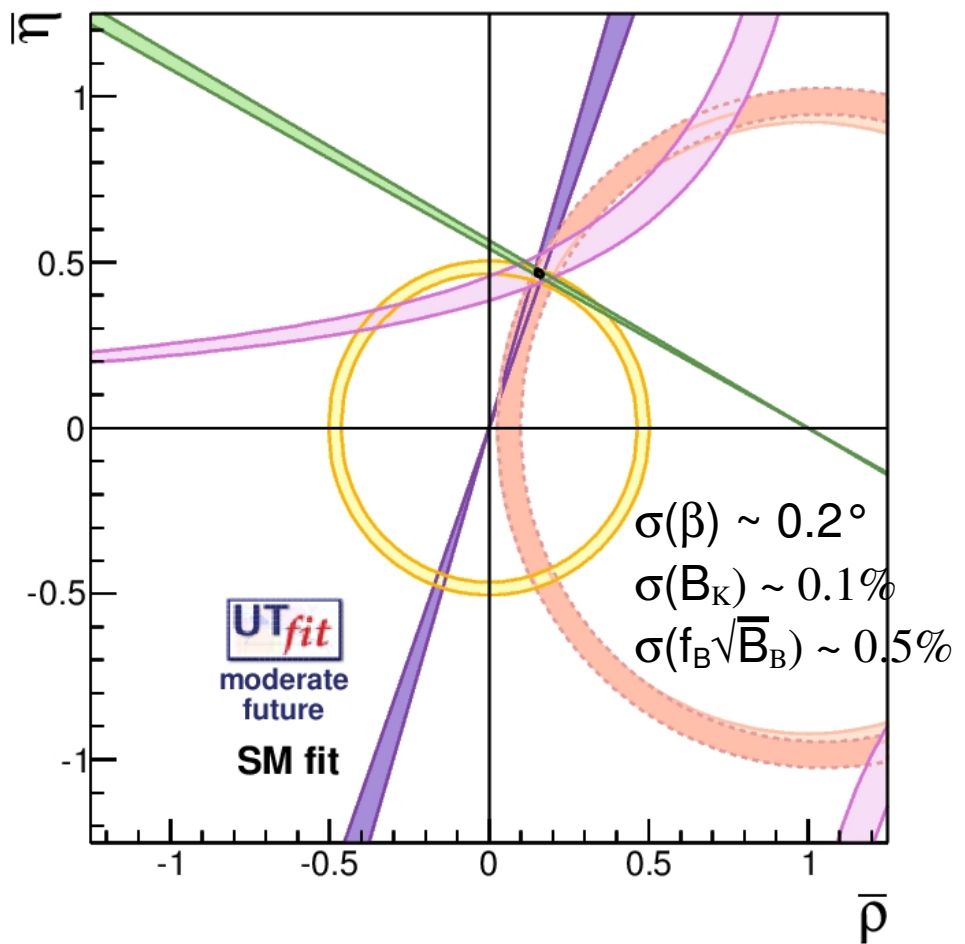
Improving  $\gamma$ ,  $\alpha$  and  $|V_{cb}|$  &  $|V_{ub}|$  crucial!

# Look at the future



errors from tree-only fit on  $\rho$  and  $\eta$ :  
 $\sigma(\rho) = 0.008$  [currently 0.040]  
 $\sigma(\eta) = 0.010$  [currently 0.037]

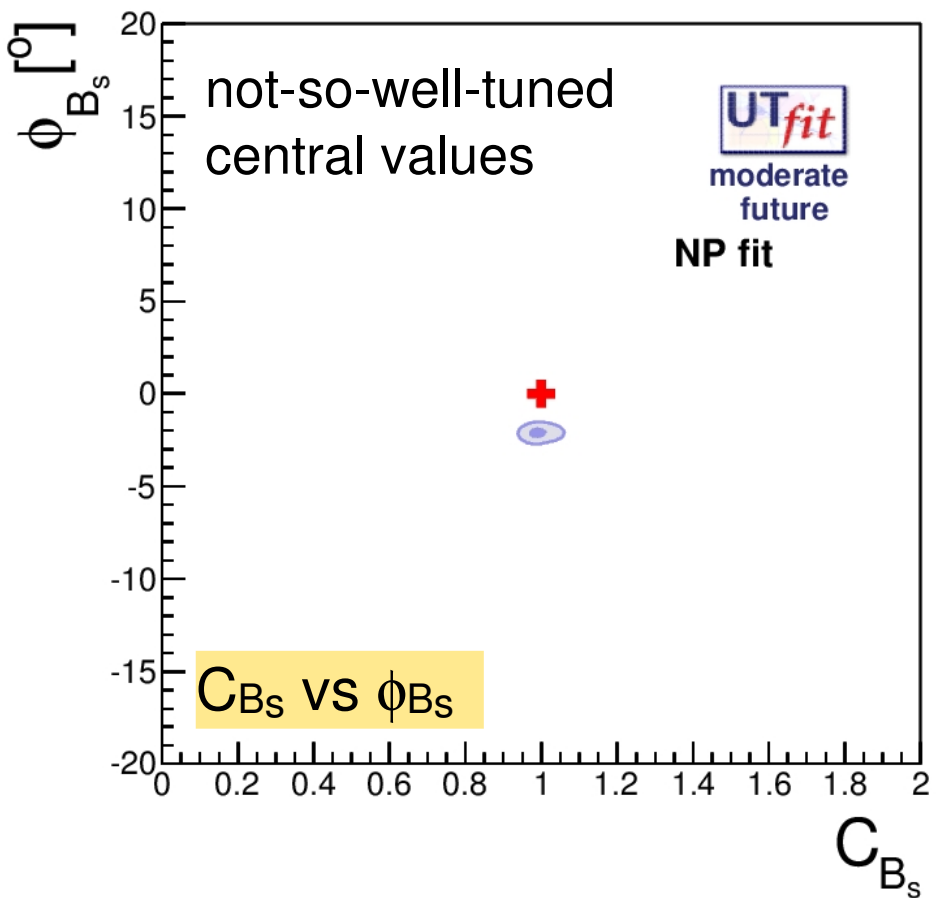
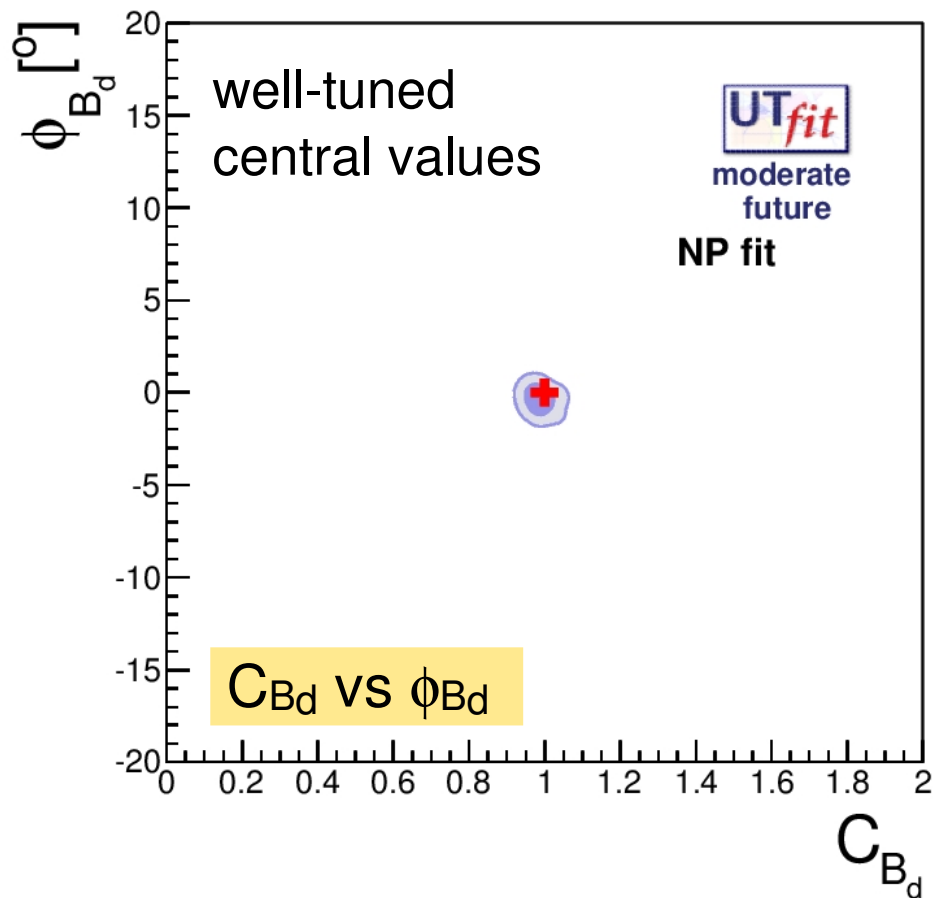
errors predicted from  
 Belle II + LHCb upgrade



errors from 5-constraint fit on  $\rho$  and  $\eta$ :  
 $\sigma(\rho) = 0.005$  [currently 0.018]  
 $\sigma(\eta) = 0.004$  [currently 0.012]

# Look at the future

errors predicted from  
Belle II + LHCb upgrade



errors on general NP parameters:

$$\sigma(C_{B_d}) = 0.03 \text{ [currently } 0.16\text{]}$$

$$\sigma(\phi_{B_d}) = 0.7 \text{ [currently } 3.2\text{]}$$

$$\sigma(C_{B_s}) = 0.03 \text{ [currently } 0.08\text{]}$$

$$\sigma(\phi_{B_s}) = 0.6 \text{ [currently } 2.0\text{]}$$

Parameter	Error				
	Now	50/fb	300/fb	1000/fb	3000/fb
$\bar{\rho}$ (SM fit)	0.002	0.0039	0.0023	0.0013	0.00064
$\bar{\eta}$ (SM fit)	0.021	0.0037	0.0019	0.0013	0.00068
$\gamma$ [°] (SM fit)	6.5	0.6	0.35	0.2	0.09
$\alpha$ [°] (SM fit)	5.5	0.6	0.37	0.2	0.1
$\beta$ [°] (SM fit)	4	0.2	0.10	0.07	0.04
$\beta_s$ [°] (SM fit)	4	0.011	0.057	0.004	0.0023
$\bar{\rho}$ (NP fit)	0.002	0.006	0.0034	0.0028	0.0022
$\bar{\eta}$ (NP fit)	0.021	0.006	0.0053	0.0061	0.0052
$\gamma$ [°] (NP fit)	6.5	0.9	0.4	0.2	0.09
$\alpha$ [°] (NP fit)	5.5	1	0.5	0.45	0.36
$\beta$ [°] (NP fit)	4	0.8	0.7	0.7	0.7
$\beta_s$ [°] (NP fit)	4	0.017	0.016	0.016	0.016
$C_{\varepsilon_K}$	0.14	0.065	0.065	0.065	0.064
$C_{B_d}$	0.15	0.024	0.024	0.024	0.022
$\Phi_{B_d}$	2.8	0.48	0.36	0.36	0.35
$C_{B_s}$	0.087	0.02	0.02	0.02	0.02
$\Phi_{B_s}$	0.96	0.26	0.11	0.063	0.038
$\Phi_{M_{12}}$ [°]	2.5	0.4	0.1	0.08	0.04
$\Phi_{\Gamma_{12}}$ [°]	—	1.2	0.4	0.24	0.12

←

←

Crucial to improve SM predictions of rare decays!

←

←

Need

← progress in

←  $|V_{ub}|$  and

←  $|V_{cb}|$

← Steady

← improvement

Very preliminary!!!



# conclusions

- ▶ SM analysis displays good overall consistency
- ▶ Still open discussion on semileptonic inclusive vs exclusive
- ▶ UTA provides determination also of NP contributions to  $\Delta F=2$  amplitudes. It currently leaves space for NP at the level of 30-40%
- ▶ So the scale analysis points to high scales for the generic scenario and at the limit of LHC reach for weak coupling. Indirect searches still essential.
- ▶ Even if we don't see relevant deviations in the down sector, we might still find them in the up sector.



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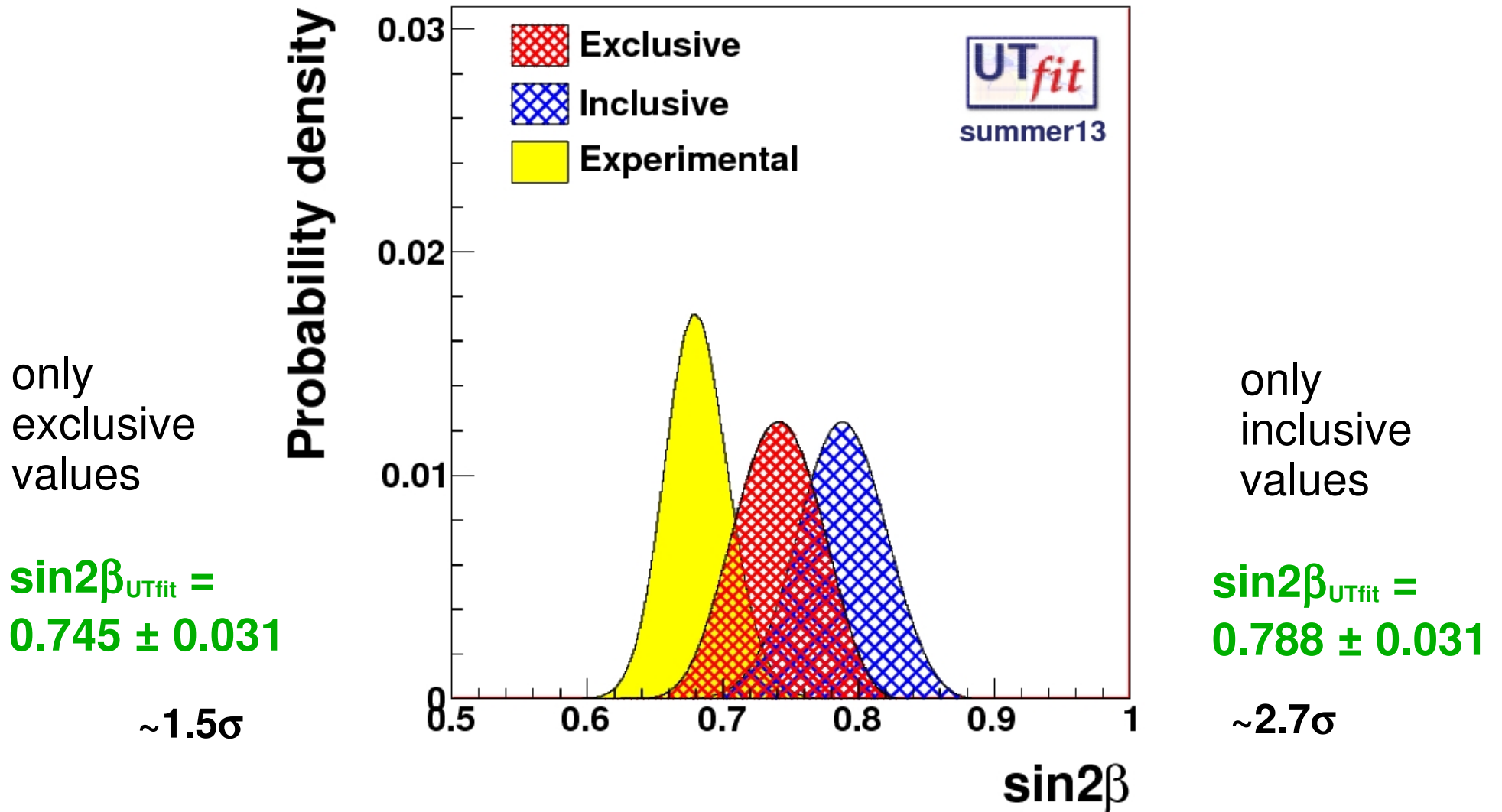
**BRUICHLADDICH**



**Back up slides**



# inclusives vs exclusives



# new-physics-specific constraints

**B meson mixing matrix element NLO calculation**  
 Ciuchini et al. JHEP 0308:031,2003.

$C_{\text{pen}}$  and  $\phi_{\text{pen}}$  are  
 parameterize possible  
 NP contributions from  
 $b \rightarrow s$  penguins

$$\frac{\Gamma_{12}^q}{A_q^{\text{full}}} = -2 \kappa C_{B_q} \left\{ e^{2\phi_{B_q}} \left( n_1 + \frac{n_6 B_2 + n_{11}}{B_1} \right) - \frac{e^{(\phi_q^{\text{SM}} + 2\phi_{B_q})}}{R_t^q} \left( n_2 + \frac{n_7 B_2 + n_{12}}{B_1} \right) \right. \\
 + \frac{e^{2(\phi_q^{\text{SM}} + \phi_{B_q})}}{R_t^{q^2}} \left( n_3 + \frac{n_8 B_2 + n_{13}}{B_1} \right) + e^{(\phi_q^{\text{Pen}} + 2\phi_{B_q})} C_q^{\text{Pen}} \left( n_4 + n_9 \frac{B_2}{B_1} \right) \\
 \left. - e^{(\phi_q^{\text{SM}} + \phi_q^{\text{Pen}} + 2\phi_{B_q})} C_q^{\text{Pen}} \left( n_5 + n_{10} \frac{B_2}{B_1} \right) \right\}$$

$\phi_s = 2\beta_s$  vs  $\Delta\Gamma_s$  from  $B_s \rightarrow J/\psi\phi$

angular analysis as a function of proper time  
 and b-tagging  
 additional sensitivity from the  $\Delta\Gamma_s$  terms

$\phi_s$  and  $\Delta\Gamma_s$ :  
 2D experimental likelihood from CDF and D0

$\phi_s$  and  $\Delta\Gamma_s$ :  
 central values with  
 gaussian errors from LHCb

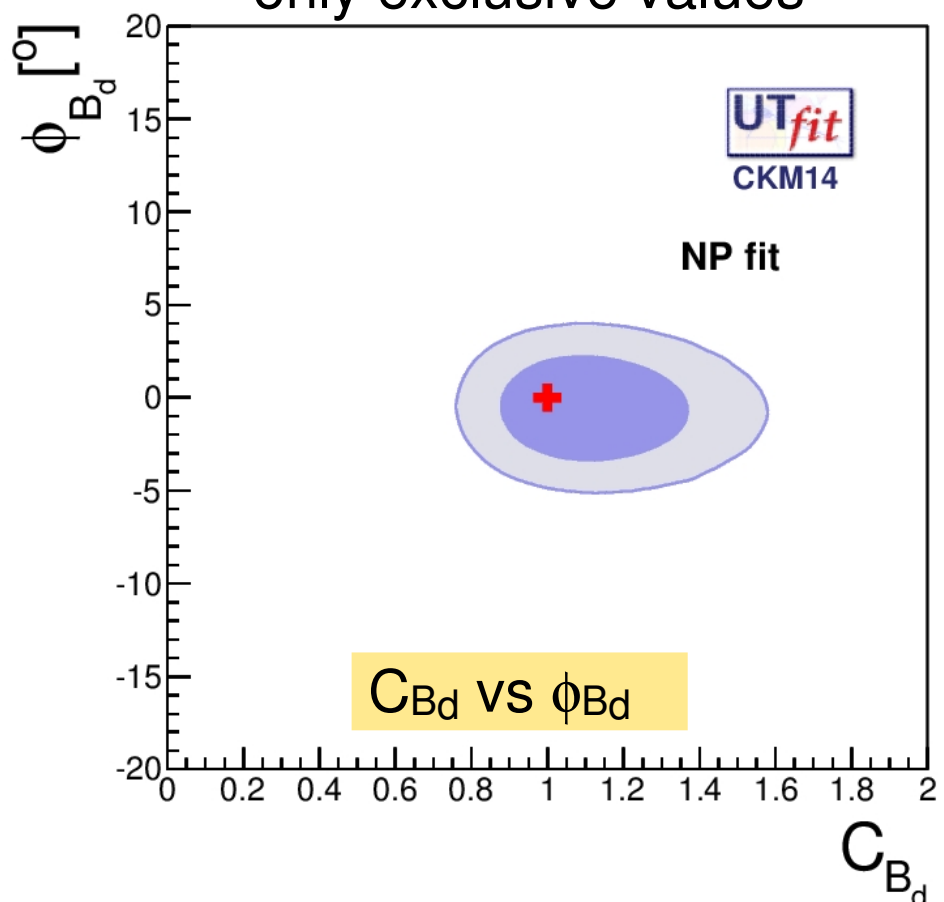
# NP parameter results: exclusives vs inclusives

dark: 68%

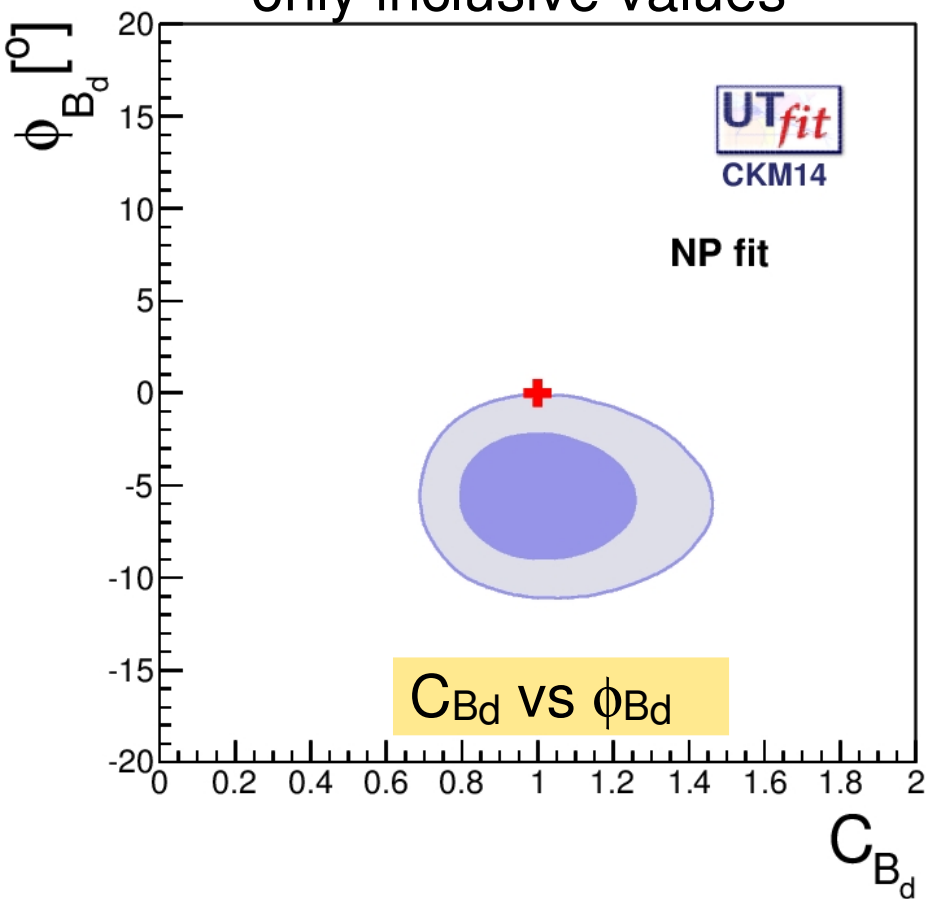
light: 95%

SM: red cross

### only exclusive values



### only inclusive values



## contribution to the mixing amplitudes

analytic expression for the contribution to the mixing amplitudes

$$\langle \bar{B}_q | \mathcal{H}_{\text{eff}}^{\Delta B=2} | B_q \rangle_i = \sum_{j=1}^5 \sum_{r=1}^5 \left( b_j^{(r,i)} + \eta c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda) \langle \bar{B}_q | Q_r^{bq} | B_q \rangle$$

Lattice QCD

*arXiv:0707.0636: for "magic numbers"  $a, b$  and  $c$ ,  $\eta = \alpha_s(\Lambda)/\alpha_s(m_t)$   
(numerical values updated last in summer'12)*

analogously for the K system

$$\langle \bar{K}^0 | \mathcal{H}_{\text{eff}}^{\Delta S=2} | K^0 \rangle_i = \sum_{j=1}^5 \sum_{r=1}^5 \left( b_j^{(r,i)} + \eta c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda) R_r \langle \bar{K}^0 | Q_1^{sd} | K^0 \rangle$$

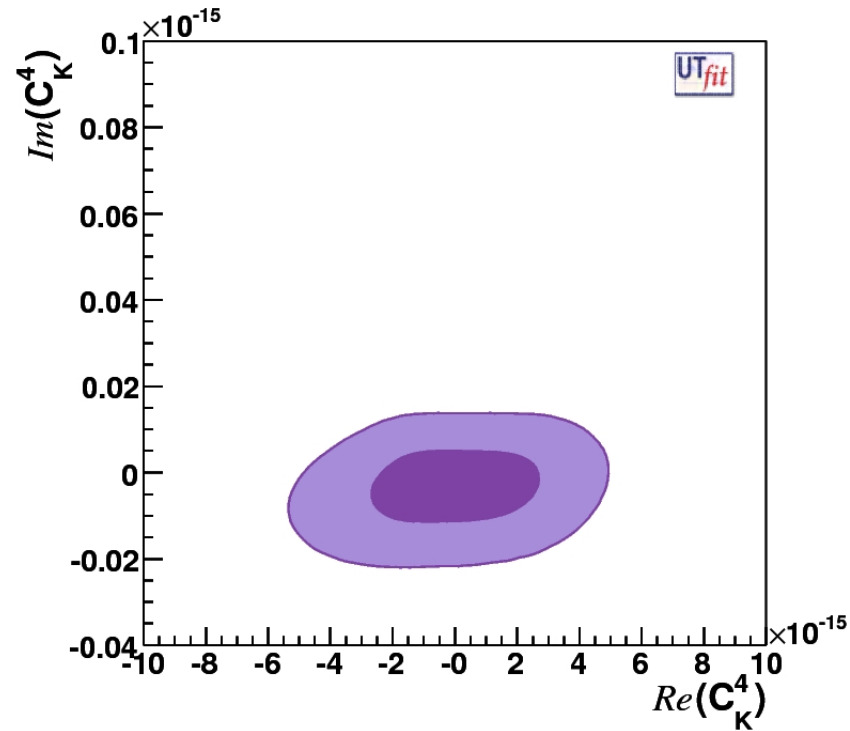
To obtain the p.d.f. for the Wilson coefficients  $C_i(\Lambda)$  at the new-physics scale, we switch on **one coefficient at a time** in each sector and calculate its value from the result of the NP analysis.

# Results from the Wilson coefficients

the results obtained for the flavour scenarios:

In deriving the lower bounds on the NP scale, we assume  $L_i = 1$ , corresponding to strongly-interacting and/or tree-level NP.

Parameter	95% allowed range	Lower limit on $\Lambda$ (TeV)	Lower limit on $\Lambda$ (TeV)
	( $\text{GeV}^{-2}$ )	for arbitrary NP	for NMFV
$\text{Re}C_K^1$	$[-9.6, 9.6] \cdot 10^{-13}$	$1.0 \cdot 10^3$	0.35
$\text{Re}C_K^2$	$[-1.8, 1.9] \cdot 10^{-14}$	$7.3 \cdot 10^3$	2.0
$\text{Re}C_K^3$	$[-6.0, 5.6] \cdot 10^{-14}$	$4.1 \cdot 10^3$	1.1
$\text{Re}C_K^4$	$[-3.6, 3.6] \cdot 10^{-15}$	$17 \cdot 10^3$	4.0
$\text{Re}C_K^5$	$[-1.0, 1.0] \cdot 10^{-14}$	$10 \cdot 10^3$	2.4
$\text{Im}C_K^1$	$[-4.4, 2.8] \cdot 10^{-15}$	$1.5 \cdot 10^4$	5.6
$\text{Im}C_K^2$	$[-5.1, 9.3] \cdot 10^{-17}$	$10 \cdot 10^4$	28
$\text{Im}C_K^3$	$[-3.1, 1.7] \cdot 10^{-16}$	$5.7 \cdot 10^4$	19
$\text{Im}C_K^4$	$[-1.8, 0.9] \cdot 10^{-17}$	$24 \cdot 10^4$	62
$\text{Im}C_K^5$	$[-5.2, 2.8] \cdot 10^{-17}$	$14 \cdot 10^4$	37



To obtain the lower bound for loop-mediated contributions, one simply multiplies the bounds by  $\alpha_s \sim 0.1$  or by  $\alpha_w \sim 0.03$ .

# The future of CKM fits

LHCb reach from:  
O. Schneider, 1<sup>st</sup> LHCb  
Collaboration Upgrade  
Workshop



2015

10/fb (5 years)  
0.07%(+0.5%)

?

0.01+syst

0.010

2.4°

4.5°

no

no



SuperB reach from:  
SuperB Conceptual  
Design Report,  
arXiv:0709.0451

1/ab (1 month  
no at Y(5S))

0.006

0.14

75/ab (5 years)

0.005

1-2°

1-2°

< 1%

1-2%

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Hadronic matrix element	Current lattice error	60 TFlop Year [2011 LHCb]	1-10 PFlop Year [2015 SuperB]
$f_+^{K\pi}(0)$	0.9% (22% on $1-f_+$ )	0.4% (10% on $1-f_+$ )	< 0.1% (2.4% on $1-f_+$ )
$\hat{B}_K$	11%	3%	1%
$f_B$	14%	2.5 - 4.0%	1 - 1.5%
$f_{B_s} B_{B_s}^{1/2}$	13%	3 - 4%	1 - 1.5%
$\xi$	5% (26% on $\xi-1$ )	1.5 - 2% (9-12% on $\xi-1$ )	0.5 - 0.8% (3.4% on $\xi-1$ )
$\mathcal{F}_{B \rightarrow D/D^*1\nu}$	4% (40% on $1-\mathcal{F}$ )	1.2% (13% on $1-\mathcal{F}$ )	0.5% (5% on $1-\mathcal{F}$ )
$f_+^{B\pi}, \dots$	11%	4 - 5%	2 - 3%
$T_1^{B \rightarrow K^*/\rho}$	13%	----	3 - 4%

S. Sharpe @ Lattice QCD: Present and Future, Orsay, 2004  
and report of the U.S. Lattice QCD Executive Committee

