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Paul Scherrer Institut
Jiachao Chen

**A New Mechanism
of Loop Formation and Transformation
in bcc Iron without Dislocation Reaction**

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Outline

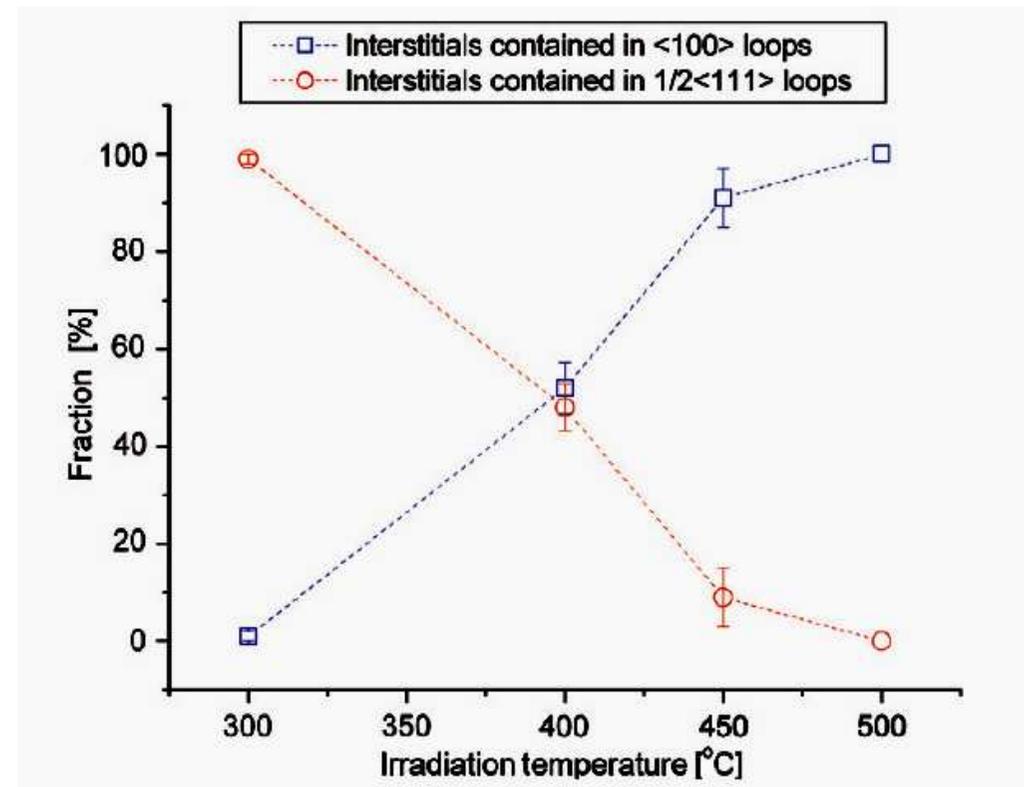
- 1. Introduction**
- 2. Experimental Evidence**
- 3. Proposed Mechanism**
- 4. MD simulation and discussion**
- 5. Conclusion**

Dislocation loops formed by irradiation damage in Fe-Cr alloys

Coexistence of $\frac{1}{2}\langle 111 \rangle$ and $\langle 100 \rangle$ loops

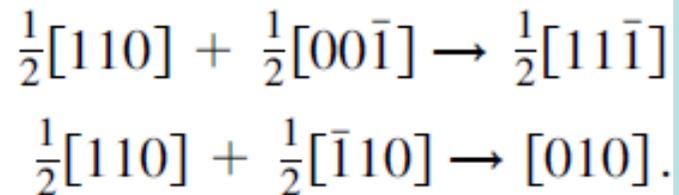
Configurations:

b	$\frac{1}{2}\langle 111 \rangle$	$\langle 100 \rangle$
habit planes	$\{111\}$ $\{110\}$	$\{100\}$
nature	I	I

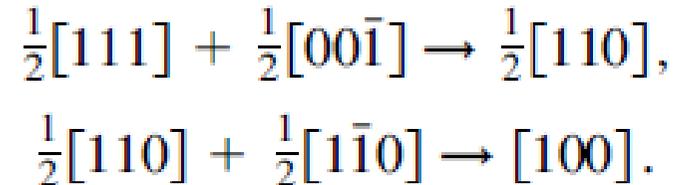


Previously proposed mechanisms for formation/transformation of $\frac{1}{2}\langle 111 \rangle$ and $\langle 100 \rangle$ loops in Fe-Cr alloys

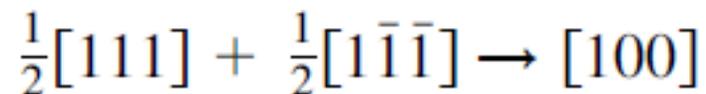
Eyre and Bullough
[Philos.Mag. 12 (1965) 31]



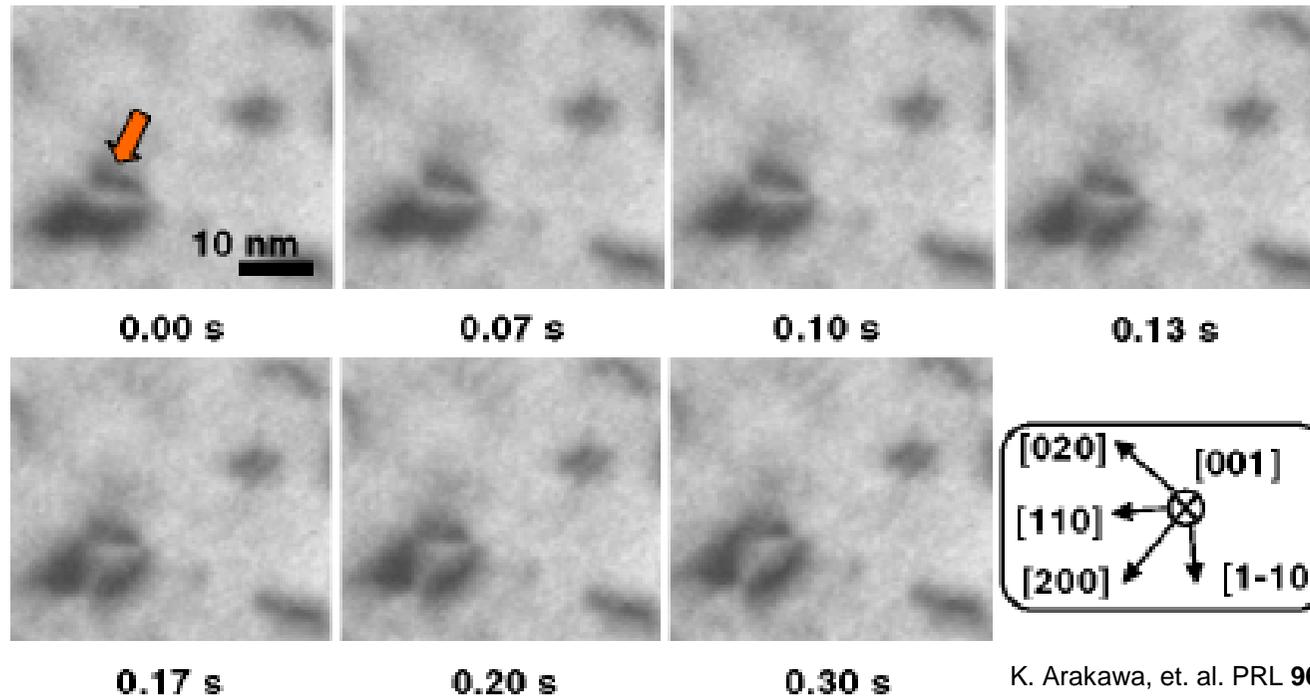
Marian et al.
[PRL 88 (2002) 255507-1]



B. C. Masters
[Philos. Mag. 11 (1965) 881]



Contradiction to recent experiments



Change of loop
Burgers vectors
from
 $\frac{1}{2}\langle 111 \rangle$ to $\langle 100 \rangle$
and vice versa
without
loop/dislocation
interaction

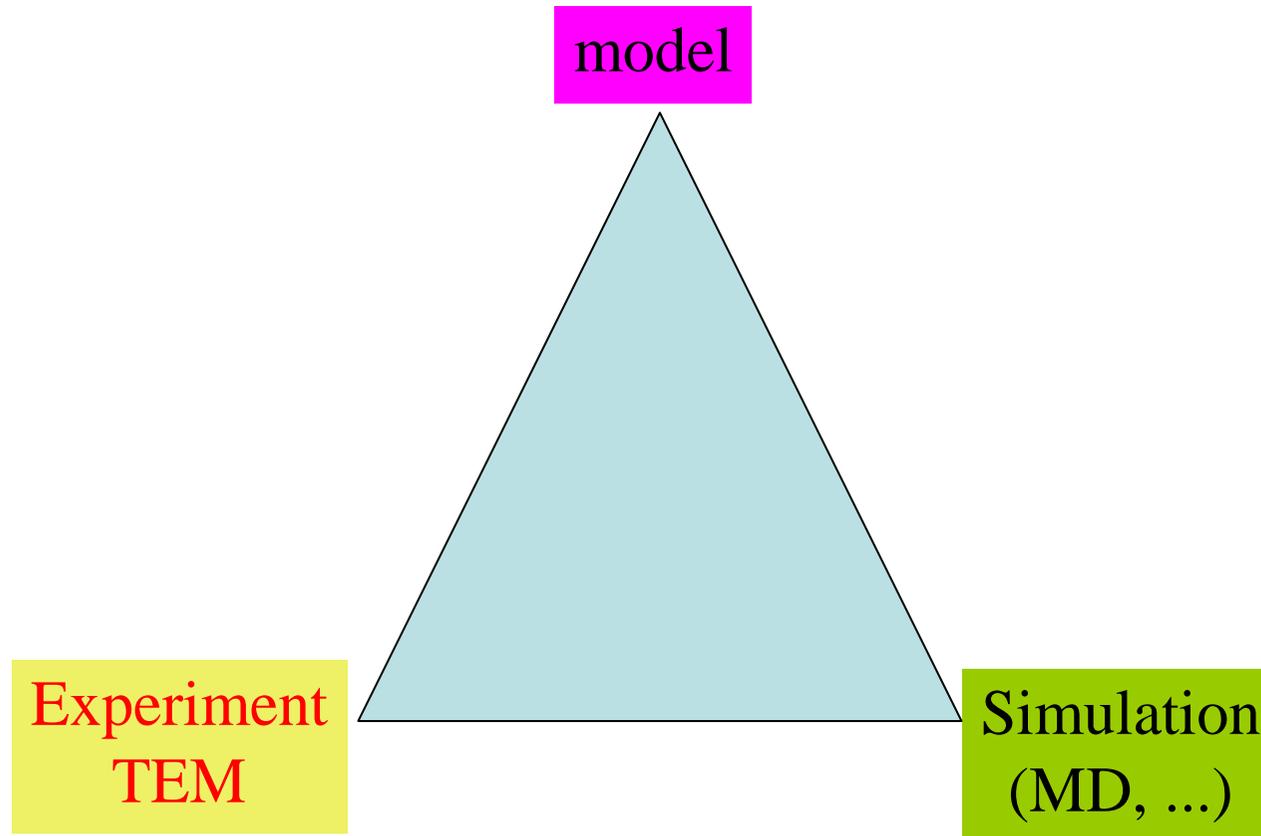
K. Arakawa, et. al. PRL **96** (2006) 125506

In another experiment: K. Arakawa, T. Amino, H. Mori, Acta Mater. 59 (2011) 141

large loops absorb small ones with different b by an absorption reaction
without changing its own b .

This is also realized in MD simulation.

Study methods



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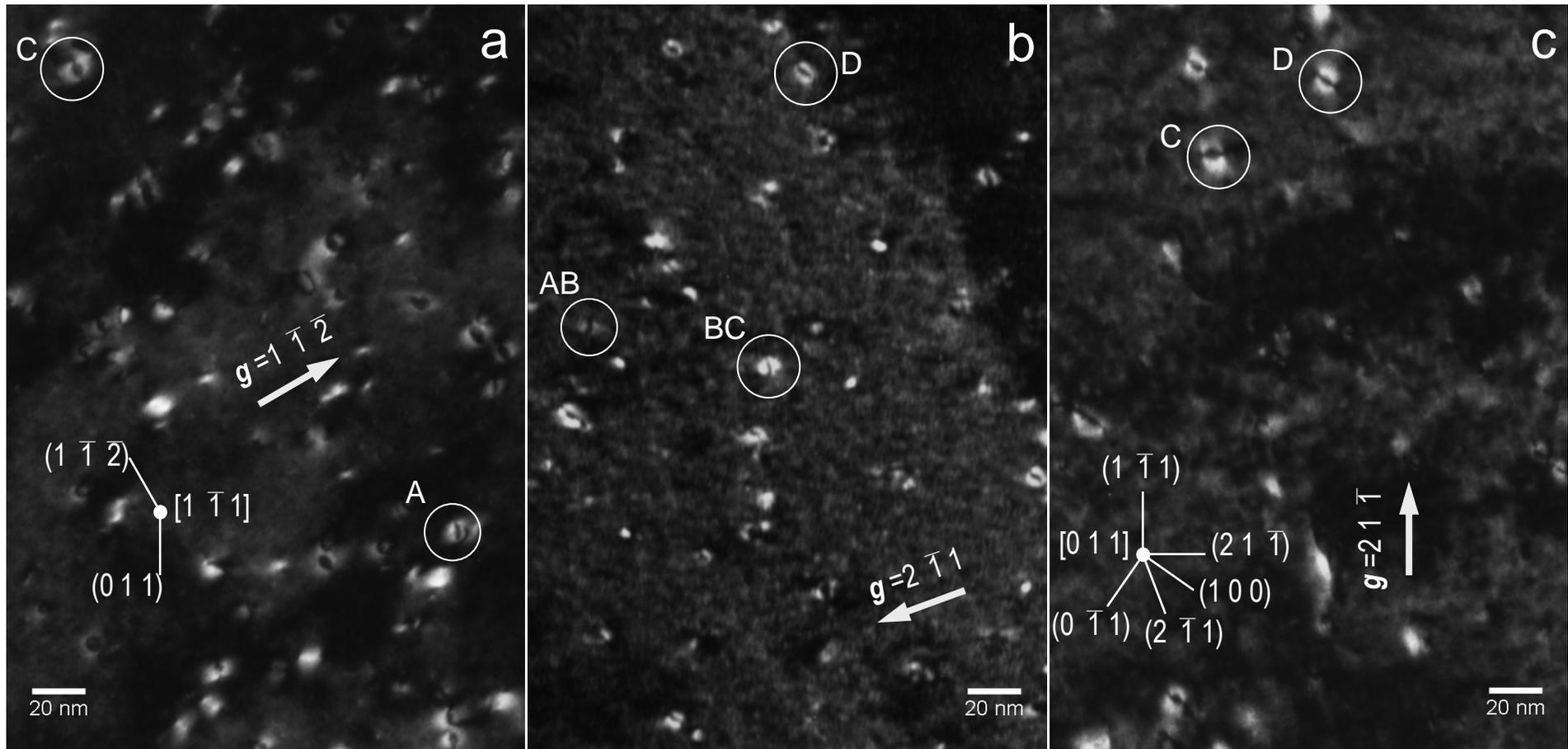
Experimental details

Materials:

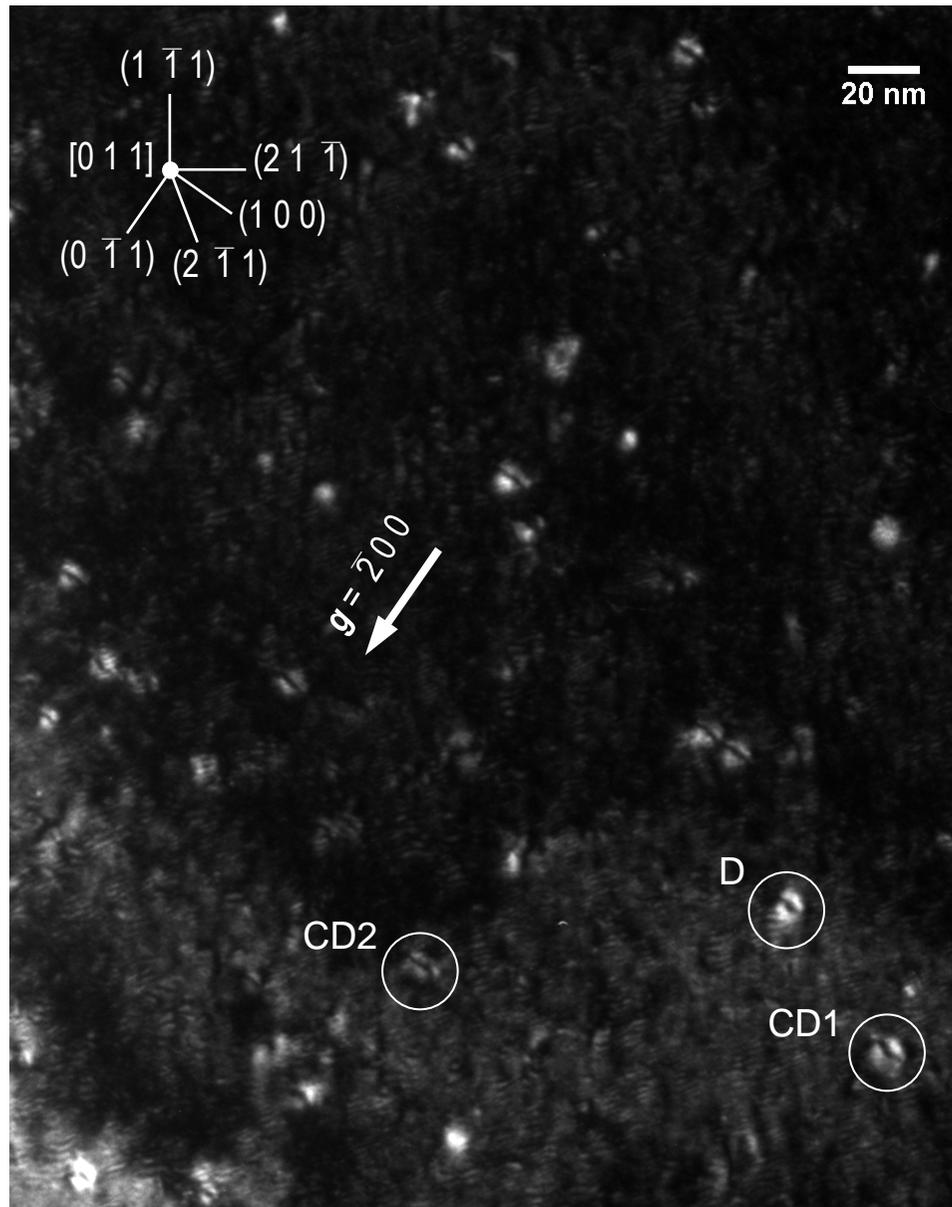
polycrystalline Fe (99.99+%), supplied by Goodfellow

Experiments Performed:

- Specimens of 0.1 mm thick annealed isothermally at 1073 K for 1 hour
- 3D-homogeneous Irradiation/implantation with 25 MeV α -particles up to 0.13 dpa with damage rate of 5.4×10^{-6} dpa/s at 573 K
- 1 mm diameter discs technique is used for TEM sample preparation
- Loop was analyzed by edge-on view and standard $g \cdot b$ techniques



Loops “A”, “B”, “C” and “D” have the configuration of $\frac{1}{2}\langle 111 \rangle \{110\}$, $\frac{1}{2}\langle 111 \rangle \{111\}$, $\frac{1}{2}\langle 111 \rangle \{211\}$ and $\langle 100 \rangle \{100\}$, respectively, double-letter loops have composite habit planes, e.g. “AB”: $(0\bar{1}1) + (1\bar{1}1)$, “BC”: $(1\bar{1}1) + (2\bar{1}1)$. Edge-on loops in Fig. a have habit planes $(1\bar{1}2)$ and (011) , in Fig. b+c: $(0\bar{1}1)$, $(1\bar{1}1)$, $(2\bar{1}1)$, $(\bar{1}00)$, $(2\bar{1}1)$, and $(1\bar{1}1)$.

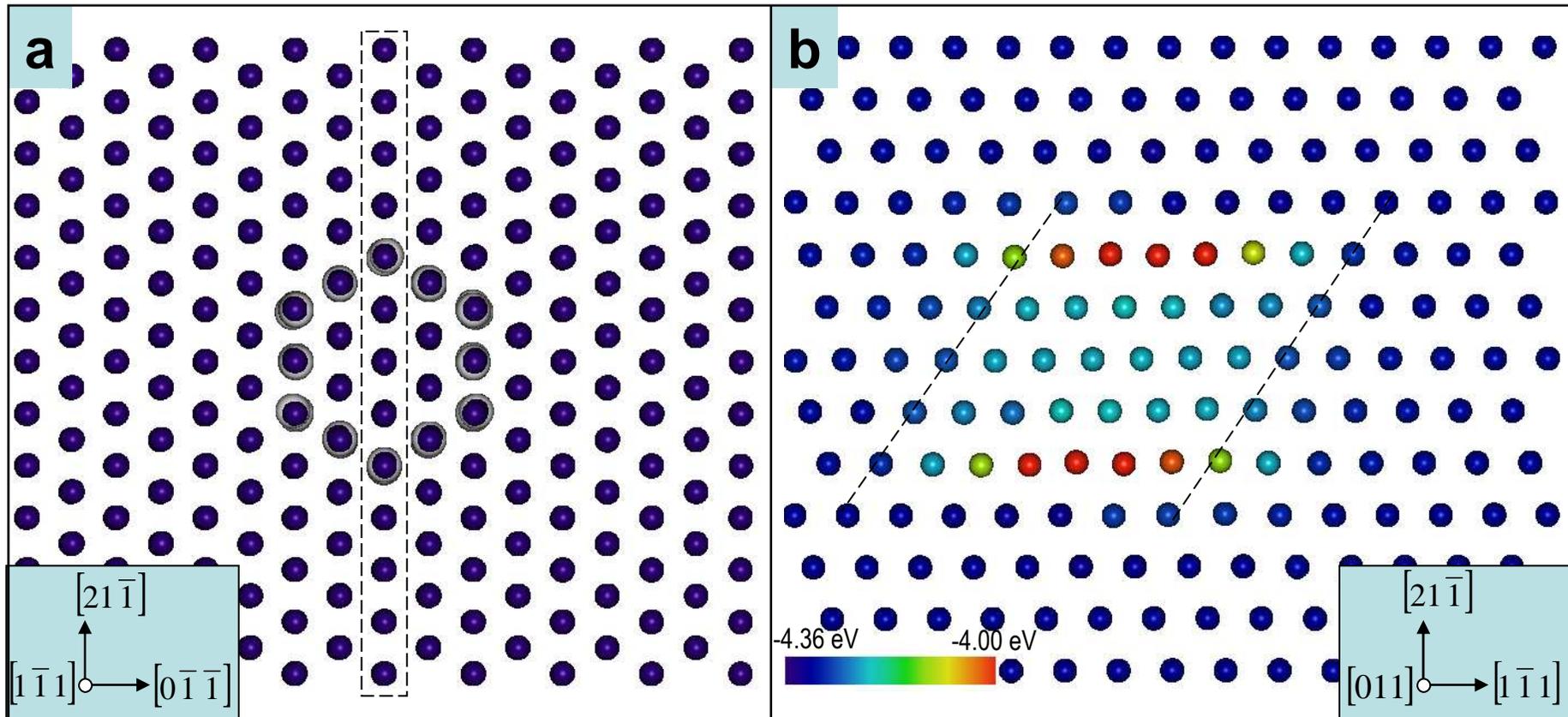


Loops composed of $\frac{1}{2}\langle 111 \rangle\{2\bar{1}1\}$ and $\langle 100 \rangle\{100\}$ components are marked by “CD1” and “CD2”.

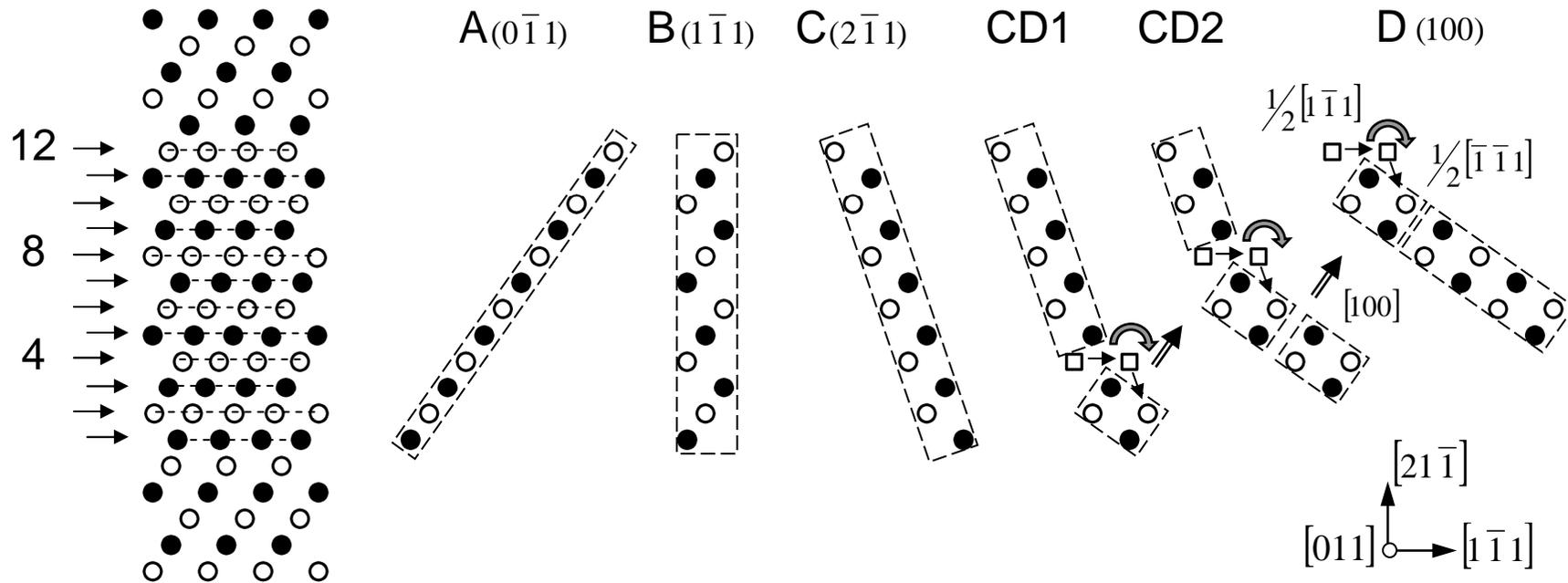
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Loop formed by closely arranged parallel $\frac{1}{2}\langle 111 \rangle$ crowdions



Proposed model in which the transformation steps are depicting the loops observed by TEM



In a (4+1)-step process (pathway) also a transition of a $\frac{1}{2}[1\bar{1}\bar{1}](2\bar{1}\bar{1})$ loop to $\langle 100 \rangle(100)$ is possible: (1) the $\langle 111 \rangle$ crowdions glide a step $\frac{1}{2}[1\bar{1}\bar{1}]$; (2) rotate from a $[1\bar{1}\bar{1}]$ to a $[1\bar{1}\bar{1}]$ crowdion; (3) the $[1\bar{1}\bar{1}]$ crowdions glide a step $\frac{1}{2}[1\bar{1}\bar{1}]$; (4) the SIAs jump to a $[100]$ orientation and form segments of a $\langle 100 \rangle(100)$ loop; (5) a segment of the $\langle 100 \rangle(100)$ loop glides a step along $[100]$. Steps (1) to (5) are repeated until the whole loop is transformed.

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MD Simulation Parameters

Ensemble: NPT

Approach: Parrinello-Rahman algorithm, nudged elastic band method

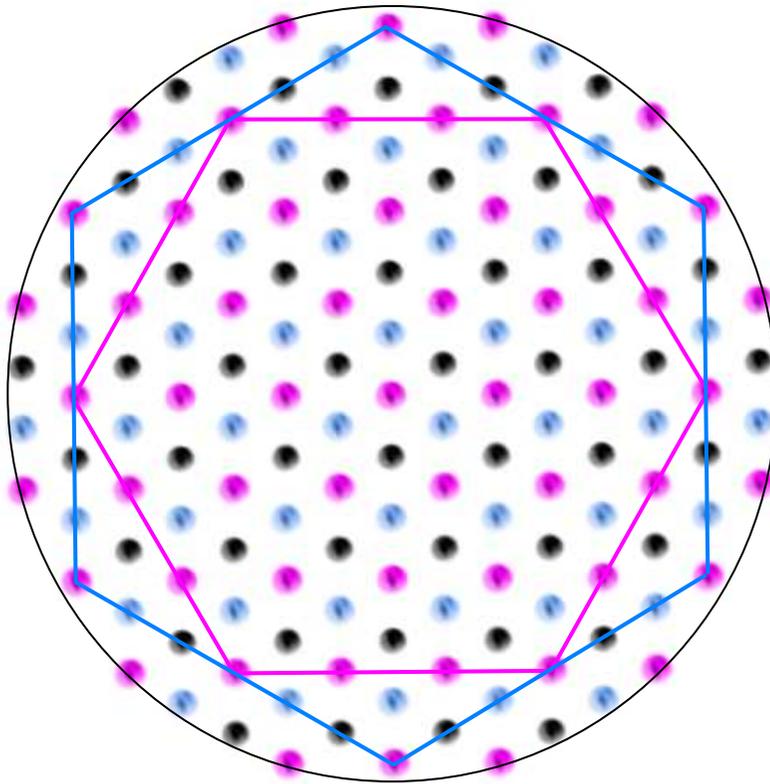
Fe-Fe potential: Ackland [1997]; Mendeleev [2003]; Dudarev-Derlet [2008]

Cell box: up to 40x15x25 cells

Loop sizes: up to 500(~ 5 nm) atoms

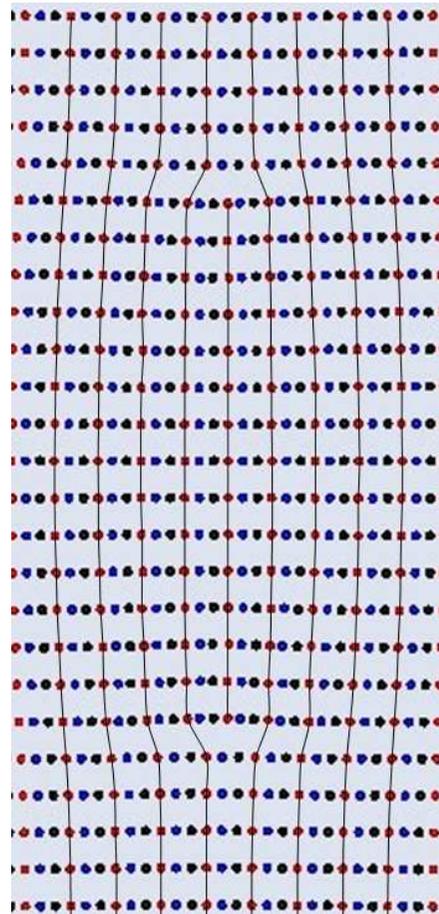
$1/2\langle 111\rangle\{111\}$ loop

Face-on view

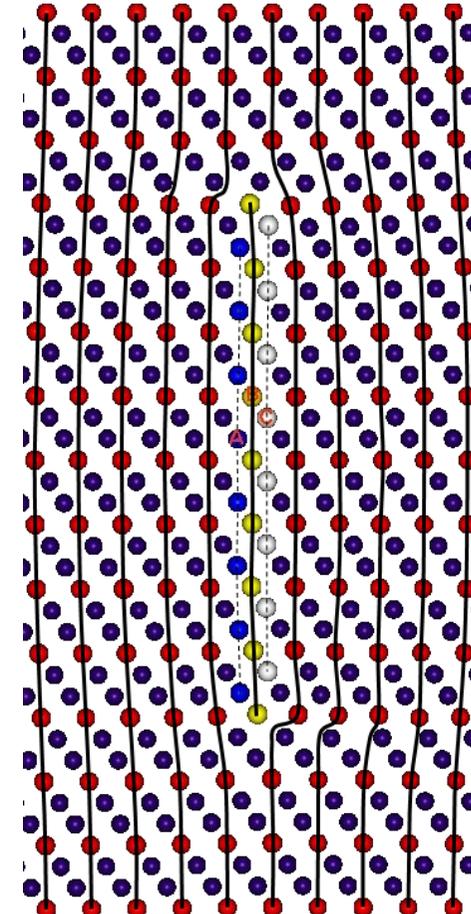


Edge-on view

Line $\langle 211 \rangle$

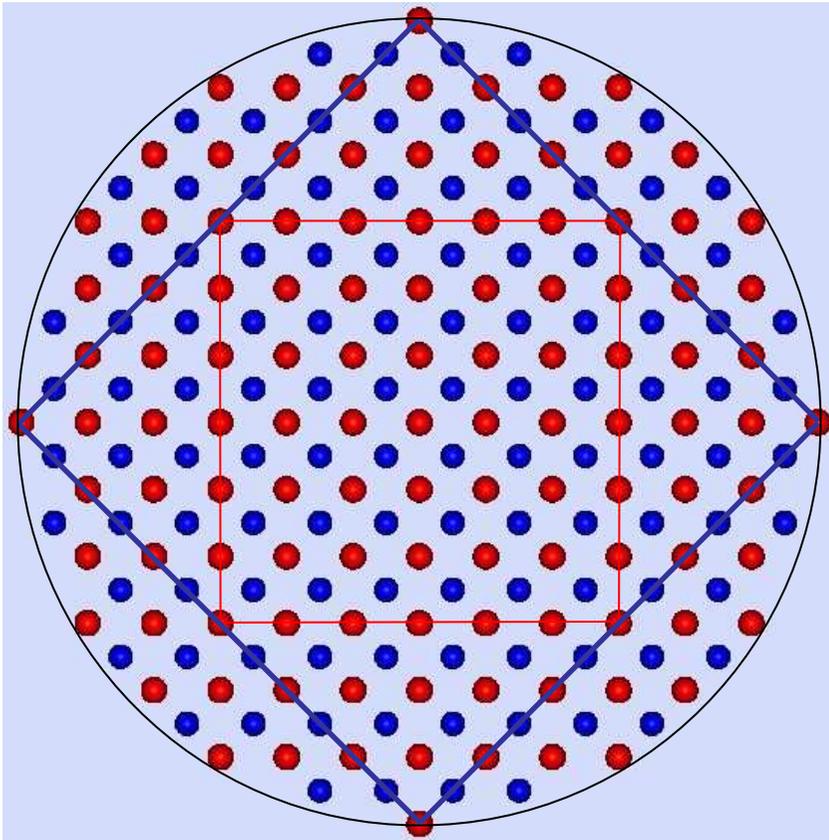


Line $\langle 110 \rangle$



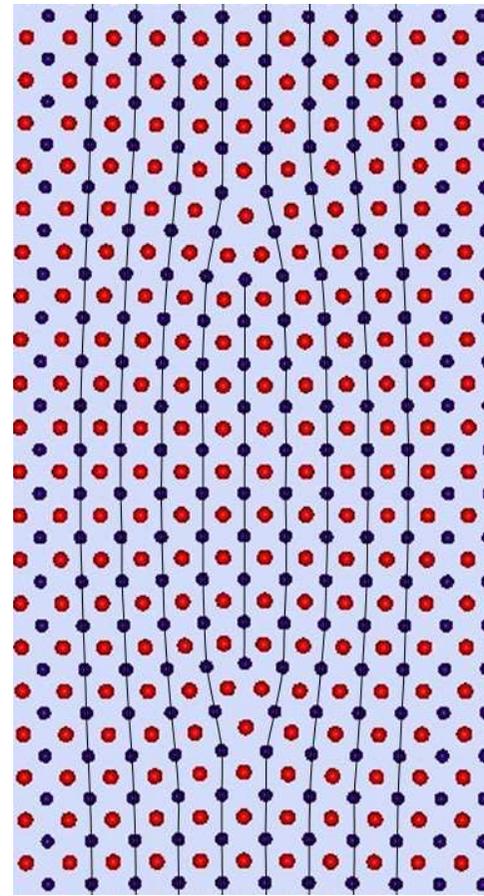
$\langle 100 \rangle \{100\}$ loop

Face-on view

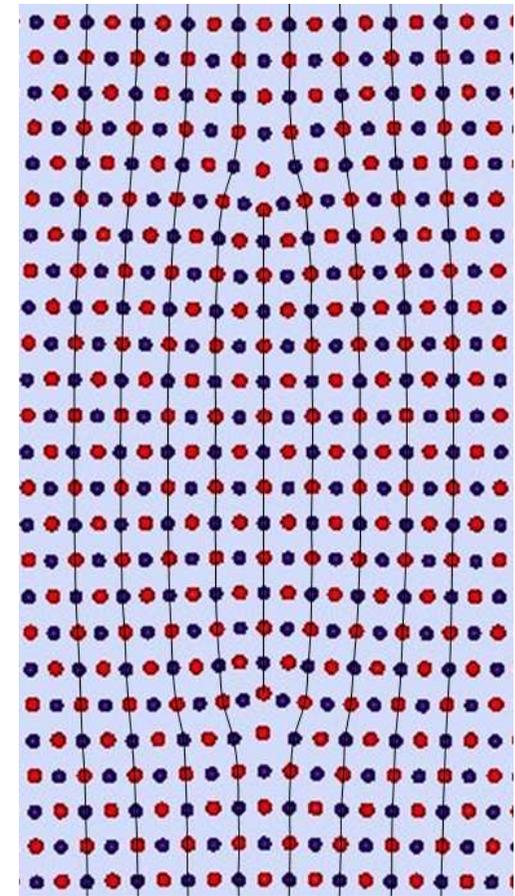


Edge-on view

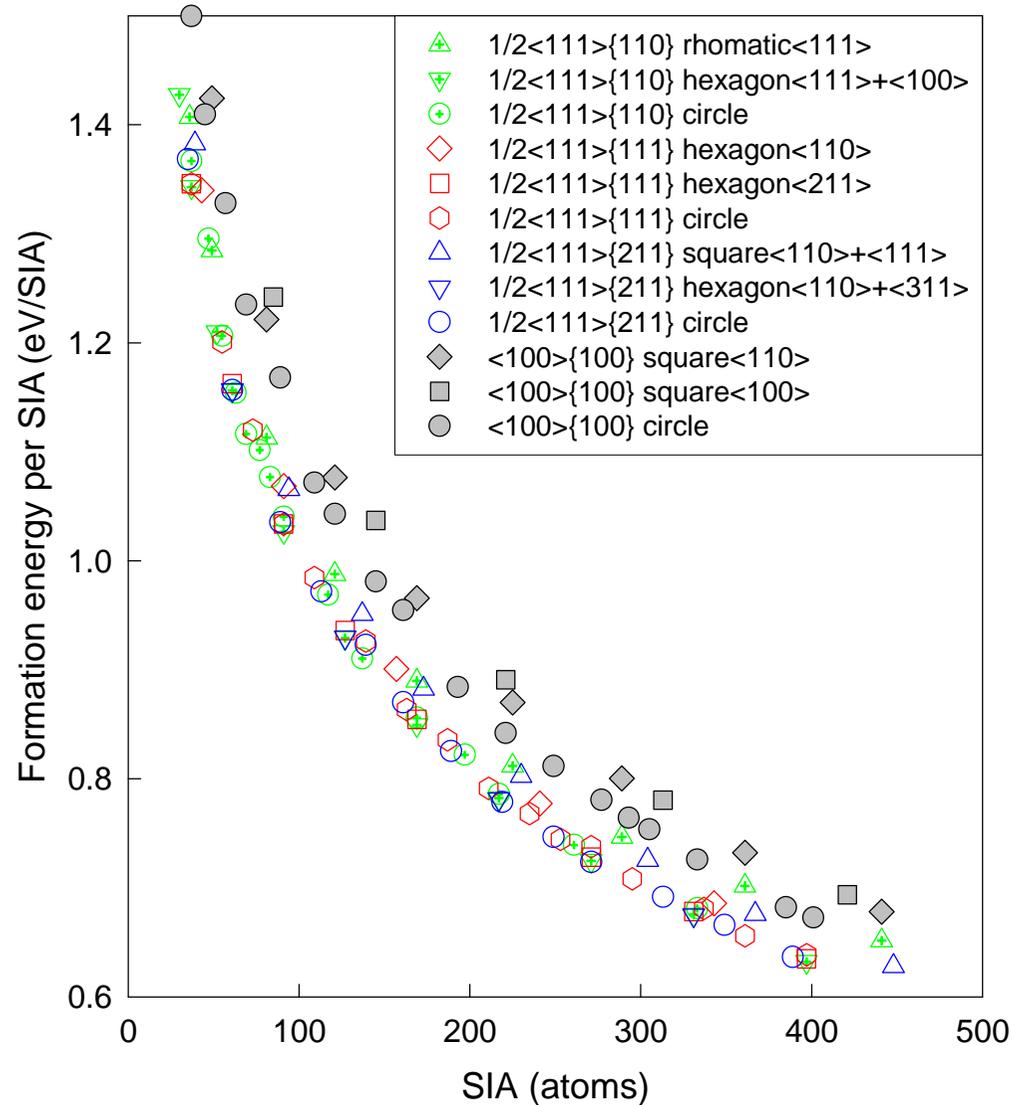
Line $\langle 100 \rangle$



Line $\langle 110 \rangle$



Formation energies per SIA of loops

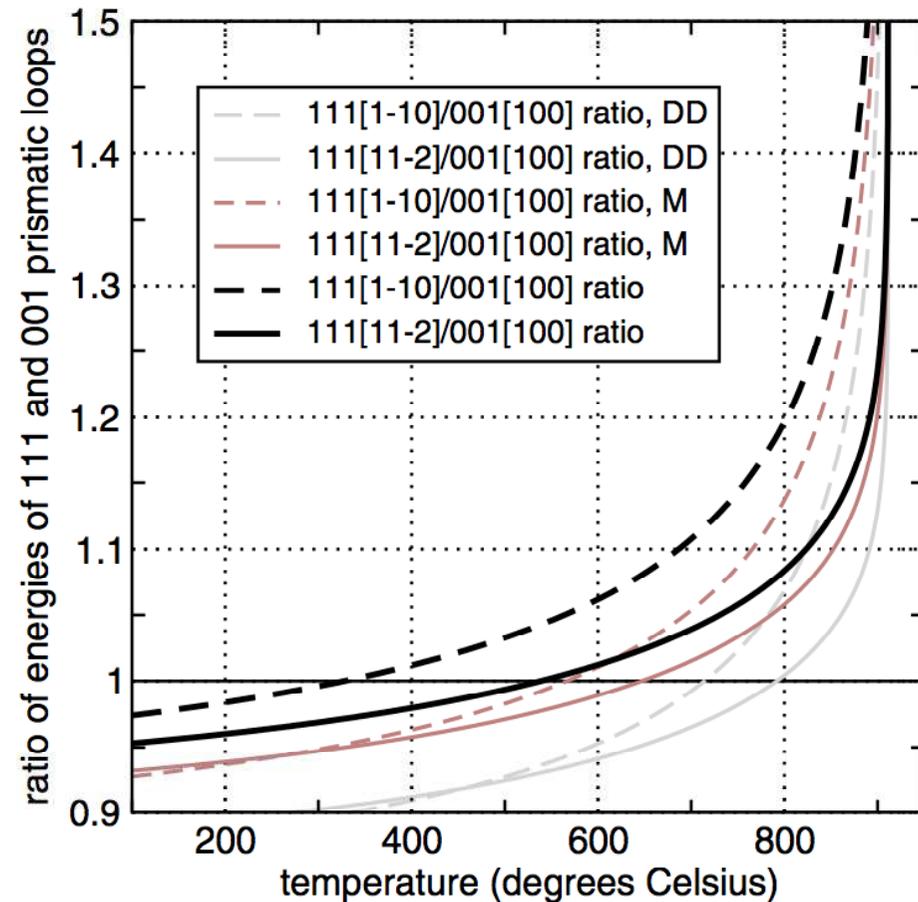


Up to now, MD simulation with various potentials shows that $1/2[111]$ loops are always more stable up to 1000 K. Only the considerations on temperature dependence of elastic constants, which affect the anisotropic elastic self-energies of dislocation in iron, could account for the experimental observation of $\langle 100 \rangle$ loops at 573 K and the increase portion of $\langle 100 \rangle \{100\}$ loop with increasing temperature up to 773 K. To clarify this point, **a new potential, valid at elevated temperatures is demanded**

Thermodynamic reasoning

Ratio of the total free energies calculated using elasticity of the continuum of 10 nm prismatic loops with $1/2 \langle 111 \rangle$ to $\langle 100 \rangle$ Burgers vector as a function of temperature. The two sets of lighter lines correspond to using a dislocation core energy found using the empirical potentials for α -Fe, whereas the darker curves represent the best fit of the dislocation core energy to experimental observations. The temperature dependence is the one of the elastic constants, which was measured experimentally.

S.L. Dudarev, R. Bullough and P.M. Derlet, Phys. Rev. Lett. 100(2008)135503



kinetic processes

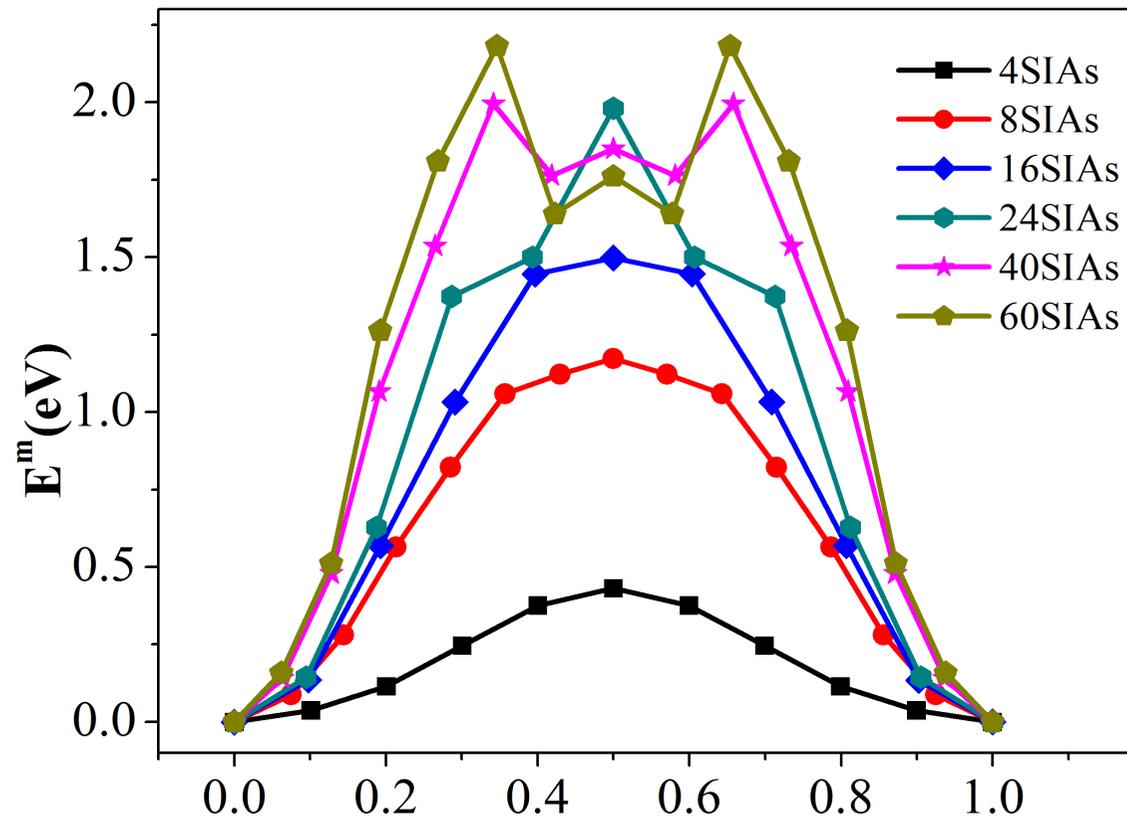
The kinetic processes occurring in the above loop transformation

- A) glide of $\langle 111 \rangle$ crowdions along $\langle 111 \rangle$;
- B) glide of segments of $\{100\}$ loops along $\langle 100 \rangle$;
- C) jumping of $\langle 111 \rangle$ crowdions between different $\langle 111 \rangle$ directions, and rotation into $\langle 100 \rangle$ orientation.

The loop glides by forming and moving double kinks as described in text books. In this case, the apparent activation energy will be the sum of kink formation energy and migration energy of a dislocation segment.

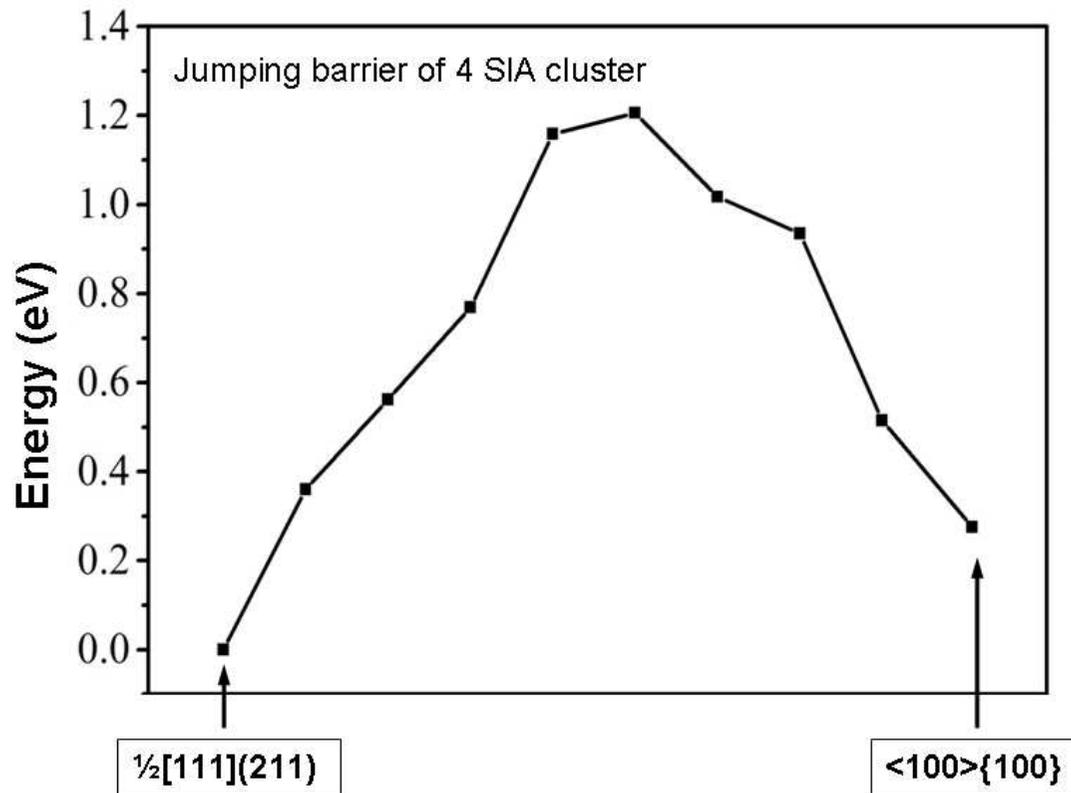
The activation energy calculated by diffusivity calculation method for the $\frac{1}{2}\langle 111 \rangle$ loop is in all cases less than 0.1 eV which is lower than the experimental value of roughly 0.4 eV for “naked” loops

glide barrier of $\langle 100 \rangle$ loop



The activation energy for gliding of a loop along $\langle 100 \rangle$ is estimated by migration energy of a dislocation segment containing 4 SIA (0.4 eV) plus a double kink formation energy (~ 0.5 eV) resulting ~ 0.9 eV.

Reorientation barrier of SIA cluster



The energy barrier for an isolated 4 SIA cluster jumping from a $\langle 111 \rangle$ crowdion arrangement to a $\langle 100 \rangle$ orientation via the proposed pathway resulting in a minimum size of a meta-stable $\langle 100 \rangle \{100\}$ loop. Calculation of the additional energy for conjunction of an original $\frac{1}{2}\langle 111 \rangle$ loop and its transformation to a $\langle 100 \rangle$ segment is still not possible. However, if we assume this additional energy to be in the order of that of a double kink, i.e. ~ 0.5 eV. Then total barrier for gradual reorientation will be in the order of **~ 1.7 eV**

Summarizing the above discussion:

The transformation barrier for the proposed mechanism is ~ 1.7 eV, corresponding to a transition temperature of 640 K, which is only slightly lower than the reported experimental value of 673 K.

Conclusions

A new mechanism of loop formation and transformation by self-interstitial atoms aggregation is proposed.

It shows how the change of Burgers vector of small loops from $\frac{1}{2}\langle 111 \rangle$ to $\langle 100 \rangle$ in Fe-materials can possibly result from the re-arrangement and reorientation of crowdion clusters, without reactions with other dislocations.