# Accelerating global fits with Machine Learning

#### Sebastian Liem





Outline

**Part I** Global scan of scalar DM EFT operators. Part II

Fast LHC Signal Prediction using Machine Learning

Real & Complex scalar DM

CMB, Direct detection, and Fermi dSph.

Proof-of-concept: reconstruction

**Gaussian Processes** 

Galactic Centre Excess

# Global analysis of scalar DM EFT operators

#### Effective field theory of dark matter: a global analysis

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#### JHEP09(2016)077

#### **Effective field theory**



#### The idea – combine all operators

If the EFT operators span the 'theory space' then

scan all EFT op  $\square$  scan the 'theory space'.

 $\mathcal{L} \supset \sum_i \lambda_i \mathcal{O}_i$ 

### **DM-parton EFT operators**

Real scalar DM operators								
Label	Coefficient	Operator	$\sigma_{ m SI}$	$\langle \sigma_{ m ann} v  angle$				
R1	$\lambda_1 \sim \frac{1}{2M^2}$	$m_q \chi^2 ar q q$	$\checkmark$	s-wave				
$\mathbf{R2}$	$\lambda_2 \sim rac{-1}{2M^2}$	$im_q\chi^2ar q\gamma^5 q$		s-wave				
$\mathbf{R3}$	$\lambda_3 \sim rac{lpha_s}{4M^2}$	$\chi^2 G_{\mu u} G^{\mu u}$	$\checkmark$	s-wave				
$\mathbf{R4}$	$\lambda_4 \sim rac{lpha_s}{4M^2}$	$i\chi^2 G_{\mu u} ilde{G}^{\mu u}$		s-wave				
Complex scalar DM operators								
Label	Coefficient	Operator	$\sigma_{ m SI}$	$\langle \sigma_{ m ann} v  angle$				
C1	$\lambda_1 \sim rac{1}{M^2}$	$m_q \chi^\dagger \chi ar q q$	$\checkmark$	s-wave				
C2	$\lambda_2 \sim rac{1}{M^2}$	$im_q\chi^\dagger\chiar q\gamma^5 q$		s-wave				
C3	$\lambda_3 \sim \frac{m_1}{M^2}$	$\chi^\dagger ar{\partial}_\mu \chi ar{q} \gamma^\mu q$	$\checkmark$	p-wave				
C4	$\lambda_4 \sim rac{1}{M^2}$	$\chi^\dagger \partial_\mu \chi ar q \gamma^\mu \gamma^5 q$		p-wave				
C5	$\lambda_5\simrac{lpha_s}{8M^2}$	$\chi^\dagger \chi G_{\mu u} G^{\mu u}$	$\checkmark$	s-wave				
C6	$\lambda_6 \sim rac{lpha_s}{8M^2}$	$i\chi^\dagger\chi G_{\mu u} ilde{G}^{\mu u}$		s-wave				

Goodman et al. arXiv:1008.1783

### Data and likelihoods

Modified SuperBayeS: MultiNest, FeynRules, micrOMEGAs, PPPC 4 DM ID

 $\ln L = \ln L_{\Omega h^2}$  Gaussian. Planck determination of relic density. (1502.01589)

- $+\ln L_{
  m LUX}$  LUXCalc (1502.02667). LUX SI x-sec (1310.8214)
- $+\ln L_{
  m dSph}$  Official likelihood. Fermi **dSph** Pass 8 (1503.02641)
- $+ \ln L_{\rm CMB} \, \stackrel{\rm Planck \, {\bf CMB \, anisotropies \, (1502.01589).}}{\rm Cline \,\&\, Scott; \, Slayter \, (1301.5908, \, 1506.03811)}$

#### Log prior

 $\log_{10} m_{\chi} \sim \text{Uniform}(0,3)$  $\log_{10} \lambda_i \sim \text{Uniform}(-20,0)$ 

#### **Dirichlet prior**

 $\log_{10} m_{\chi} \sim \text{Uniform}(0,3)$  $\lambda_i = f_i A$  $\log_{10} A \sim \text{Uniform}(-20, 0)$  $\sum f_i = 1$  $\bar{f} \sim \text{Dir}(\alpha = 0.1) = \frac{\Gamma(\alpha K)}{\Gamma(K)} \prod_{i=1}^{K} f_i^{\alpha - 1}$ 

### **Two priors**

### **Real scalar DM**





### Complex scalar DM

#### **P-wave operators!**

	Real sca	alar DM operato	ors		
Label	Coefficient	Operator	$\sigma_{ m SI}$	$\langle \sigma_{\rm ann} v \rangle$	
R1	$\lambda_1 \sim \frac{1}{2M^2}$	$m_q \chi^2 \bar{q} q$	$\checkmark$	s-wave	
R2	$\lambda_2 \sim \frac{1}{2M^2}$	$im_q\chi^2ar q\gamma^5 q$		s-wave	
$\mathbf{R3}$	$\lambda_3 \sim \frac{\alpha_s}{4M^2}$	$\chi^2 G_{\mu u} G^{\mu u}$	$\checkmark$	s-wave	
R4	$\lambda_4 \sim rac{lpha_s}{4M^2}$	$i\chi^2 G_{\mu u}  ilde{G}^{\mu u}$		s-wave	
Complex scalar DM operators					
Label	Coefficient	Operator	$\sigma_{ m SI}$	$\langle \sigma_{\rm ann} v \rangle$	
C1	$\lambda_1 \sim rac{1}{M^2}$	$m_q \chi^\dagger \chi ar q q$	$\checkmark$	s-wave	
Co	$M^2$	$im_q \chi' \chi q \gamma q$		a wave	
C3	$\lambda_3 \sim \frac{m_1}{M^2}$	$\chi^\dagger ar{\partial}_\mu \chi ar{q} \gamma^\mu q$	$\checkmark$	p-wave	
C4	$\lambda_4 \sim rac{1}{M^2}$	$\chi^\dagger \partial_\mu \chi ar q \gamma^\mu \gamma^5 q$		p-wave	
C0	$\alpha_s$	$\chi^{\dagger}\chi C C^{\mu\nu}$		s-wave	
C6	$\lambda_6 \sim \frac{\alpha_s}{2M^2}$	$i\chi^\dagger\chi G_{\mu u} ilde{G}^{\mu u}$		s-wave	



## Galactic centre excess (Galactic Bulge Emission)

# Null results is not (interesting) enough.

Additional measurements required to identify operators.



- <sup>16</sup> Likelihood from
- $\sim$  Calore et al.

arxiv:1409.0042

Bartels et al. Phys. Rev. Lett. 116, 051102





### So far

Simplest possible combined EFT approach using cosmology, direct, and indirect detection. Measurements are key for minimal theoretical assumption approaches to be interesting in global analysis context.

# Alright, cool. But what about the LHC?

The plan was to include the monojet search.

But the debate on the validity of DM EFT at LHC happened which lead to DM simplified models.

Possible workaround was suggested: Make LHC low-energy experiment by introducing a Mcut = gM and demanding that all events have Ecm < Mcut.

Decent but not satisfying. We want to be all-encompassing.

Considering the cost of simulation and we deemed it not worth it.

### Cost is a limitation

Profile likelihood maps of a 15-dimensional MSSM

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Effective field theory of dark matter: a global analysis

Sebastian Liem,<sup>a</sup> Gianfranco Bertone,<sup>a</sup> Francesca Calore,<sup>a</sup> Roberto Ruiz de Austri,<sup>b</sup> Tim M.P. Tait,<sup>c</sup> Roberto Trotta<sup>d,e</sup> and Christoph Weniger<sup>a</sup>

120M points 400 CPU-years required

#### 4.2M points

Dark matter interpretations of ATLAS searches for the electroweak production of supersymmetric particles in  $\sqrt{s} = 8$  TeV proton–proton collisions

The ATLAS Collaboration

500k points

LHC cannot currently be treated as yet another observation in global fits. We need to be clever.

# Part II Fast LHC Signal Prediction using Machine Learning

#### Collaborators

Physics:

Machine Learning:

**Gianfranco Bertone** 

Marc Peter Deisenroth

Roberto Ruiz de Austri

Max Welling

Jong Soo Kim

(iDark project - Sascha Caron, Tom Heskes, etc. visualisation, scanning, ML techniques, + actual computer scientists)



Do we really have to do the calculation from scratch for all points in the sample?

Not if we assume that points in parameter space that are close to each other have similar predictions. (We do.)

Iffy for cross-sections, probable for SR efficiencies.

#### **Model parameters**

#### Surrogate function

Replace the expensive MC calculation with a **surrogate function**.

 $\epsilon = f(\theta) + \xi$ 

In ML language this a supervised regression problem. Many algorithms, we will use Gaussian processes.



#### Gaussian Processes

**Non-parametric:** No need to assume a functional form.

**Probabilistic:** Produces posteriors, i.e. estimates its own error.

**Bayesian:** Needs a prior on the type of functions. Kernel/covariance function.

 $k_y(x_p, x_q) = \sigma_f^2 \exp\left(-\frac{1}{2\ell^2}(x_p - x_q)^2\right) + \sigma_n^2 \delta_{pq}.$ 

Picking hyperparameters is the *learning* task in Gaussian processes.

To deal with large training dataset, we use Distributed Gaussian Processes Deisenroth & Ng, arXiv:1502.02843



Figure 1.1: Panel (a) shows four samples drawn from the prior distribution. Panel (b) shows the situation after two datapoints have been observed. The mean prediction is shown as the solid line and four samples from the posterior are shown as dashed lines. In both plots the shaded region denotes twice the standard deviation at each input value x.

"Gaussian Processes for Machine Learning" Rasmussen & Williams, 2006 www.gaussianprocess.org

Online demo, http://www.tmpl.fi/gp/

# THE GOAL: Predict SR efficiencies from theory parameters.

#### INPUT: Natural SUSY

Natural because it stabilizes the electroweak scale without fine-tuning.

Only few SUSY states needs to be light.

$$heta = \{ aneta, \ \mu, \ M_3, \ m_{ ilde{Q}_t}, \ m_{ ilde{t}_R}, \ A_t \}$$

Low-dimensional yet realistic theory.

We already had the training data from previous paper.



FIG. 1. The minimal natural SUSY mass spectrum on the left while the remaining supersymmetric particles are decoupled on the right.

"Natural SUSY: Now or Never?" Kim et al. arXiv:1606.06738

### **OUTPUT: Two Signal Regions**

Defined in ATLAS-PHYS-PUB-2013-011

Looking direct stop production with HL-LHC, 14 TeV with 3000 fb<sup>-1</sup>

Stops decay typically to top or b quarks, W/Z or Higgs bosons, and a LSP. Multiple jets, b-jets, large MET, possibly leptons.

1-lepton

p

0-lepton

MET > 750 GeV m<sub>T</sub>(lepton, MET) > 550 GeV MET > 800 GeV m<sub>T</sub>(b-jet, MET) > 400 GeV

Total bkg: 21.1 ± 5.9

Total bkg: 12.2 ± 3.9

### Some Machine Learning Lingo

We have a dataset with inputs and outputs. Which is randomly split into

**TRAINING DATA** which is used to inform the Gaussian process about the functional relationship between our theory parameter and SR efficiency.

**TESTING DATA** which is **not** use to train the Gaussian process. We use predict the output from the inputs of this data and compare with the true value as defined by the MC.

### **Training the Gaussian Processes**

18647 models split into 16647 models for training and 2000 models for testing.

O(10 min) to train per signal region.

The lunch is not free, just cheaper!

SPheno, Pythia, NLLFAST, CheckMATE, Delphes etc. still needed to generate training data. Models uniformly sampled from these ranges:

```
\begin{array}{ll} 0.1 \, {\rm TeV} \leq \ |\mu| \ \leq 1.0 \, {\rm TeV}, \\ 0.1 \, {\rm TeV} \leq m_{\tilde{Q}_t} \leq 2.0 \, {\rm TeV}, \\ 0.1 \, {\rm TeV} \leq m_{\tilde{t}_R} \leq 2.0 \, {\rm TeV}, \\ 0.1 \, {\rm TeV} \leq |M_3| \leq 3.0 \, {\rm TeV}, \\ |A_t| \leq 3.0 \, {\rm TeV}, \\ 1 \leq \tan\beta \leq 20. \end{array}
```

All models avoid LEP-II chargino limit, and have reasonable Higgs boson mass. Is it fast?

# 0.06 sec/pred

4 GHz single core, unoptimized Python

Is it right?

#### How to evaluate the trained GPs?





0-lepton signal region

1-lepton signal region

### Proof-of-concept: Reconstruction



Schematic view of our reconstruction exercise



#### No MC performed! No event gen. No detector sim. ~154k likelihood eval ~18h on 10 CPUs

Reconstruction, 4D scan

### Outlook

Active learning

Multi-output prediction

Extend to MSSM and other theories

More SRs

**Cross-sections** 

Robust public code library

#### To conclude...

Using LHC data is a bottleneck.

Gaussian processes makes it really really fast.

Bypasses event generation, detector simulation.

The speed enables a fast reconstruction of theory parameters, and eventually dark matter properties.

The method is completely generic, with many possible applications.

# Backup slides

### Proof-of-concept Reconstruction

Assume one particular model is true, and measured at the LHC.

Try to reconstruct the parameters using our two signal regions.

Reasonable Higgs mass, LEP chargino limit.

Benchmark model has
63.2 events in 0-lepton SR
161.8 events in 1-lepton SR
(Bkg: 12.2 ± 3.9 and 21.1 ± 5.9)

#### Distributed Gaussian Processes

The standard Gaussian process scales badly with **N** the size of the training dataset. It involves inverting **N**x**N** matrices.

We use **distributed** Gaussian processes to avoid this. The data is randomly partitioned and on each partition a Gaussian process is defined. Predictions from each process is then combined.



Figure 1. Computational graphs of hierarchical PoE models. Main computations are at the leaf nodes (GP experts, black). All other nodes recombine computations from their direct children. The top node (blue) computes the overall prediction.

"Distributed Gaussian Processes" Deisenroth & Ng, arXiv:1502.02843

#### Benchmark

$$\mu = -192.5$$
$$M_3 = 2100$$
$$m_{\tilde{Q}_3} = 1300$$
$$m_{\tilde{t}_R} = 1200$$
$$\tan \beta = 8.43$$
$$A_t = 2000$$

