Daniel Grumiller

Institute for Theoretical Physics TU Wien

Beyond Lorentzian Geometry II, University of Edinburgh, February 2023



based on work in progress w. F. Ecker, J. Hartong, A. Perez, S. Prohazka, and R. Troncoso



## Outline

Motivation for Carrollian gravity

Carrollian dilaton gravity in  $1\!+\!1$  dimensions

Carrollian extremal surfaces

Outlook towards Carrollian black holes

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Motivation for Carrollian gravity

Carrollian dilaton gravity in 1+1 dimensions

Carrollian extremal surfaces

Outlook towards Carrollian black holes

### Symmetries ubiquitious in constraining physics





Daniel Grumiller - Carrollian black holes

Symmetries ubiquitious in constraining physics

- Kinematics & Dynamics
- Correlations functions



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- Decay channels



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- Kinematics & Dynamics
- Correlations functions
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- Density of states

$$S_{\rm BH} = S_{\rm Cardy}$$

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#### Carrollian symmetries arise in various contexts

 $\blacktriangleright$  Formally  $c \rightarrow 0$  limit of Poincaré

### Carrollian Archeology Jean-Marc Lévy-Leblond Université de Nice ....notwithstanding the sagacious advice by Lewis Carroll, who wrote : "It's no use going back to yesterday, because I was a different person then."

12/08/33

JHLL Caral Markhop, Ne

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- Symmetries of tensionless strings

Tensile closed string String grows longer and fills out spacetime as tension decreases

Space-filling D-brane



tension =  $\frac{1}{2\pi a'}$ 



Eð







tension = 0

Decreasing string tension

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### Following history from SR to GR: natural to gauge Carroll algebra

Gravity actions (but with Carroll boost invariance)

$$\tilde{\mathcal{L}} = E \left[ \frac{1}{4} \Pi^{\mu\nu} \Pi^{\rho\sigma} T_{\mu\rho} T_{\nu\sigma} + \sigma \Pi^{\mu\nu} \overset{(C)}{R}_{\mu\nu} - \sigma^2 T^{\mu} T^{\nu} \overset{(C)}{R}_{\mu\nu} \right]$$

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google "carroll black hole" images; 17th result is song 'Black Hole' by Mackin Carroll

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### Can some of the latter be regarded as "Carrollian black holes?"

### Carrollian gravity

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- Cosmology aplications
- Carrollian holography
- Cond-mat applications
- because it is there

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- metric-like variables
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What?

- generic statements about whole model space vs. specific examples
- focus on simple model
- Iowest dimension possible: 1+1
- consider limit from 2d (dilaton) gravity

No lightcones in Carroll gravity!

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Thanks for your attention! Questions?

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This talk pursues last option: Carrollian extremal surfaces (CES)

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### Carrollian dilaton gravity in $1\!+\!1$ dimensions

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Outlook towards Carrollian black holes

▶ In first order form:  $sl(2, \mathbb{R})$  BF theory

$$I_{\rm JT} \sim \int \langle \mathcal{X} F \rangle \qquad \qquad F = \mathrm{d}A + A \wedge A$$

Expand connection 1-form as

$$A = \omega J + e_a P^a$$

generators  $J, P_a$  obey (A)dS<sub>2</sub> algebra

$$[P_a, P_b] = \epsilon_{ab} \Lambda \qquad [P_a, J] = \epsilon_a{}^b P_b$$

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Expanding in components more familiar Cartan formulation

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with  $R = d\omega$ ,  $T_a = de_a + \epsilon_{ab} \omega \wedge e^b$ , and  $\epsilon = \epsilon_{ab} e^a \wedge e^b$ Interpretation of fields:

- ω: (dualized) Lorentz connection
- ▶ e<sub>a</sub>: zweibein (dyad)
- X: dilaton

X<sup>a</sup>: Lagrange multipliers for torsion constraints

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Constant curvature solutions

$$R=\epsilon\,\Lambda$$

depending on sign( $\Lambda$ )  $\neq$  0: (A)dS<sub>2</sub>
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depending on  $\mathsf{sign}(\Lambda) \neq 0$ : (A)dS\_2

- ► JT/SYK correspondence (Schwarzian, chaos bound, ...)
- JT gravity has black hole solutions

Consider most general consistent deformation of JT

- maintain number of field degrees of freedom
- maintain number of gauge degrees of freedom
- maintain gravity interpretation (Lorentzian boosts, 2d diffeos)

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$$I_{\rm dil} = \frac{k}{2\pi} \int \left( XR + X^a T_a - \epsilon \, \mathcal{V}(X, \, X^a X_a) \right)$$

JT recovered for  $\mathcal{V}=\Lambda X$ 

 $\boldsymbol{k}$  is gravitational coupling constant

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notation:  $A_I = (\omega, e_a) = (\omega, e_+, e_-)$  and  $X^I = (X, X^a) = (X, X^+, X^-)$   
 $X^I$  can be interpreted as (target space) coordinates of a Poisson manifold  
 $I_{\text{dil}} = I_{\text{PSM}} \sim \int \left( X^I \, \mathrm{d}A_I + P^{IJ}(X^K) \, A_I \wedge A_J \right) \qquad P^{IJ} = \begin{pmatrix} 0 & -X^+ & X^- \\ X^+ & 0 & Y \\ -X^- & -Y & 0 \end{pmatrix}$ 

Ikeda '93; Schaller, Strobl '94

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Have infinite family of (toy) models available!

Selected list of models  $\mathcal{V} = V(X) + X^a X_a U(X)$  (see review hep-th/0604049)

Black holes in (A)dS<sub>2</sub>, asymptotically flat or arbitrary spaces (Wheeler property)

Model	U(X)	V(X)
1. Schwarzschild (1916)	$-\frac{1}{2X}$	$-\lambda^2$
2. Jackiw–Teitelboim [JT] (1984)	0	$\Lambda X$
3. Witten Black Hole (1991)	$-\frac{1}{X}$	$-2b^2X$
4. CGHS (1992)	0	$-2\Lambda$
5. (A)dS <sub>2</sub> ground state (1994)	$-\frac{a}{X}$	BX
6. Rindler ground state (1996)	$-\frac{a}{X}$	$BX^a$
7. Black Hole attractor (2003)	0	$BX^{-1}$
8. Spherically reduced gravity $(N > 3)$	$-\frac{N-3}{(N-2)X}$	$-\lambda^2 X^{(N-4)/(N-2)}$
9. All above: <i>ab</i> -family (1997)	$-\frac{a}{X}$	$BX^{a+b}$
10. Liouville gravity	a	$be^{\alpha X}$
11. Reissner–Nordström (1916)	$-\frac{1}{2X}$	$-\lambda^2 + \frac{Q^2}{X}$
12. Schwarzschild-(A)dS	$-\frac{21}{2X}$	$-\lambda^2 - \ell X$
13. Katanaev–Volovich (1986)	α	$\beta X^2 - \Lambda$
14. BTZ/Achucarro–Ortiz (1993)	0	$\frac{Q^2}{X} - \frac{J}{4X^3} - \Lambda X$
15. KK reduced CS (2003)	0	$\frac{1}{2}X(c-X^2)$
16. KK red. conf. flat (2006)	$-\frac{1}{2} \tanh{(X/2)}$	$A \sinh X$
17. 2D type 0A string Black Hole	$-\frac{1}{X}$	$-2b^2X + \frac{b^2q^2}{8\pi}$
18. exact string Black Hole (2005)	lengthy	lengthy

Daniel Grumiller — Carrollian black holes

Carrollian dilaton gravity in 1+1 dimensions

Change to sl(2) basis convenient for Carrollian contraction

$$[H, J] = -\delta P \qquad [P, J] = -H \qquad [H, P] = -\Lambda J$$

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Write down corresponding BF action

$$I_{\rm CJT} \sim \int \left( XR + X_{\rm H}T + X_{\rm P} \,\mathrm{d}e - \epsilon \,\Lambda X \right)$$

with curvature  $R = d\omega$ , torsion  $T = d\tau + \omega \wedge e$ , and volume  $\epsilon = \tau \wedge e$ 

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$$egin{aligned} \delta_\lambda X &= 0 & & \delta_\lambda X_{
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m H} \,\lambda \ \delta_\lambda \omega &= {
m d}\lambda & & & \delta_\lambda au &= -e \,\lambda & & & \delta_\lambda e &= 0 \end{aligned}$$

invariance of dilaton  $X_{\rm I}$  auxiliary scalar  $X_{\rm H}$  , and spatial vielbein e

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Opportunities for students to carve out interesting corners of model space!

Selected list of Carrollian models  $\mathcal{V} = V(X)$ 

ModelV(X)1. Carrollian Jackiw–Teitelboim [CJT] (2020) $\Lambda X$ 

Opportunities for students to carve out interesting corners of model space!

EOM (with their suggested names\*):

Carrollian curvature: Carrollian torsion: No intrinsic torsion: Carrollian expansion: Carrollian Casimir: Boost non-invariant:

 $\begin{aligned} \mathrm{d}\omega &+ \partial_X V(X, \, X_\mathrm{H}) \, \tau \wedge e = 0 \\ \mathrm{d}\tau &+ \omega \wedge e + \partial_{X_\mathrm{H}} V(X, \, X_\mathrm{H}) \, \tau \wedge e = 0 \\ \mathrm{d}e &= 0 \\ \mathrm{d}X &+ X_\mathrm{H} \, e = 0 \\ \mathrm{d}X_\mathrm{H} &+ V(X, \, X_\mathrm{H}) \, e = 0 \\ \mathrm{d}X_\mathrm{P} &- V(X, \, X_\mathrm{H}) \, \tau - X_\mathrm{H} \, \omega = 0 \end{aligned}$ 

\* please speak up now if you object to some of these names

1. Constant dilaton vacua

$$X_{\rm H} = 0$$
  $X = \text{const.}$  s.t.  $\mathcal{V}(X, 0) = 0$ 

constant curvature,  $R = \partial_X \mathcal{V}$ ; slightly boring sector

Follow Lorentzian algorithm DG, Kummer, Vassilevich '02

- 1. Constant dilaton vacua
- 2. Linear dilaton vacua: follow essentially standard algorithm

Follow Lorentzian algorithm DG, Kummer, Vassilevich '02

1. Constant dilaton vacua

#### 2. Linear dilaton vacua: follow essentially standard algorithm

• assume  $X_{\rm H} \neq 0$  and write  $e = - \,\mathrm{d} X / X_{\rm H}$ 

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- 1. Constant dilaton vacua
- 2. Linear dilaton vacua: follow essentially standard algorithm
  - assume  $X_{\rm H} \neq 0$  and write  $e = \,\mathrm{d}X/X_{\rm H}$
  - integration of EOM

$$\frac{1}{2} \mathrm{d}X_{\mathrm{H}}^2 = \mathcal{V}(X, \, X_{\mathrm{H}}) \mathrm{d}X$$

establishes conserved Casimir/mass

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• for simplicity: 
$$\mathcal{V} = V(X)$$
; define  $U(X) = \int^X V(y) \, dy$ ; yields mass

$$dM = 0$$
  $M = U(X) - \frac{1}{2}X_{\rm H}^2$ 

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- assume  $X_{\rm H} \neq 0$  and write  $e = \,\mathrm{d} X / X_{\rm H}$
- integration of EOM

$$\frac{1}{2} \, \mathrm{d}X_{\mathrm{H}}^2 = \mathcal{V}(X, \, X_{\mathrm{H}}) \, \mathrm{d}X$$

establishes conserved Casimir/mass

• for simplicity: 
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• curvature  $R = \partial_X V$  not necessarily constant

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# Global structure? Singularities? Special surfaces? (horizon, trapping, extremal, ...)

### Outline

Motivation for Carrollian gravity

Carrollian dilaton gravity in 1+1 dimensions

Carrollian extremal surfaces

Outlook towards Carrollian black holes

#### Lorentzian extremal surfaces in target space

classify co-dimension-2 surfaces according to their null expansions

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 $X^- = 0$  | marginally anti-trapped marginally trapped

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$$\begin{array}{c|c} \text{gns} & X^+ > 0 & X^+ < 0 & X^+ = 0 \\ \hline > 0 & \text{anti-trapped} & \text{anti-normal} & \text{marginally anti-trapped} \\ \hline < 0 & \text{normal} & \text{trapped} & \text{marginally trapped} \end{array}$$

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Extremal surfaces are boost invariant loci!

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Carrollian classification of co-dimension-2 surfaces simple

signs	$X_{\rm H} > 0$	$X_{\rm H} < 0$	$X_{\scriptscriptstyle \mathrm{H}}=0$
	normal	anti-normal	extremal

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definition above conceivably generalizes to higher dimensions

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#### Address these issues in remainder of talk

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- implement solution algorithm for CJT:

$$X = \frac{1}{2} e^{r/\ell} + M\ell^2 e^{-r/\ell} \qquad \qquad \omega = -\frac{X}{\ell^2} dt$$
$$X_{\rm H} = \frac{1}{2} e^{r/\ell} - M\ell^2 e^{-r/\ell} \qquad \qquad \tau = X_{\rm H} dt$$
$$X_{\rm P} = 0 \qquad \qquad e = dr$$

 $r \to \pm \infty$ : asymptotic regions;  $r \to \frac{\ell}{2} \ln(2M\ell^2)$ : CES (for M > 0) note similarities to AdS<sub>2</sub> black holes! suggestive to impose Brown-Henneaux-like boundary conditions

$$\begin{split} X &= \frac{1}{2} e^{r/\ell} + \mathcal{O}(e^{-r/\ell}) & \omega = \left( -\frac{1}{2\ell^2} e^{r/\ell} + \mathcal{O}(e^{-r/\ell}) \right) \,\mathrm{d}t \\ X_\mathrm{H} &= \frac{1}{2\ell} e^{r/\ell} + \mathcal{O}(e^{-r/\ell}) & \tau = \left( \frac{1}{2\ell} e^{r/\ell} + \mathcal{O}(e^{-r/\ell}) \right) \,\mathrm{d}t \\ X_\mathrm{P} &= 0 & e = \mathrm{d}r \end{split}$$

CJT black holes require positive mass  ${\cal M}$ 

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- $1_{\cdot}$  defined using covariant phase space methods

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  - ► apply Lorentzian PSM definition to Carrollian PSM  $2\pi T := P^{IJ} * (A_I \land A_J)|_{\text{extremal}} \checkmark$
  - define Carrollian surface gravity  $2\pi T := \kappa$  with  $\nabla^{\mu} (e^{\nu} \partial_{\nu} X) |_{\text{extremal}} =: \kappa e^{\mu} \checkmark$
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  - exploit holonomy condition
- 3. know the result is Wald entropy (so that dE = T dS)

$$S = kX|_{\text{extremal}}$$

but not yet how to derive this from first principles ?

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Apply definitions to CJT black holes  $_{\text{set }\ell\,=\,1}$ 

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## Thermal properties

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$$\mathrm{d} E = T \ \mathrm{d} S$$

evidence suggests: CJT black holes are thermal states with (large) entropy

Assessment of Carrollian extremal surface definition

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No final verdict on Carrollian black holes but incentive to continue research

