



# SUSY, Dark Matter and Dark Sector at future $e^+e^-$ collider (CEPC/FCC-ee)

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Peking University

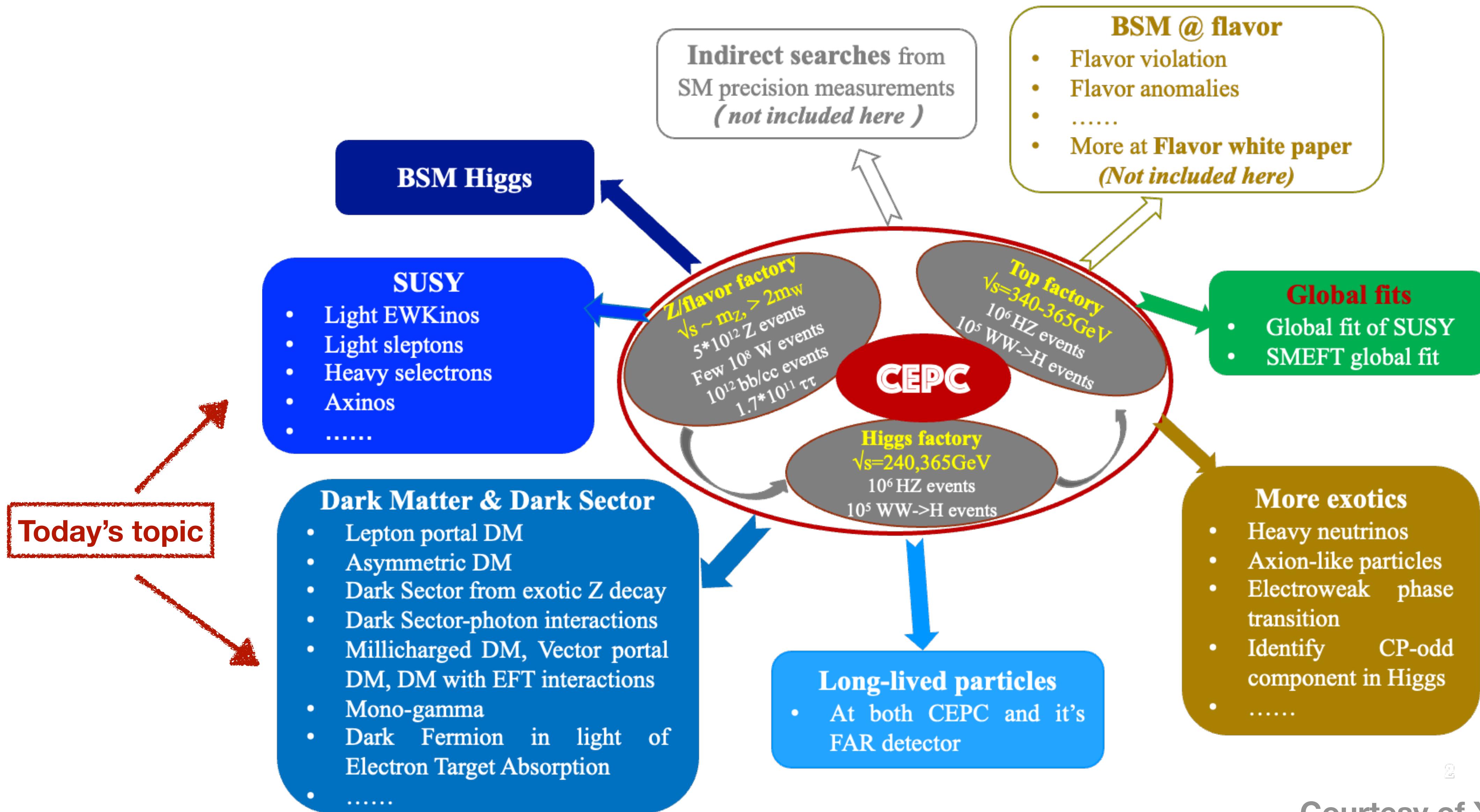
Bias warning: mostly on recent findings of the CEPC study, but work for FCC-ee as well.  
Apologize for omitting papers, can not cover all...

The 2023 international workshop on the Circular Electron Positron Collider  
@ the University of Edinburgh  
2023-07-04

# Outline

- CEPC and New Physics Searches
  - SUSY
  - Dark Matter and Dark Sector
    - Fermion portal – lepton portal
    - Higgs portal
    - Vector portal
    - EFT models
  - Summary

# CEPC BSM Physics Program



# BSM inputs and status

■ BSM Higgs (1709.06103; 1808.02037; 1912.01431; 2008.05492; 2011.04540)

■ SUSY Searches

- Direct SUSY Searches (CPC46(2022)013106; 2101.12131; 2203.10580; 2202.11011, 2211.08132)
- Indirect search of SUSY (2010.09782)

■ Dark Matter and Dark Sector searches

- Lepton portal DM (JHEP 06 (2021) 149 )
- Asymmetric DM (PRD 104(2021)055008 )
- Dark Sector from exotic Z decay (1712.07237)
- DM (Millicharged DM, Vector portal DM, DM with EFT interactions): 1903.1211
- Mono-gamma (2205.05560), Dark Sector-photon interactions (2208.08142)
- Dark Fermion in light of Electron Target Absorption (2306.00657)

■ Long-lived particles (1904.10661, 1911.06576, 2201.08960 )

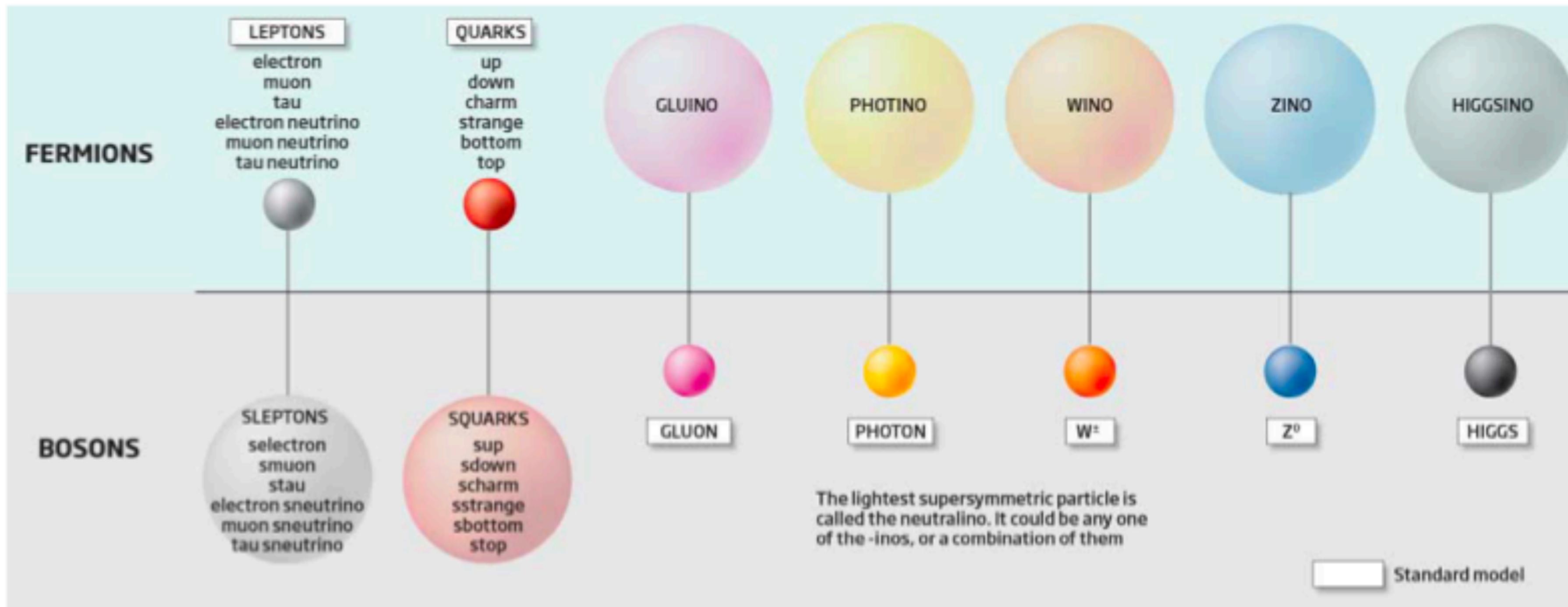
■ More exotics:

- Heavy neutrinos (2102.12826, 2201.05831);
- Axion-like particles (2103.05218, 2204.04702, 2210.09335, [J. Phys. G](#))
- Electroweak phase transition (1911.10210, 1911.10206, 2011.04540, 2204.05085)
- Identify CP-odd component in Higgs study (2212.05390)
- .....

■ Global fits:

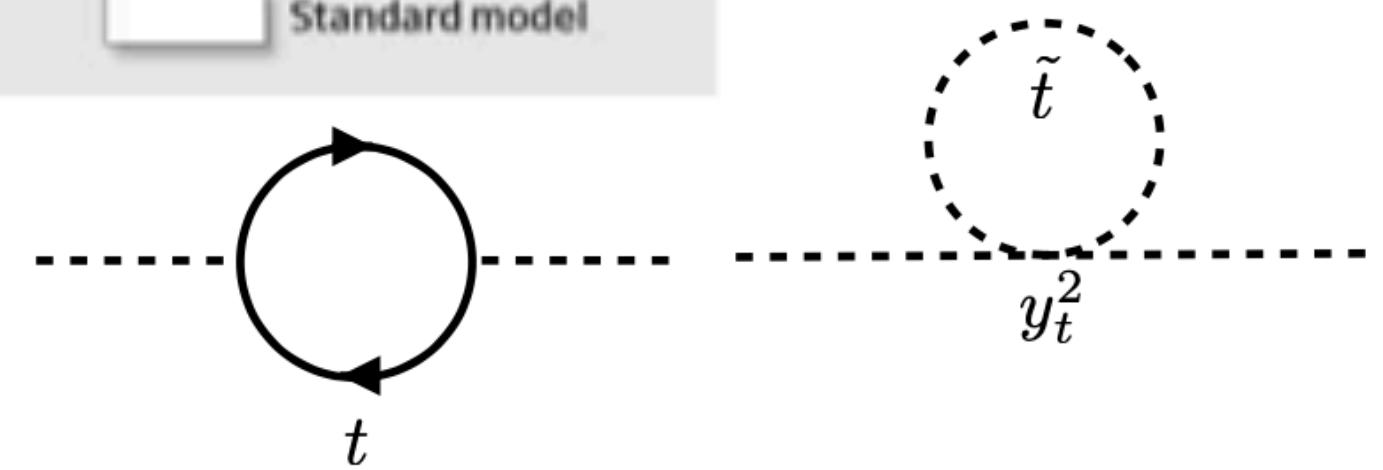
- Global fit of SUSY (2203.04828, 2203.07883)
- SMEFT global fit (2206.08326)

# The motivation for SUSY



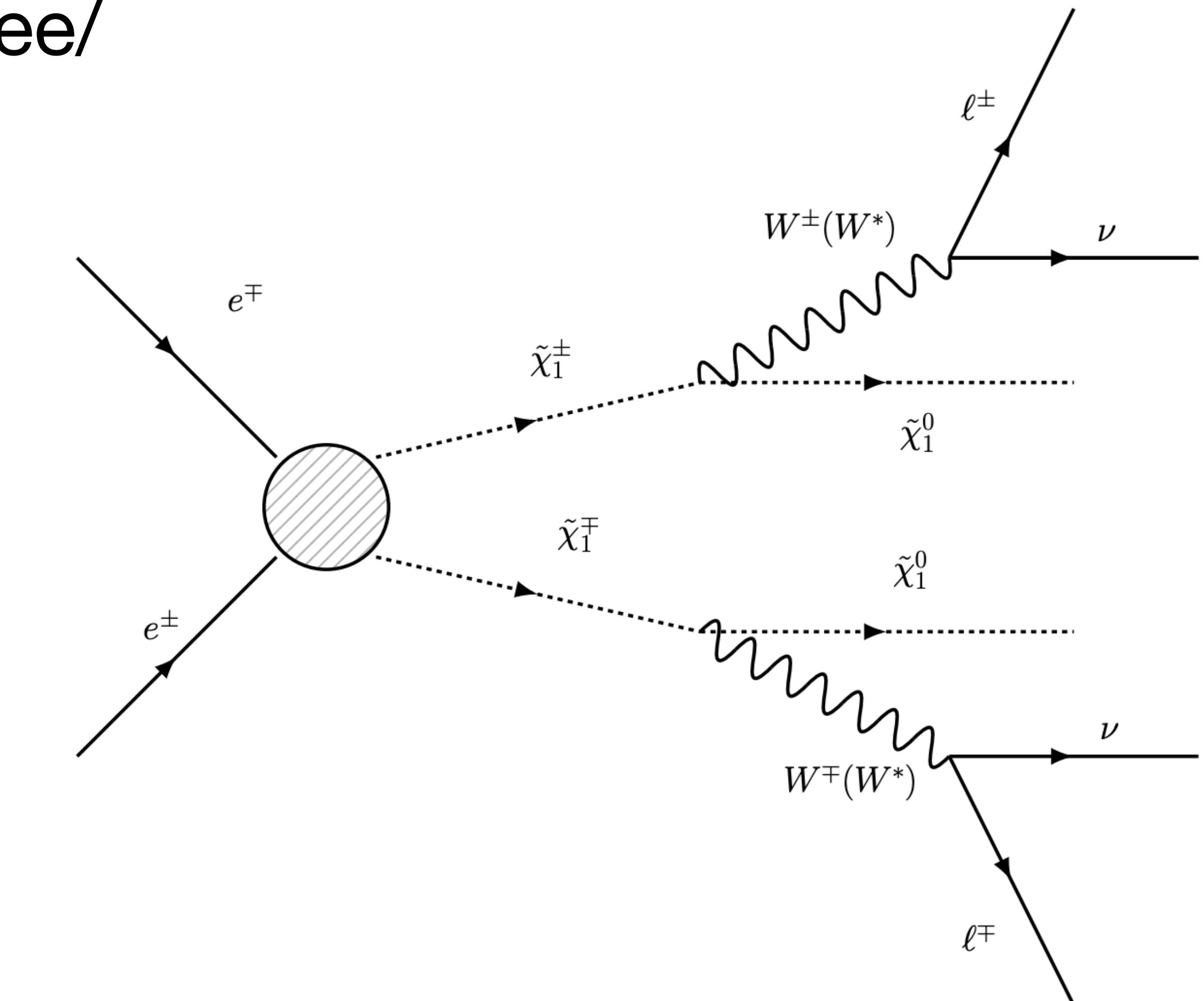
Credit:  
New Scientist

- Stabilize the weak scale from Planck scale via a new space-time symmetry
- MSSM leads to unification of gauge couplings
- EW breaking is induced radiatively
- Neutralino as an excellent candidate for cold dark matter

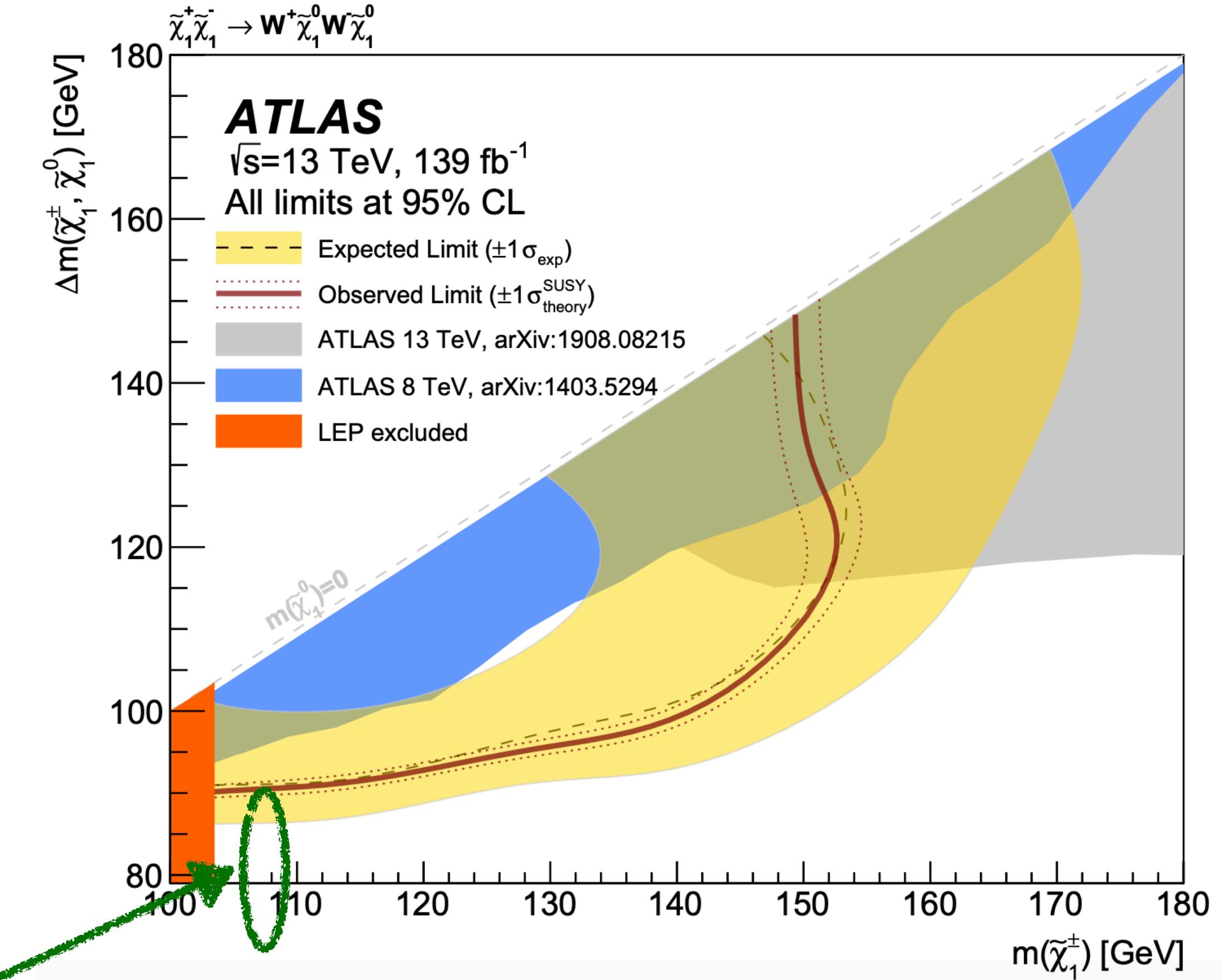
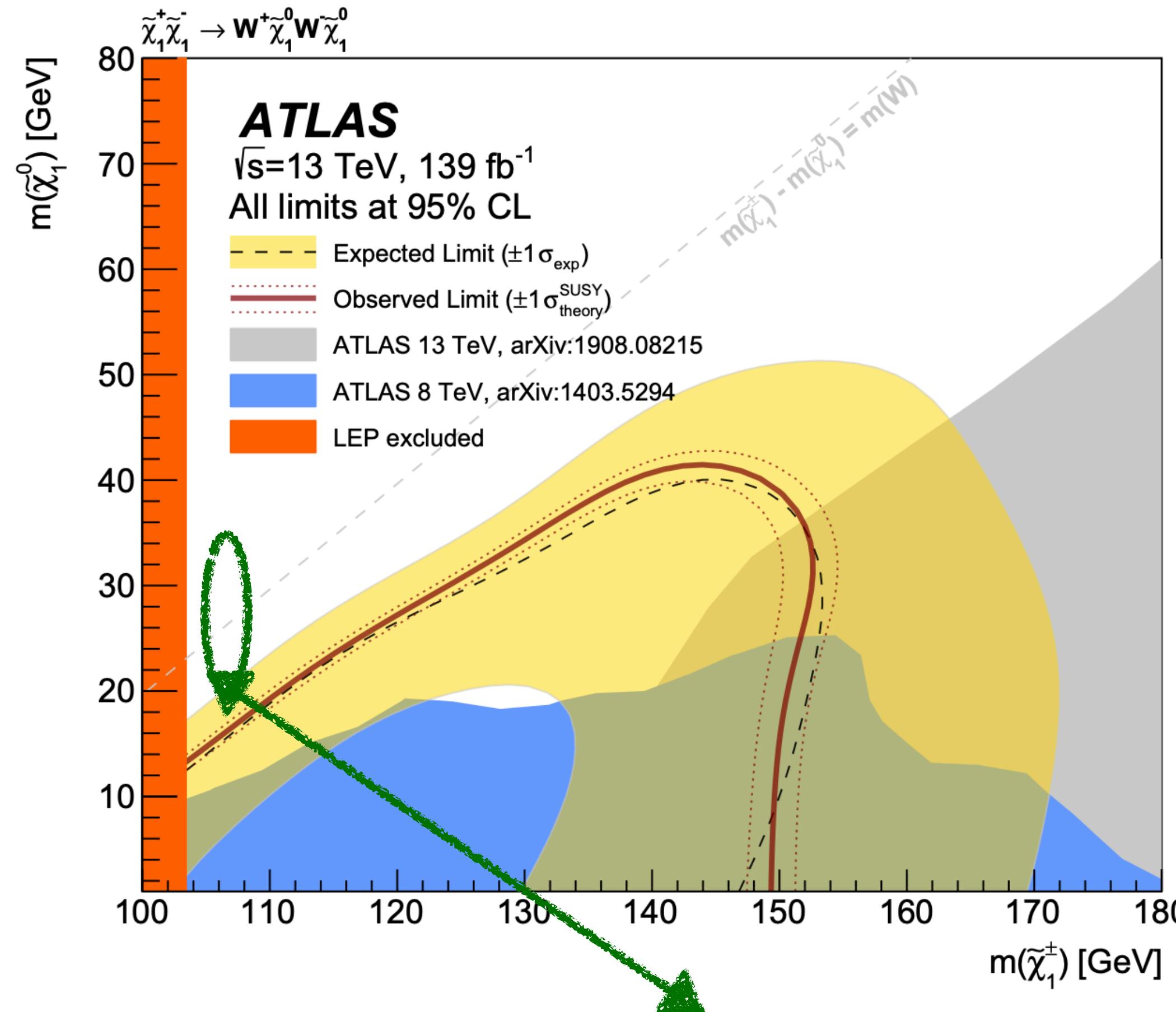


# SUSY at future $e^+e^-$ collider

- Why future  $e^+e^-$  collider CEPC/FCC-ee/ILC?
- Could be complementary with LHC:
  - Soft energy region
  - Lower mass region
- Study light EWKinos, sleptons
- Indirect searches through precision measurements

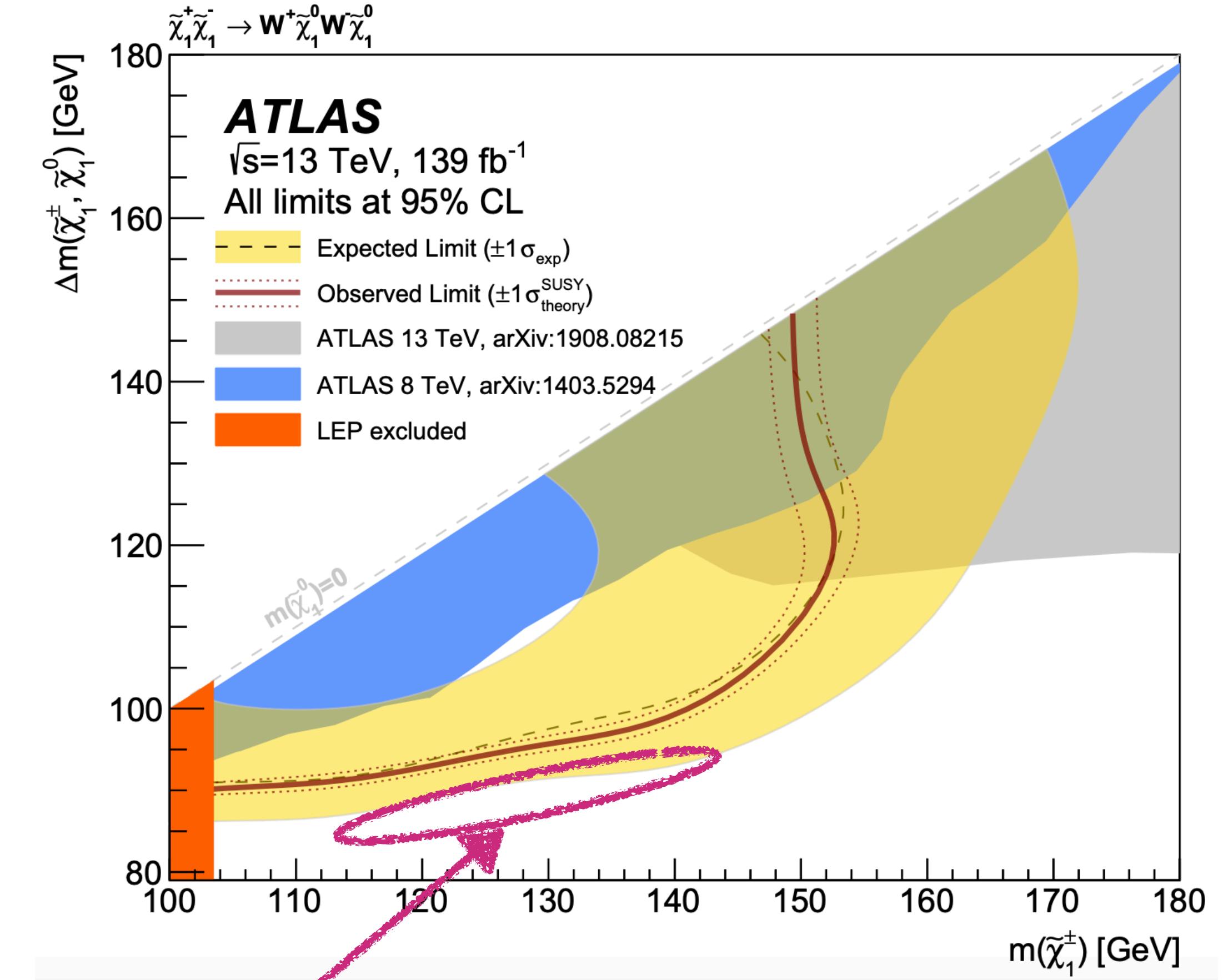
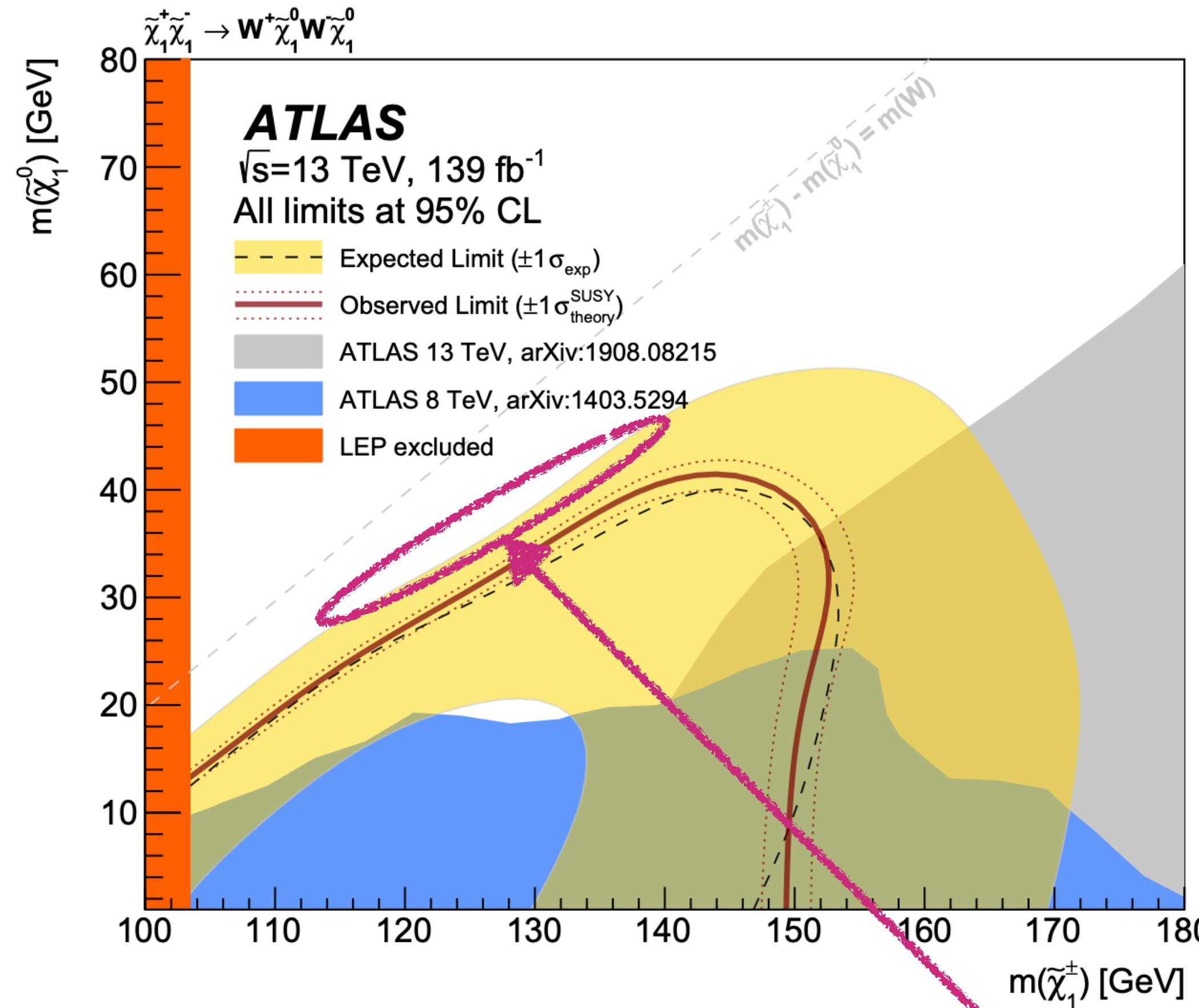


# Comparing with recent LHC results: where to go



- Improving low mass boundary: naively up to  $\sqrt{s}/2$ , but can be better than that
- Improving the soft lepton region

# Comparing with recent LHC results: where to go



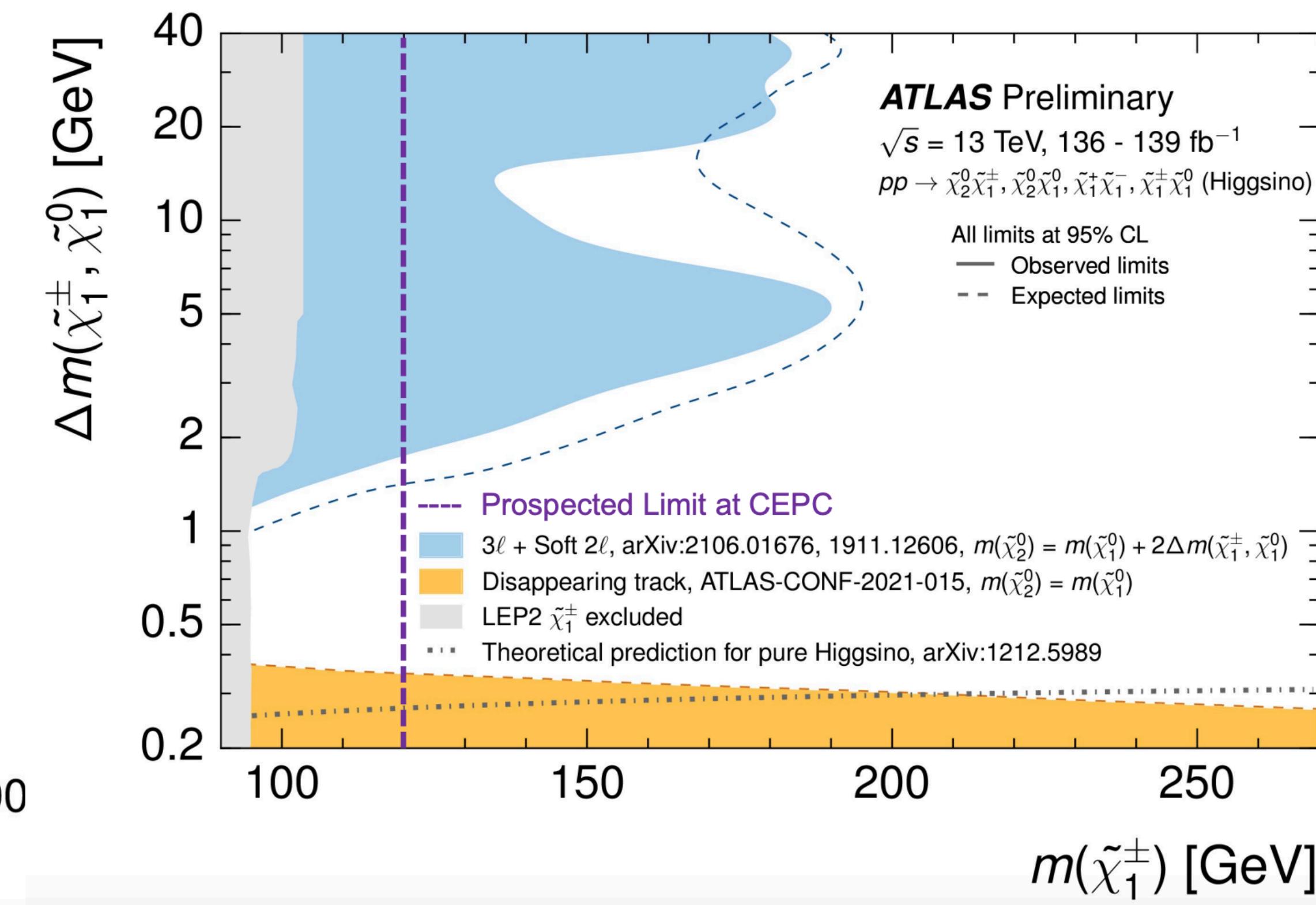
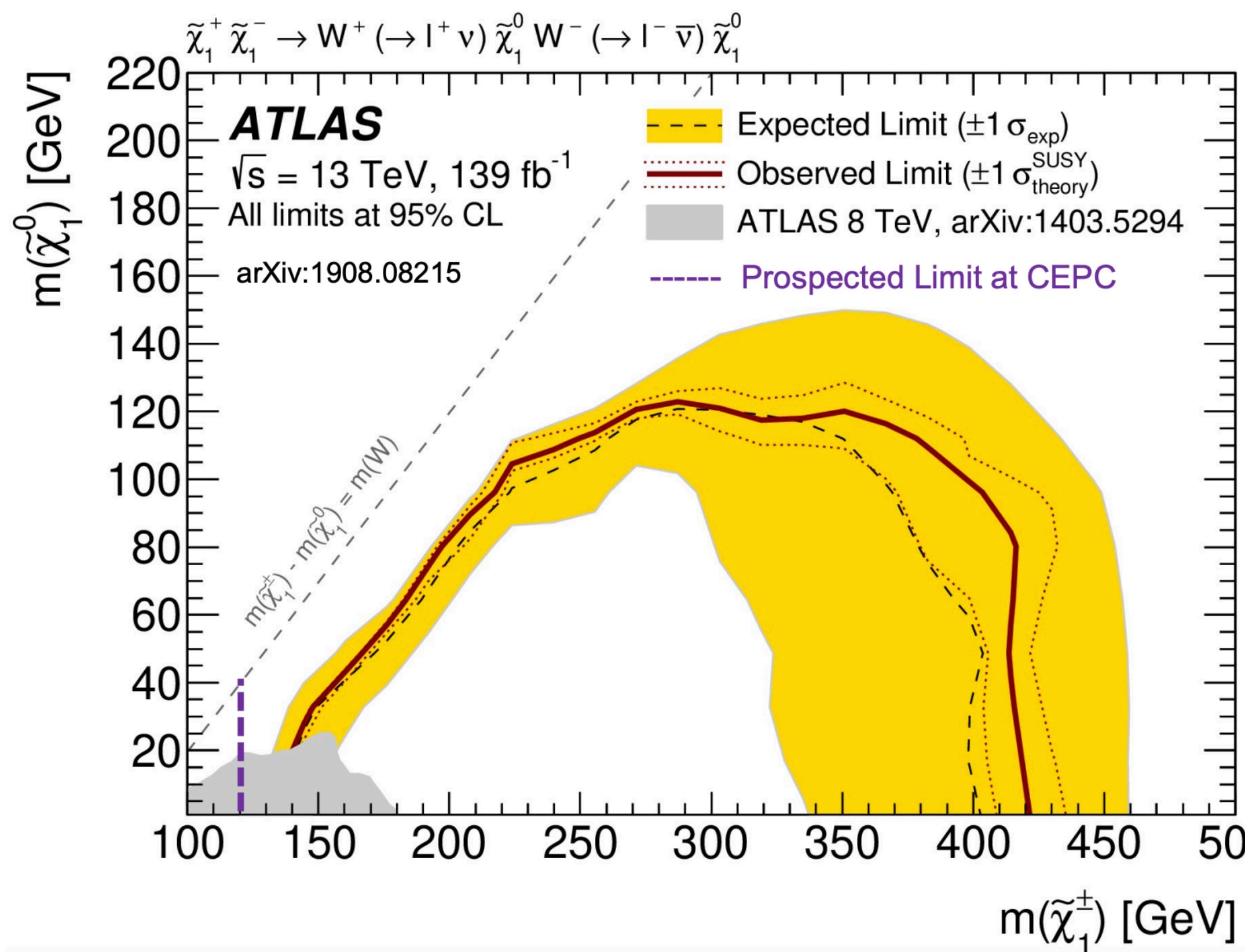
- Improving low mass boundary
- Improving the soft lepton region

# Dilepton + Missing energy signal region

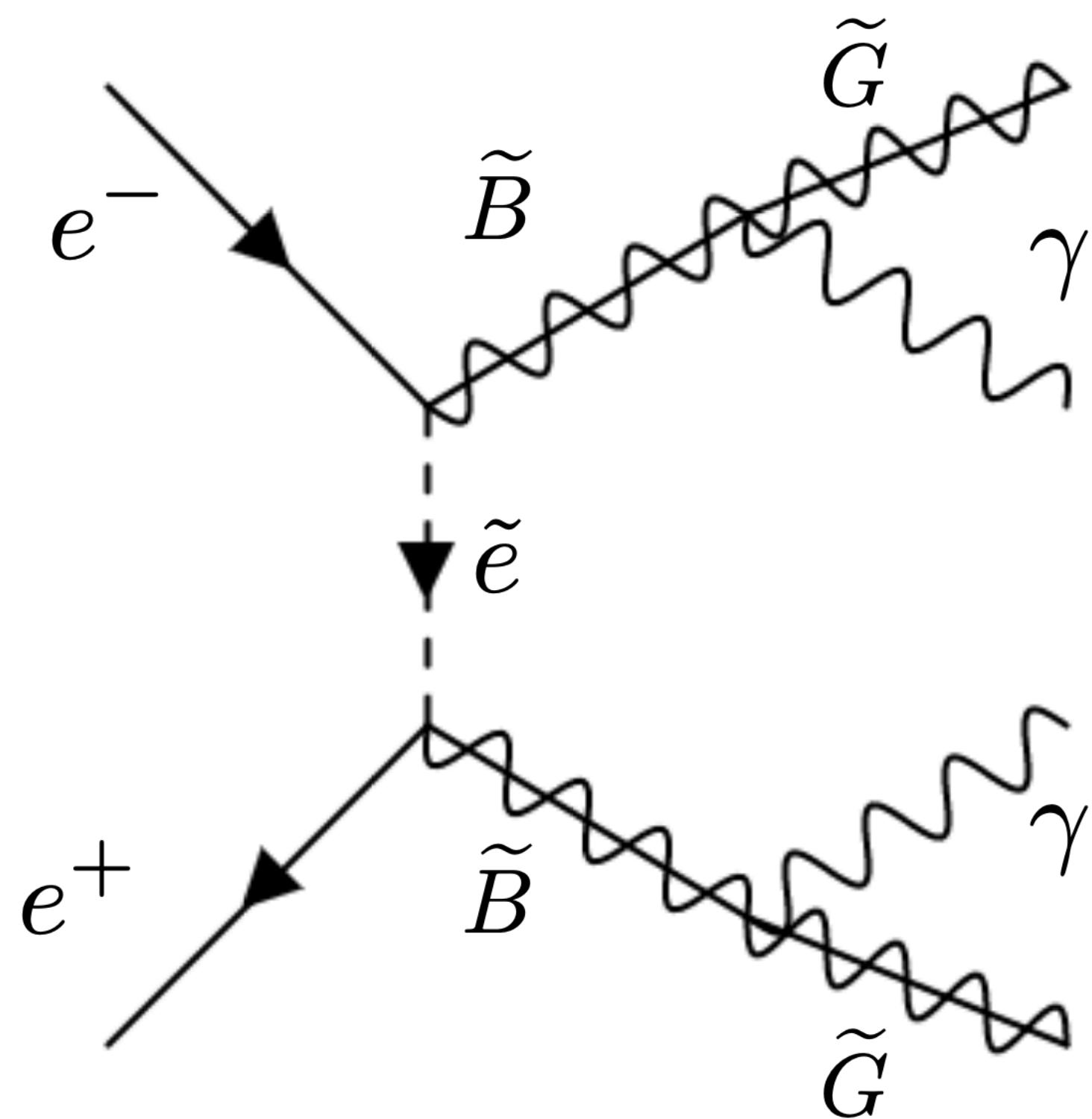
- Focus on opposite-sign dimuon
- Large recoil mass  $\sim$  Large missing energy
- Up to mass =  $\sqrt{s}/2$

Xuai Zhuang  
et al, 2105.06135 (CPC)

Signal Region
$\equiv$ 2 muons (OS)
$E_{\mu^\pm} > 10 \text{ GeV}$
$0.4 < \Delta R(\mu^+, \mu^-) < 1.6$
$P_T^{\mu^\pm} > 30 \text{ GeV}$
$M_{recoil} > 130 \text{ GeV}$

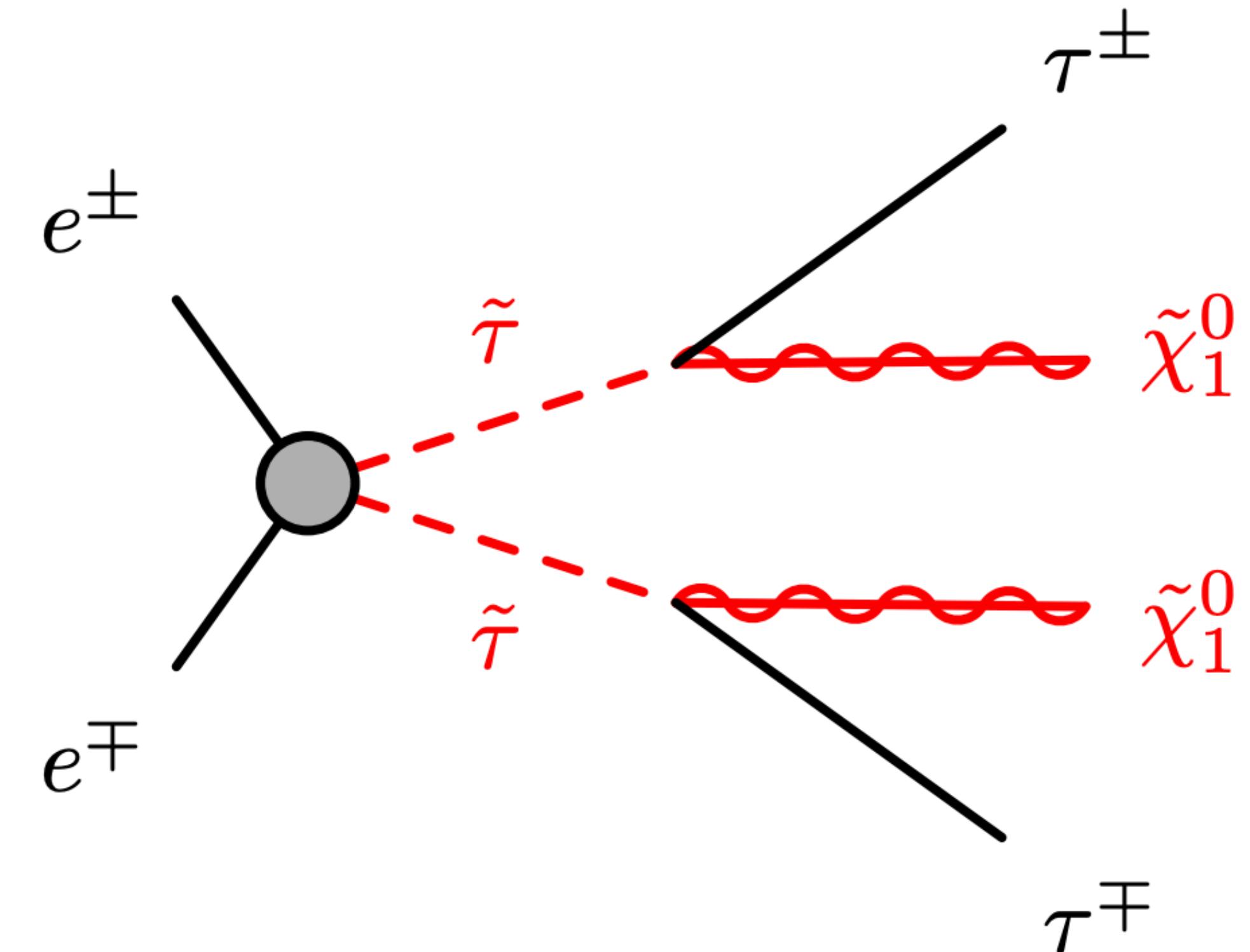


# Other possible complementary searches



- Low mass Bino decay  $\sim O(100)$  GeV:  
 $\tilde{B} \rightarrow \tilde{G} + \gamma$

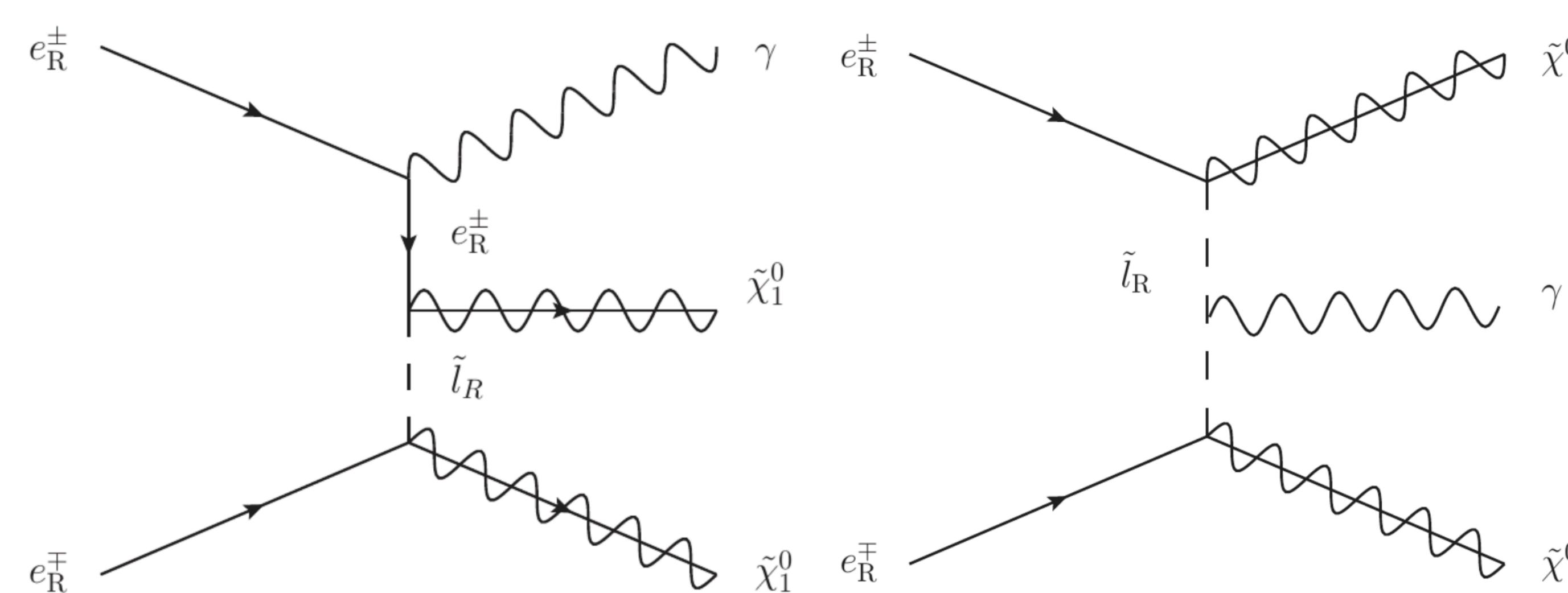
Junmou Chen et al, 2101.12131 (PRD)



- Low mass stau/smuon  $\sim O(100)$  GeV:  $\tilde{\tau} \rightarrow \tilde{\chi}_1^0 + \tau$
- Cleaner background for tau final states comparing with LHC

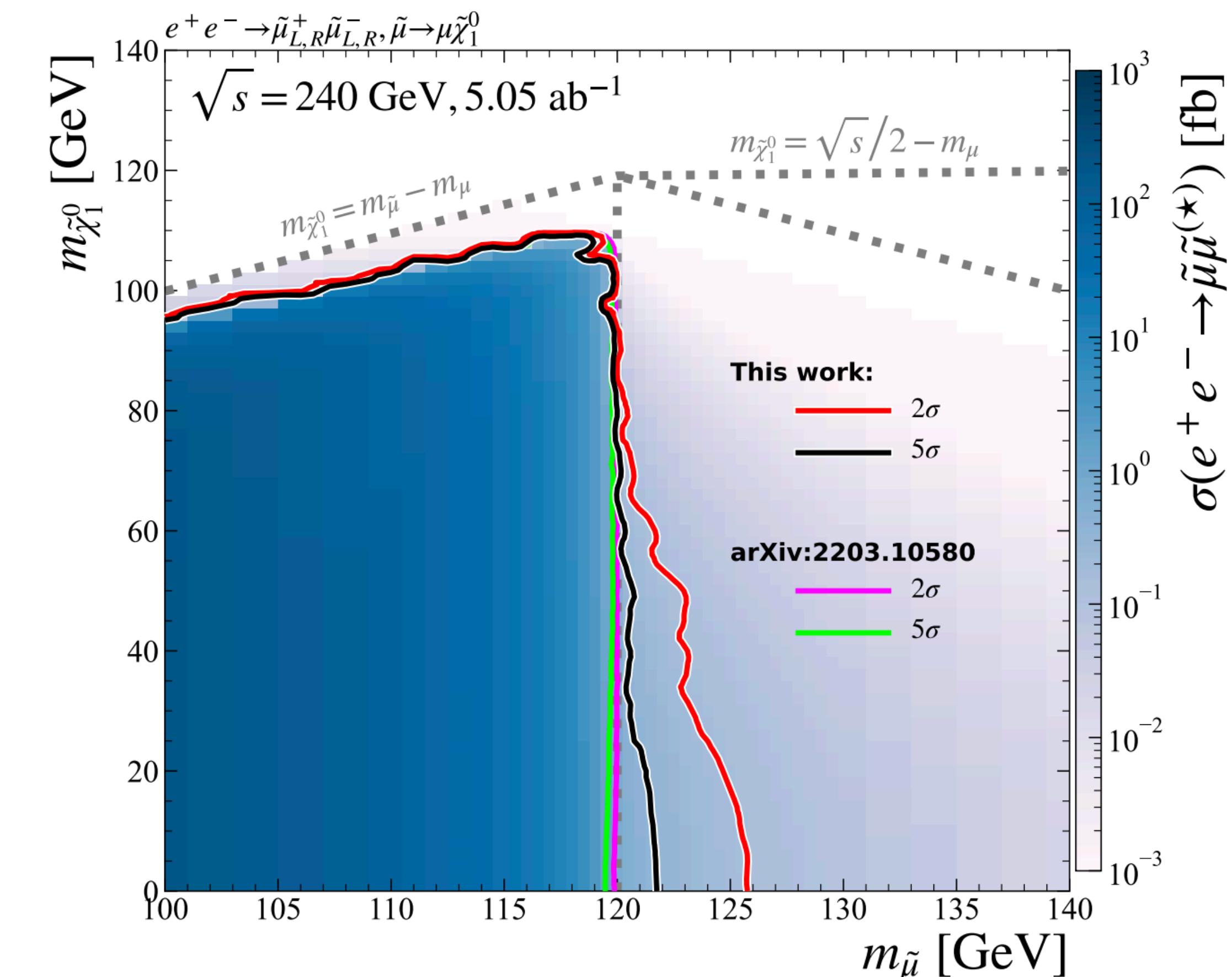
Xuai Zhuang et al, 2203.10580

# Other possible complementary searches



- Mono-photon search:  $\gamma \tilde{\chi}_0^1 \tilde{\chi}_0^1$
- Cleaner background comparing with LHC

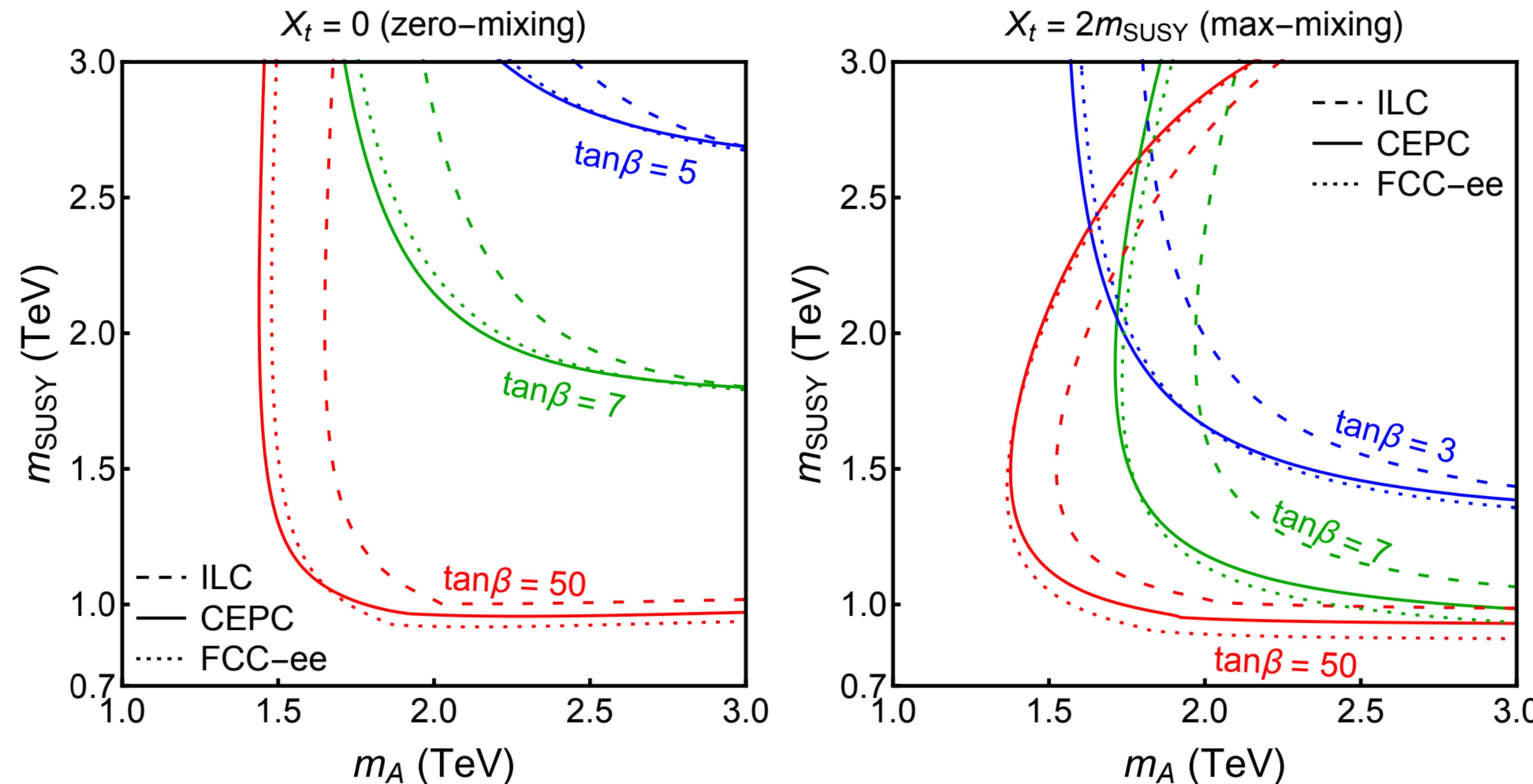
Waqas Ahmed et al, 2202.11011 (PLB)



- Trying to cover higher slepton mass as best as one could

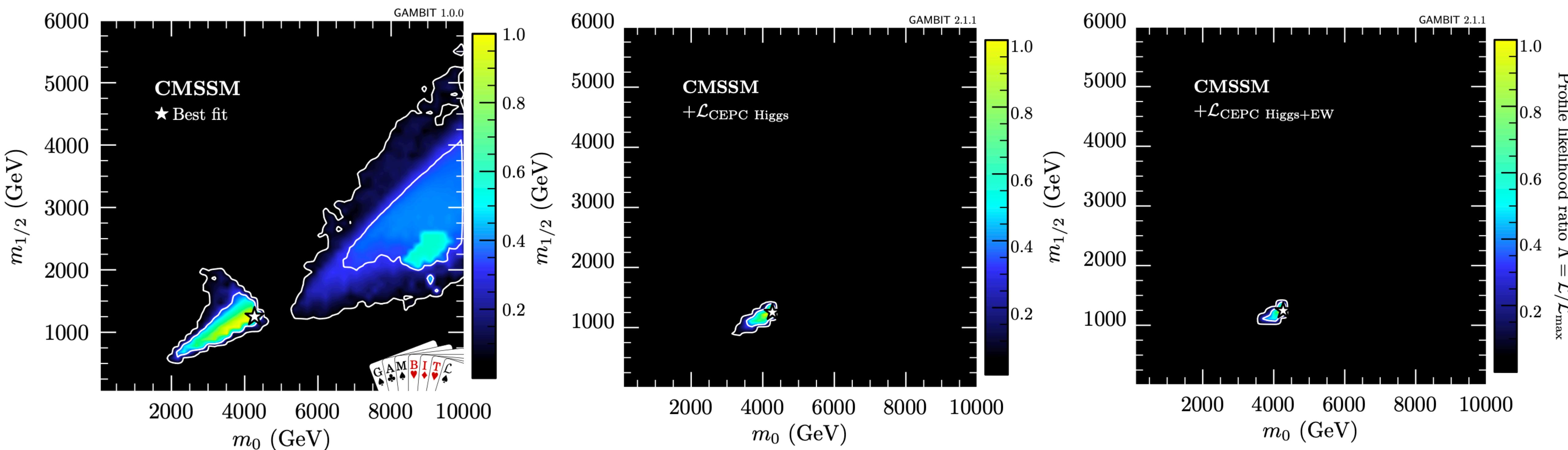
Jinmin Yang et al, 2211.08132

# Other possible complementary searches



- Probing SUSY parameters **without producing SUSY particles**
- Indirect search at Higgs factories through Higgs precision measurements

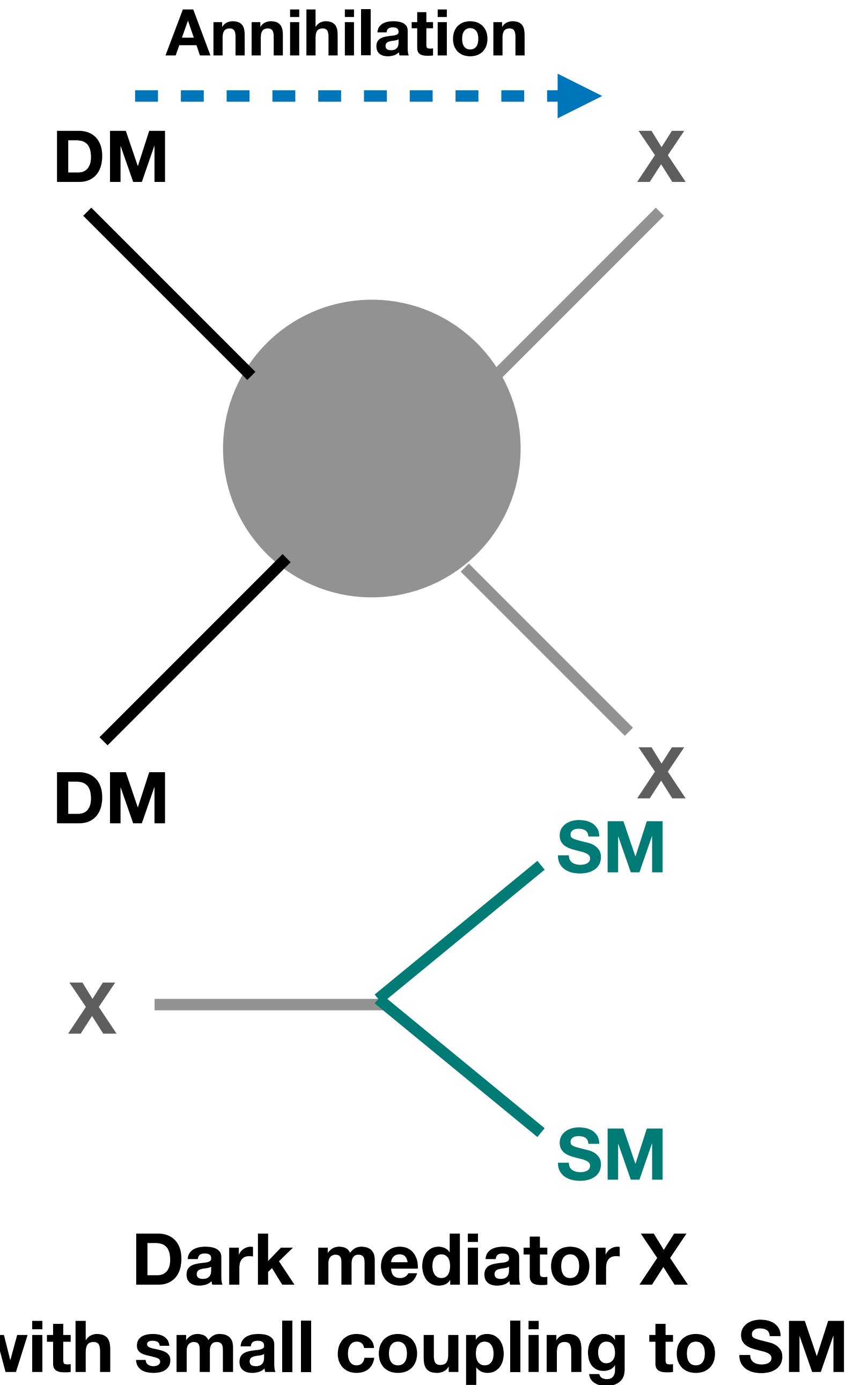
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    - Higgs portal
    - Vector portal
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# Fermion portal

JL, S. Brian, N. Weiner, I. Yavin, 1303.4404 (JHEP)  
Y. Bai, J. Berger, 1308.0612 (JHEP)

- The Fermion portal to DM
- The Interaction:  $\lambda \bar{\psi} \phi f_{\text{SM}}$
- Dark Matter and Dark Sector
  - Fermion portal – lepton portal
  - $f$  can be quark-lepton, L/R-handed

	DM	SM charge	$Z_2$
$\psi$	Yes	No	-1
$\phi$	No	Yes	-1
$f_{\text{SM}}$	No	Yes	1

$\phi$  can be DM as well!

# Searching lepton portal dark sector at CEPC

$$\mathcal{L}_\chi = \frac{1}{2}\bar{\chi}i\not{\partial}\chi - \frac{1}{2}m_\chi\bar{\chi}\chi + y_\ell (\bar{\chi}_L S^\dagger \ell_R + \text{h.c.}) ,$$

JL, XP Wang, KP Xie, 2104.06421 (JHEP)

$$\mathcal{L}_S = (D^\mu S)^\dagger D_\mu S - V(H, S),$$

$$V(H, S) = \mu_H^2 |H|^2 + \mu_S^2 |S|^2 + \lambda_H |H|^4 + \lambda_S |S|^4 + 2\lambda_{HS} |H|^2 |S|^2$$

- DM ( $\chi$ ) couples to SM via the lepton portal
- Mediated by charged particle ( $S$ ), similar to slepton
- DM thermal relic requirements
- Lepton collider production of dark sector particle S
- Higgs precision test on the model
- Gravitational Wave signal and its complementary with ee collider

# Searching lepton portal dark sector at CEPC

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- DM ( $\chi$ ) couples to SM via the lepton portal
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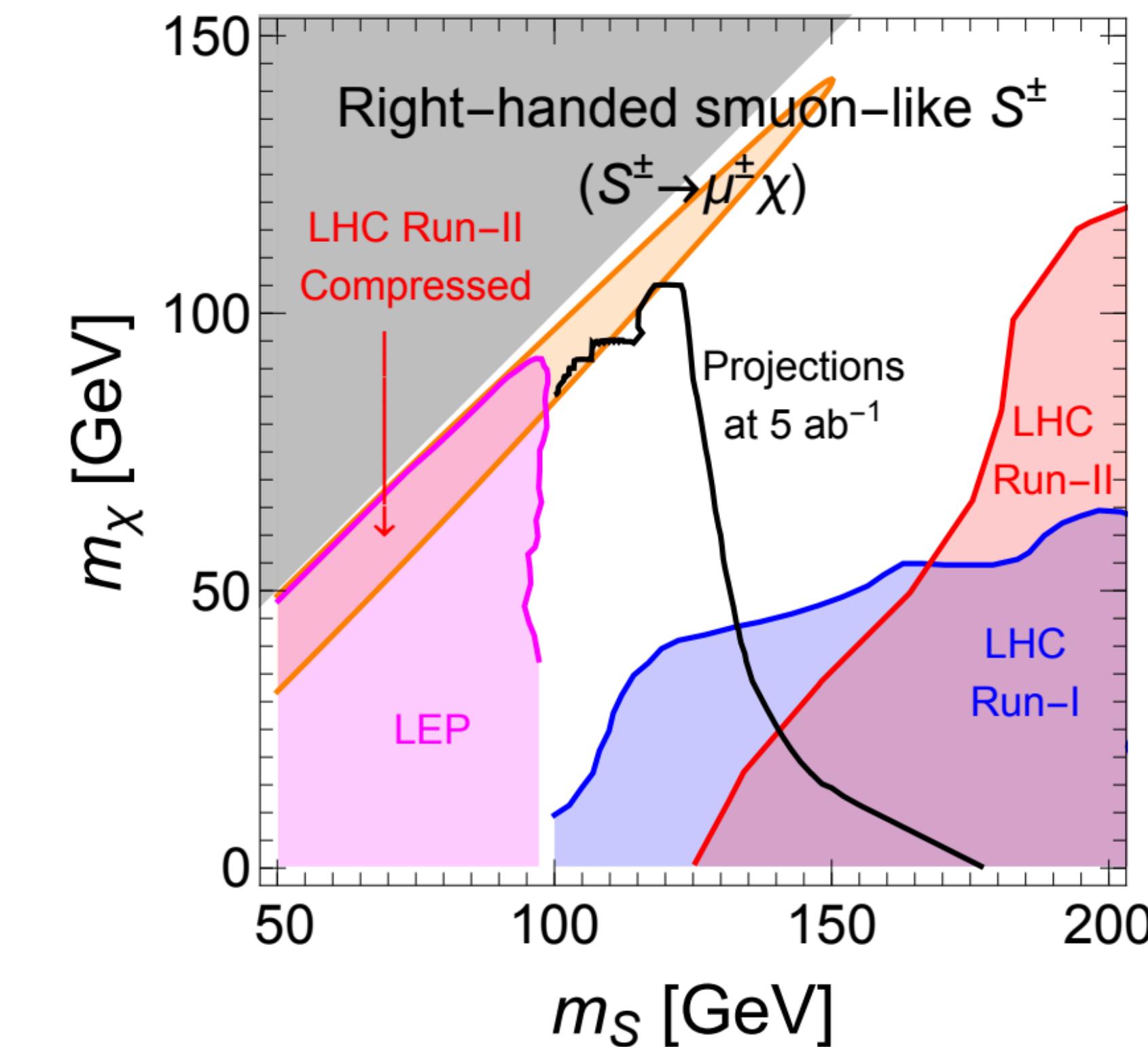
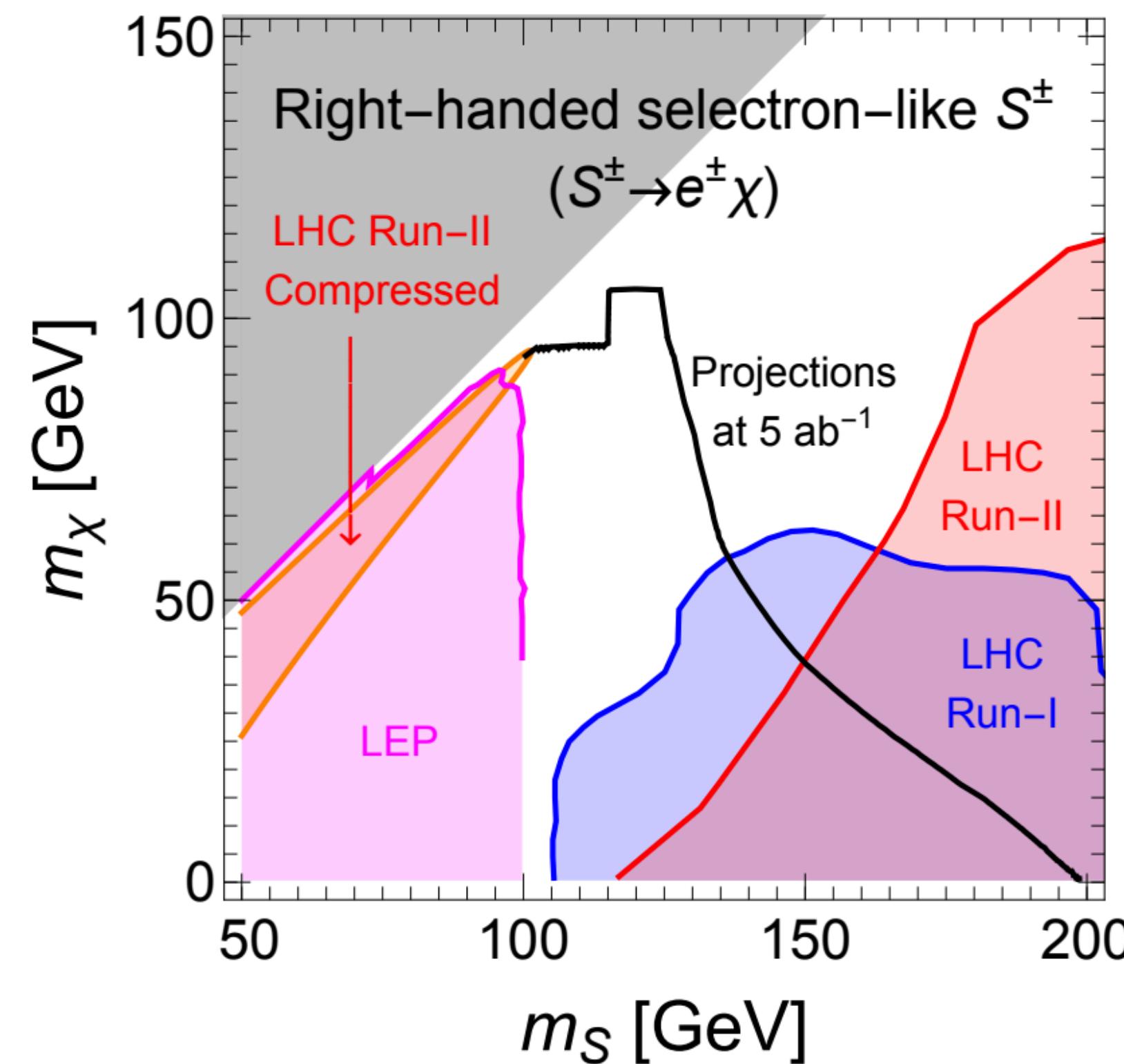
$\chi\chi \rightarrow \ell^+\ell^-$ , p-wave annihilation,  
less constrained by indirect detection

# Dark sector particle production in 3-body final state

- 3-body final state:  $e^+e^- \rightarrow S^+S^{-*} \rightarrow S^+\ell^-\chi \rightarrow (\ell^+\chi)\ell^-\chi$

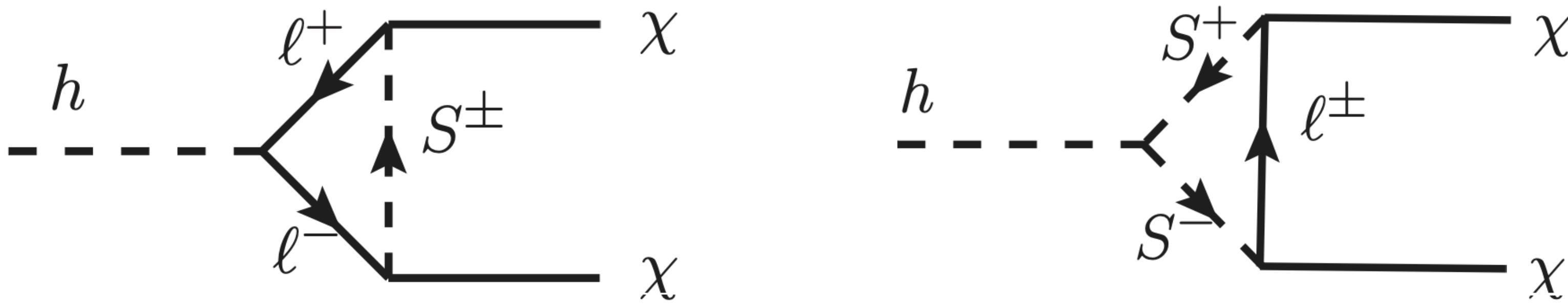
JL, XP Wang, KP Xie, 2104.06421 (JHEP)

- Reaching higher mass:  $m_S \gtrsim \sqrt{s}/2 = 120$  GeV

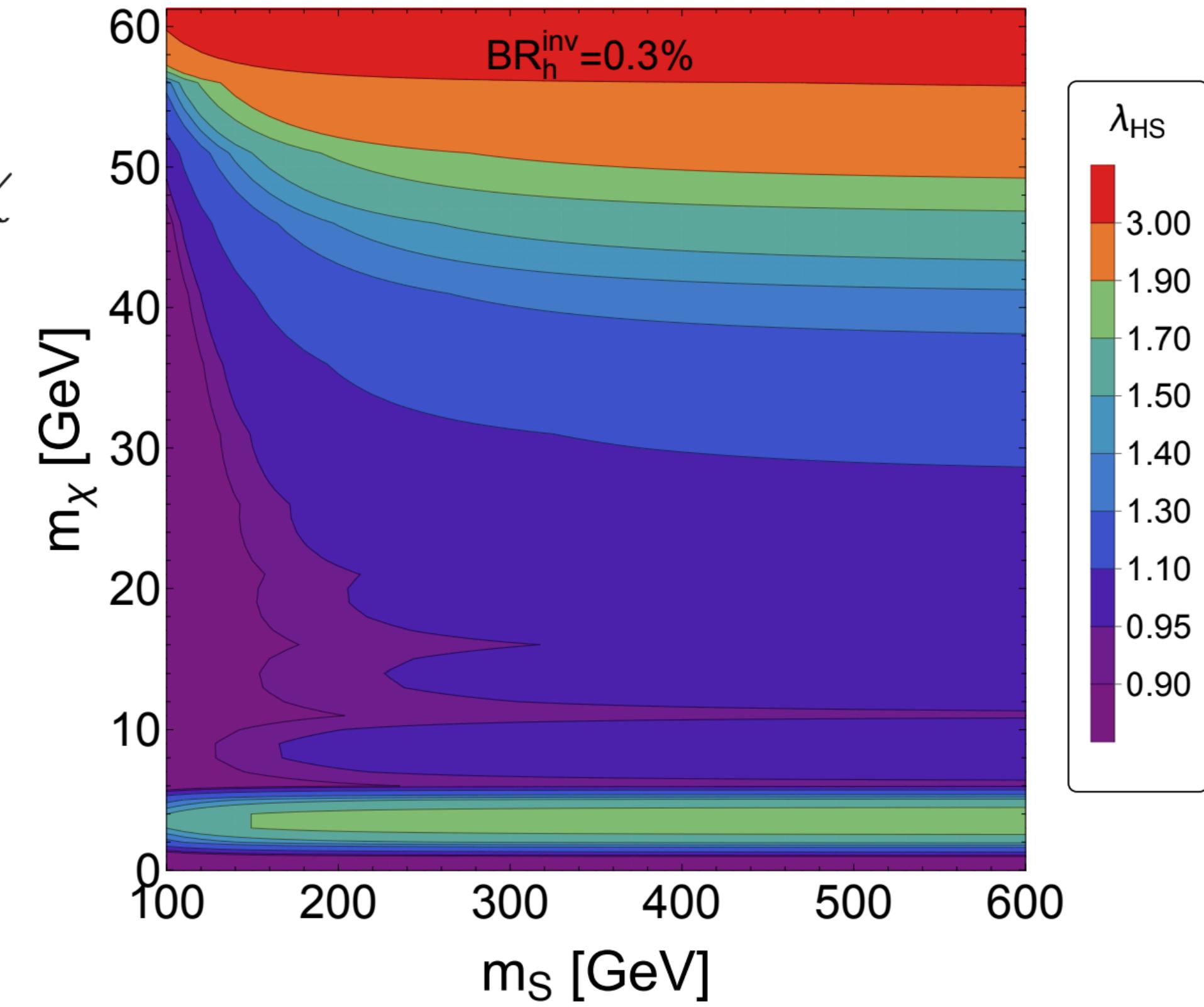


# Exotic Higgs decay constraints

- Invisible Higgs decay at 1-loop



- The Higgs-dark sector coupling is constrained



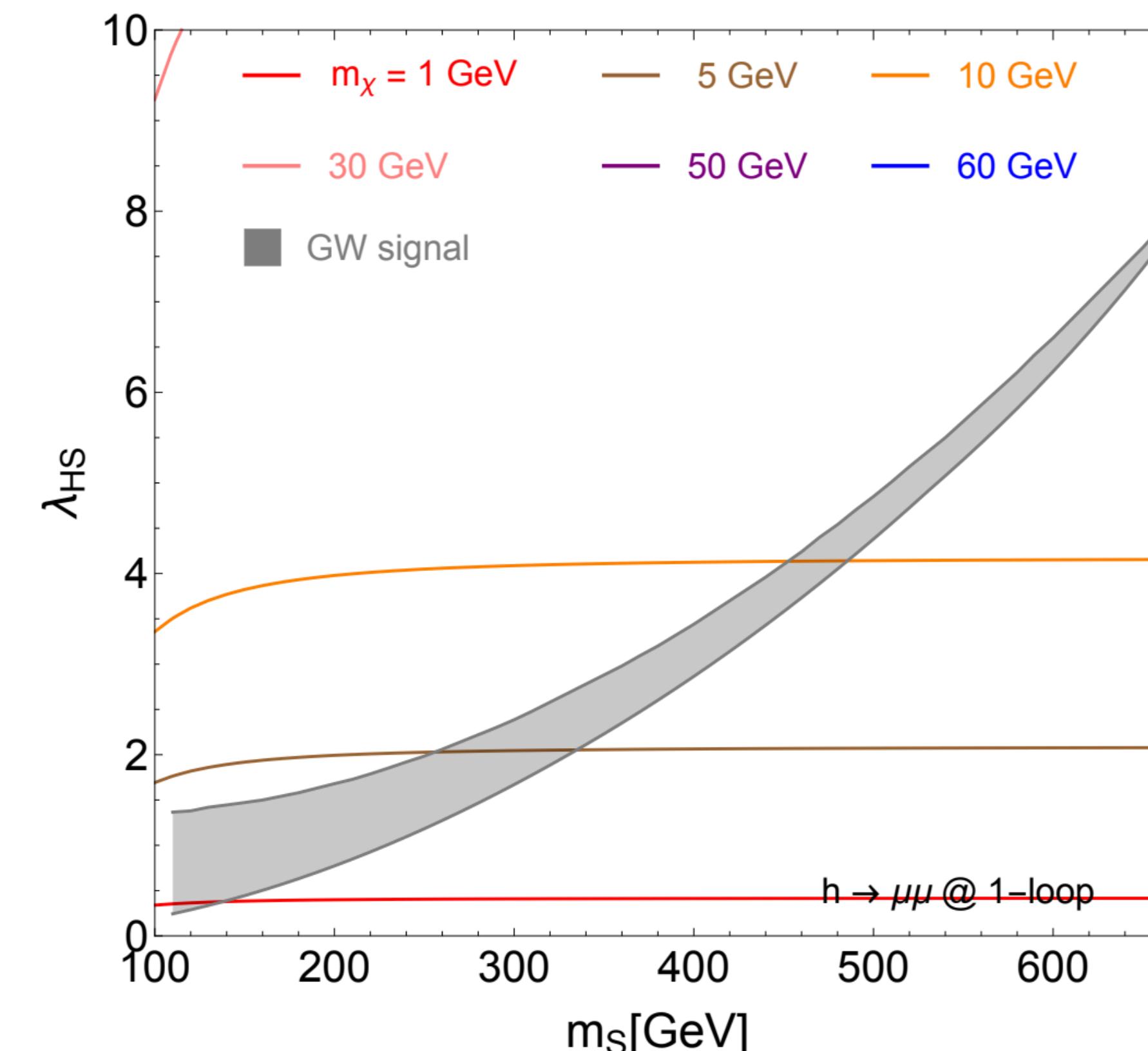
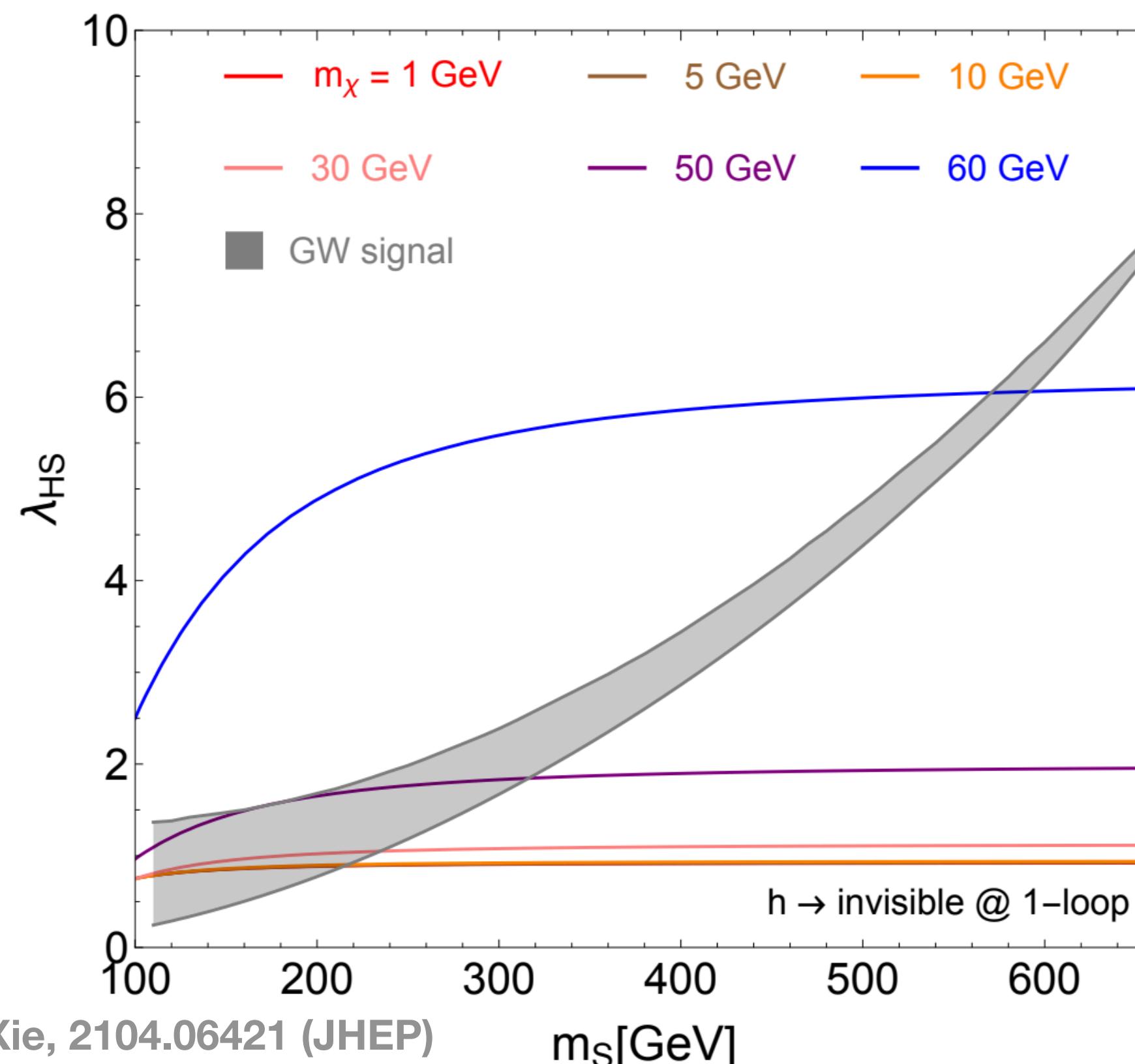
# Future ee collider search and GW complementarity

- Higgs precision measurement can cover 1st order phase transition region,  $m_S$ , up to hundred of GeV

$$\mathcal{L}_\chi = \frac{1}{2}\bar{\chi}i\not{\partial}\chi - \frac{1}{2}m_\chi\bar{\chi}\chi + y_\ell (\bar{\chi}_L S^\dagger \ell_R + \text{h.c.}),$$

$$\mathcal{L}_S = (D^\mu S)^\dagger D_\mu S - V(H, S),$$

$$V(H, S) = \mu_H^2 |H|^2 + \mu_S^2 |S|^2 + \lambda_H |H|^4 + \lambda_S |S|^4 + 2\lambda_{HS} |H|^2 |S|^2$$



# Asymmetric dark matter connecting to lepton portal

Mengchao Zhang, 2104.06988 (PRD)

- A model for (dark) baryogenesis via leptogenesis
- The dark matter is dark baryon and asymmetric

	$SU(3)'$	$SU(3)$	$U_Y(1)$	Spin	$L$	$B$	$B'$
$N_1/N_2$	1	1	0	1/2	0	0	0
$\Phi$	3	1	1	0	-1	0	1/3
$\chi$	3	1	1	1/2	-1	0	1/3
$q'$	3	1	0	1/2	0	0	1/3
$l_R$	1	1	-1	1/2	1	0	0
$d_R$	1	3	-1/3	1/2	0	1/3	0
$u_R$	1	3	2/3	1/2	0	1/3	0

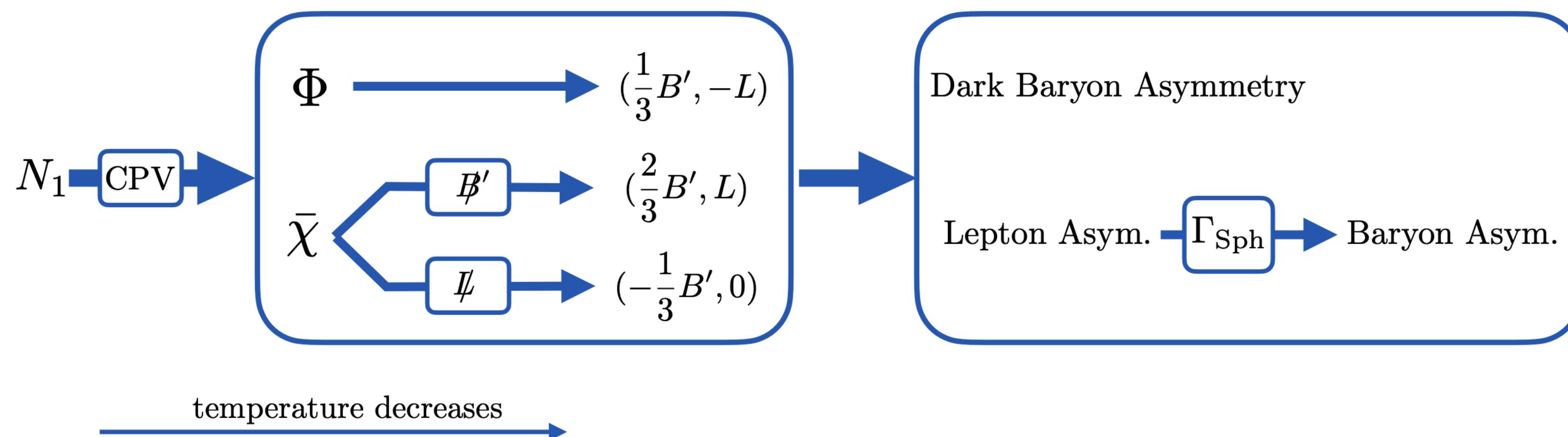
$$\begin{aligned} \mathcal{L} = & \mathcal{L}_{\text{SM}} - \frac{1}{2} \sum_{i=1,2} M_{N_i} \bar{N}_i N_i^C - m_\Phi^2 \Phi^\dagger \Phi - m_\chi \bar{\chi} \chi - m_{q'} \bar{q}' q' + \mathcal{L}_{\text{kinetic}} \\ & - \sum_{i=1,2} \lambda_i \bar{N}_i \chi \Phi^\dagger - \kappa \Phi \bar{q}'_L l_R - \frac{1}{\Lambda_1^2} (\bar{q}'^C \chi) (\bar{q}'^C_L l_R) - \frac{1}{\Lambda_2^2} (\bar{\chi} \gamma^\mu q') (\bar{d}_R \gamma_\mu u_R) + h.c. \end{aligned}$$

# Dark matter and Baryon asymmetry

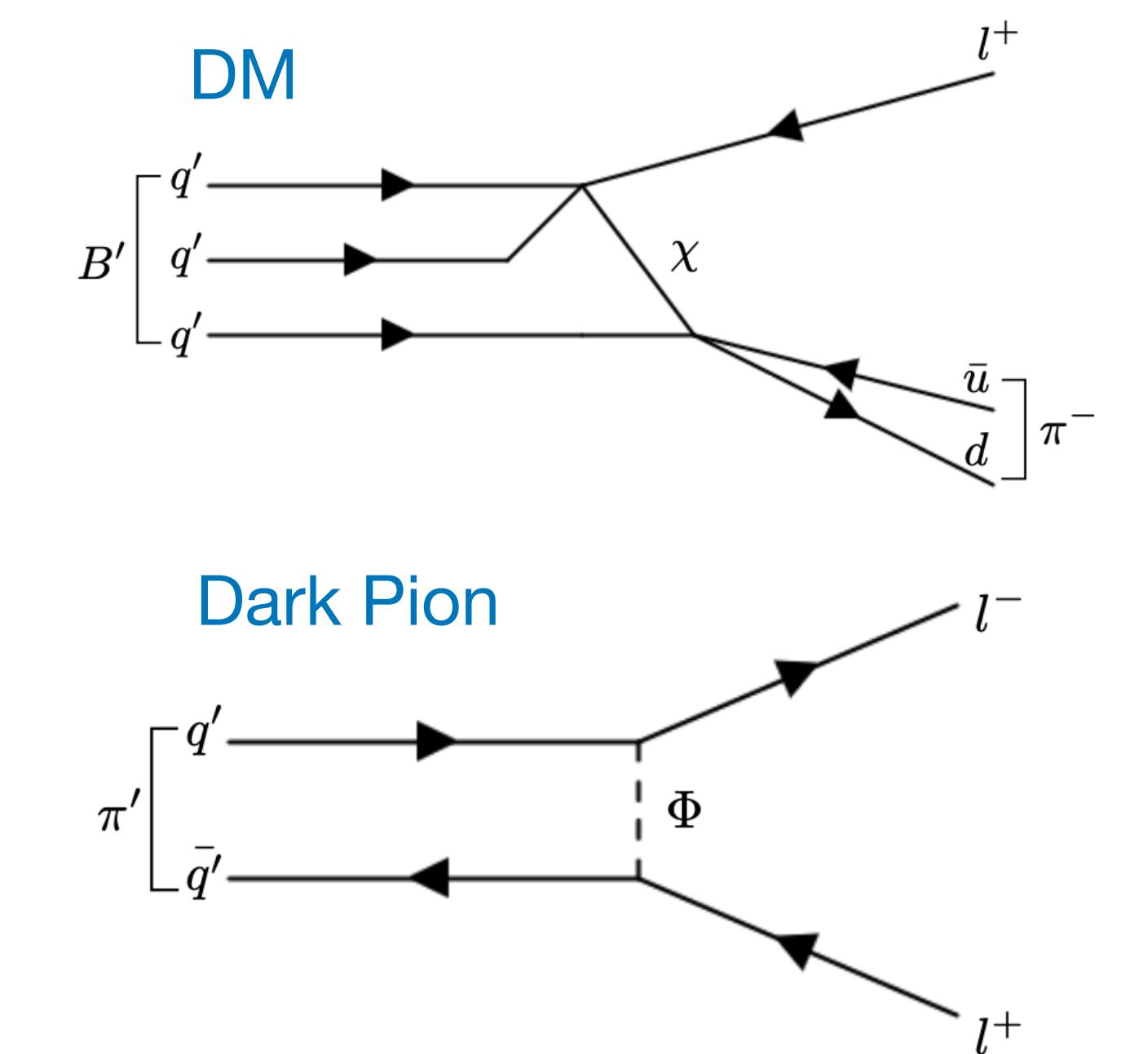
$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{2} \sum_{i=1,2} M_{N_i} \bar{N}_i N_i^C - m_\Phi^2 \Phi^\dagger \Phi - m_\chi \bar{\chi} \chi - m_{q'} \bar{q}' q' + \mathcal{L}_{\text{kinetic}}$$

$$- \sum_{i=1,2} \lambda_i \bar{N}_i \chi \Phi^\dagger - \kappa \Phi \bar{q}'_L l_R - \frac{1}{\Lambda_1^2} (\bar{q}'^C \chi) (\bar{q}'_L l_R) - \frac{1}{\Lambda_2^2} (\bar{\chi} \gamma^\mu q') (\bar{d}_R \gamma_\mu u_R) + h.c.$$

$$\Delta B' \neq 0 \quad \Delta L \neq 0$$



	$SU(3)'$	$SU(3)$	$U_Y(1)$	Spin	$L$	$B$	$B'$
$N_1/N_2$	1	1	0	$1/2$	0	0	0
$\Phi$	3	1	1	0	-1	0	$1/3$
$\chi$	3	1	1	$1/2$	-1	0	$1/3$
$q'$	3	1	0	$1/2$	0	0	$1/3$
$l_R$	1	1	-1	$1/2$	1	0	0
$d_R$	1	3	- $1/3$	$1/2$	0	$1/3$	0
$u_R$	1	3	$2/3$	$1/2$	0	$1/3$	0



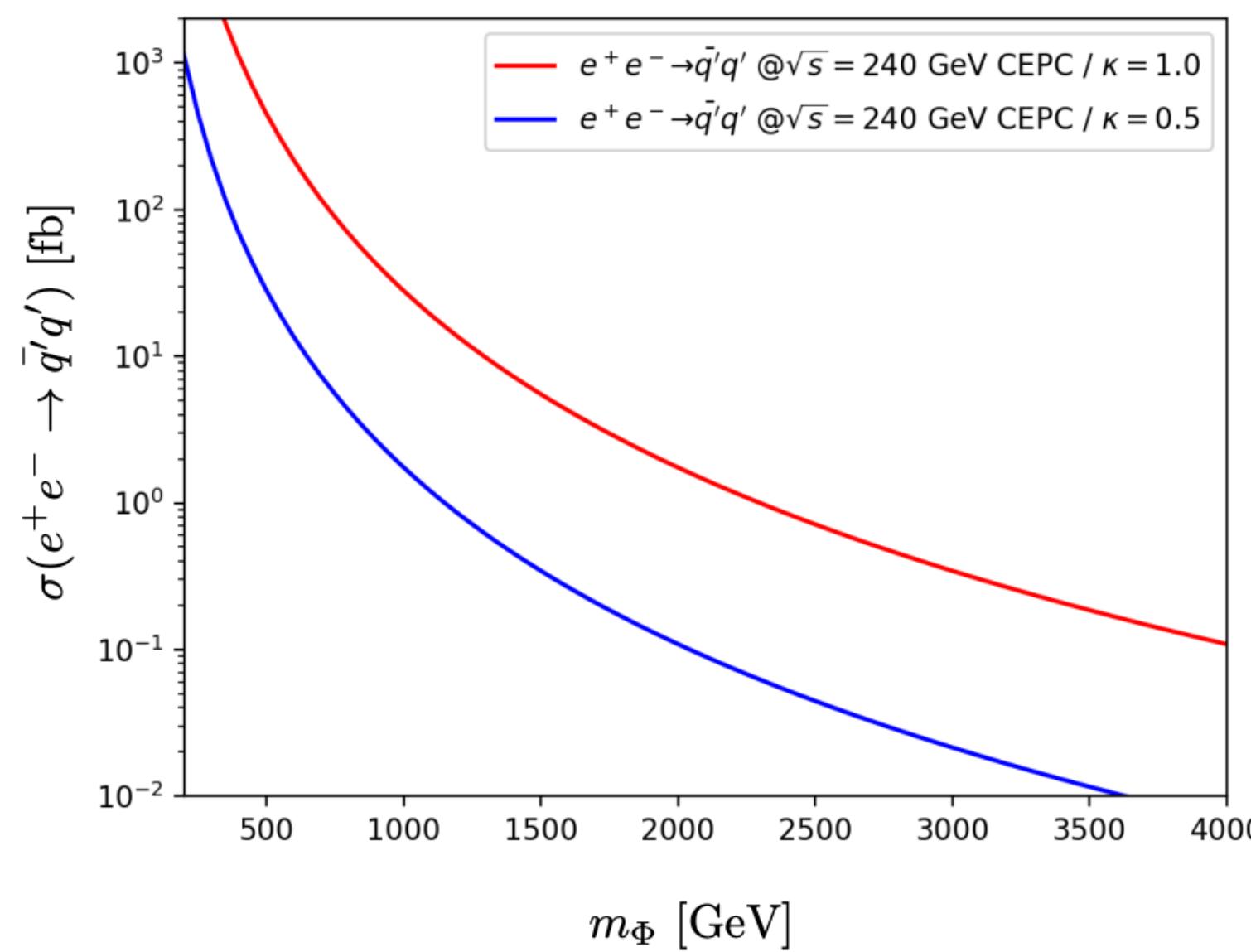
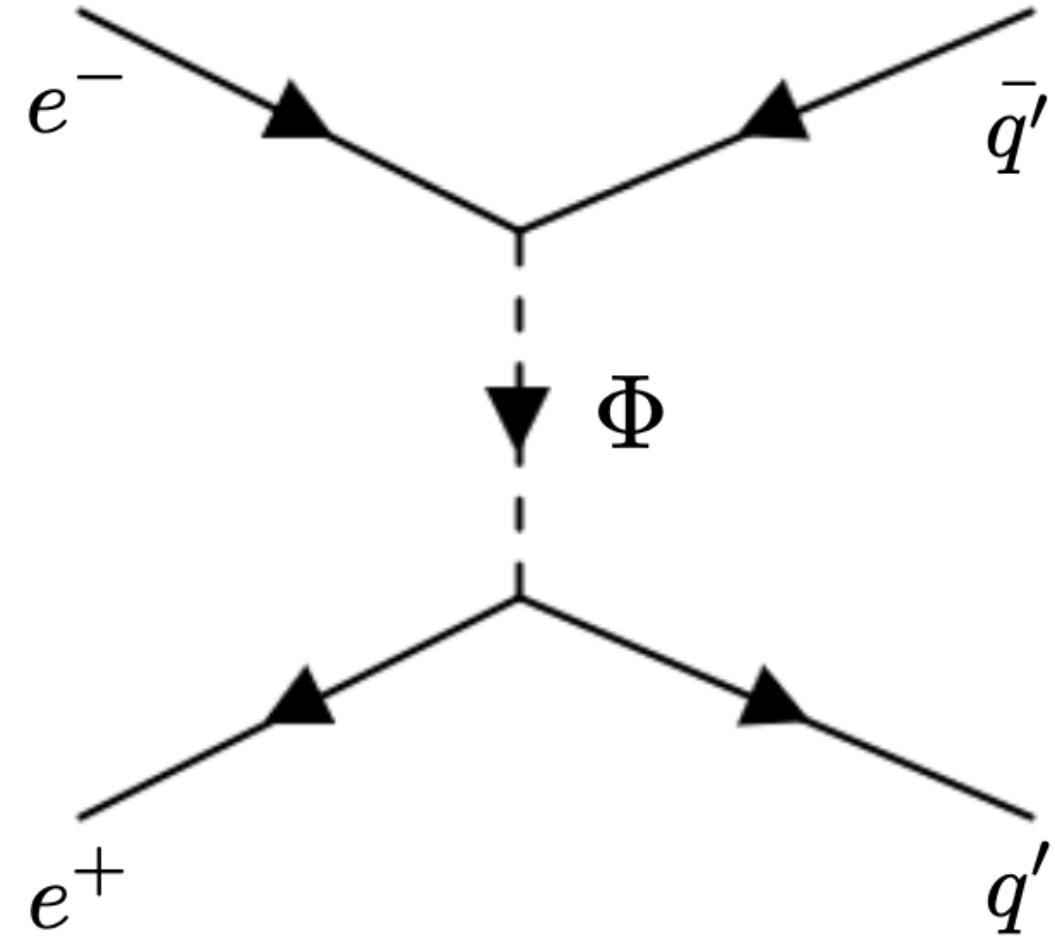
# Dark jets at ee collider

$$\begin{aligned} \mathcal{L} = & \mathcal{L}_{\text{SM}} - \frac{1}{2} \sum_{i=1,2} M_{N_i} \bar{N}_i N_i^C - m_\Phi^2 \Phi^\dagger \Phi - m_\chi \bar{\chi} \chi - m_{q'} \bar{q}' q' + \mathcal{L}_{\text{kinetic}} \\ & - \sum_{i=1,2} \lambda_i \bar{N}_i \chi \Phi^\dagger - \boxed{\kappa \Phi \bar{q}'_L l_R} - \frac{1}{\Lambda_1^2} (\bar{q}'^C \chi) (\bar{q}'^C_L l_R) - \frac{1}{\Lambda_2^2} (\bar{\chi} \gamma^\mu q') (\bar{d}_R \gamma_\mu u_R) + h.c. \end{aligned}$$

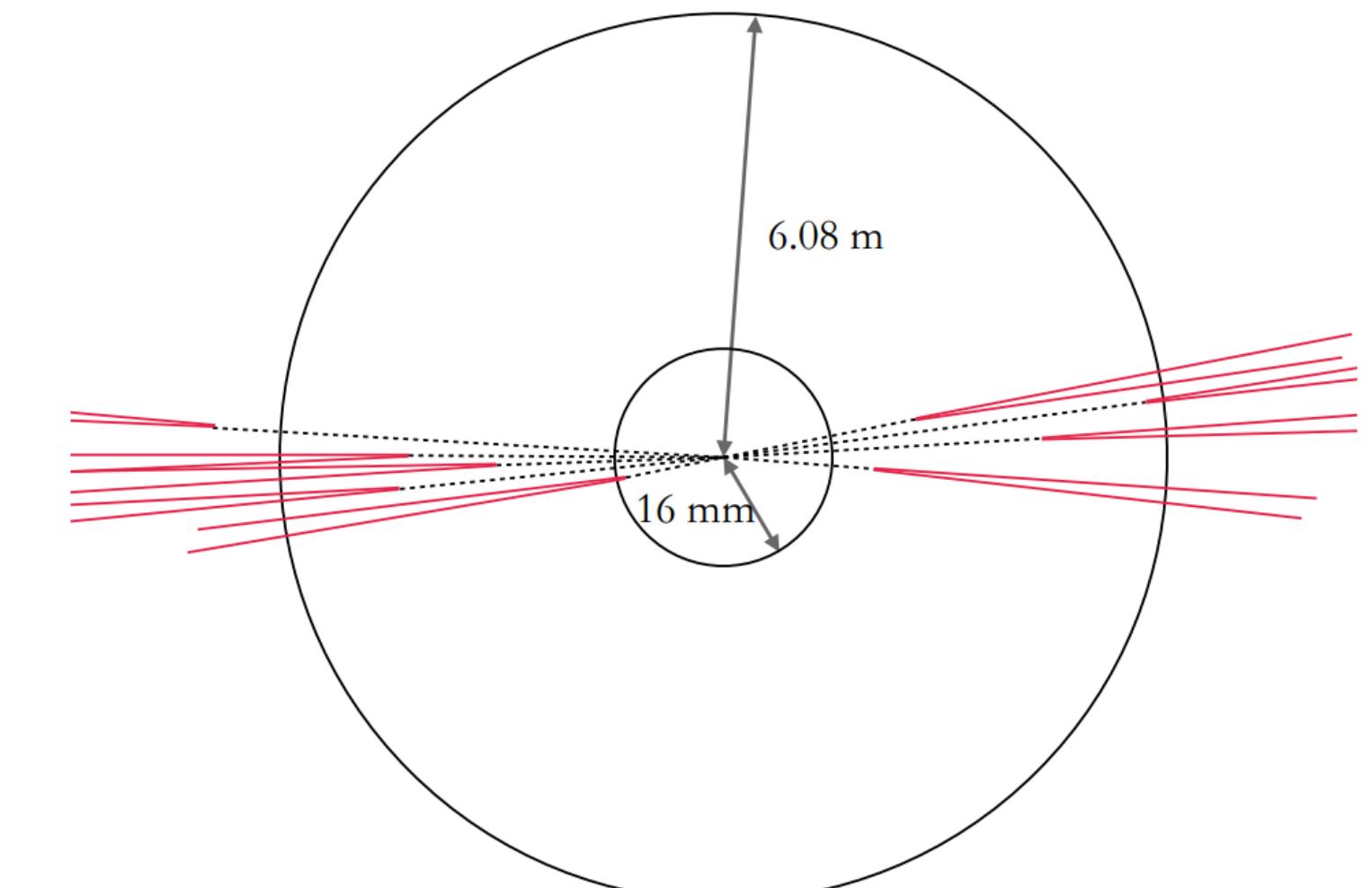
- Relevant Lagrangian

P. Schwaller et al, 1502.05409 (JHEP)

$$\mathcal{L} \supset \bar{q}'(\not{D} - m_{q'}) q' + (D_\mu \Phi)^\dagger (D^\mu \Phi) - m_\Phi^2 \Phi^\dagger \Phi - \frac{1}{4} G'^{\mu\nu} G'_{\mu\nu} - (\kappa \Phi \bar{q}'_L l_R + h.c.)$$

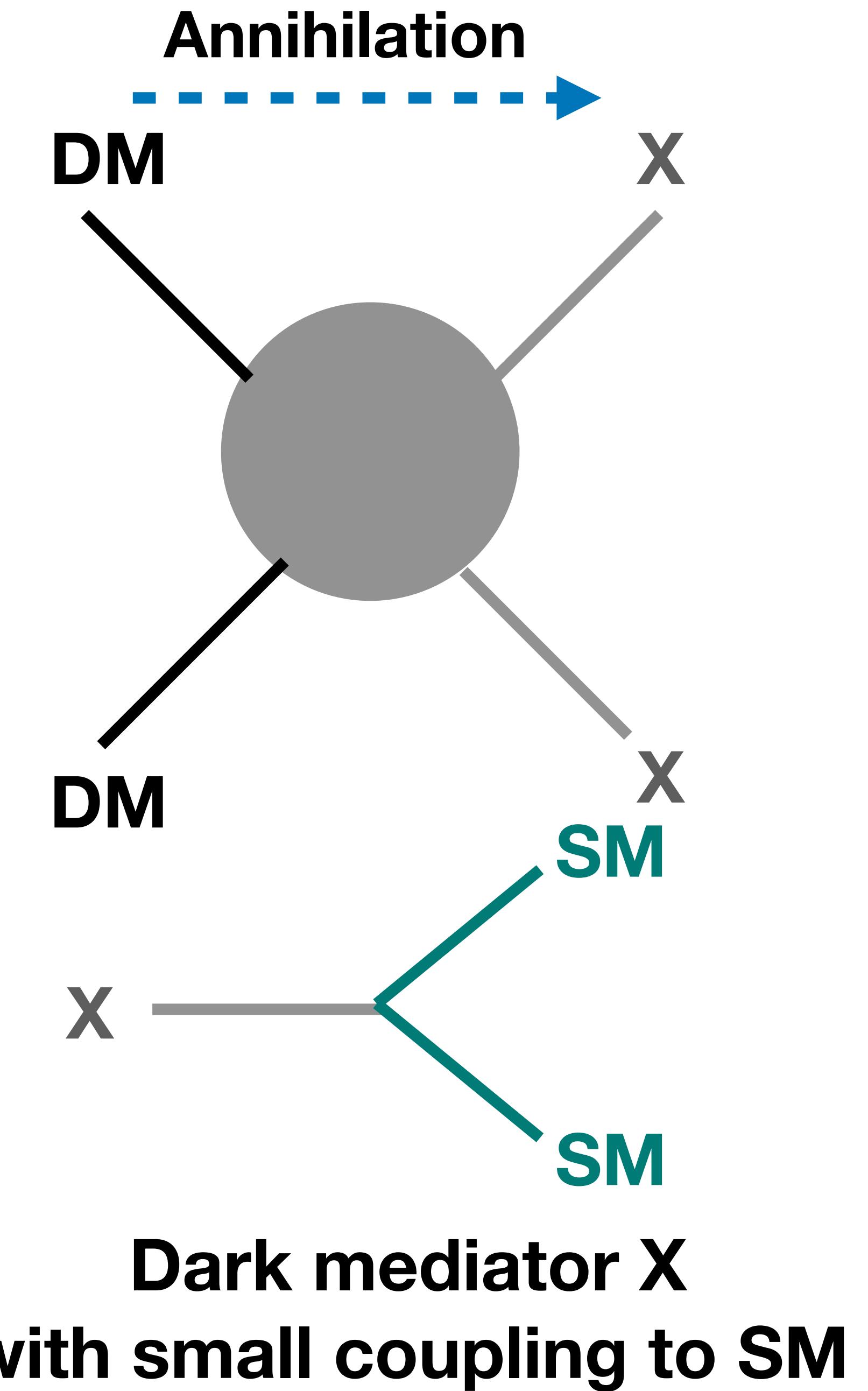


- $e^+ e^- \rightarrow \bar{q}'q'$  followed by dark hadronization
- Displaced dark meson decay  $\pi_d = (\bar{q}'q')$



# Outline

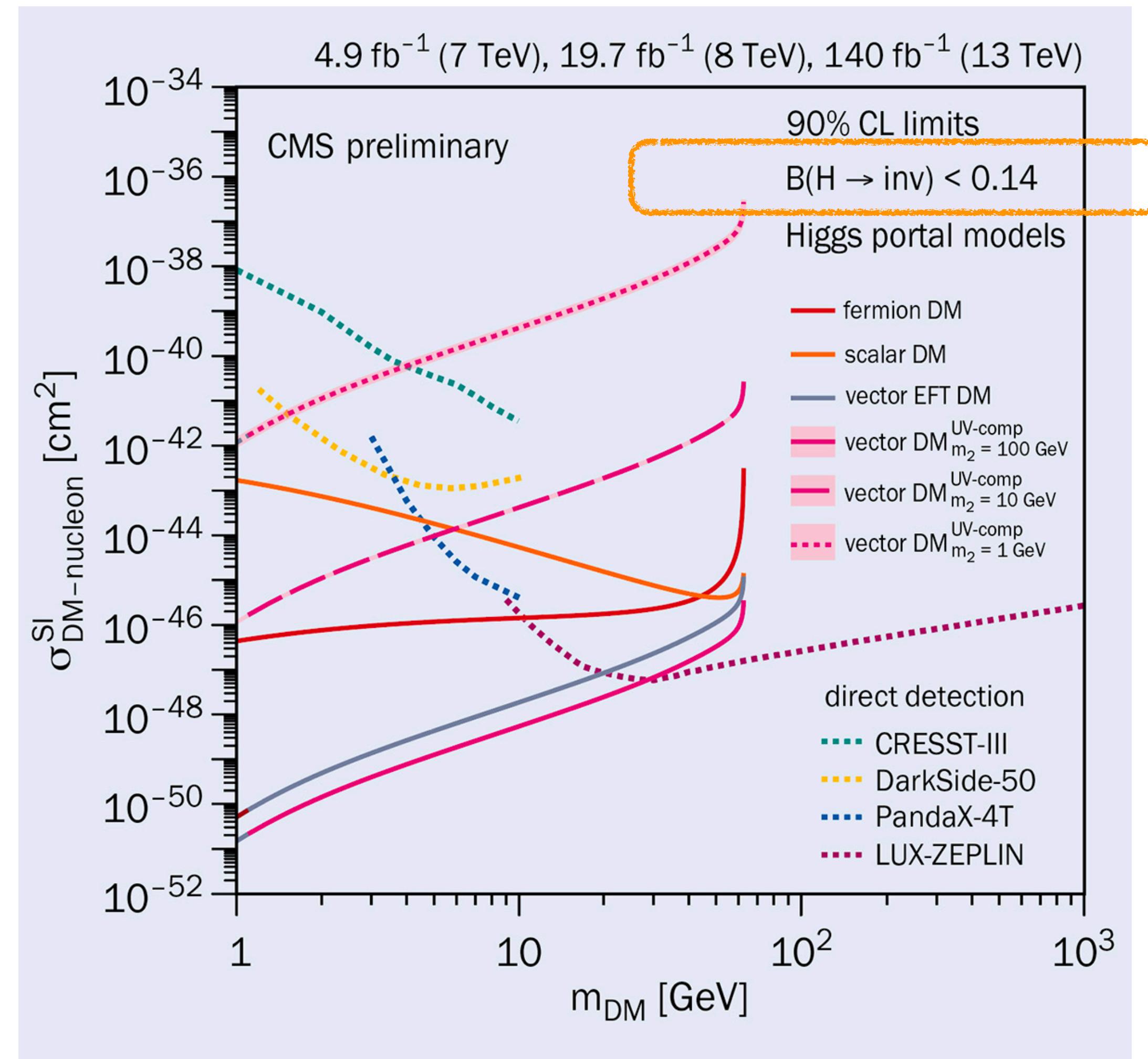
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# Higgs portal DM: invisible Higgs decay

- Precision measurement of Higgs
  - Invisible Higgs test to **0.07%@CEPC**

	240 GeV, 20 ab <sup>-1</sup>		360 GeV, 1 ab <sup>-1</sup>		
	ZH	vvH	ZH	vvH	eeH
inclusive	<b>0.26%</b>		<b>1.40%</b>	\	\
H→bb	<b>0.14%</b>	<b>1.59%</b>	<b>0.90%</b>	<b>1.10%</b>	<b>4.30%</b>
H→cc	<b>2.02%</b>		<b>8.80%</b>	<b>16%</b>	<b>20%</b>
H→gg	<b>0.81%</b>		<b>3.40%</b>	<b>4.50%</b>	<b>12%</b>
H→WW	<b>0.53%</b>		<b>2.80%</b>	<b>4.40%</b>	<b>6.50%</b>
H→ZZ	<b>4.17%</b>		<b>20%</b>	<b>21%</b>	
H → ττ	<b>0.42%</b>		<b>2.10%</b>	<b>4.20%</b>	<b>7.50%</b>
H → γγ	<b>3.02%</b>		<b>11%</b>	<b>16%</b>	
H → μμ	<b>6.36%</b>		<b>41%</b>	<b>57%</b>	
H → Zγ	<b>8.50%</b>		<b>35%</b>		
Br <sub>upper</sub> (H → inv.)	<b>0.07%</b>				
Γ <sub>H</sub>	<b>1.65%</b>		<b>1.10%</b>		

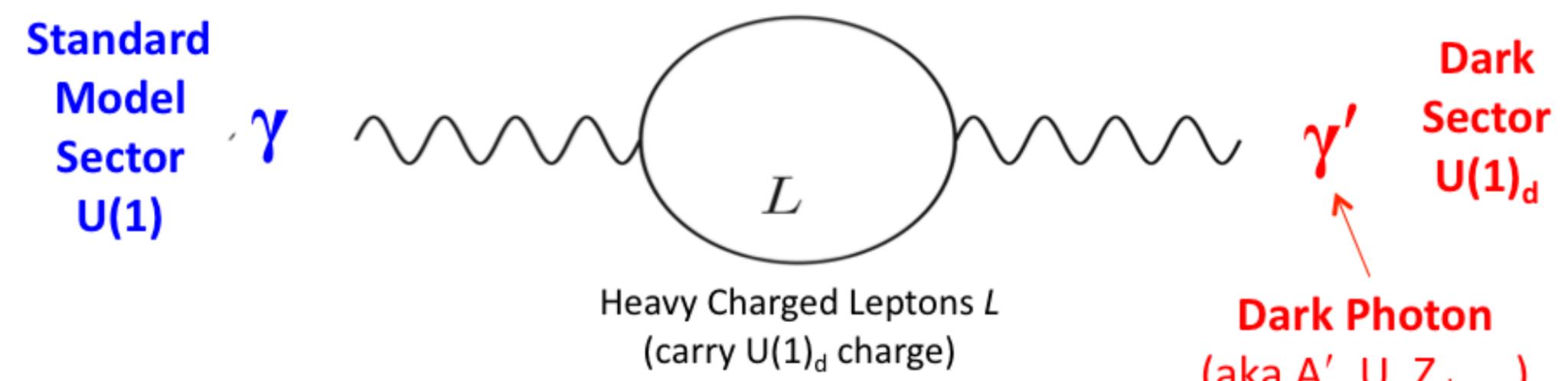
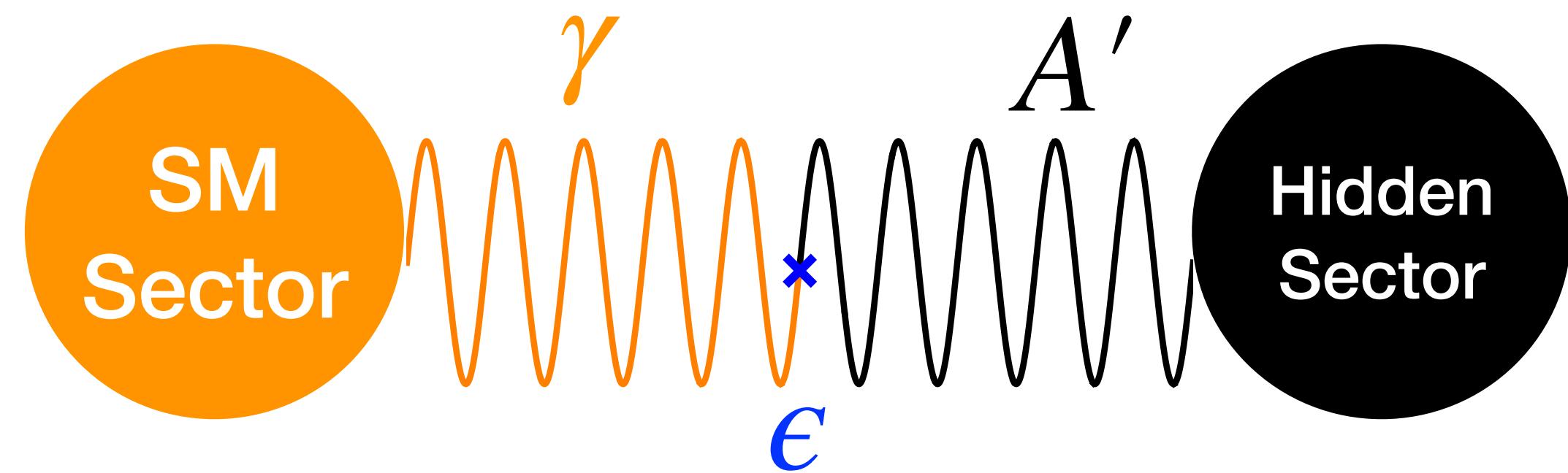


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# Vector portal

- The vector portal to DM
- The Interaction:  $g' \bar{\chi} \gamma^\mu \chi Z'_\mu$
- Millicharged DM: EM with  $\epsilon e$
- Kinetic mixing portal
- Effective  $Z'$  through fermion mass mixing

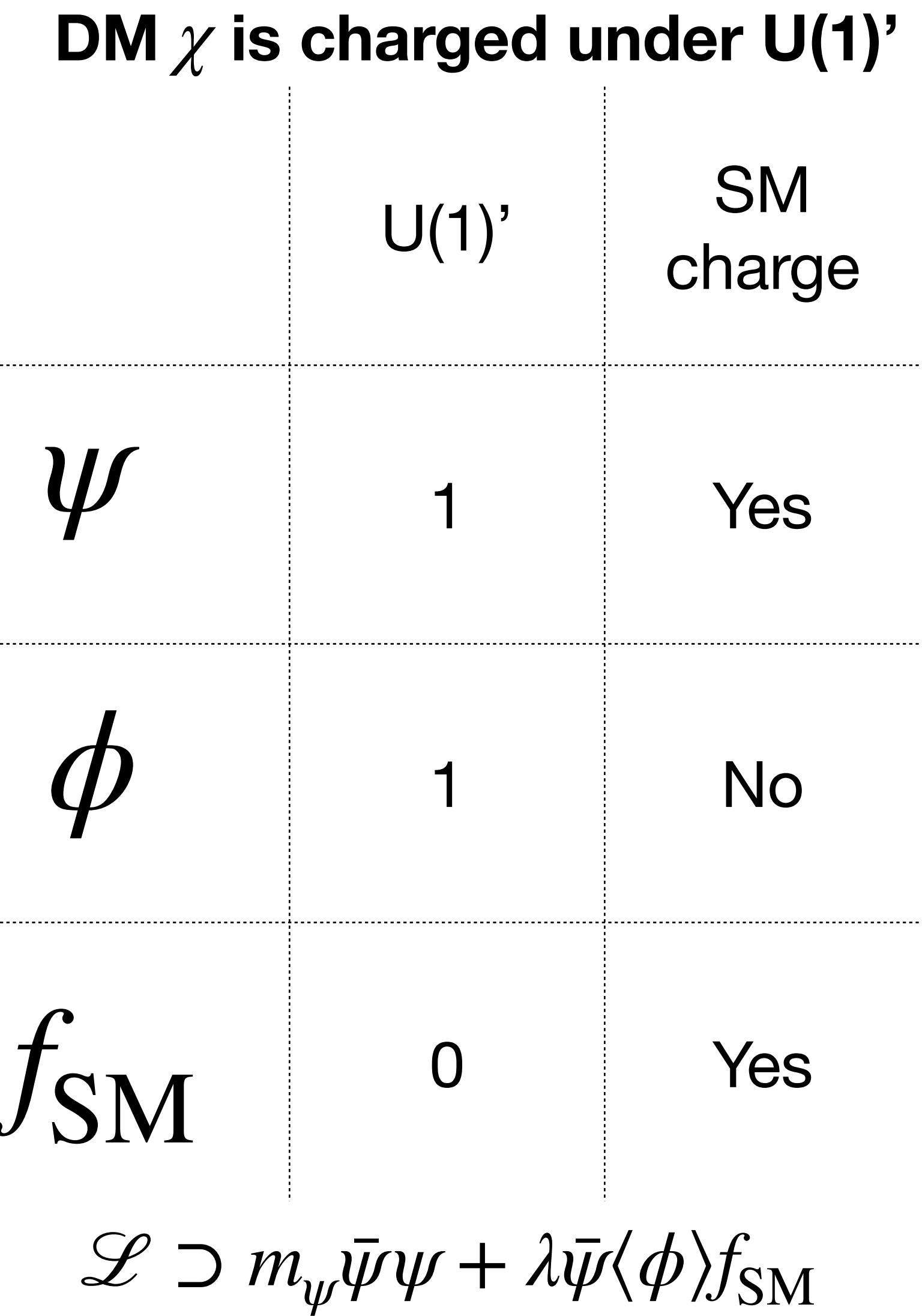


$$\epsilon \sim -\frac{gg'}{16\pi^2} \log \left( \frac{m_L^2}{\mu^2} \right)$$

$$\mathcal{L} = -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{1}{2} m_{A'}^2 A'^{\mu} A'^{\nu} - \frac{1}{2} \epsilon F'_{\mu\nu} F^{\mu\nu}$$

# Vector portal

- The vector portal to DM
- The Interaction:  $g' \bar{\chi} \gamma^\mu \chi Z'_\mu$
- Millicharged DM: EM with  $\epsilon e$
- Kinetic mixing portal
- Effective  $Z'$  through fermion mass mixing ( $\psi/f_{\text{SM}}$ )  $g' \bar{f} \gamma^\mu f Z'_\mu$

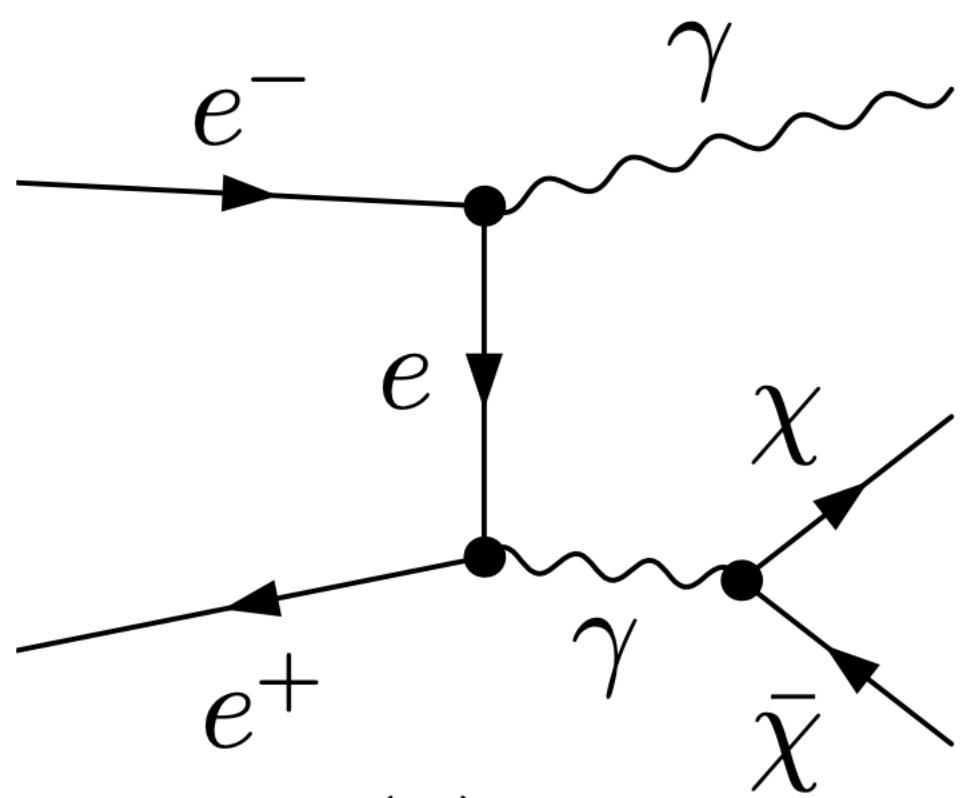


# Searching DM via mono-photon at CEPC

- DM searches via mono-photon final states

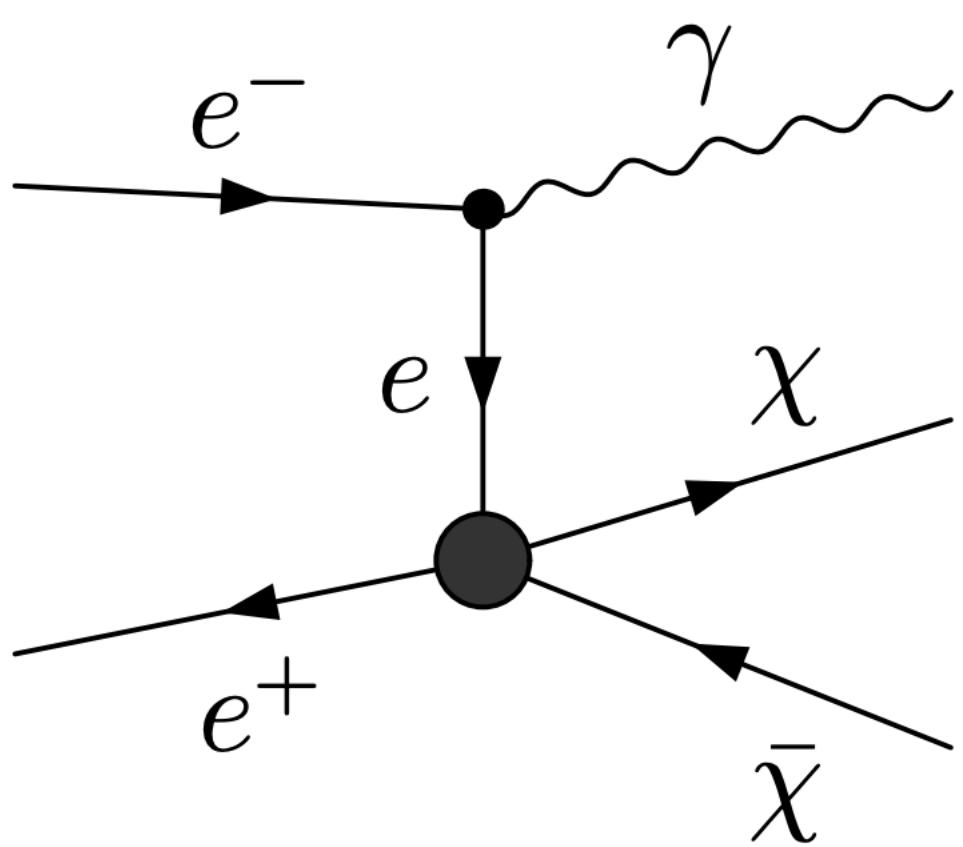
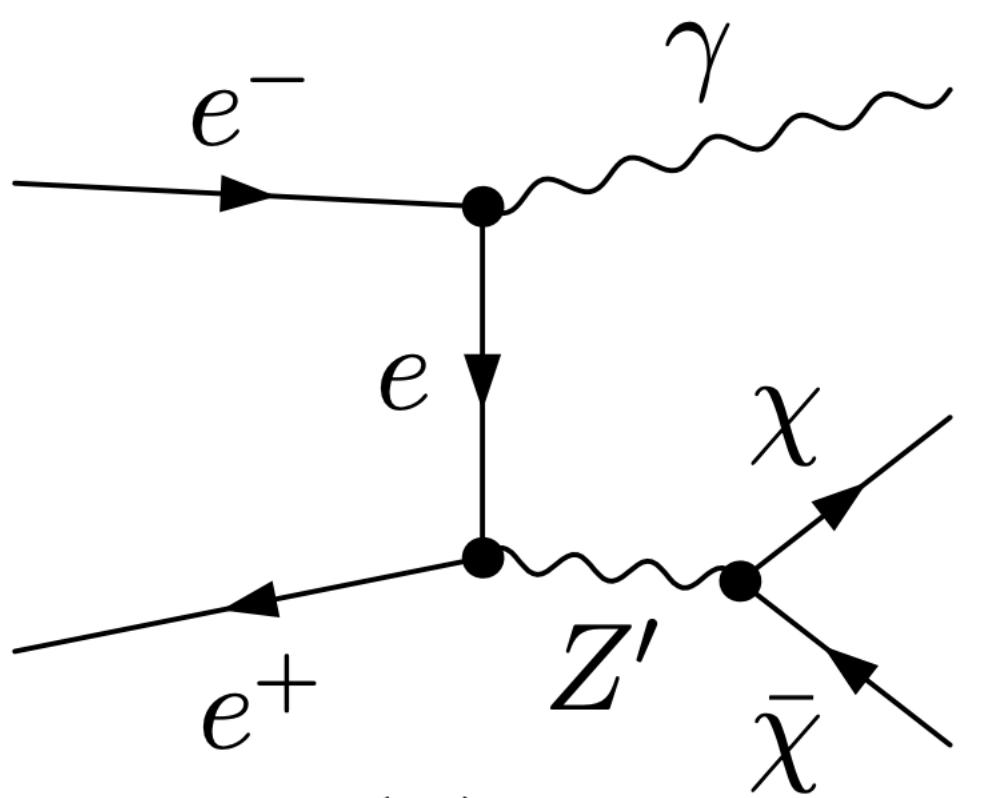
Millicharged DM

$$\mathcal{L} = e\varepsilon A_\mu \bar{\chi} \gamma^\mu \chi,$$



$Z'$  portal DM

$$\mathcal{L} = Z'_\mu \bar{\chi} \gamma^\mu (g_V^\chi - g_A^\chi \gamma_5) \chi + Z'_\mu \bar{f} \gamma^\mu (g_V^f - g_A^f \gamma_5) f,$$



DM EFTs

$$\mathcal{L} = \frac{1}{\Lambda_V^2} \bar{\chi} \gamma_\mu \chi \bar{\ell} \gamma^\mu \ell,$$

$$\mathcal{L} = \frac{1}{\Lambda_s^2} \bar{\chi} \chi \bar{\ell} \ell,$$

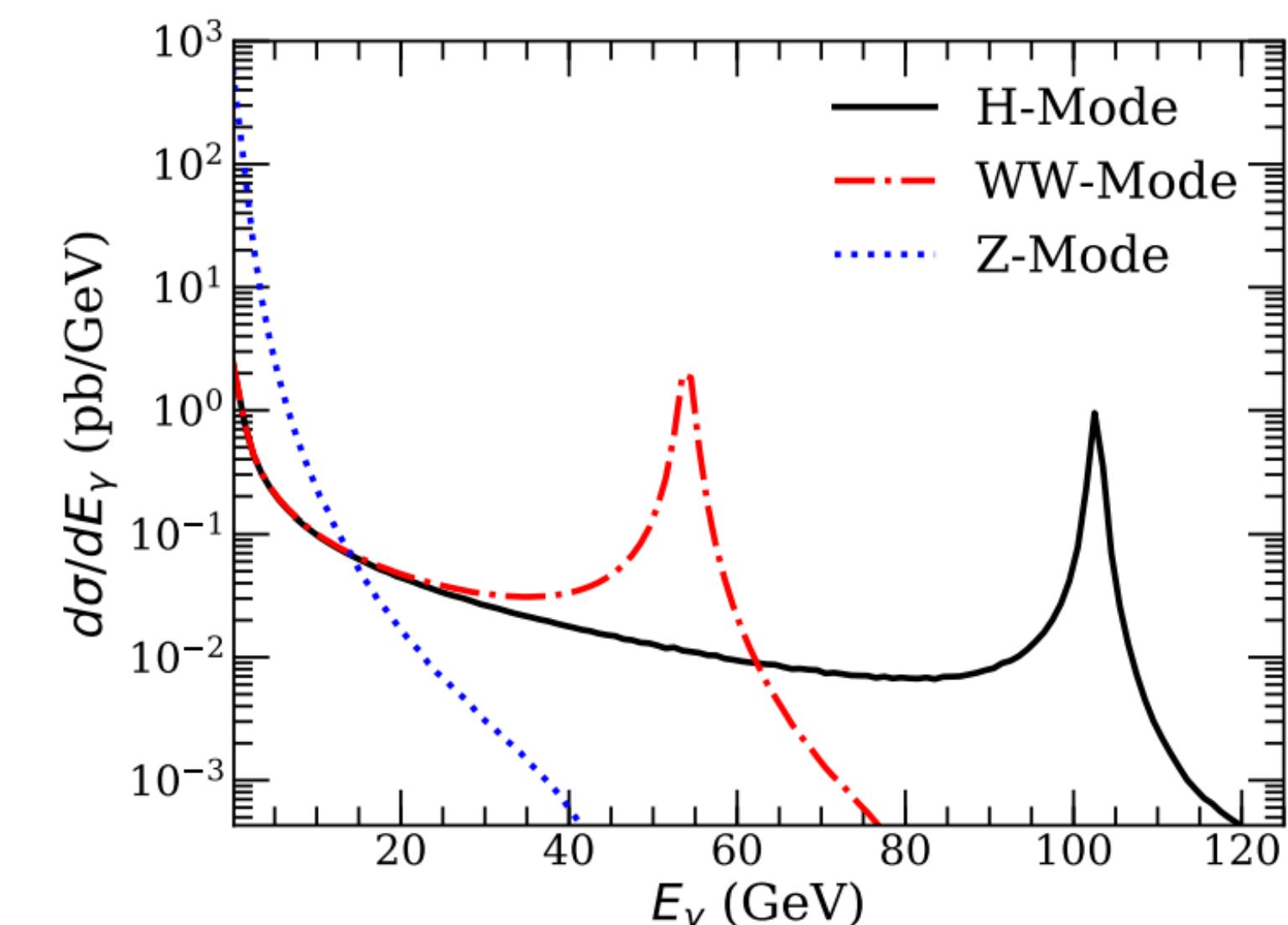
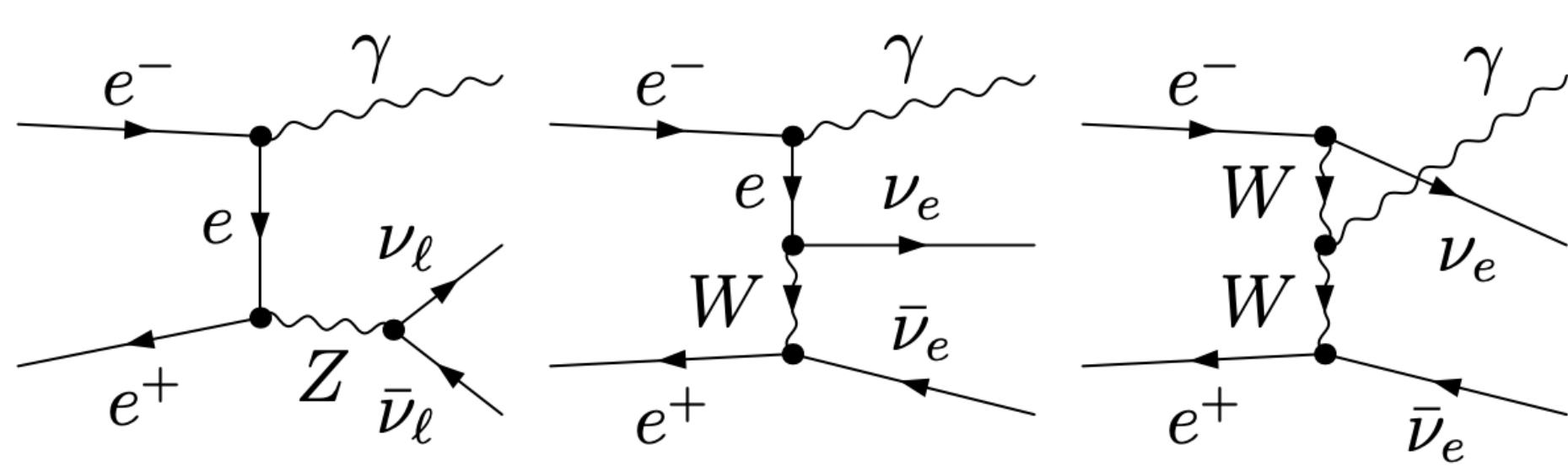
$$\mathcal{L} = \frac{1}{\Lambda_A^2} \bar{\chi} \gamma_\mu \gamma_5 \chi \bar{\ell} \gamma^\mu \gamma_5 \ell,$$

$$\mathcal{L} = \frac{1}{\Lambda_t^2} \bar{\chi} \ell \bar{\ell} \chi$$

# Mono-photon background at CEPC

- Irreducible background

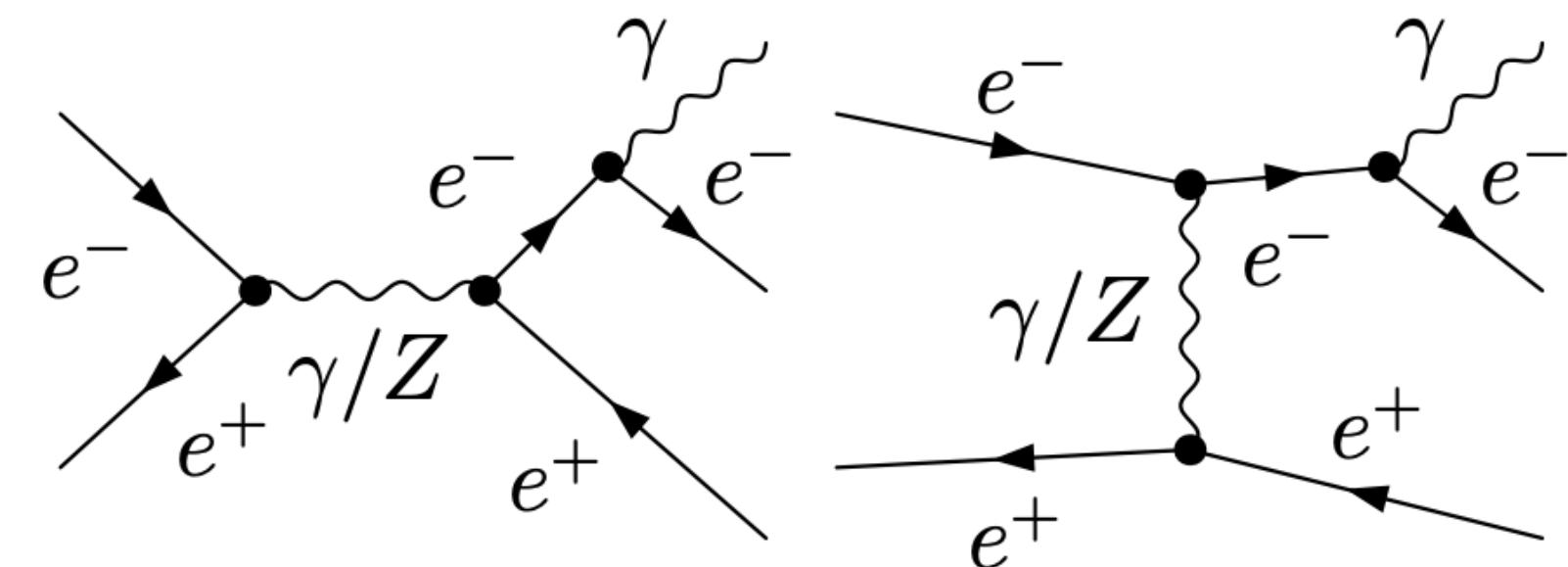
$$e^+ e^- \rightarrow \gamma \nu \bar{\nu}$$



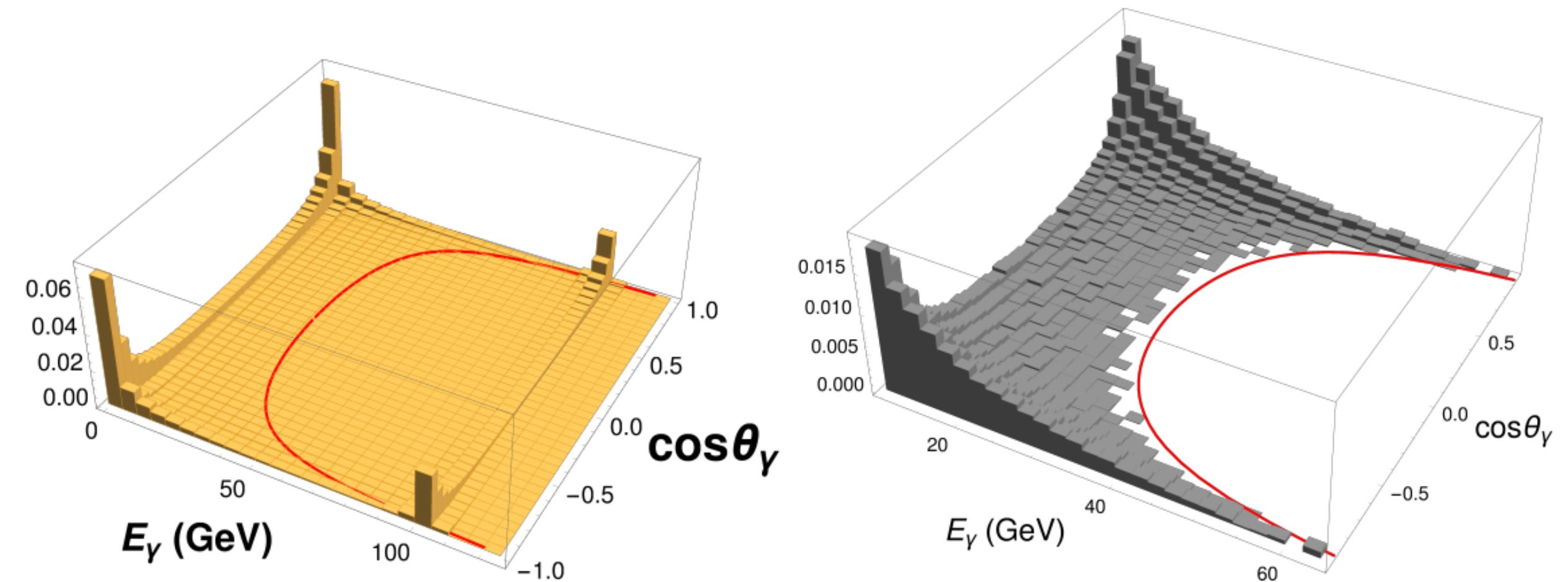
Zuowei Liu et al, 1903.12114 (JHEP)

- Reducible background

$$e^+ e^- \rightarrow \gamma e^+ e^-$$

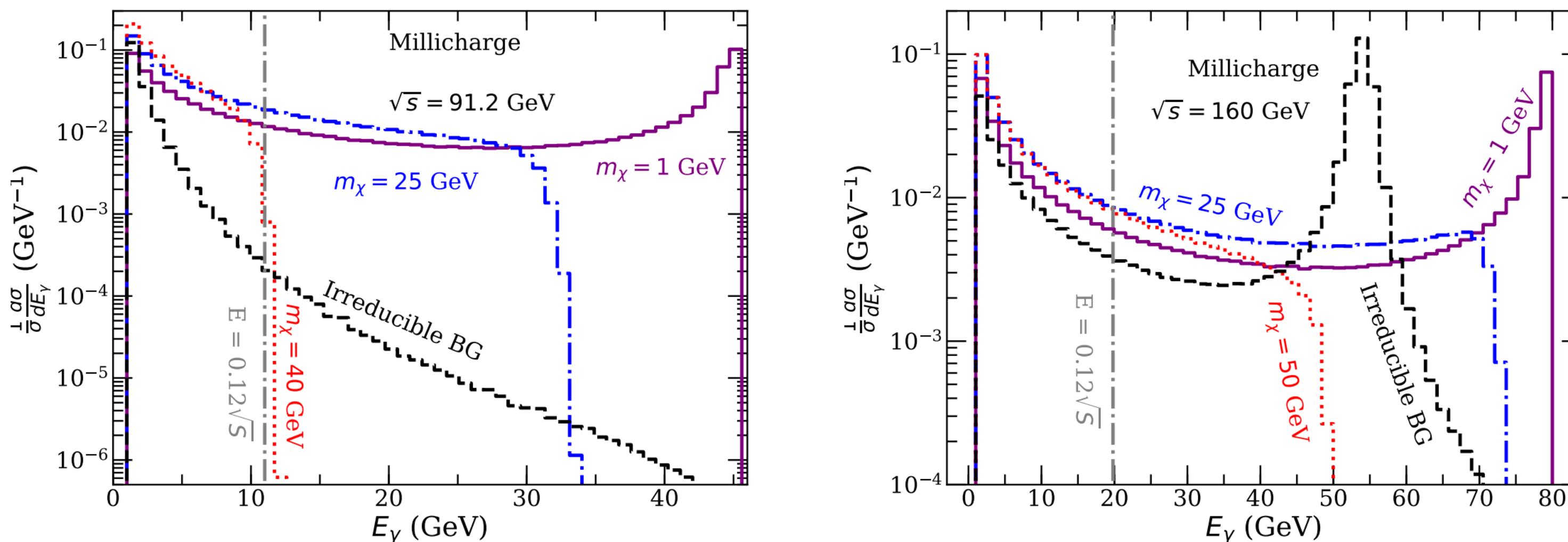


- Red curve shows the signal region

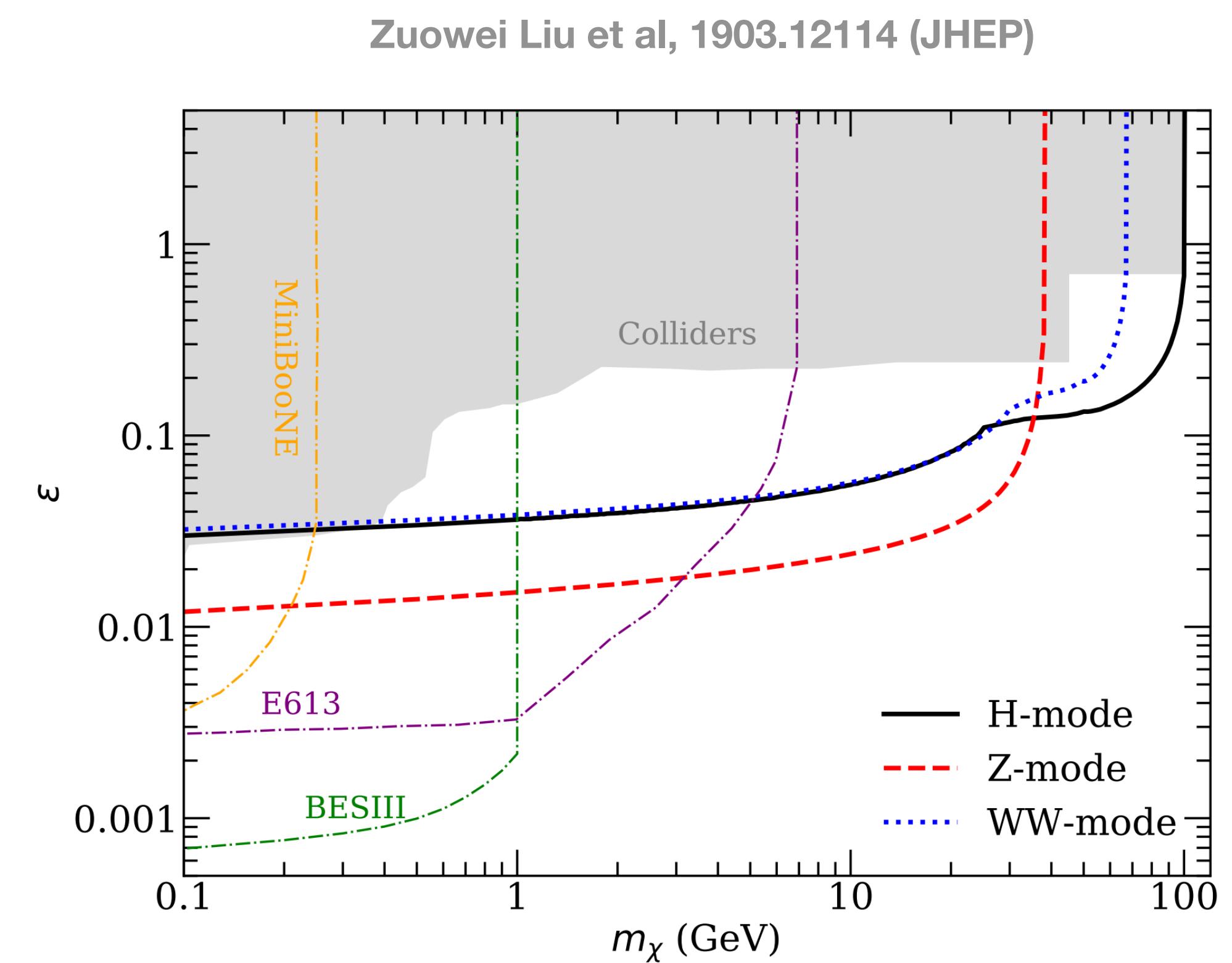


# Constraints on millicharged DM

- The signal distribution at Z ( $2.6 \text{ ab}^{-1}$ ), H ( $5.6 \text{ ab}^{-1}$ ), WW modes ( $16 \text{ ab}^{-1}$ )

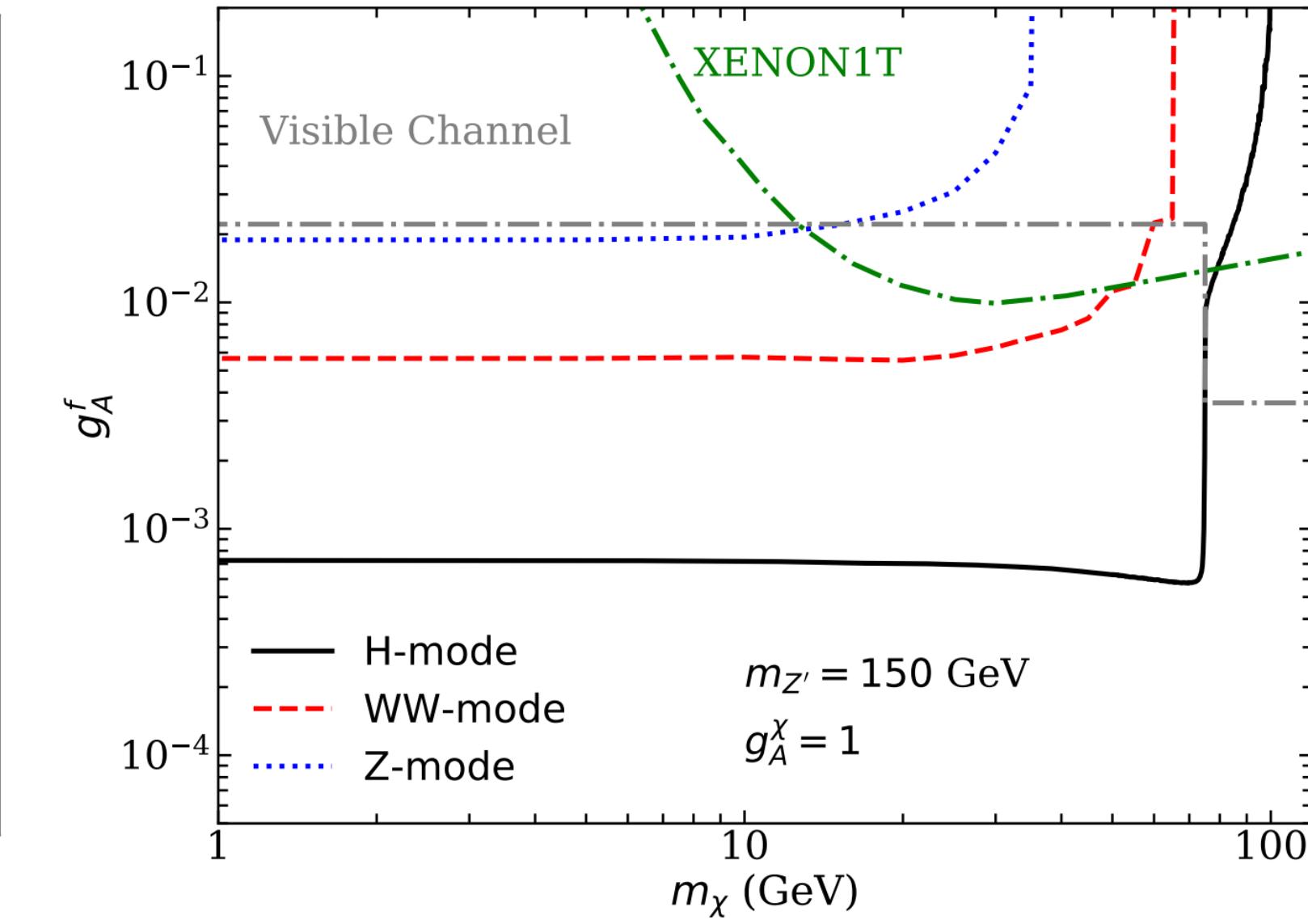
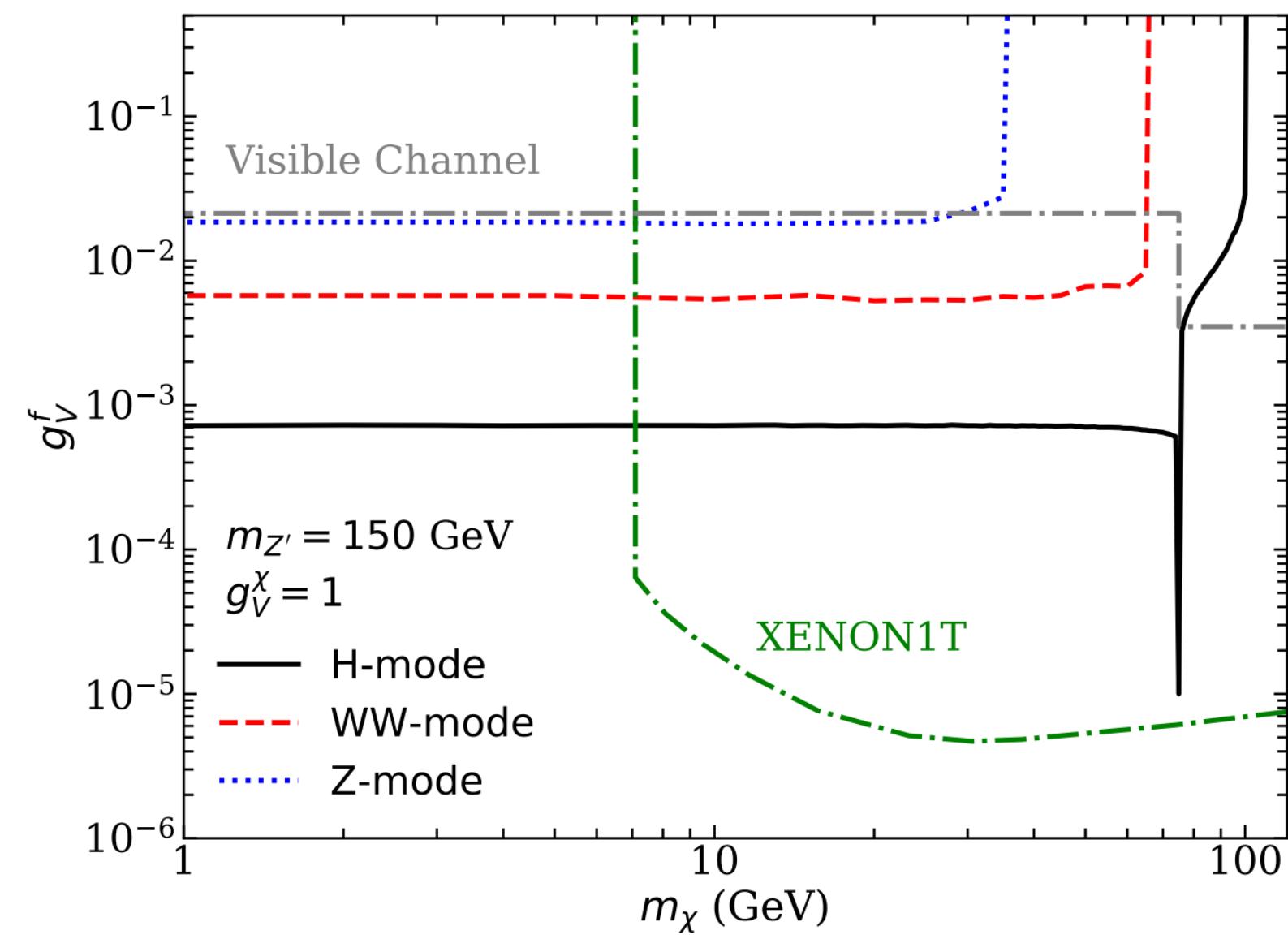
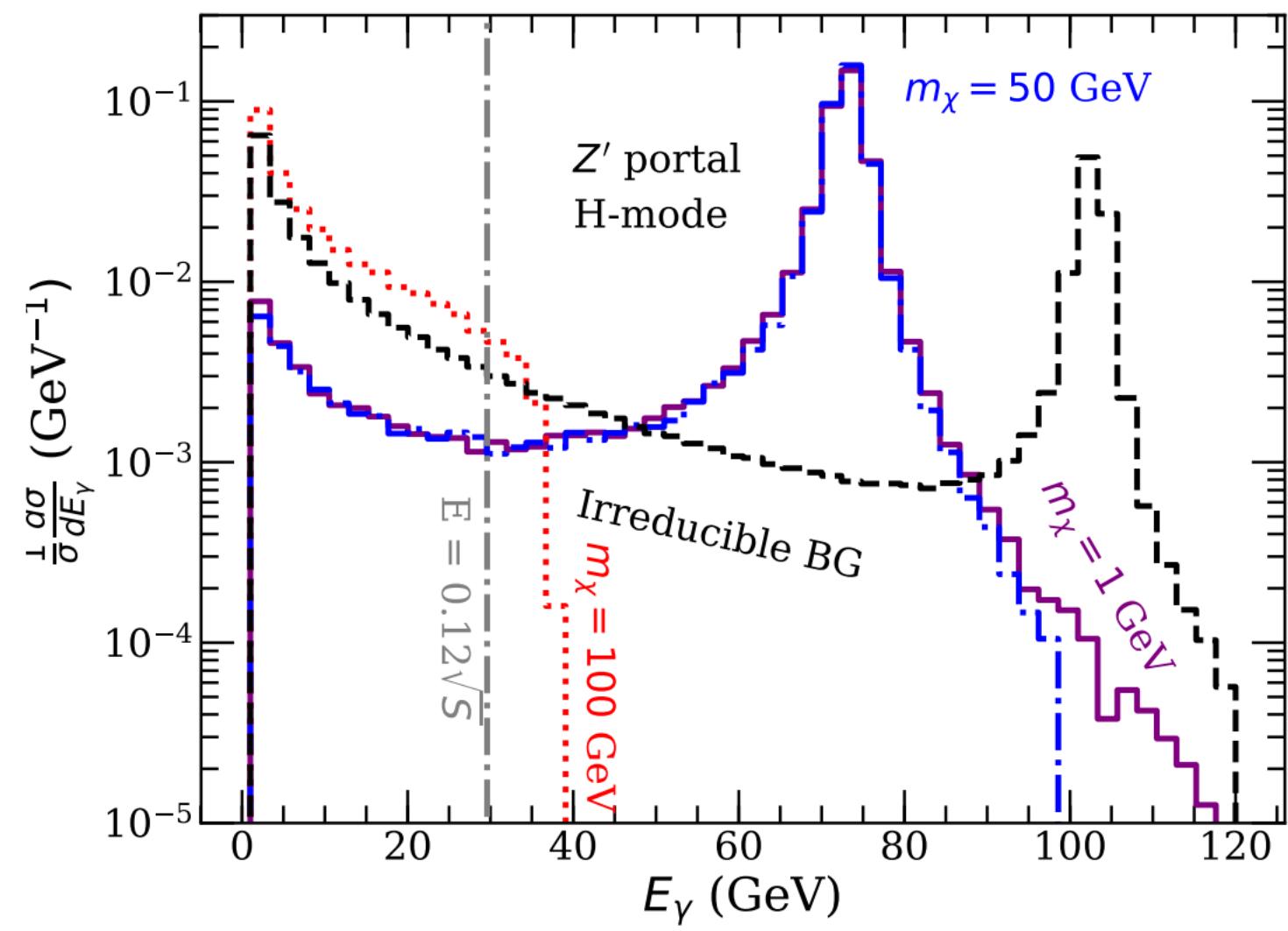


- The Z/Higgs factory modes can provide competitive sensitivity comparing with existing colliders



# Constraints on Z' portal DM model

- The signal distribution at Z ( $2.6 \text{ ab}^{-1}$ ), H ( $5.6 \text{ ab}^{-1}$ ), WW modes ( $16 \text{ ab}^{-1}$ )



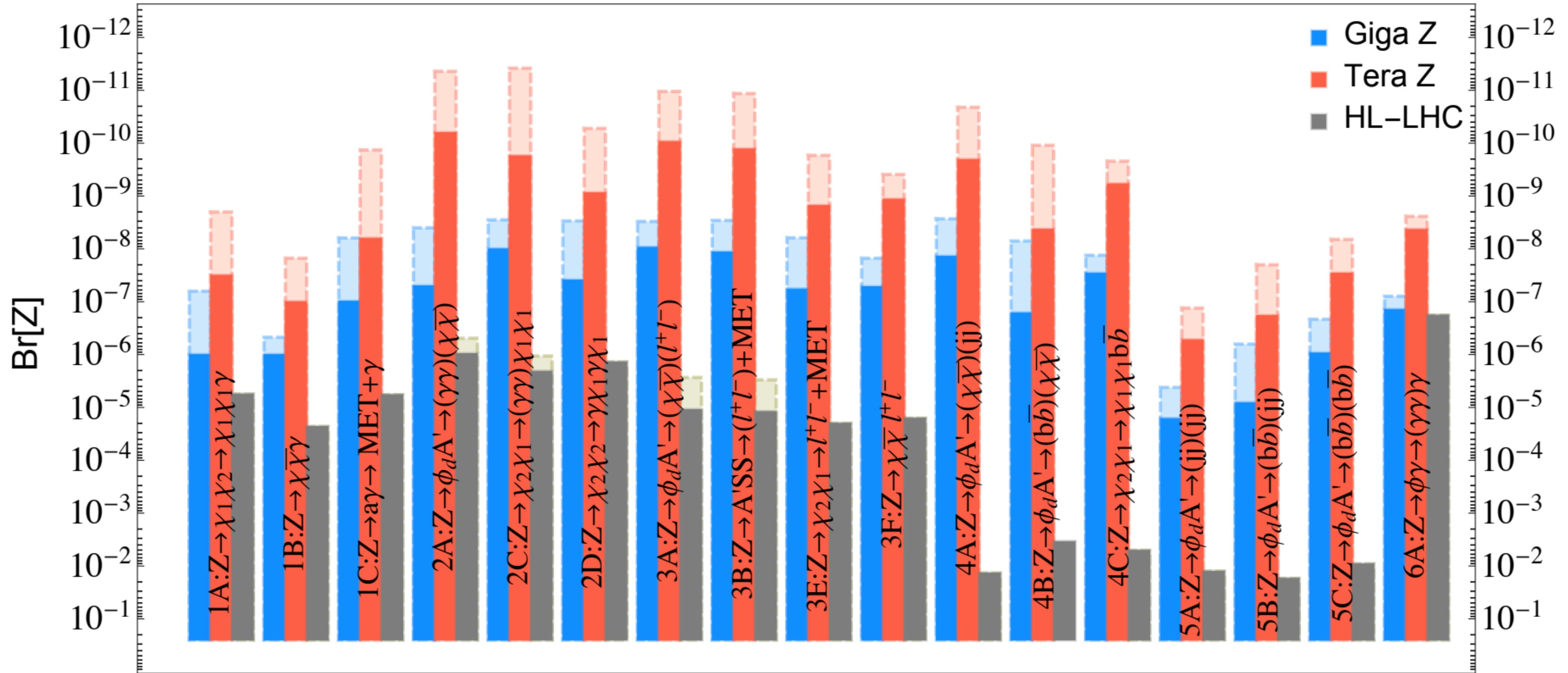
- The Z/Higgs factory modes can provide competitive sensitivity comparing with existing colliders and direct detection searches

# Probing dark sector models at future Z-factory

exotic decays	topologies	$n_{res}$	models
$Z \rightarrow \cancel{E} + \gamma$	$Z \rightarrow \chi_1 \chi_2, \chi_2 \rightarrow \chi_1 \gamma$	0	1A: $\frac{1}{\Lambda_{1A}} \bar{\chi}_2 \sigma^{\mu\nu} \chi_1 B_{\mu\nu}$ (MIDM)
	$Z \rightarrow \chi \bar{\chi} \gamma$	0	1B: $\frac{1}{\Lambda_{1B}^3} \bar{\chi} \chi B_{\mu\nu} B^{\mu\nu}$ (RayDM)
	$Z \rightarrow a \gamma \rightarrow (\cancel{E}) \gamma$	1	1C: $\frac{1}{4\Lambda_{1C}} a B_{\mu\nu} \tilde{B}^{\mu\nu}$ (long-lived ALP)
	$Z \rightarrow A' \gamma \rightarrow (\bar{\chi} \chi) \gamma$	1	1D: $\epsilon^{\mu\nu\rho\sigma} A'_\mu B_\nu \partial_\rho B_\sigma$ (WZ terms)
$Z \rightarrow \cancel{E} + \gamma\gamma$	$Z \rightarrow \phi_d A', \phi_d \rightarrow (\gamma\gamma), A' \rightarrow (\bar{\chi} \chi)$	2	2A: Vector portal
	$Z \rightarrow \phi_H \phi_A, \phi_H \rightarrow (\gamma\gamma), \phi_A \rightarrow (\bar{\chi} \chi)$	2	2B: 2HDM extension
	$Z \rightarrow \chi_2 \chi_1, \chi_2 \rightarrow \chi_1 \phi, \phi \rightarrow (\gamma\gamma)$	1	2C: Inelastic DM
	$Z \rightarrow \chi_2 \chi_2, \chi_2 \rightarrow \gamma \chi_1$	0	2D: MIDM

$Z \rightarrow \cancel{E} + \ell^+ \ell^-$	$Z \rightarrow \phi_d A', A' \rightarrow (\ell^+ \ell^-), \phi_d \rightarrow (\bar{\chi} \chi)$	2	3A: Vector portal
	$Z \rightarrow A' SS \rightarrow (\ell\ell) SS$	1	3B: Vector portal
	$Z \rightarrow \phi(Z^*/\gamma^*) \rightarrow \phi \ell^+ \ell^-$	1	3C: Long-lived ALP, Higgs portal
	$Z \rightarrow \chi_2 \chi_1 \rightarrow \chi_1 A' \chi_1 \rightarrow (\ell^+ \ell^-) \cancel{E}$	1	3D: Vector portal and Inelastic DM
	$Z \rightarrow \chi_2 \chi_1, \chi_2 \rightarrow \chi_1 \ell^+ \ell^-$	0	3E: MIDM, SUSY
$Z \rightarrow \cancel{E} + JJ$	$Z \rightarrow \bar{\chi} \chi \ell^+ \ell^-$	0	3F: RayDM, slepton, heavy lepton mixing
	$Z \rightarrow \phi_d A' \rightarrow (\bar{\chi} \chi)(jj)$	2	4A: Vector portal
	$Z \rightarrow \phi_d A' \rightarrow (bb)(\bar{\chi} \chi)$	2	4B: Vector portal + Higgs portal
$Z \rightarrow (JJ)(JJ)$	$Z \rightarrow \chi_2 \chi_1 \rightarrow bb \chi_1 + \chi_1 \rightarrow bb \cancel{E}$	0	4C: MIDM
	$Z \rightarrow \phi_d A', \phi_d \rightarrow jj, A' \rightarrow jj$	2	5A: Vector portal + Higgs portal
	$Z \rightarrow \phi_d A', \phi_d \rightarrow b\bar{b}, A' \rightarrow jj$	2	5B: vector portal + Higgs portal
	$Z \rightarrow \phi_d A', \phi_d \rightarrow b\bar{b}, A' \rightarrow b\bar{b}$	2	5C: vector portal + Higgs portal
$Z \rightarrow \gamma\gamma\gamma$	$Z \rightarrow \phi \gamma \rightarrow (\gamma\gamma)\gamma$	1	6A: ALP, Higgs portal

# Probing dark sector models at future Z-factory



JL, LT Wang, XP Wang, W. Xue, 1712.07237 (PRD)

# Outline

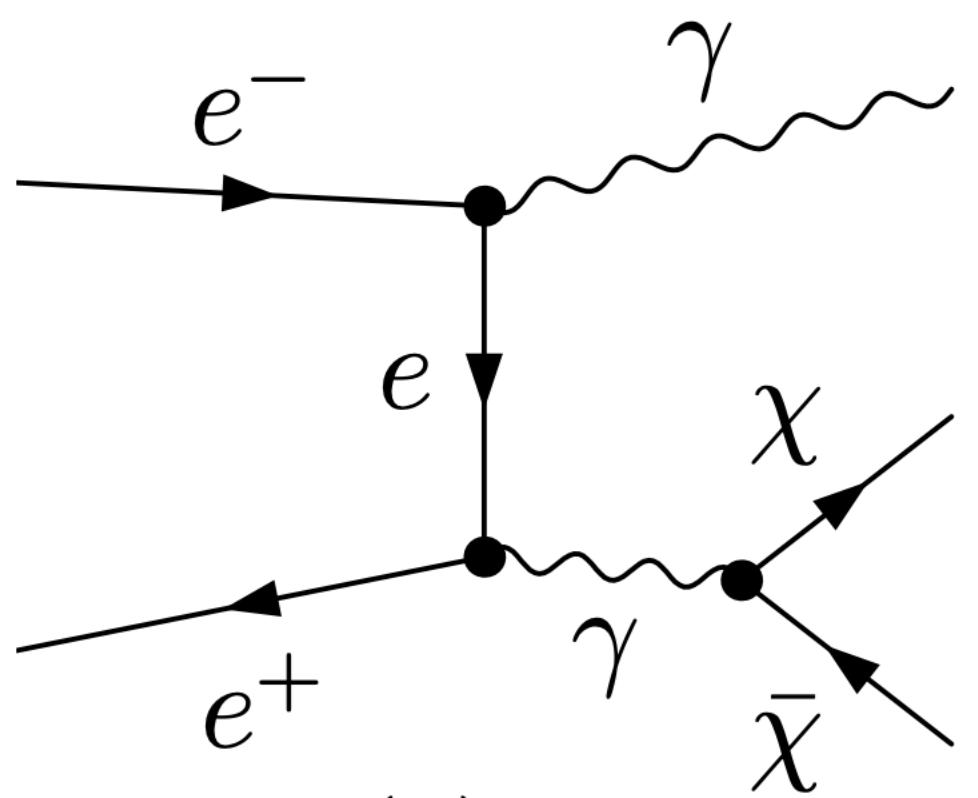
- CEPC and New Physics Searches
  - SUSY
  - Dark Matter and Dark Sector
    - Fermion portal – lepton portal
    - Higgs portal
    - Vector portal
    - EFT models
  - Summary

# Searching DM via mono-photon at CEPC

- DM searches via mono-photon final states

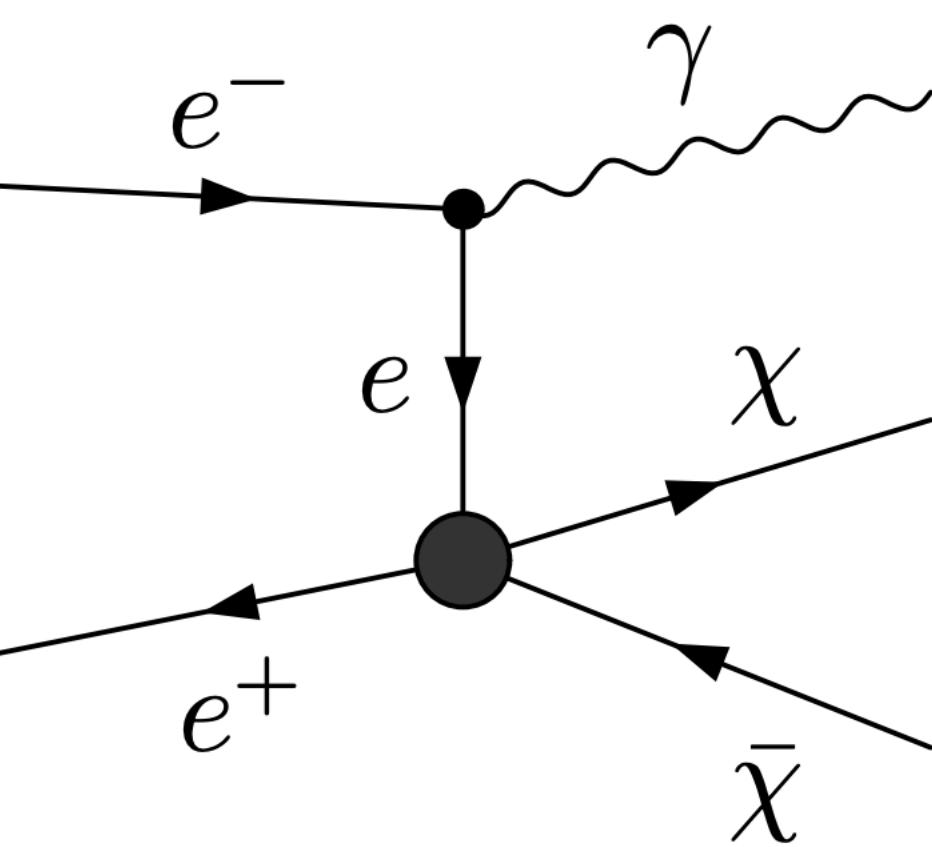
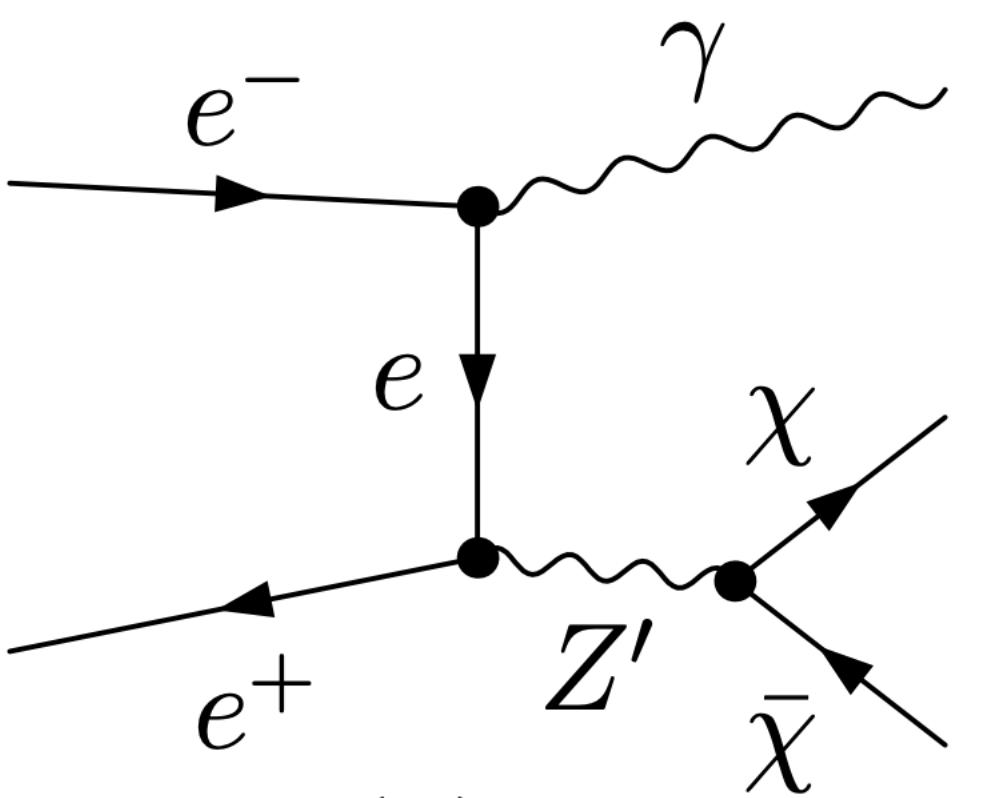
Millicharged DM

$$\mathcal{L} = e\varepsilon A_\mu \bar{\chi} \gamma^\mu \chi,$$



$Z'$  portal DM

$$\mathcal{L} = Z'_\mu \bar{\chi} \gamma^\mu (g_V^\chi - g_A^\chi \gamma_5) \chi + Z'_\mu \bar{f} \gamma^\mu (g_V^f - g_A^f \gamma_5) f,$$



DM EFTs

$$\mathcal{L} = \frac{1}{\Lambda_V^2} \bar{\chi} \gamma_\mu \chi \bar{\ell} \gamma^\mu \ell,$$

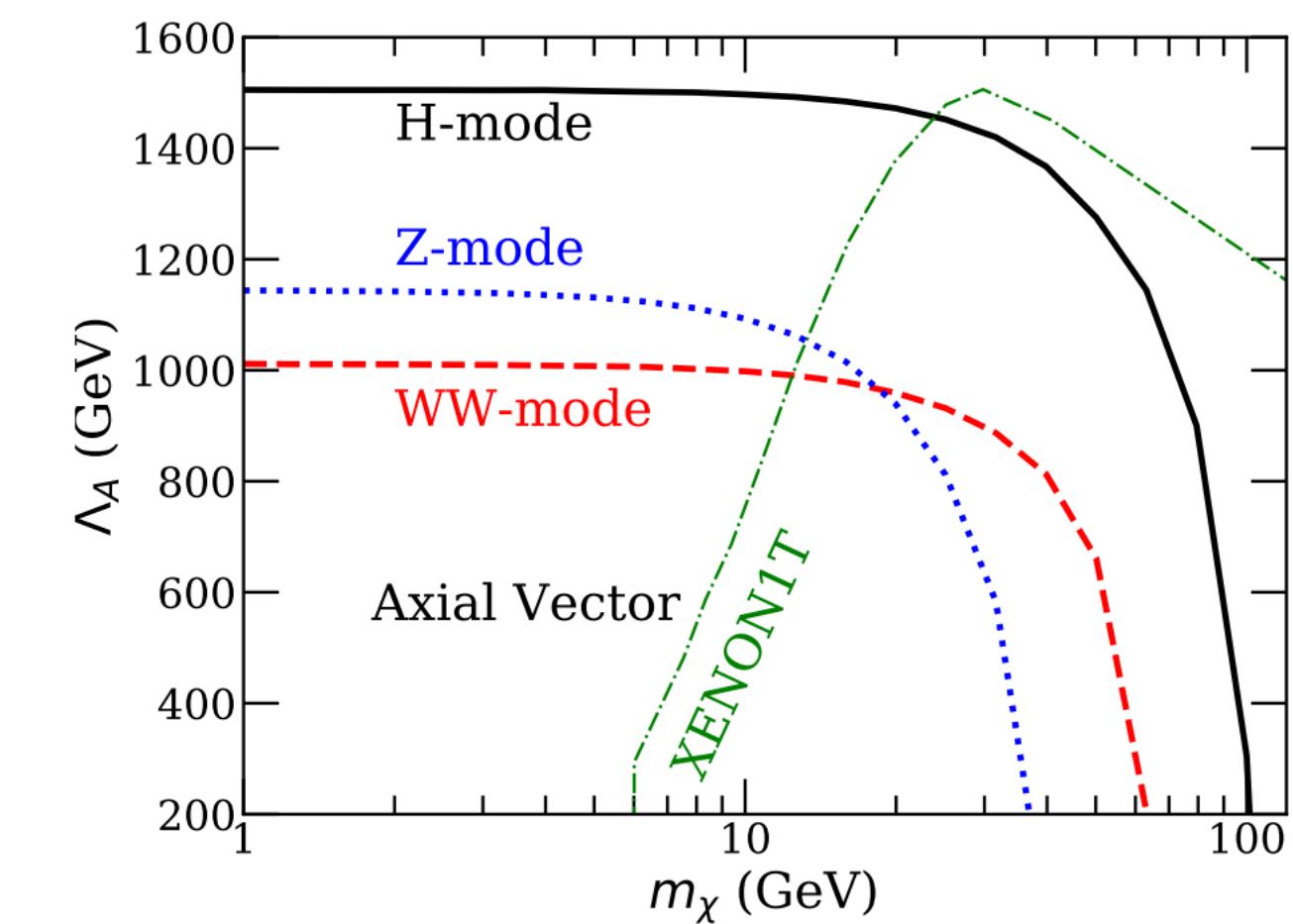
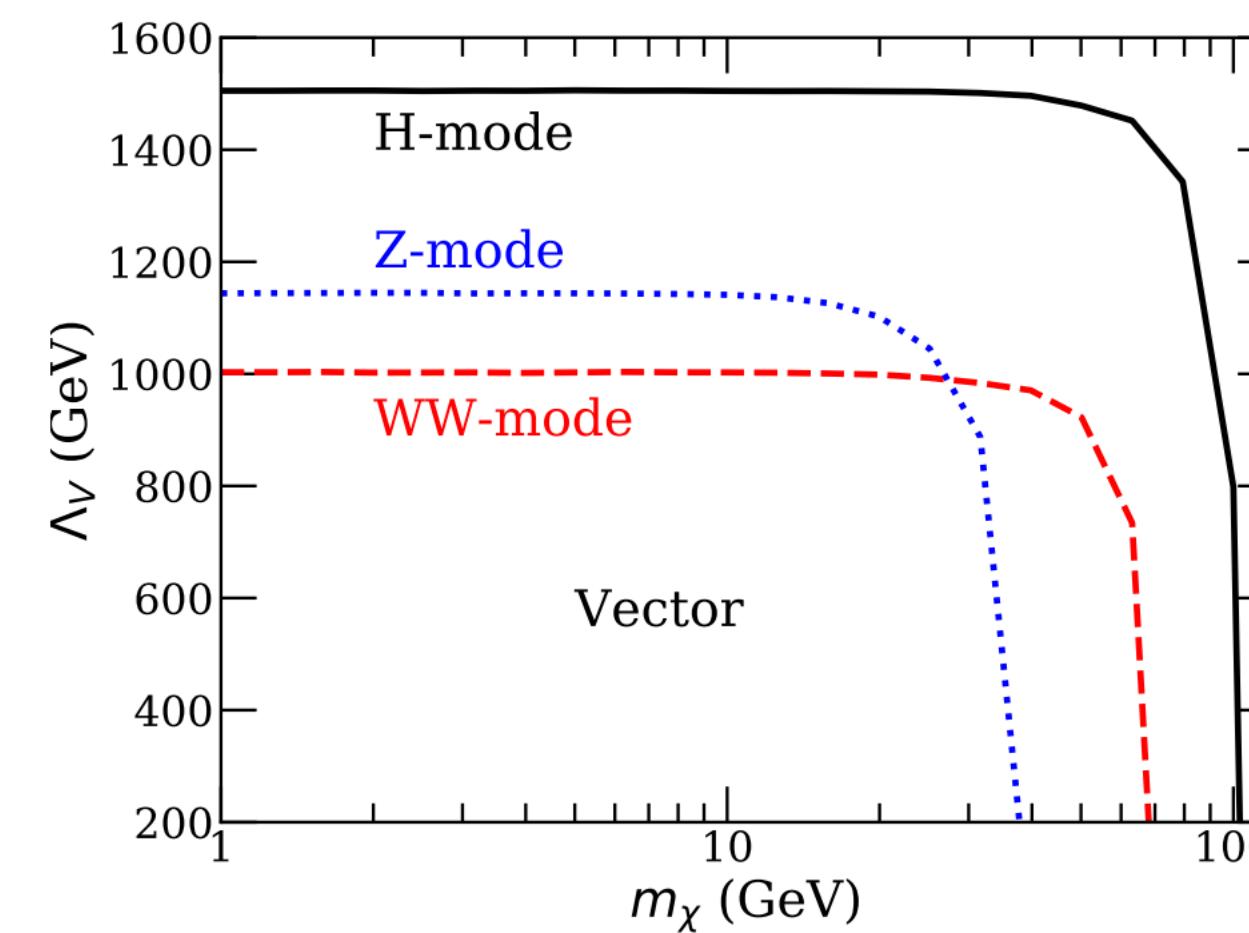
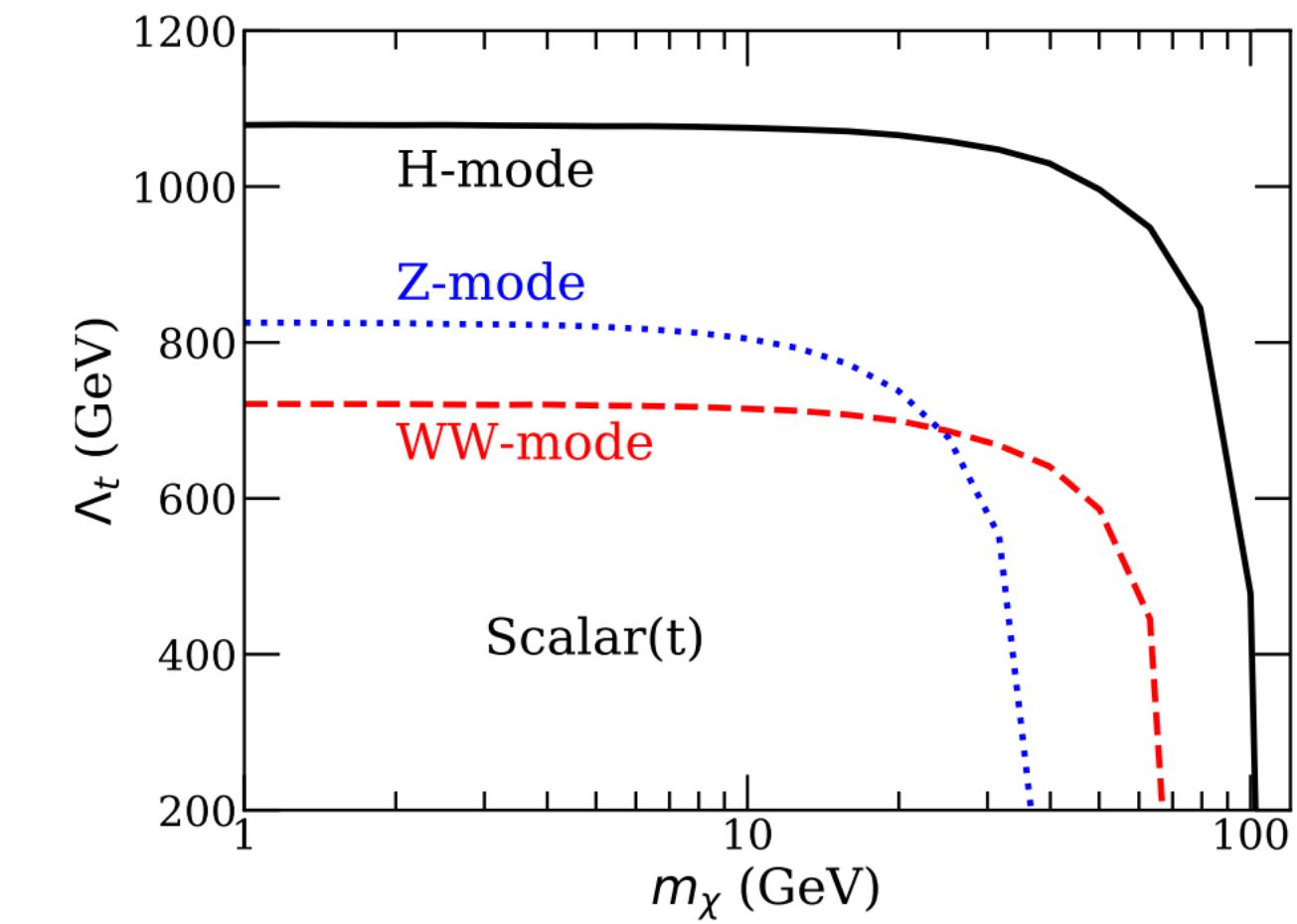
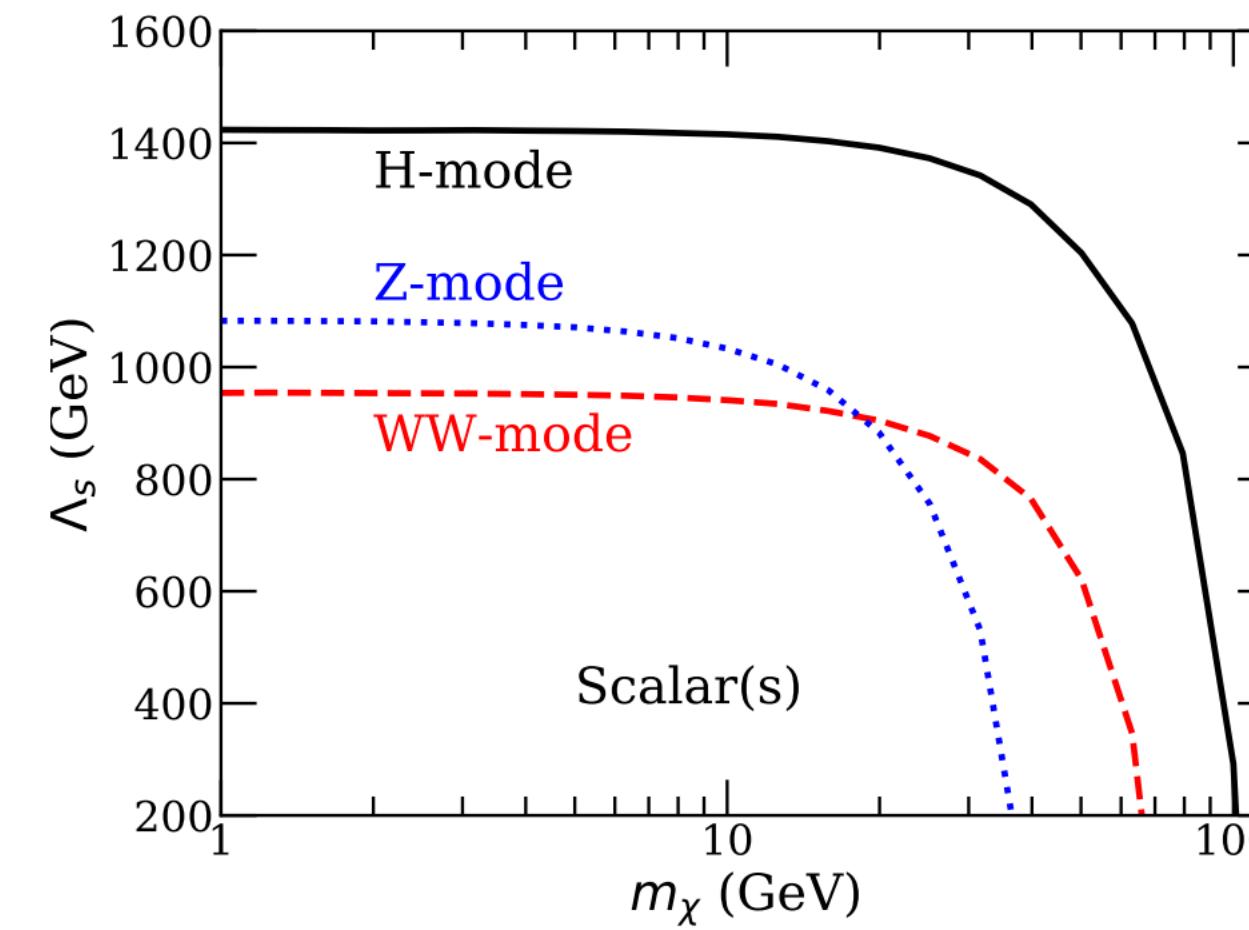
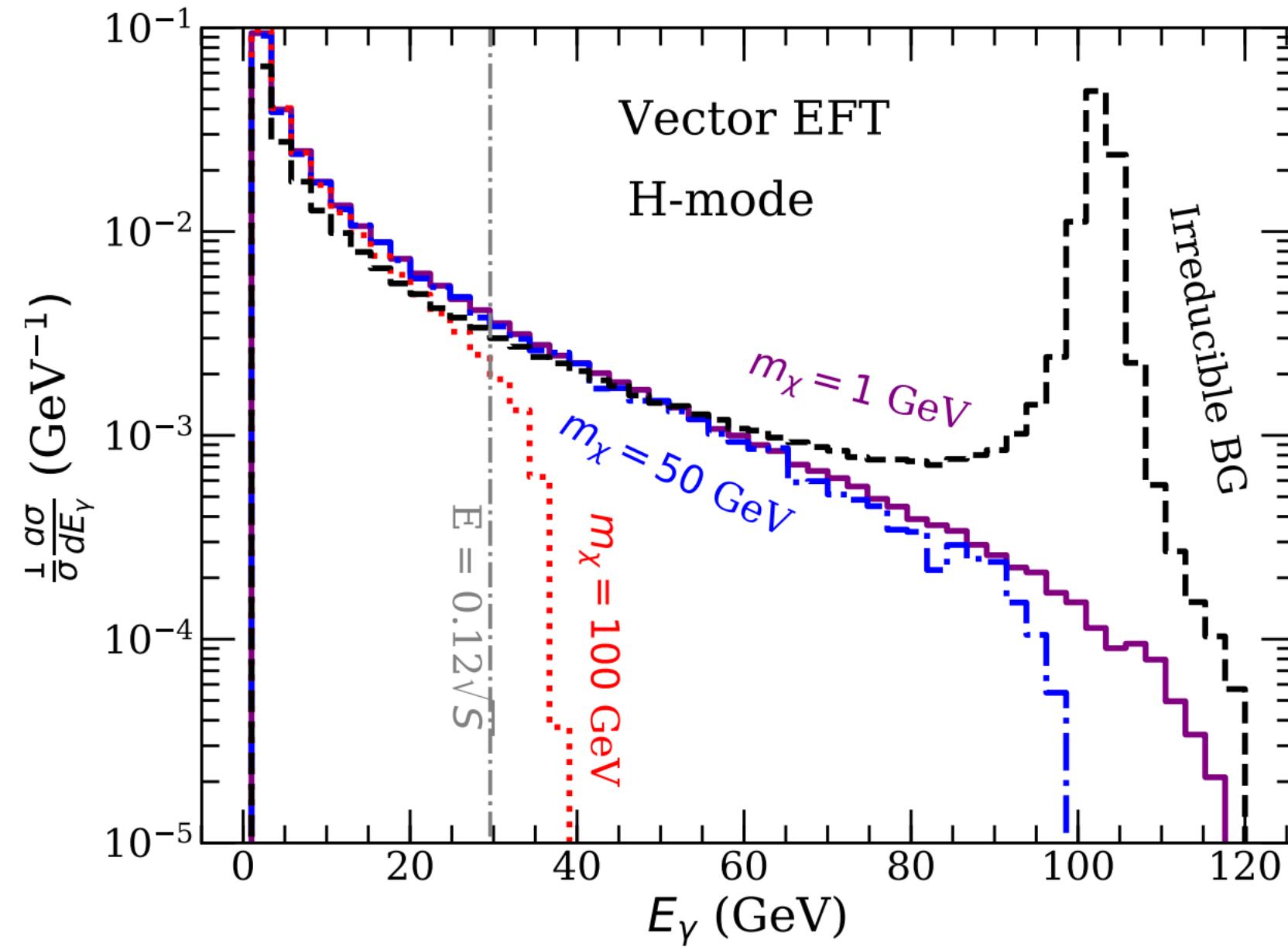
$$\mathcal{L} = \frac{1}{\Lambda_s^2} \bar{\chi} \chi \bar{\ell} \ell,$$

$$\mathcal{L} = \frac{1}{\Lambda_A^2} \bar{\chi} \gamma_\mu \gamma_5 \chi \bar{\ell} \gamma^\mu \gamma_5 \ell,$$

$$\mathcal{L} = \frac{1}{\Lambda_t^2} \bar{\chi} \ell \bar{\ell} \chi$$

# Constraints on DM EFTs

- The signal distribution at Z ( $2.6 \text{ ab}^{-1}$ ), H ( $5.6 \text{ ab}^{-1}$ ), WW modes ( $16 \text{ ab}^{-1}$ )



- The Z/Higgs factory modes can provide competitive sensitivity comparing with existing colliders and direct detection searches

# Searching fermionic DM absorption at CEPC

- Light DM can provide enough energy to direct detection via down-scattering (fermionic)/absorption (bosonic)
- Inelastic DM  $\chi_2 q \rightarrow \chi_1 q$ , Luminous DM long-lived  $\chi_2 \rightarrow \chi_1 \gamma$
- Fermionic DM down-scattering:  $\chi e \rightarrow \nu e$ , no  $Z_2$  protection

Shaofeng Ge et al, 2201.11497 (JHEP)

$$\mathcal{O}_{e\nu\chi}^S \equiv (\bar{e}e)(\bar{\nu}_L\chi_R),$$

$$\mathcal{O}_{e\nu\chi}^P \equiv (\bar{e}i\gamma_5 e)(\bar{\nu}_L\chi_R),$$

$$\mathcal{O}_{e\nu\chi}^V \equiv (\bar{e}\gamma_\mu e)(\bar{\nu}_L\gamma^\mu\chi_L),$$

$$\mathcal{O}_{e\nu\chi}^A \equiv (\bar{e}\gamma_\mu\gamma_5 e)(\bar{\nu}_L\gamma^\mu\chi_L),$$

$$\mathcal{O}_{e\nu\chi}^T \equiv (\bar{e}\sigma_{\mu\nu} e)(\bar{\nu}_L\sigma^{\mu\nu}\chi_R),$$

- DM is stable, because

$$m_\chi < 2m_e$$

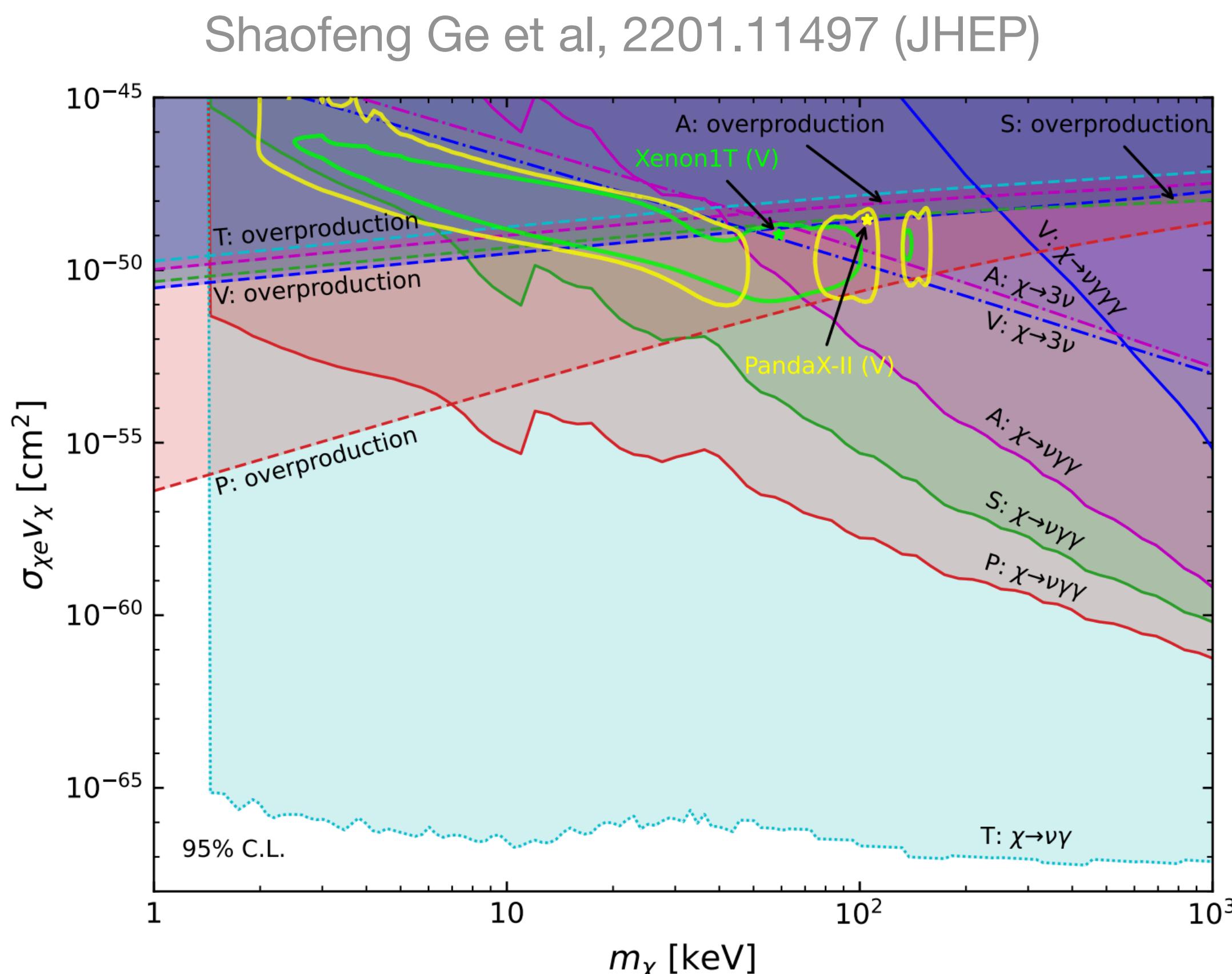
- Radiative decay to  $\nu + \gamma^n$  is small enough

$$(\bar{\nu}_L\chi_R)F^2, (\bar{\nu}_L\chi_R)F\tilde{F}$$

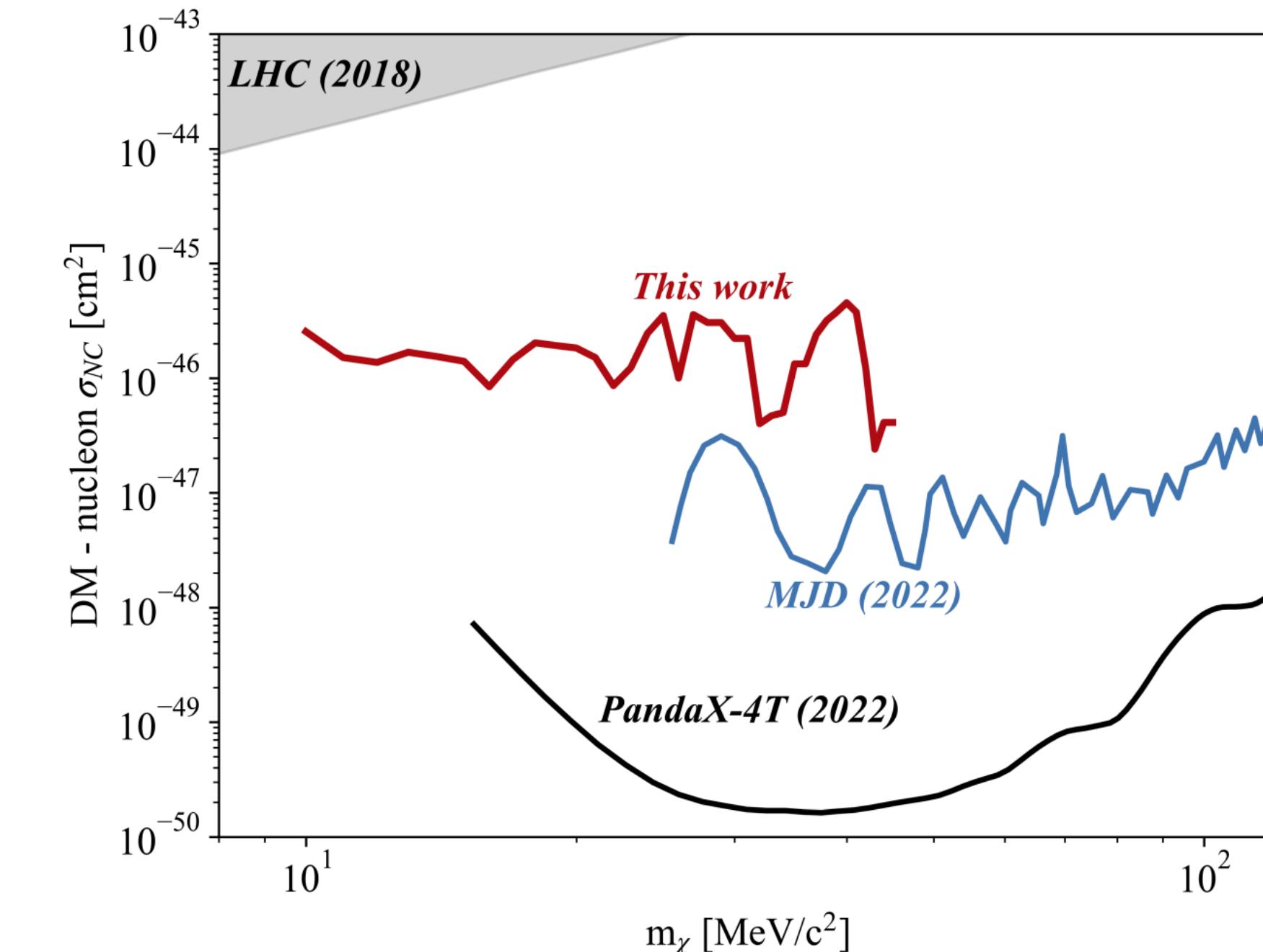

Operator \ Process	$\chi \rightarrow \nu\gamma$	$\chi \rightarrow \nu\gamma\gamma$	$\chi \rightarrow \nu\gamma\gamma\gamma$	$\chi \rightarrow 3\nu$
S: $\mathcal{O}_{e\nu\chi}^S$	✗	✓	✗	✗
P: $\mathcal{O}_{e\nu\chi}^P$	✗	✓	✗	✗
V: $\mathcal{O}_{e\nu\chi}^V$	✗	✗	✓	✓
A: $\mathcal{O}_{e\nu\chi}^A$	✗	✓	✗	✓
T: $\mathcal{O}_{e\nu\chi}^T$	✓	✗	✗!	✗!

# Searching fermionic DM absorption at CEPC

- Light DM can provide enough energy to direct detection via fermionic DM down-scattering:  $\chi e \rightarrow \nu e$



- Actively searched by DM direct detection experiments



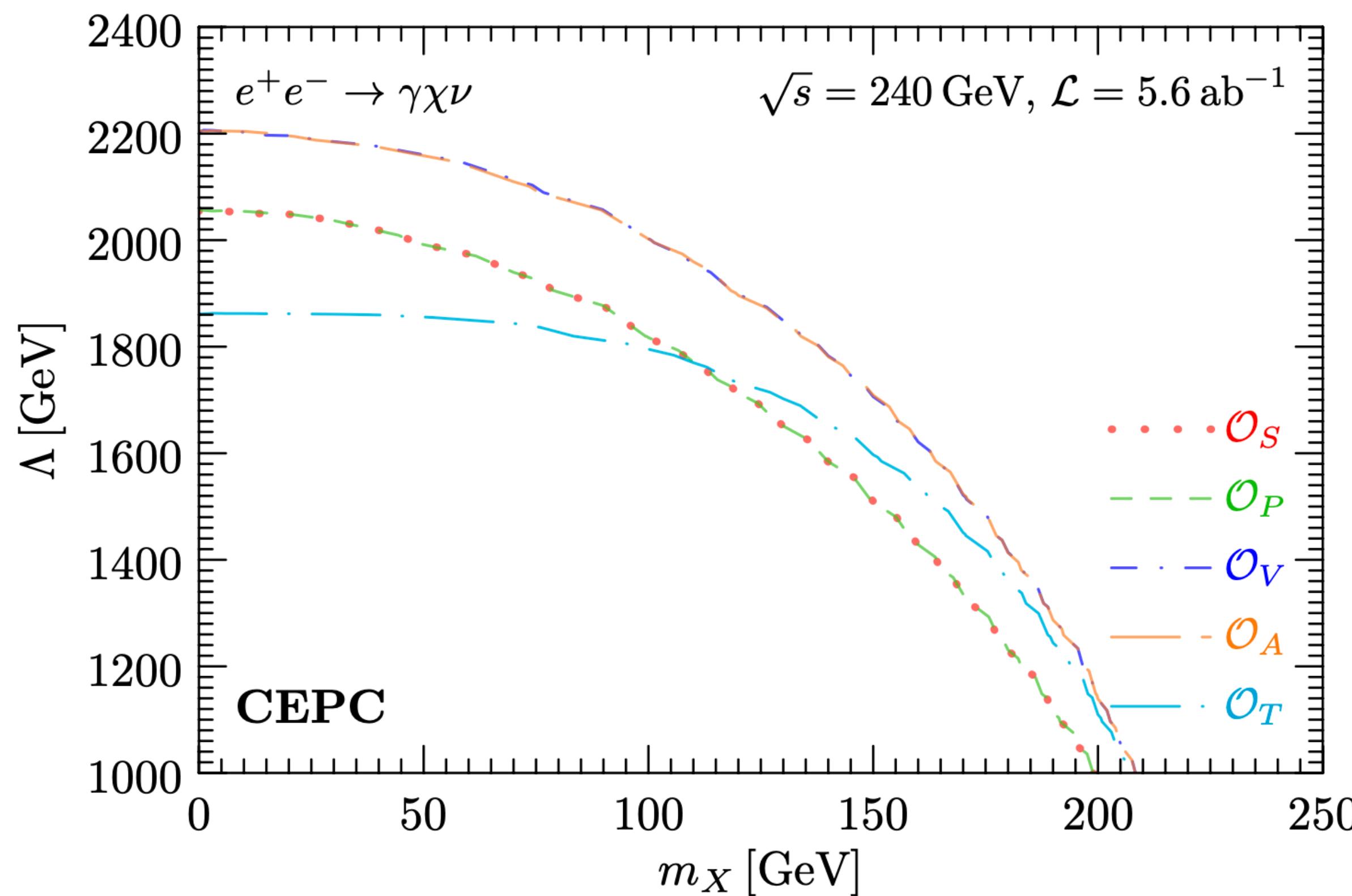
PANDAX 2206.02339 (PRL)  
CDEX 2209.00861(PRL)

# Searching fermionic DM absorption at CEPC

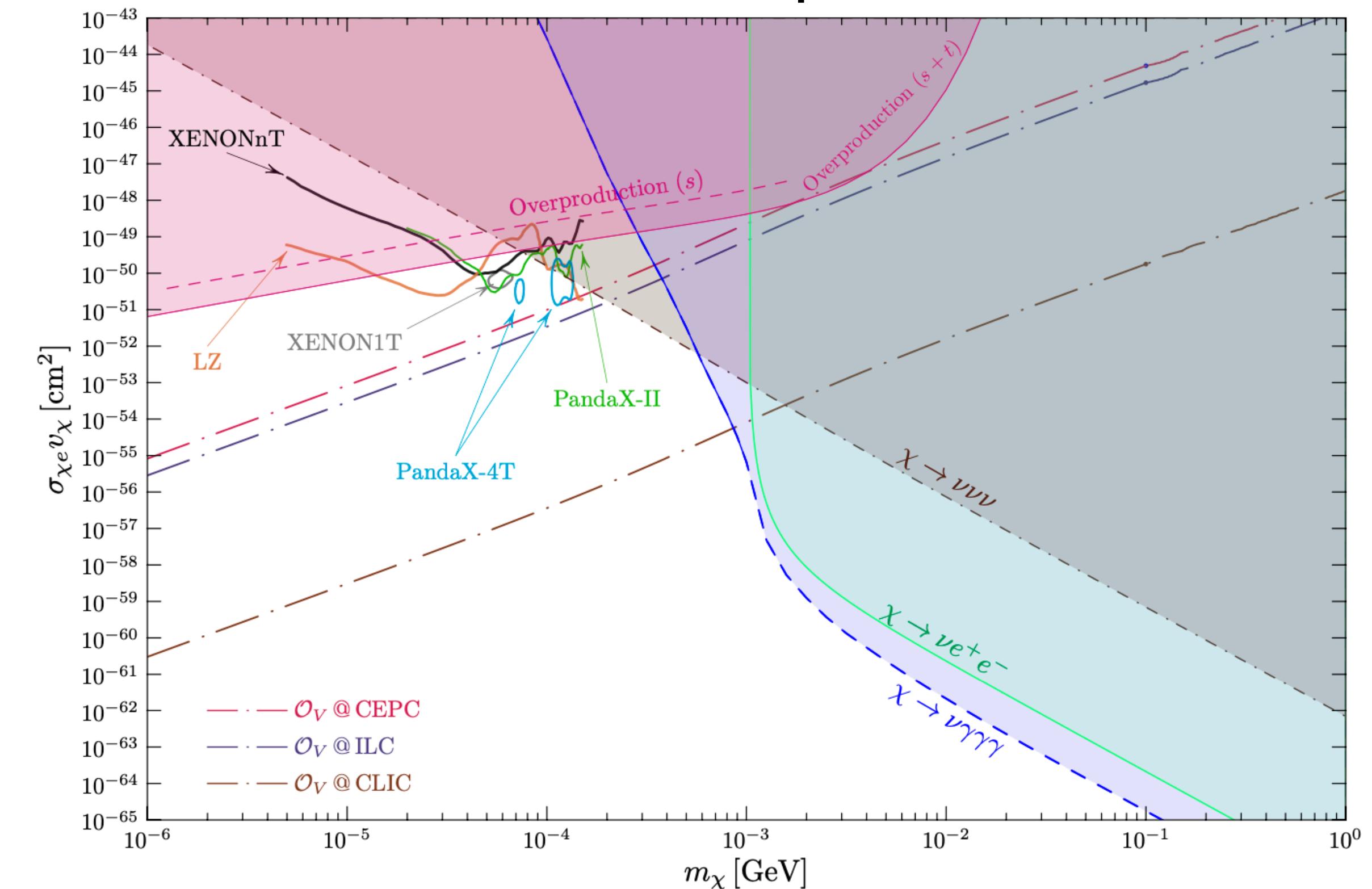
- The future ee collider can also test this scenario via

- $e^+e^- \rightarrow \gamma\nu\bar{\chi}, e^+e^-\nu\bar{\chi}, e^+e^-\nu\bar{\nu}$ , in multi-particle final states

Shaofeng Ge et al, 2306.00657



- Collider complementary between DM direct and indirect experiments



# Summary

- Future ee collider (CEPC/FCC-ee etc) provides valuable opportunities to test Beyond the Standard Model physics
  - Interactions with leptons and photons can be generally covered
  - Light particles are covered, even larger than  $\sqrt{s}/2$  through 3-body final states
  - Much heavier particle can be tested indirectly through precision measurements
  - Complementarity with other dark matter studies: direct detection, indirect detection
  - Suitable for searching for light dark sector bosons and its complementarity with gravitational wave studies
- CEPC offers a diverse and extensive research program covering various areas of fundamental physics

*Thank you!*

# Backup slides

