

中山大學天琴中心

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Precise calculations of electroweak phase transition dynamics

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Sun Yat-sen University

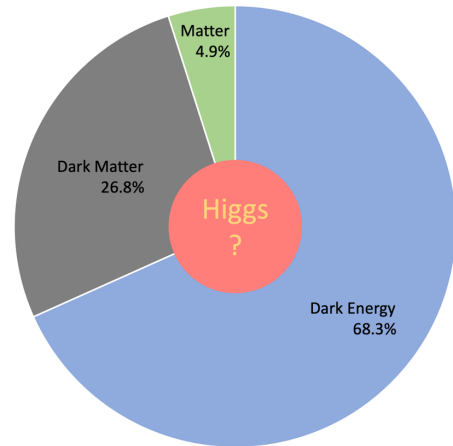
The 2023 international workshop on the Circular Electron Positron Collider (CEPC)
@University of Edinburgh, UK, 3-6 July, 2023



Motivation

What is the role of Higgs in the early universe ?

After the discovery of 125 GeV scalar, it becomes a realistic portal to study the fundamental physics and its deep connections to cosmology



How to test?

Hi, Higgs !
Let us help you to explore your confusion!

Future lepton colliders
(ILC, CEPC, FCC-ee)



Future Gravitational wave experiments (LISA)

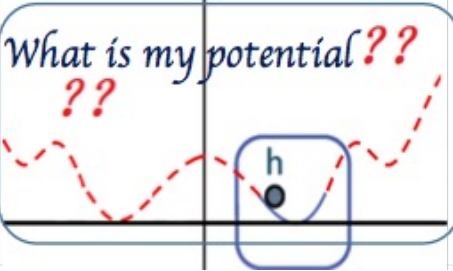
Hi, I am Higgs boson!
I have many questions?



What is the my role in
Baryon asymmetry of
the universe??
Electroweak
baryogenesis??

What is the my role in
dark matter models??

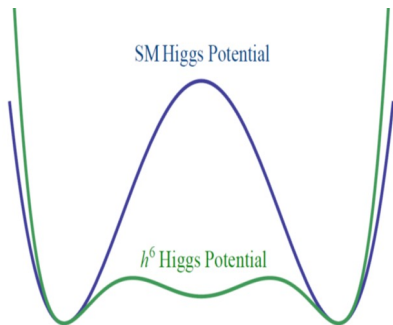
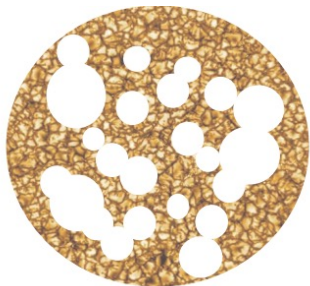
Why my mass is so light
compared to Planck mass??
Cosmological relaxion??



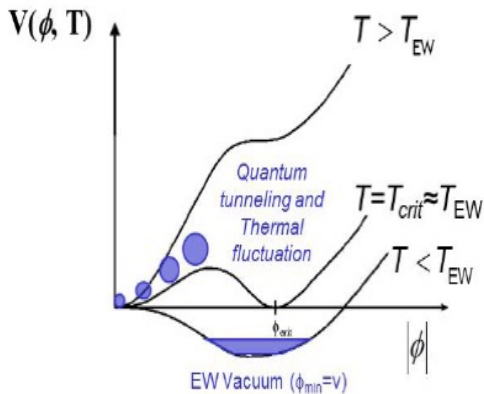


Motivation

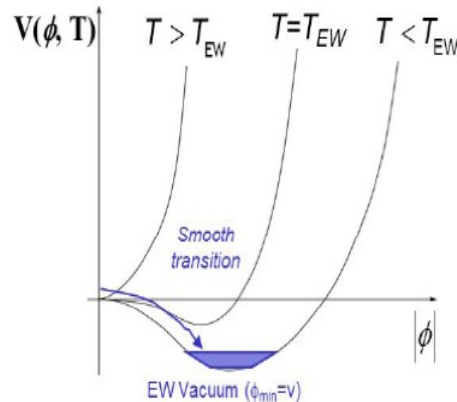
From
lattice
simulation



SFOPT for $m_H < 75$ GeV



Cross over for $m_H > 75$ GeV



Extension of the Higgs sector is needed to SFOPT for 125 GeV Higgs boson.

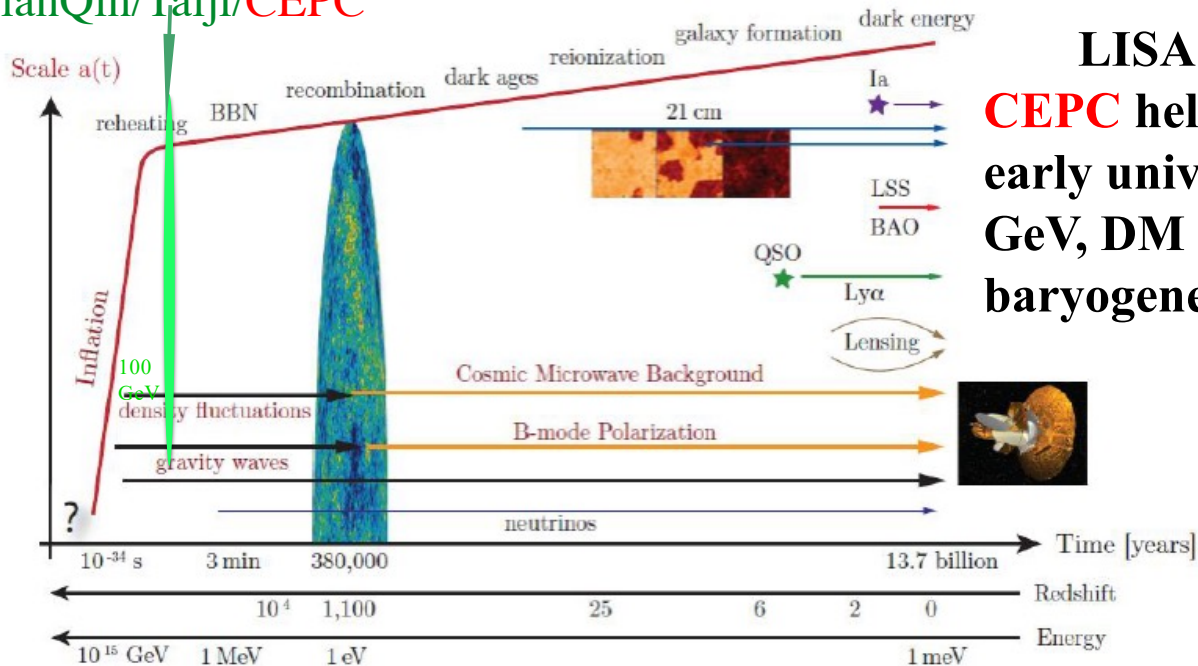
We discuss well-motivated extensions (baryogenesis, DM...) of Higgs section to realize strong first-order phase transition (SFOPT) with abundant cosmological effects.

EW phase transition and its GW signals becomes realistic after the discovery of Higgs by LHC and GW by LIGO.



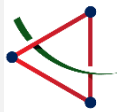
Motivation

EW phase transition/
baryogenesis:
LISA/TianQin/Taiji/CEPC

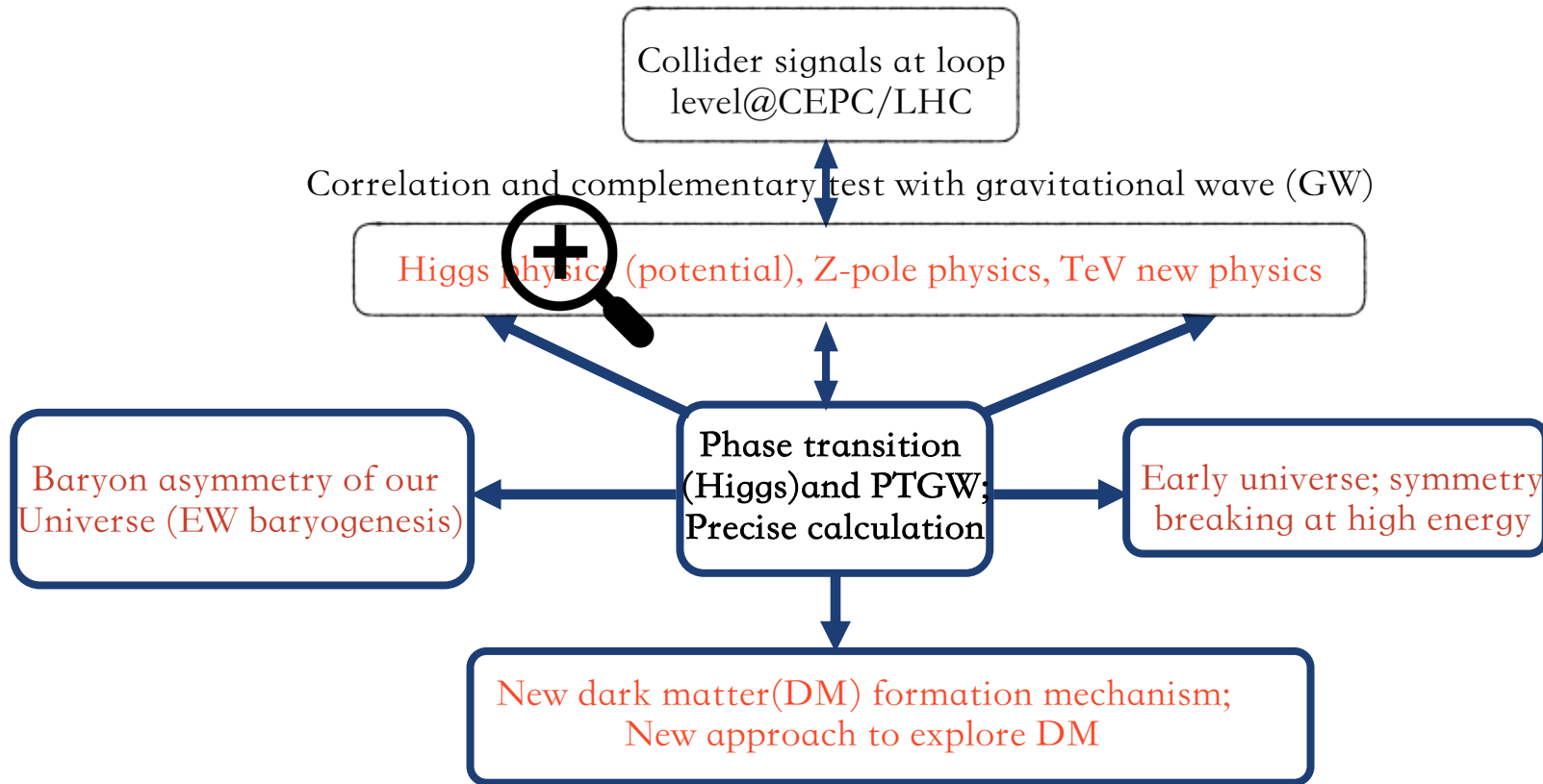


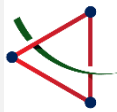
LISA in synergy with
CEPC helps to explore the
early universe around 100
GeV, DM and
baryogenesis.

credit:D.Baumann



Motivation

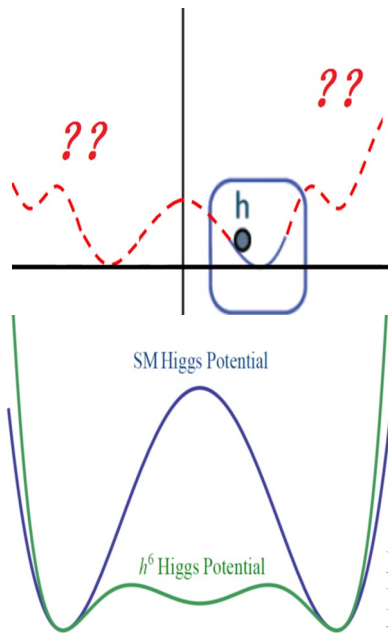




Higgs potential

What is the shape of Higgs potential?

Current data tells us nothing but the quadratic oscillation around the VEV 246 GeV with 125 GeV mass.



$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4$$



or
$$V(h) = \frac{1}{2}\mu^2 h^2 - \frac{\lambda}{4}h^4 + \frac{1}{\Lambda^2}h^6$$

Produce a SFOPT, large deviation of Higgs trilinear coupling and GW

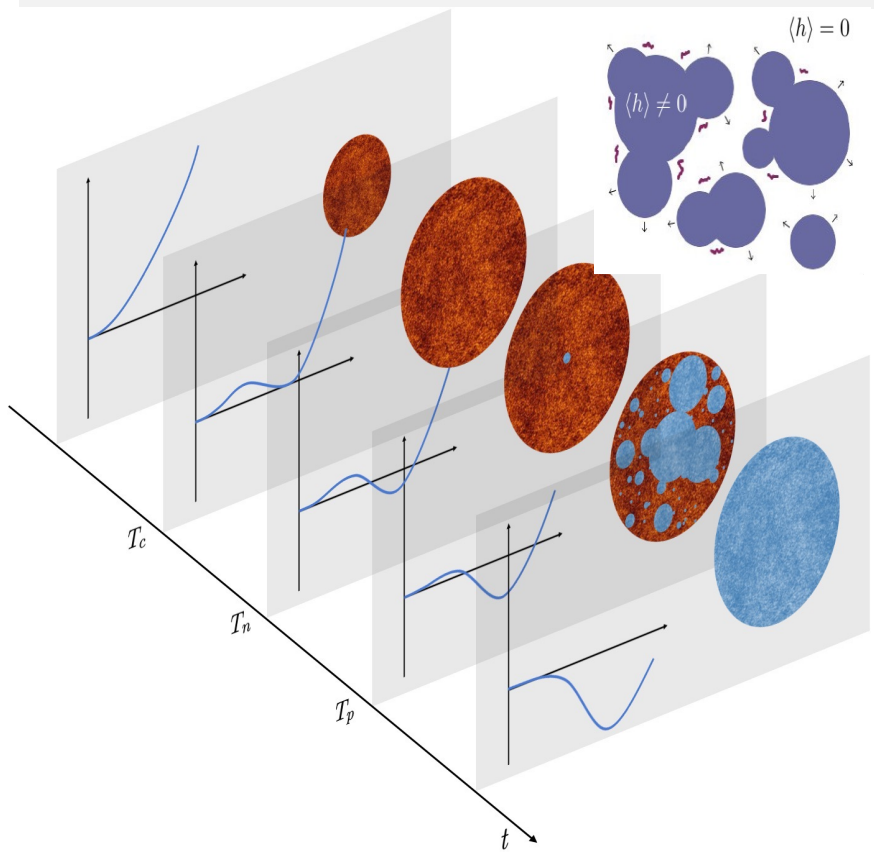
Xinmin Zhang Phys.Rev. D47 (1993) 3065-3067; C. Grojean, G. Servant, J. Well PRD71(2005)036001
D.J.H. Chung, A. J. Long, Lian-tao Wang Phys.Rev. D87(2013) 023509

FPH, et.al, Phys.Rev.D94(2016)no.4,041702 ; FPH, et.al, Phys.Rev.D93 (2016) no.10,103515

arXiv:1511.06495, Nima Arkani-Hamed et. al.; PreCDR of CEPC; arXiv: [1811.10545](https://arxiv.org/abs/1811.10545), CDR of CEPC



Phase transition GW in a nutshell



calculate the finite-temperature effective potential using the thermal field theory: free energy density.

$$V_{\text{eff}}^{(1)}(\bar{\phi}) = \sum_i n_i \left[\int \frac{d^D p}{(2\pi)^D} \ln(p^2 + m_i^2(\bar{\phi})) + J_{\text{B,F}} \left(\frac{m_i^2(\bar{\phi})}{T^2} \right) \right]$$

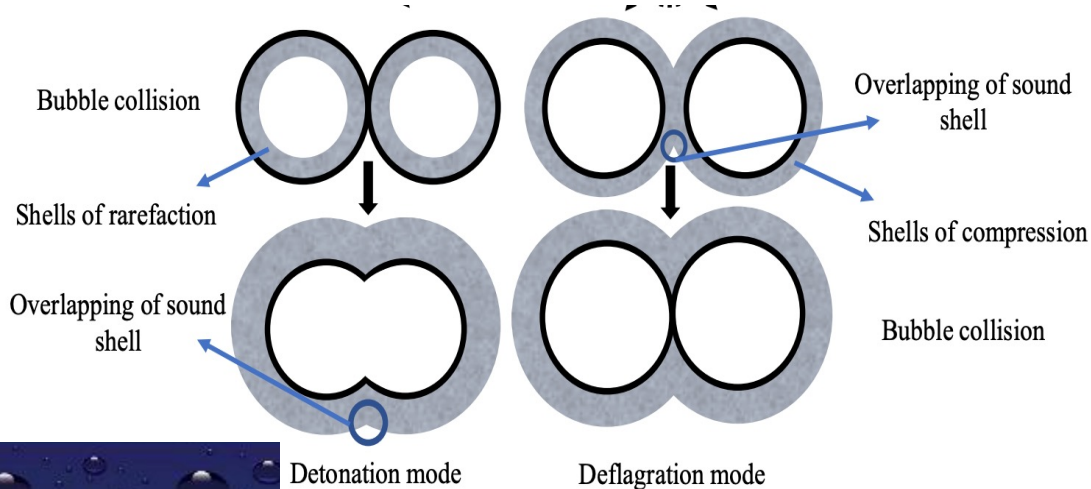
$$S(T) = \int d^4 x \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 + V_{\text{eff}}(\phi, T) \right]$$

$$\Gamma = \Gamma_0 e^{-S(T)}$$

Xiao Wang, **FPH**, Xinmin Zhang, JCAP05(2020)045



Phase transition GW in a nutshell



$$R_{\mu\nu} - \frac{1}{2}R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$h_{ij} \simeq \frac{2G}{c^4 r} \ddot{Q}_{ij}^{TT}(t - r/c)$$

E. Witten, Phys. Rev. D 30, 272 (1984)

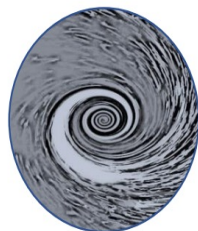
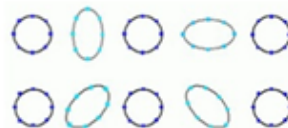
C. J. Hogan, Phys. Lett. B 133, 172 (1983);

M. Kamionkowski, A. Kosowsky and M. S. Turner, Phys. Rev. D 49, 2837 (1994))

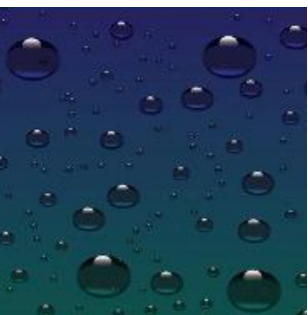
EW phase transition

GW becomes more interesting and realistic after the discovery of

Higgs by LHC and GW by LIGO.



Turbulence



Xiao Wang, **FPH**, Xinmin Zhang, JCAP05(2020)045

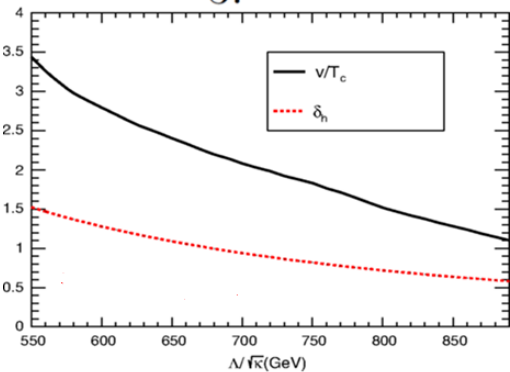


Higgs potential

Correlate collider and GW: double test on Higgs potential from particle to wave

SFOPT leads to obvious deviation of the tri-linear Higgs coupling

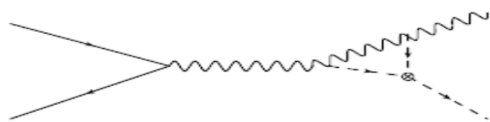
$$\mathcal{L}_{hhh} = -\frac{1}{3!}(1 + \delta_h)A_h h^3$$



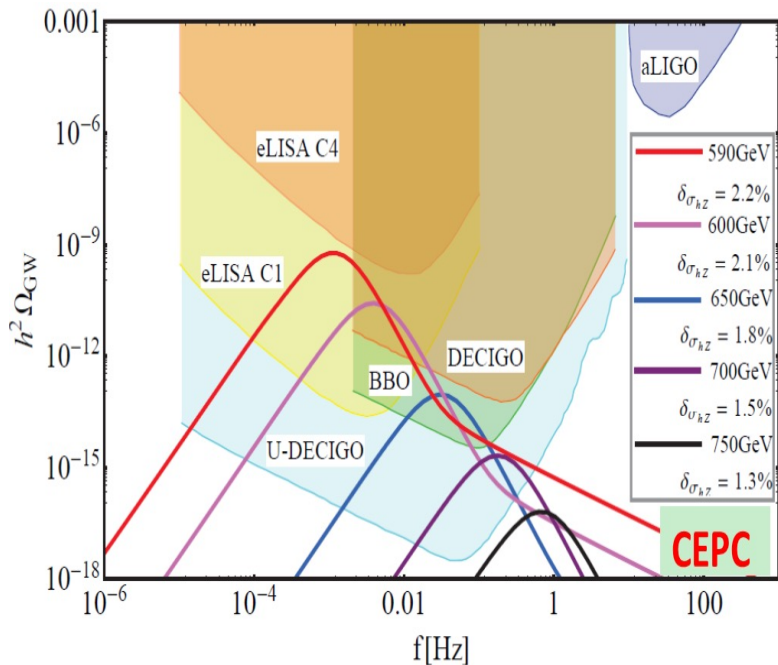
At one-loop level, deviation of the tri-linear Higgs coupling

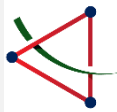
$$\delta_h \in (0.6, 1.5)$$

The Circular Electron Positron Collider (CEPC), ILC, FCC-ee can precisely test this scenario by precise measurements of the hZ cross section ($e^+e^- \rightarrow hZ$). SM NNLO hZ cross section recently by Lilin Yang, et al 2016, Yu Jia et al 2016



$$\delta_\sigma = \frac{\sigma_{hz, \delta_h \neq 0}}{\sigma_{hz, SM}} - 1$$





Higgs potential

SM EFT

$$\mathcal{L} \supset -\mu^2 |H|^2 - \lambda |H|^4 + c_6 |H|^6$$

$$+ c_T \mathcal{O}_T + c_{WW} \mathcal{O}_{WW} + \text{other dimension-six operators}$$

$$\delta_{\sigma(hZ)} \approx (0.26c_{WW} + 0.01c_{BB} + 0.04c_{WB} - 0.06c_H - 0.04c_T + 0.74c_L^{(3)\ell} + 0.28c_{LL}^{(3)\ell} + 1.03c_L^\ell - 0.76c_R^e) \times 1 \text{ TeV}^2 + 0.016\delta_h,$$

SFOPT produces large modification of trilinear Higgs coupling

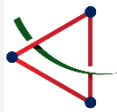
δ_h

c_6

dominates the hZ cross section deviation

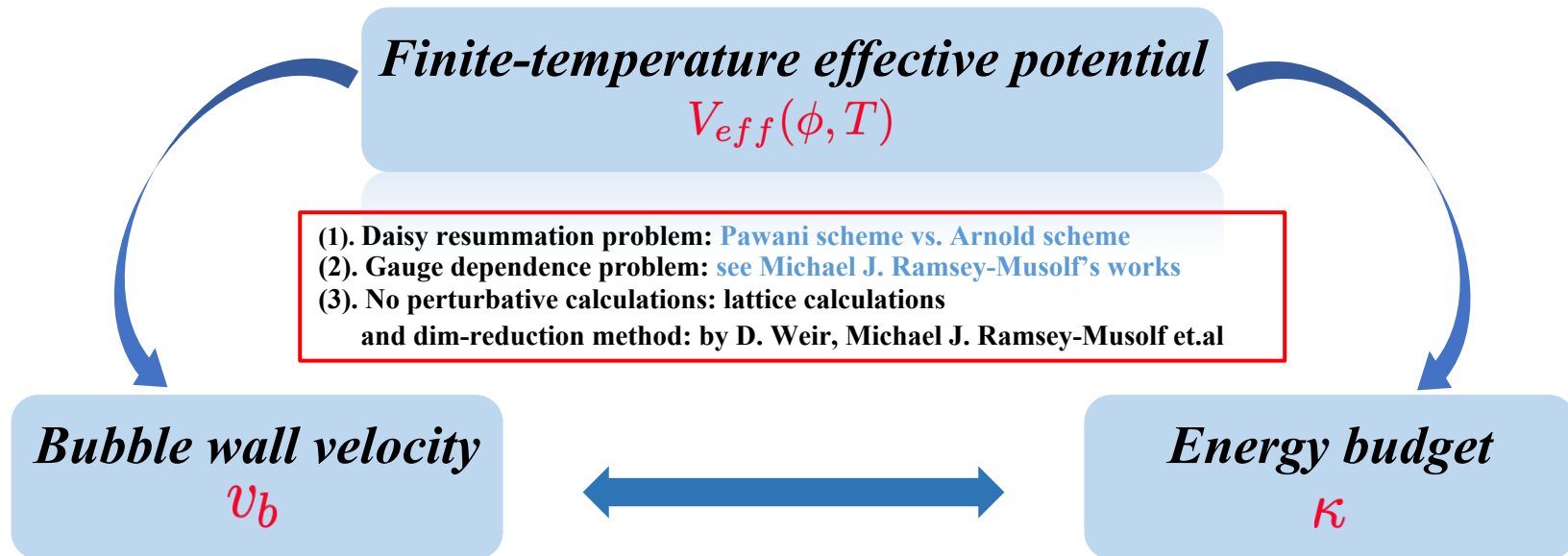
Taking a general study of the scalar extended models and the composite Higgs model as examples, we find that the Higgs sextic scenario still works well after considering all the dim-6 operators and the precise measurements.

Qing-Hong Cao, **FPH**, Ke-Pan Xie, Xinmin Zhang, arXiv:1708.0473.



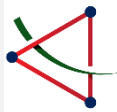
Phase transition Dynamics

Precise predictions on the phase transition dynamics and its GW signals

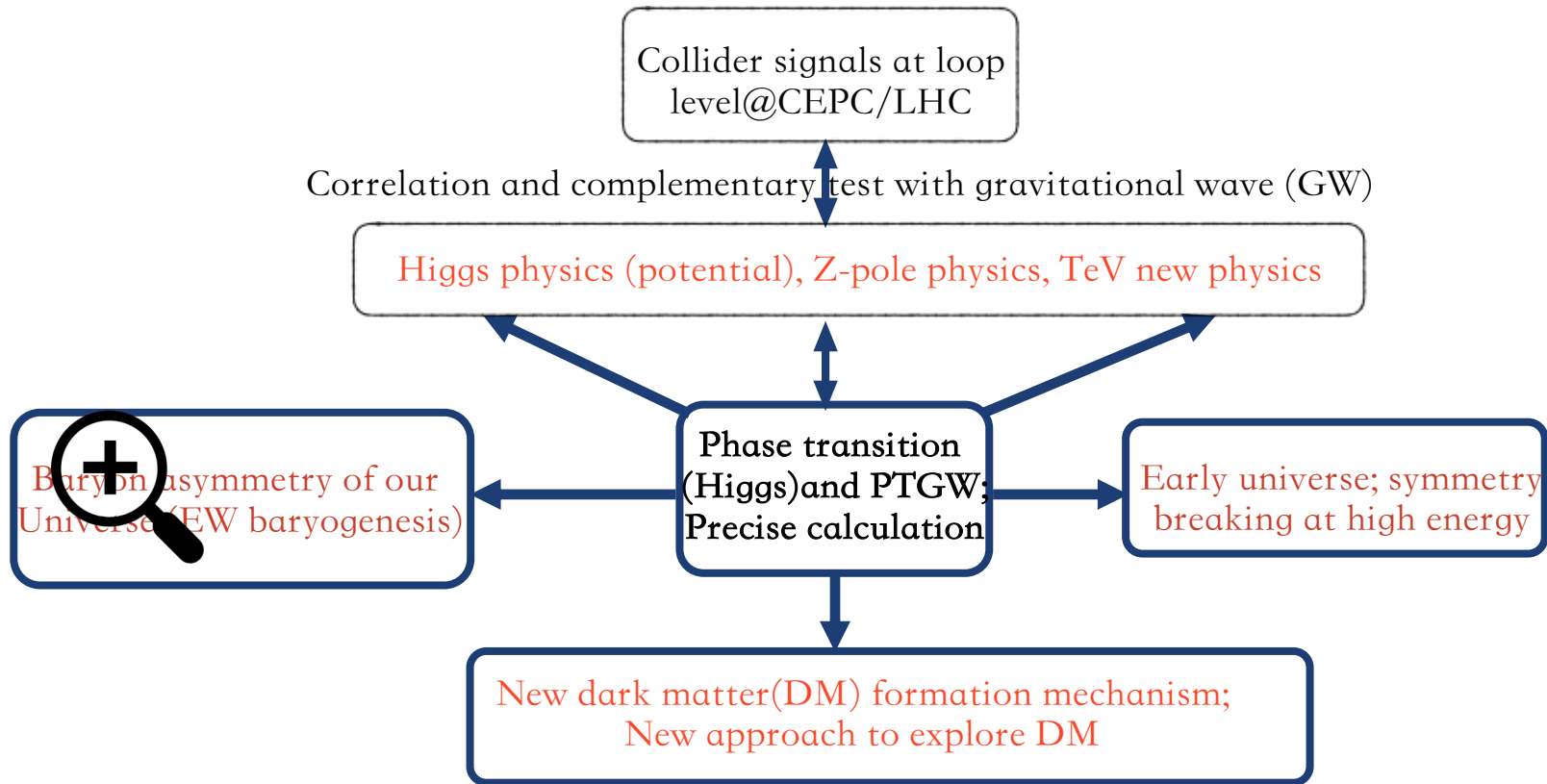


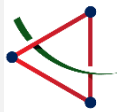
S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang, arXiv:2007.10343,
Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith,
arXiv:2009.14295v2
Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903
Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005

F. Giese, T. Konstandin, K. Schmitz and J. van de Vis, arXiv:2010.09744
Xiao Wang, **FPH** and Xinmin Zhang, Phys.Rev.D 103 (2021) 10, 103520
Xiao Wang, Chi Tian, **FPH**, arXiv: 2301.12328



Motivation





SFOPT and EW baryogenesis

A long standing problem in particle cosmology is the origin of baryon asymmetry of the universe.

After discovery of Higgs boson@LHC & GW @aLIGO, EW baryogenesis becomes a timely and testable scenario.

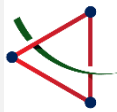
SM technically has all the 3 elements for baryogenesis(Sakharov conditions)

- B violation from anomaly in B+L current;
- C and CP-violation: CKM matrix, but too weak, need new CP-violating sources;
- **Departure from thermal equilibrium: SFOPT with expanding Higgs bubble wall**



$$\eta_B = n_B/n_\gamma = 5.8 - 6.6 \times 10^{-10}$$

(CMB, BBN)



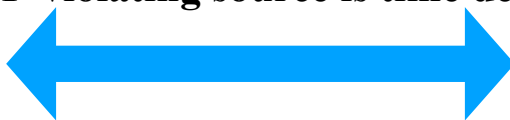
SFOPT and EW baryogenesis

Question: How to alleviate the tension of EW baryogenesis?

Answer: Dynamical CP-violating source

**Large enough
CP-violating source
for successful
EW baryogenesis**

Alleviate by assuming the
CP-violating source is time dependent



Dynamical/cosmological evolve

**Negligible
CP-violating source
at current time
to avoid strong EDM
constraints**

Effective field theory: **FPH**, Zhuoni Qian, Mengchao Zhang, Phys.Rev. D98 (2018) no.1, 015014, **FPH**, Chong Sheng Li, Phys. Rev. D 92, 075014 (2015);

Renormalizable model: Complex 2HDM, Xiao Wang, **FPH**, Xinmin Zhang, arXiv: 1909.02978, work in progress with Eibun Senaha, Xiao Wang in an extended IDM model
Baltes, T. Konstandin and G. Servant, arXiv:1604.04526; I. Baltes, T. Konstandin and G. Servant, JHEP 1612, 073 (2016);

S. Bruggisser, T. Konstandin and G. Servant, JCAP 1711, no. 11, 034 (2017)

Matthew Reece recently has more deep study on the dynamical cp-violation.



SFOPT and EW baryogenesis

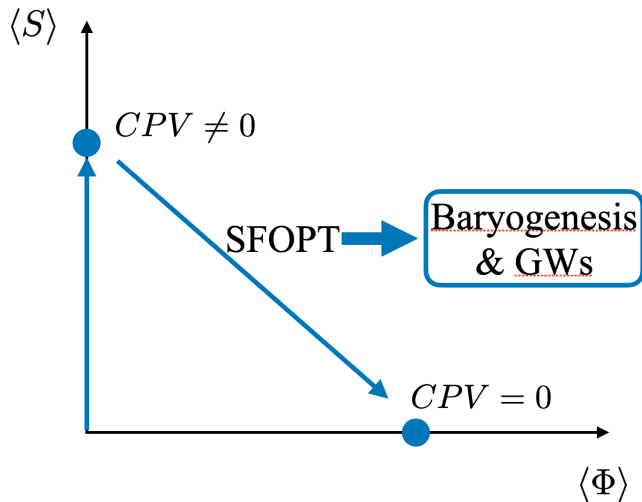
Taking the effective scenario as a representative example:

$$\mathcal{L}_{\text{SM}} - y_t \frac{\eta}{\Lambda} S \bar{Q}_L \tilde{\Phi} t_R + \text{H.c.} + \frac{1}{2} \partial_\mu S \partial^\mu S + \frac{1}{2} \mu^2 S^2 - \frac{1}{4} \lambda S^4 - \frac{1}{2} \kappa S^2 (\Phi^\dagger \Phi)$$

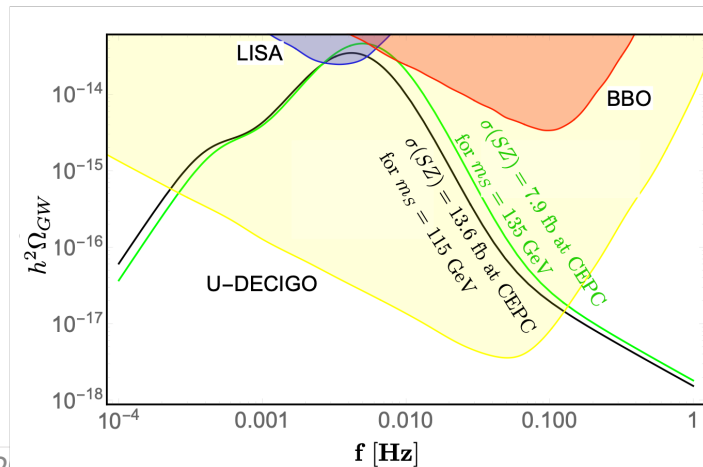
The singlet and the dim-5 operator can come From many types composite Higgs model, [arXiv:0902.1483](#), [arXiv:1703.10624](#)

[arXiv:1704.08911](#), J. M. Cline and K. Kainulainen, JCAP **1301**, 012 (2013)

J. R. Espinosa, B. Gripaios, T. Konstandin and F. Riva, JCAP **1201**, 012 (2012)



Phys.Rev. D98 (2018) no.1, 015014, **FPH**, Zhuoni Qian, Mengchao Zhang
lots of new works unmentioned here





EW baryogenesis & Z-pole physics

Renormalizable DM and EW baryogenesis with dynamical CP-violating source.

arXiv:1905.10283, Phys.Rev.D 100 (2019) 3, 035014, FPH, Eibun Senaha

and work in progress with Eibun Senaha

$$V_0(\Phi, \eta) = \mu_1^2 \Phi^\dagger \Phi + \mu_2^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\Phi^\dagger \Phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \lambda_3 (\Phi^\dagger \Phi) (\eta^\dagger \eta) + \lambda_4 (\Phi^\dagger \eta) (\eta^\dagger \Phi) + \left[\frac{\lambda_5}{2} (\Phi^\dagger \eta)^2 + \text{h.c.} \right],$$

The new lepton Yukawa interaction is

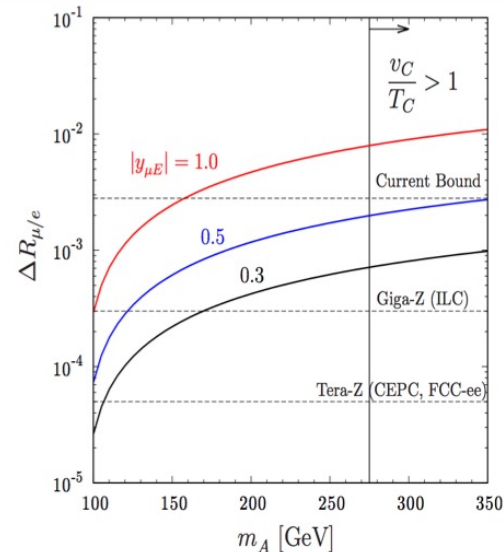
$$-\mathcal{L}_Y \ni y_{ij} \bar{\ell}_{iL} \eta E_{jR} + m_{E_i} \bar{E}_{iL} E_{iR} + \text{h.c.}$$

vector-like lepton (E_i)

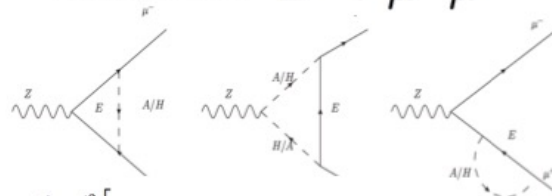
D. Borah, S. Sadhukhan and S. Sahoo, Phys. Lett. B **771**, 624 (2017).

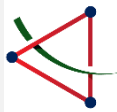
L. Calibbi, R. Ziegler and J. Zupan, JHEP **1807**, 046 (2018).

D. Borah, P. S. B. Dev and A. Kumar, Phys. Rev. D **99**, no. 5, 055012 (2019).

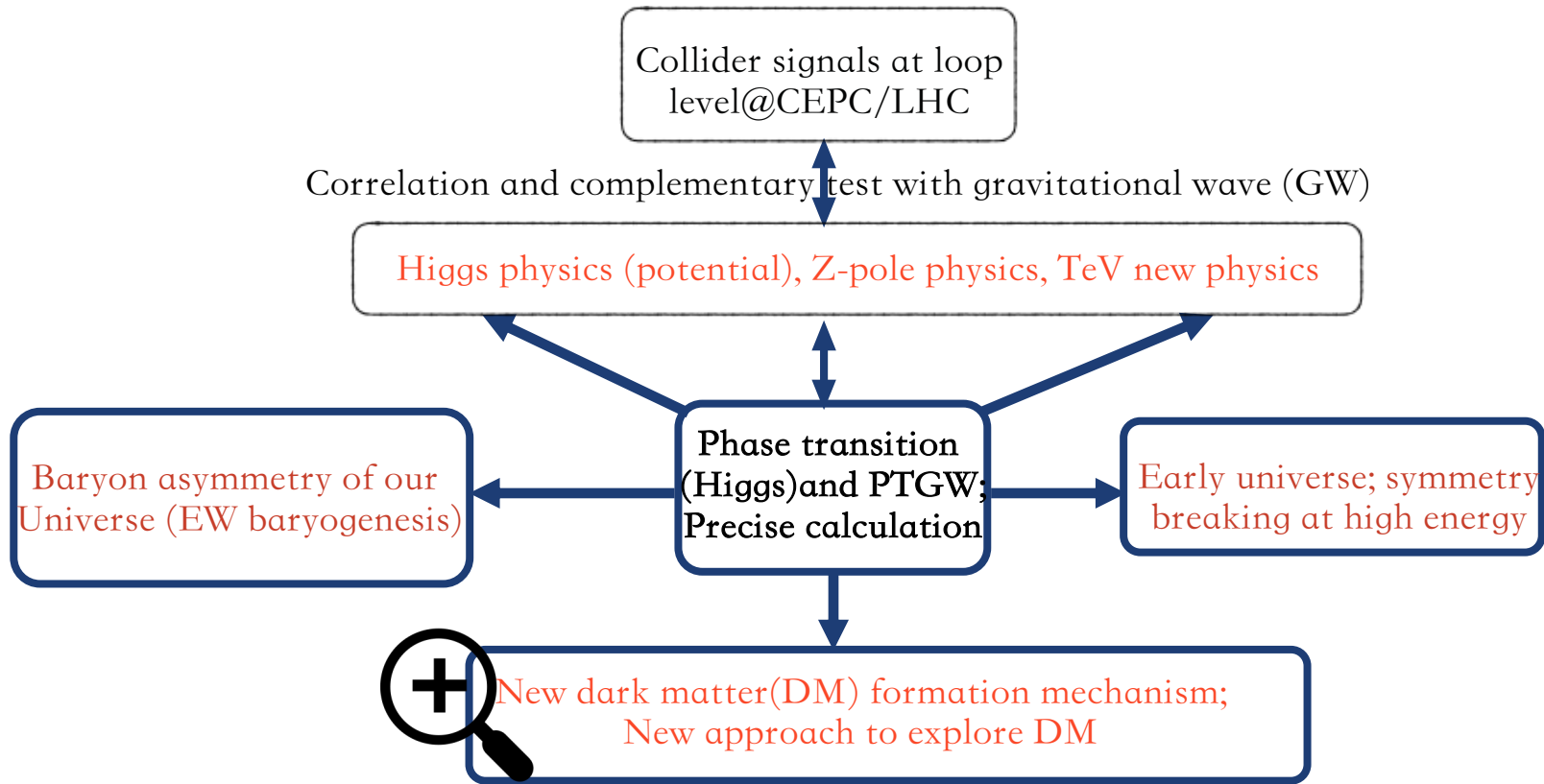


Indirect search by $Z \rightarrow \mu^+ \mu^-$





Motivation





SFOPT and new DM mechanism/signal

- **New** The observation of GW by LIGO has initiated a new era of exploring DM by GW.
- **DM can trigger a SFOPT in the early universe, which can leads to detectable GW signals.**

Hearing the signal of dark sectors with gravitational wave detectors

J.Jaeckel, V. V. Khoze, M. Spannowsky, Phys.Rev. D94 (2016) no.10, 103519

Zhaofeng Kang, et.al. arXiv:2101.03795

Zhaofeng Kang, et. al. arXiv:2003.02465

Yan Wang, Chong Sheng Li, and **FPH**, arXiv:2012.03920

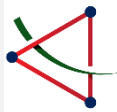
FPH, Eibun Senaha Phys.Rev. D100 (2019) no.3, 03501

FPH PoS ICHEP2018 (2019) 397

FPH, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028

FPH, Jiang-Hao Yu, Phys.Rev. D98 (2018) no.9, 095022

FPH, Xinmin Zhang, Phys.Lett. B788 (2019) 288-29



SFOPT and new DM mechanism/signal

Inert Doublet Models

$$V_0 = M_D^2 D^\dagger D + \lambda_D (D^\dagger D)^2 + \lambda_3 \Phi^\dagger \Phi D^\dagger D \\ + \lambda_4 |\Phi^\dagger D|^2 + (\lambda_5/2)[(\Phi^\dagger D)^2 + h.c.],$$

mixed singlet-doublet model

$$V_0 = \frac{1}{2} M_S^2 S^2 + M_D^2 H_2^\dagger H_2 + \frac{1}{2} \lambda_S S^2 |\Phi|^2 + \lambda_3 \Phi^\dagger \Phi H_2^\dagger H_2 \\ + \lambda_4 |\Phi^\dagger H_2|^2 + \frac{\lambda_5}{2} [(\Phi^\dagger H_2)^2 + \text{H.c.}] + A [S \Phi H_2^\dagger + \text{H.c.}].$$

mixed singlet-triplet model

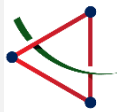
$$V_0 = \frac{1}{2} M_S^2 S^2 + M_\Sigma^2 \text{Tr}(H_3^2) + \kappa_\Sigma \Phi^\dagger \Phi \text{Tr}(H_3^2) \\ + \frac{\kappa}{2} |\Phi|^2 S^2 + \xi S \Phi^\dagger H_3 \Phi.$$

provide natural
DM candidate

produce SFOPT and phase transition
GW

FPH, Jiang-Hao Yu, Phys.Rev. D98 (2018) no.9, 095022

Yan Wang, Chong Sheng Li, and FPH, Phys.Rev.D 104 (2021) 5, 053004;



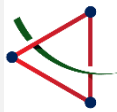
SFOPT and new DM mechanism/signal

New dark matter (DM) production scenario filtered by the bubbles of the SFOPT in the early universe.

The cosmic phase transition with Q-balls production can explain baryogenesis and DM simultaneously, where constraints on DM mass and reverse dilution are significantly relaxed. We study how to probe this scenario by GW signals and collider signals at QCD NLO.

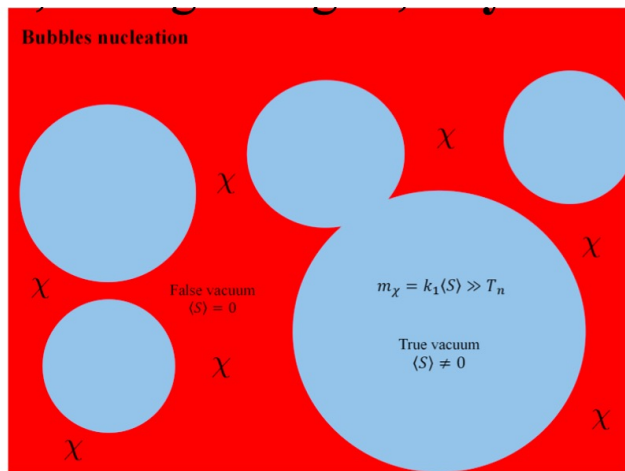
FPH, Chong Sheng Li, *Phys. Rev. D*96 (2017) no.9, 095028

$$\rho_{\text{DM}}^4 v_b^{3/4} = 73.5 (2\eta_B s_0)^3 \lambda_S \sigma^4 \Gamma^{3/4}$$

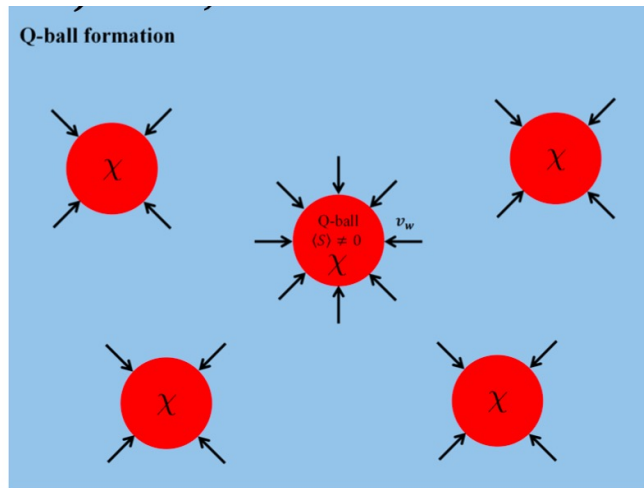


SFOPT and new DM mechanism/signal

SFOPT naturally correlates DM, baryogenesis, particle collider and GW signals.



(a) Bubble nucleation: χ particles trapped in the false vacuum due to Boltzmann suppression



(b) Q-ball formation: After the formation of Q-balls, they should be squeezed by the true vacuum

FPH, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028



SFOPT and new DM mechanism/signal

In the recent two years, this dynamical DM formed by phase transition has become a new idea and attracted more and more attentions.

Namely, bubbles in SFOPT can be the “filters” to packet your needed heavy DM.

FPH, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028

arXiv:1912.04238, Dongjin Chway, Tae Hyun Jung, Chang Sub Shin

arXiv:1912.02830, Michael J. Baker, Joachim Kopp, and Andrew J. Long

arXiv:2012.15113, Wei Chao, Xiu-Fei Li, Lei Wang

arXiv:2101.05721, Aleksandr Azatov, Miguel Vanvlasselaer, Wen Yin

arXiv:2103.09827, Pouya Asadi, Eric D. Kramer, Eric Kuflik, Gregory W. Ridgway, Tracy R. Slatyer, J. Smirnov

arXiv:2103.09822, Pouya Asadi, Eric D. Kramer, Eric Kuflik, Gregory W. Ridgway, Tracy R. Slatyer, J. Smirnov

arXiv:2008.04430 Jeong-Pyong Hong, Sunghoon Jung, Ke-pan Xie

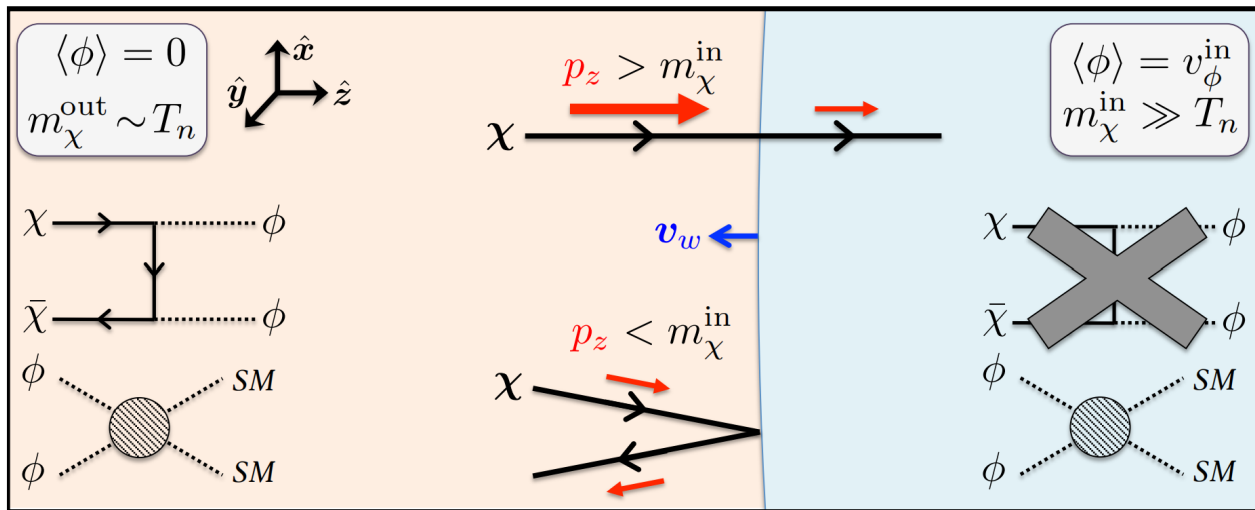
arXiv:2305.02218 Siyu Jiang, **FPH**, Chong Sheng Li



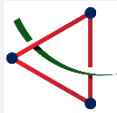


SFOPT and new DM mechanism/signal

Michael J. Baker, Joachim Kopp, Andrew J. Long, Phys.Rev.Lett. 125 (2020) 15, 151102

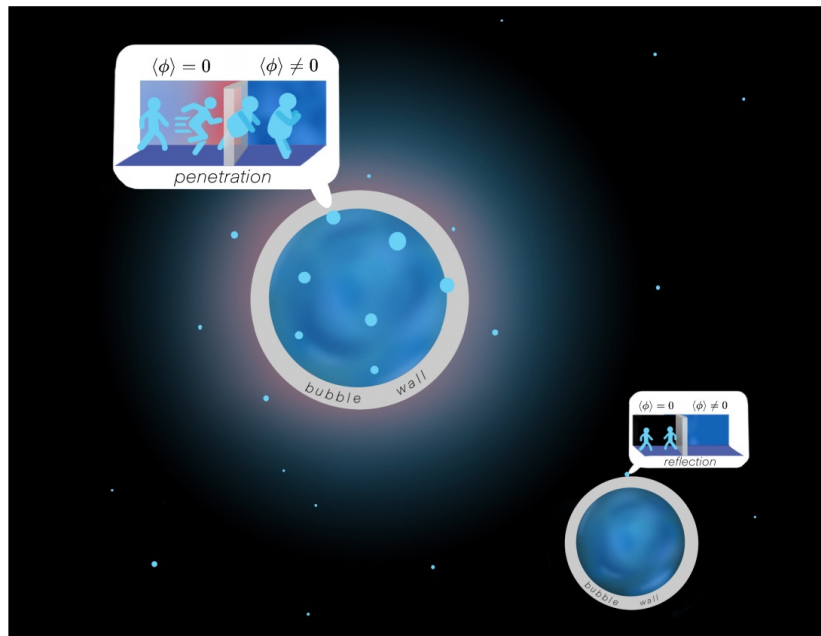


$$\Omega_{\text{DM}} h^2 \approx 0.17 \left(\frac{T_n}{\text{TeV}} \right) \left(\frac{m_\chi^\infty}{30 T_n} \right)^{-\frac{5}{2}} \exp\left(-\frac{m_\chi^\infty}{30 T_n} \right)$$

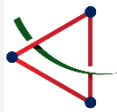


SFOPT and new DM mechanism/signal

Bubble wall dynamics play an essential role in the filtered dark matter mechanism



Siyu Jiang, **FPH**, Chong Sheng Li, arXiv:2305.02218



SFOPT and new DM mechanism/signal

outside the bubble wall $f_{\chi}^{\text{eq}} = \frac{1}{e^{\tilde{\gamma}_{\text{pl}}(\sqrt{(p^w)^2 + m_0^2} - \tilde{v}_{\text{pl}} p_z^w)/T} \mp 1}$,

$$J_{\chi}^w = g_{\chi} \int \frac{d^3 p^w}{(2\pi)^3} \frac{p_z^w}{E^w} f_{\chi}^{\text{eq}} \Theta \left(p_z^w - \sqrt{\Delta m^2} \right),$$

$$n_{\chi}^{\text{in}} = \frac{J_{\chi}^w}{\gamma_w v_w}.$$

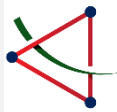
$$n_{\chi}^{\text{in}} \simeq \frac{g_{\chi} T^3}{\gamma_w v_w} \left(\frac{\tilde{\gamma}_{\text{pl}} (1 - \tilde{v}_{\text{pl}}) m_{\chi}^{\text{in}}/T + 1}{4\pi^2 \tilde{\gamma}_{\text{pl}}^3 (1 - \tilde{v}_{\text{pl}})^2} \right) e^{-\frac{\tilde{\gamma}_{\text{pl}}(1 - \tilde{v}_{\text{pl}}) m_{\chi}^{\text{in}}}{T}},$$

original works:

$$\tilde{v}_{\text{pl}} = v_w, \quad T = T' = T_n \quad \rightarrow \quad \Omega_{\text{DM}}^{(0)} h^2 [2].$$

our works:

$$\tilde{v}_{\text{pl}} = \tilde{v}_+, \quad T = T_+, \quad T' = T_- \quad (\text{this work with hydrodynamic effects}).$$



SFOPT and new DM mechanism/signal

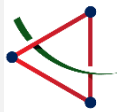
$$\mathcal{L} = -y_\chi \bar{\chi} \Phi \chi - V(\Phi) - \kappa \Phi^\dagger \Phi H^\dagger H + h.c. ,$$

toy model

$$V_{\text{eff}}(\phi, T) = \frac{\mu^2 + DT^2}{2} \phi^2 - CT\phi^3 + \frac{\lambda}{4} \phi^4 - \frac{g_\star \pi^2 T^4}{90} ,$$

$$\alpha \equiv \frac{\left(1 - \frac{T}{4} \frac{\partial}{\partial T}\right) \Delta V_{\text{eff}}}{\pi^2 g_\star T^4 / 30} ,$$

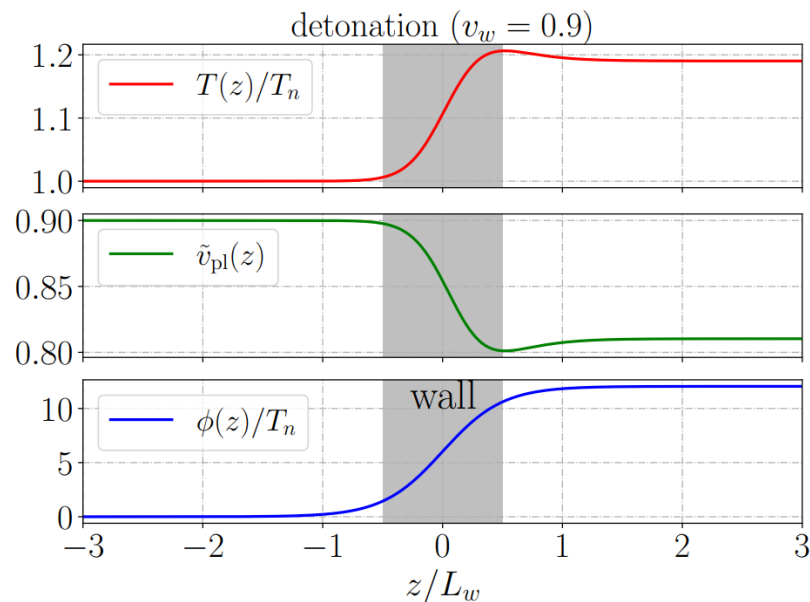
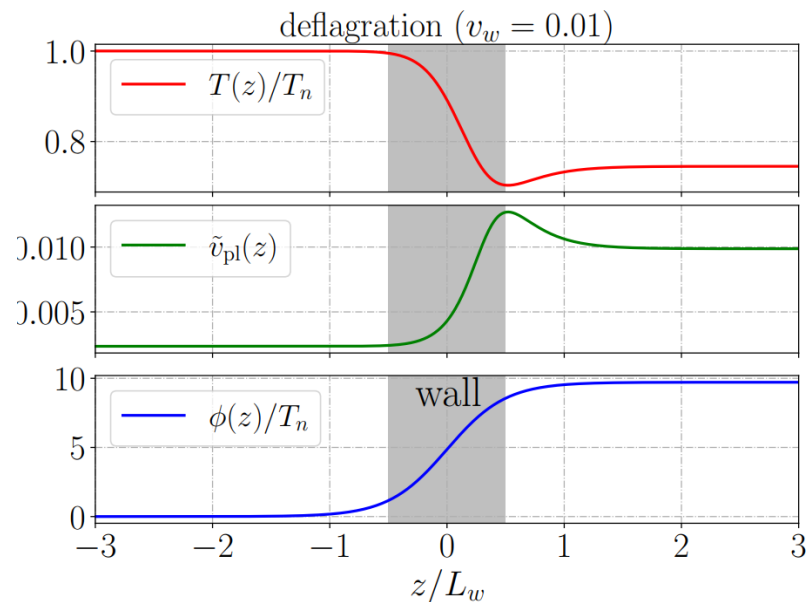
$$\beta = -\frac{d}{dt} \frac{S_3(T)}{T} = H(T) T \frac{d}{dT} \frac{S_3(T)}{T} .$$

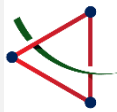


SFOPT and new DM mechanism/signal

bubble wall with non-zero width

$$\phi(z) = \frac{\phi(T_-)}{2} \left(1 + \tanh \frac{2z}{L_w} \right) .$$





SFOPT and new DM mechanism/signal

Boltzmann equation

$$\mathbf{L}[f_\chi] = \mathbf{C}[f_\chi]$$

$$f_\chi = \mathcal{A}(z, p_z) f_{\chi,+}^{\text{eq}} = \mathcal{A}(z, p_z) \exp\left(-\frac{\tilde{\gamma}_+(E - \tilde{v}_+ p_z)}{T_+}\right),$$

$$p_z^2 + m_\chi(z)^2 = \text{const.} \rightarrow \mathbf{L}[f_\chi] = \frac{p_z}{E} \frac{\partial f_\chi}{\partial z} - \frac{m_\chi}{E} \frac{\partial m_\chi}{\partial z} \frac{\partial f_\chi}{\partial p_z}.$$

$$g_\chi \int \frac{dp_x dp_y}{(2\pi)^2} \mathbf{L}[f_\chi] \approx$$

$$\left[\left(\frac{p_z}{m_\chi} \frac{\partial}{\partial z} - \left(\frac{\partial m_\chi}{\partial z} \right) \frac{\partial}{\partial p_z} - \left(\frac{\partial m_\chi}{\partial z} \right) \frac{\tilde{\gamma}_+ + \tilde{v}_+}{T_+} \right) \mathcal{A}(z, p_z) \right] \frac{g_\chi m_\chi T_+}{2\pi \tilde{\gamma}_+} e^{\tilde{\gamma}_+(\tilde{v}_+ p_z - \sqrt{m_\chi^2 + p_z^2})/T_+}.$$

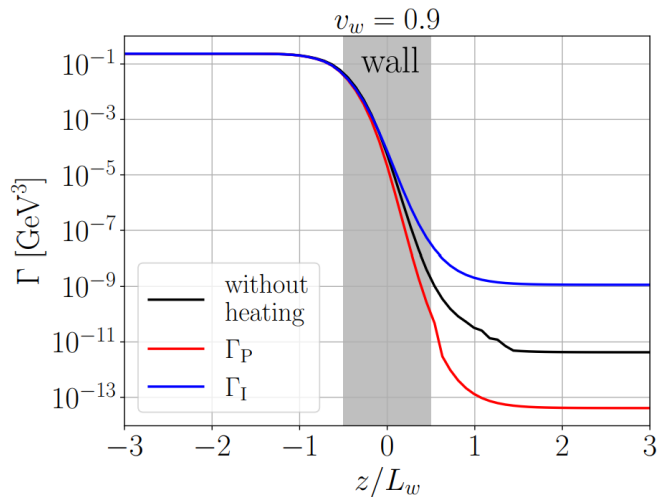
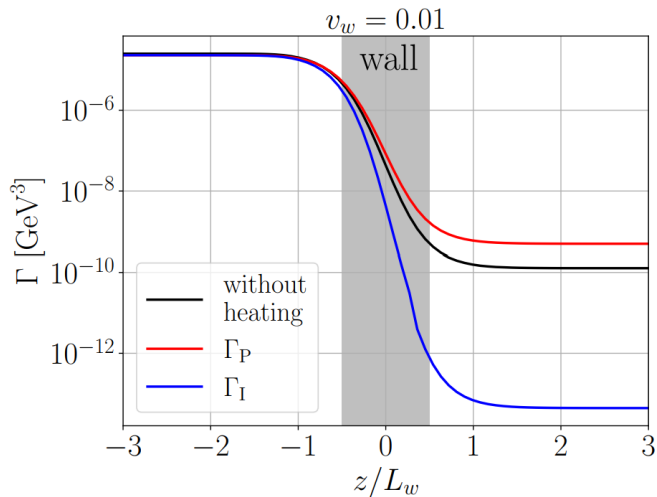
$$\mathbf{C}[f_\chi] = - \sum_{\text{spins}} d\Pi_{q^p} d\Pi_{k^p} d\Pi_{l^p} \frac{(2\pi)^4}{2E_p^p} \delta^{(4)}(p^p + q^p - k^p - l^p) |\mathcal{M}|^2 \mathcal{P}[f_\chi]$$

$$\mathcal{P}[f_\chi] = f_{\chi_p} f_{\bar{\chi}_q} (1 \pm f_{\phi_k}) (1 \pm f_{\phi_l}) - f_{\phi_k} f_{\phi_l} (1 \pm f_{\chi_p}) (1 \pm f_{\bar{\chi}_q})$$



SFOPT and new DM mechanism/signal

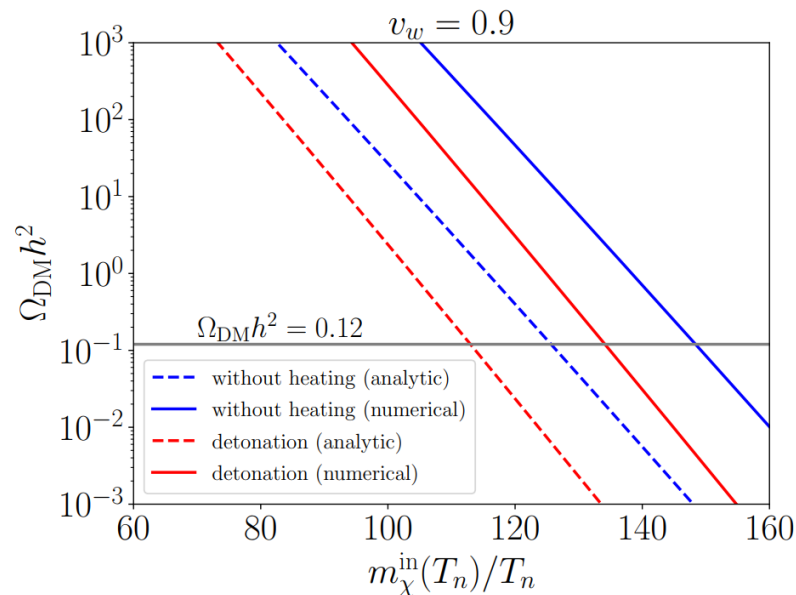
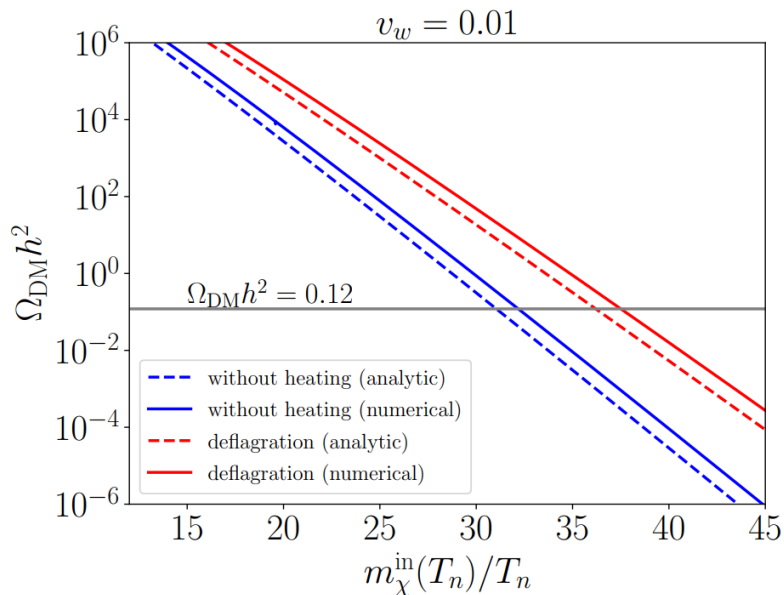
$$\begin{aligned}
 g_\chi \int \frac{dp_x dp_y}{(2\pi)^2} \mathbf{C}[f_\chi] &= -g_\chi g_{\bar{\chi}} \int \frac{dp_x dp_y}{(2\pi)^2 2E_p^{\mathcal{P}}} d\Pi_{q^{\mathcal{P}}} 4F \sigma_{\chi\bar{\chi} \rightarrow \phi\phi} \left[f_{\chi_p} f_{\bar{\chi}_q, +}^{\text{eq}} - f_{\chi_p}^{\text{eq}} f_{\bar{\chi}_q} \right] \\
 &= -g_\chi g_{\bar{\chi}} \int \frac{dp_x dp_y}{(2\pi)^2 2E_p^{\mathcal{P}}} d\Pi_{q^{\mathcal{P}}} 4F \sigma_{\chi\bar{\chi} \rightarrow \phi\phi} \left[\mathcal{A} f_{\chi_p, +}^{\text{eq}} f_{\bar{\chi}_q, +}^{\text{eq}} - f_{\chi_p}^{\text{eq}} f_{\bar{\chi}_q}^{\text{eq}} \right] \\
 &\equiv \Gamma_{\text{P}}(z, p_z) \mathcal{A}(z, p_z) - \Gamma_{\text{I}}(z, p_z) ,
 \end{aligned}$$

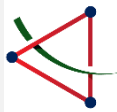




SFOPT and new DM mechanism/signal

$$n_{\chi}^{\text{in}} = \frac{T_+}{\gamma_w \tilde{\gamma}_+} \int_0^{\infty} \frac{dp_z}{(2\pi)^2} \mathcal{A}(z \gg L_w, p_z) \exp \left[\tilde{\gamma}_+ \left(\tilde{v}_+ p_z - \sqrt{p_z^2 + (m_{\chi}^{\text{in}})^2} \right) / T_+ \right] \left(\sqrt{p_z^2 + (m_{\chi}^{\text{in}})^2} + \frac{T_+}{\tilde{\gamma}_+} \right)$$





SFOPT and new DM mechanism/signal

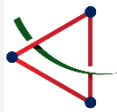
Ratio of dark matter relic density with and without hydrodynamic effects

$$v_w = 0.01$$

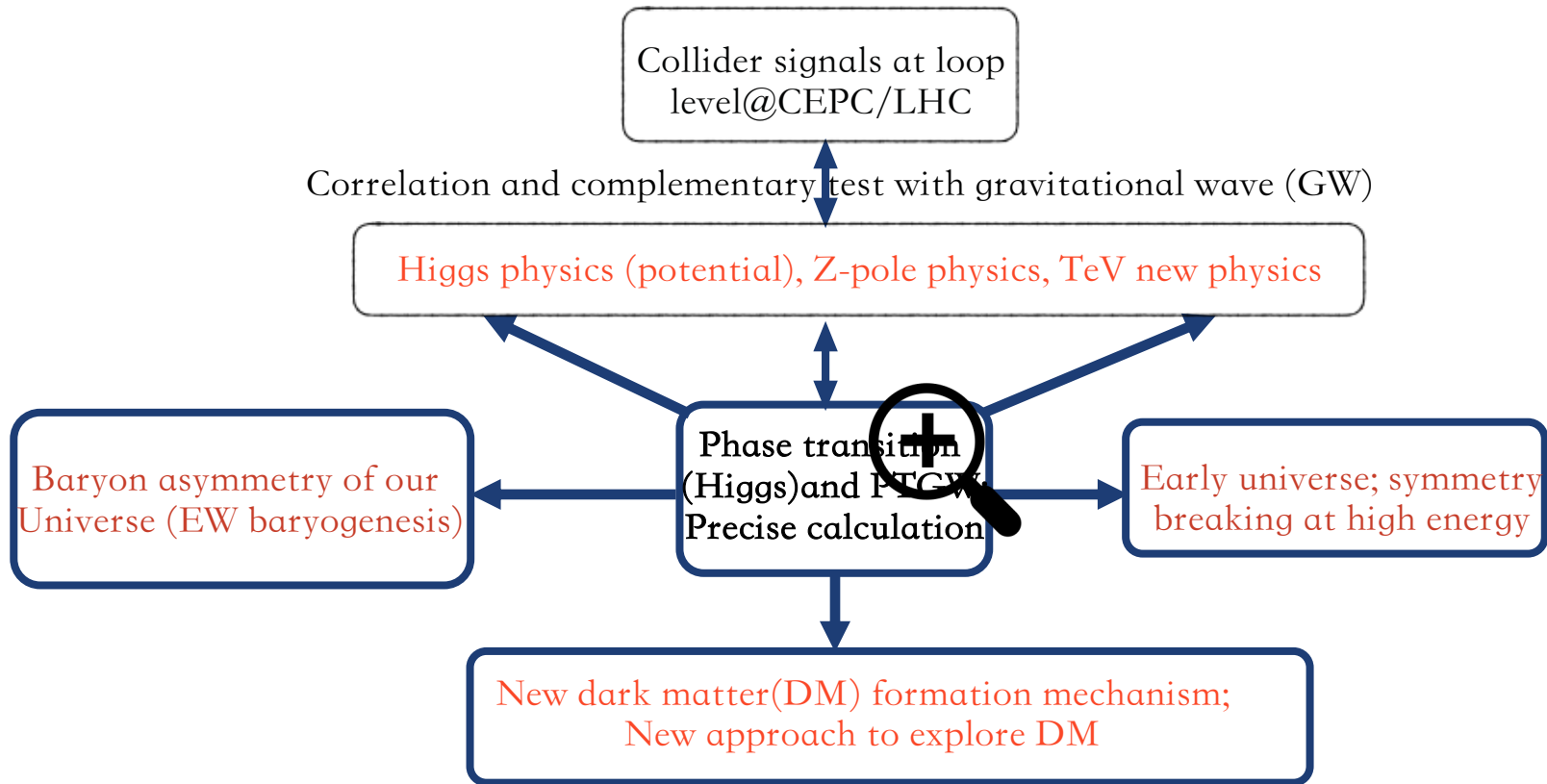
| | analytic | | numerical | |
|--------|-------------------------------|---|-------------------------------|---|
| | $m_\chi^{\text{in}}(T_n)/T_n$ | $\Omega_{\text{DM}}^{(\text{hy})} h^2 / \Omega_{\text{DM}}^{(0)} h^2$ | $m_\chi^{\text{in}}(T_n)/T_n$ | $\Omega_{\text{DM}}^{(\text{hy})} h^2 / \Omega_{\text{DM}}^{(0)} h^2$ |
| BP_1 | 31 | 66 | 32 | 71 |
| BP_2 | 31.1 | 7.9 | 32.2 | 8.1 |
| BP_3 | 30.8 | 778.8 | 31.9 | 858.5 |
| BP_4 | * | * | * | * |

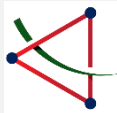
$$v_w = 0.9$$

| | analytic | | numerical | |
|--------|-------------------------------|---|-------------------------------|---|
| | $m_\chi^{\text{in}}(T_n)/T_n$ | $\Omega_{\text{DM}}^{(\text{hy})} h^2 / \Omega_{\text{DM}}^{(0)} h^2$ | $m_\chi^{\text{in}}(T_n)/T_n$ | $\Omega_{\text{DM}}^{(\text{hy})} h^2 / \Omega_{\text{DM}}^{(0)} h^2$ |
| BP_1 | 125.3 | 1/19 | 147.8 | 1/27 |
| BP_2 | 125.9 | 1/7 | 148.7 | 1/9 |
| BP_3 | 124.6 | 1/10 | 147.3 | 1/12 |
| BP_4 | 123.8 | $1/(1.2 \times 10^{13})$ | 146.5 | $1/(2.2 \times 10^{15})$ |



Motivation





Phase transition Dynamics

Precise predictions on the phase transition dynamics and its GW signals

Finite-temperature effective potential

$$V_{eff}(\phi, T)$$

- (1). Daisy resummation problem: Pawani scheme vs. Arnold scheme
- (2). Gauge dependence problem: see Michael J. Ramsey-Musolf's works
- (3). No perturbative calculations: lattice calculations
and dim-reduction method: by D. Weir, Michael J. Ramsey-Musolf et.al

Bubble wall velocity

$$v_b$$

Energy budget

$$\kappa$$

S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang, arXiv:2007.10343,
Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith,
arXiv:2009.14295v2
Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903
Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005

F. Giese, T. Konstandin, K. Schmitz and J. van de Vis, arXiv:2010.09744
Xiao Wang, **FPH** and Xinmin Zhang, Phys.Rev.D 103 (2021) 10, 103520
Xiao Wang, Chi Tian, **FPH**, arXiv: 2301.12328

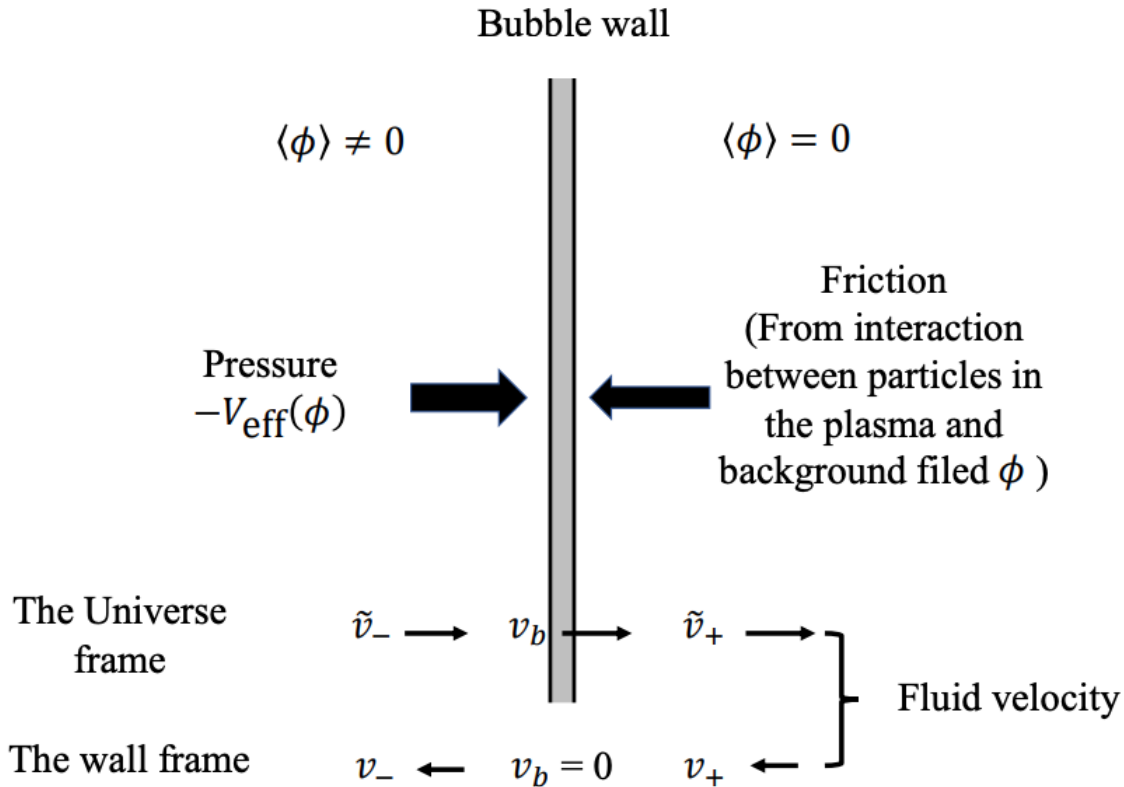


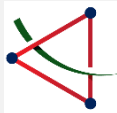
Phase transition Dynamics

bubble wall velocity:

most difficult; most important for EW baryogenesis and GW experiments

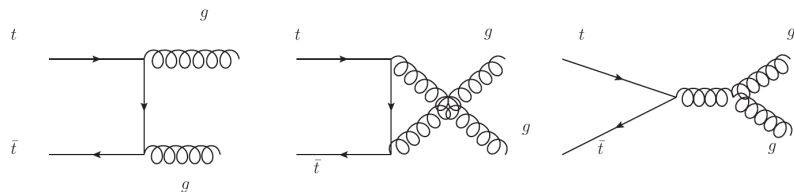
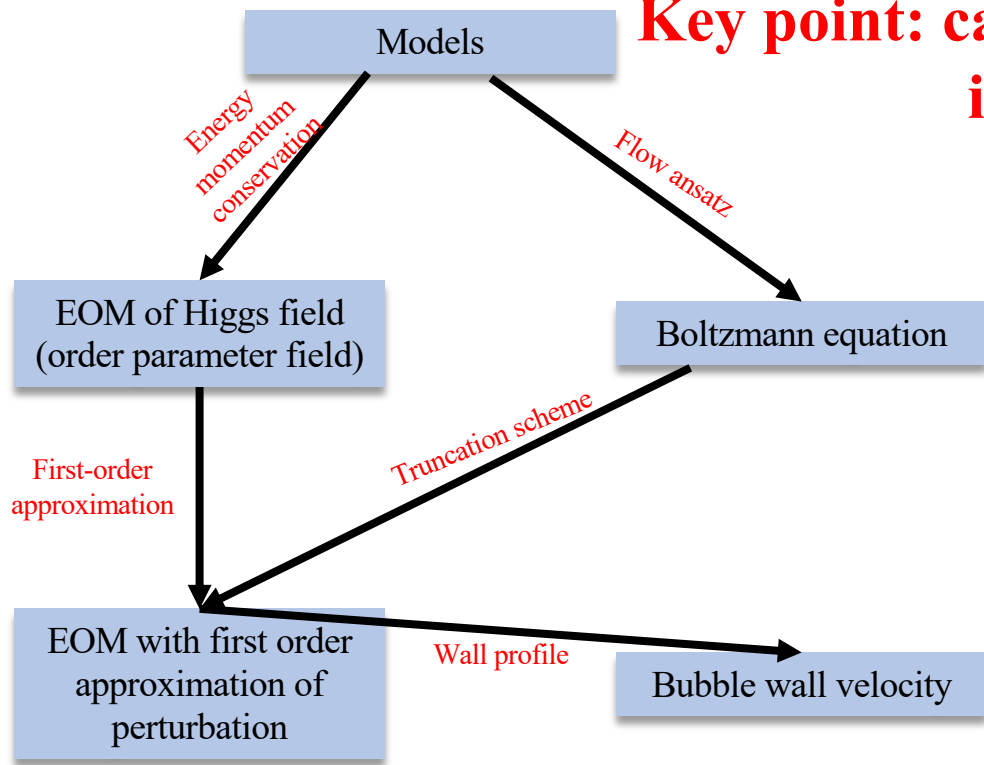
See James cline and Hindmarsh's recent work;
lots of unsolved problems, eg. resummation





bubble wall velocity in SM EFT

Key point: calculate the particle scattering in thermal plasma

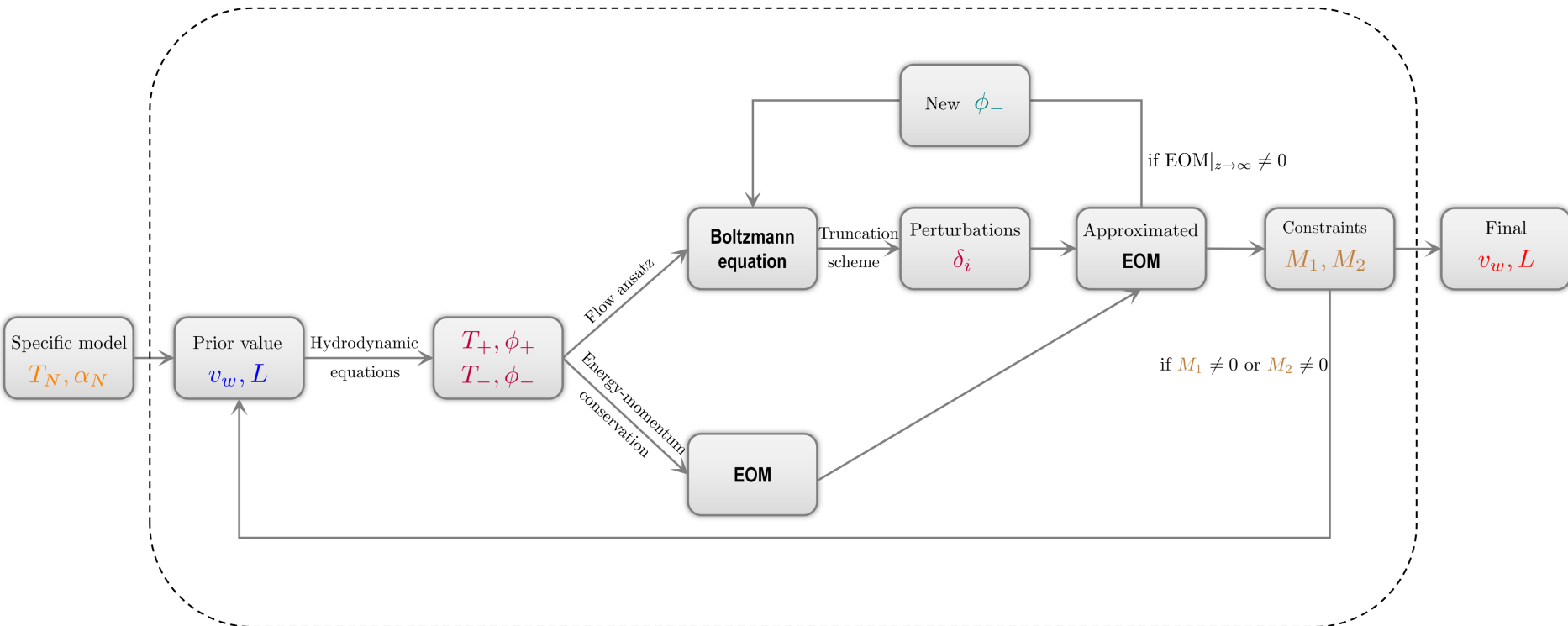


- For a specific benchmark point, the results of beyond leading-log approximation with the contribution of Higgs boson

| | $\Lambda/\sqrt{\kappa}$ [GeV] | v_w | $L_w T$ | v_c/T_c | T_n [GeV] |
|----------------|-------------------------------|--------|---------|-----------|-------------|
| Two-particle | 780 | 0.3382 | 20.1863 | 1.1044 | 100.977 |
| Three-particle | 780 | 0.2499 | 18.1759 | 1.1044 | 100.977 |



bubble wall velocity in inert 2HDM



Siyu Jiang, *FPH*, xiao wang, *Phys.Rev.D* 107 (2023) 9, 095005



Phase transition Dynamics

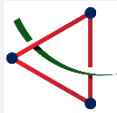
$$V_0 = \mu_1^2 |\Phi|^2 + \mu_2^2 |\eta|^2 + \frac{1}{2} \lambda_1 |\Phi|^4 + \frac{1}{2} \lambda_2 |\eta|^4 \\ + \lambda_3 |\Phi|^2 |\eta|^2 + \lambda_4 |\Phi^\dagger \eta|^2 + \frac{1}{2} \{ \lambda_5 (\Phi^\dagger \eta)^2 + \text{H.c.} \} ,$$

$$\Phi = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(h + v + iG^0) \end{pmatrix}, \quad \eta = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}}(H + iA) \end{pmatrix},$$

$$V_{\text{CW}}(\phi, T = 0) = \sum_i \frac{n_i}{64\pi^2} \left[m_i^4(\phi) \left(\ln \frac{m_i^2(\phi)}{\bar{m}_i^2} - \frac{3}{2} \right) + 2\bar{m}_i^2 \bar{m}_i^2(\phi) \right],$$

$$V_{\text{T}}(\phi, T > 0) = \sum n_i \frac{T^4}{2\pi^2} I_b \left(\frac{M_i^2}{T^2} \right),$$

$$V_{\text{eff}}(\phi, T) = V_0(\phi) + V_{\text{CW}}(\phi) + V_{\text{T}}(\phi, T) .$$

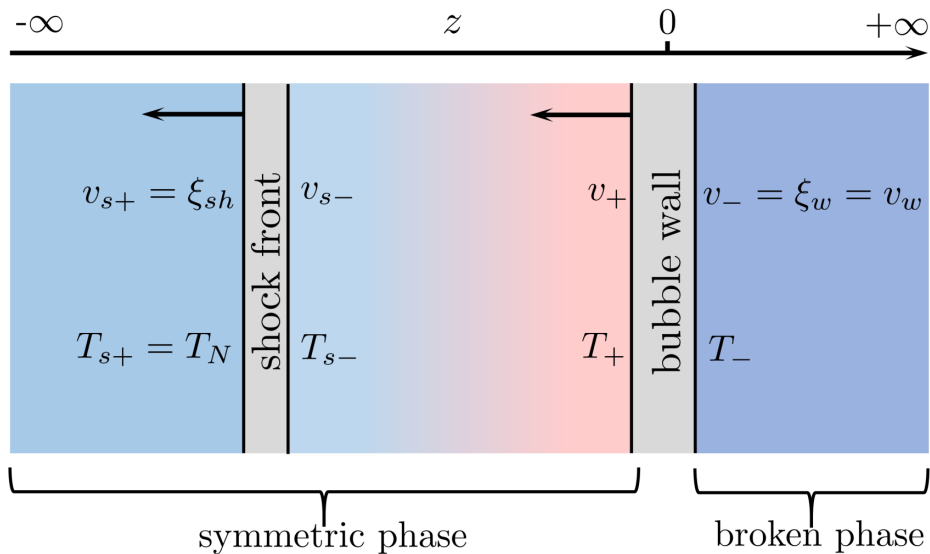


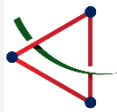
Phase transition Dynamics

Hydrodynamic equations

$$(\xi - v) \frac{\partial_\xi e_f}{\omega_f} = 2 \frac{v}{\xi} + [1 - \gamma^2 v(\xi - v)] \partial_\xi v,$$

$$(1 - v\xi) \frac{\partial_\xi p_f}{\omega_f} = \gamma^2 (\xi - v) \partial_\xi v,$$





Phase transition Dynamics

Evolution of the perturbations

$$\hat{A}\delta' + \boxed{\Gamma}\delta = \boxed{\Sigma}, \quad \text{source term}$$

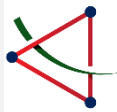
collision term

$$\delta = (\mu_t, \delta T_t, T\delta v_t, \mu_W, \delta T_W, T\delta v_W, \boxed{\mu_A, \delta T_A, T\delta v_A}),$$

New particles introduced by inert-doublet model

$$\Sigma = \frac{v_w}{2T} (c_1^t (m_t^2)', c_2^t (m_t^2)', 0, c_1^W (m_W^2)', c_2^W (m_W^2)', 0, c_1^A (m_A^2)', c_2^A (m_A^2)', 0),$$

$$\hat{A} = \begin{pmatrix} \hat{A}_t & 0 & 0 \\ 0 & \hat{A}_W & 0 \\ 0 & 0 & \hat{A}_A \end{pmatrix}, \quad \text{where} \quad \hat{A}_i = \begin{pmatrix} v_w c_2^i & v_w c_3^i & \frac{1}{3} c_3^i \\ v_w c_3^i & v_w c_4^i & \frac{1}{3} c_4^i \\ \frac{1}{3} c_3^i & \frac{1}{3} c_4^i & \frac{1}{3} v_w c_4^i \end{pmatrix},$$

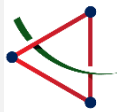


Phase transition Dynamics

Perturbations for background

$$\tilde{c}_4 \left(v_w \delta T'_{bg} + \frac{\delta v'_{bg}}{3} T \right) = N_t (\mu_t \Gamma_{\mu 2,t} + \delta T_t \Gamma_{T 2,t}) + \sum_{\text{bosons}} N_b (\mu_b \Gamma_{\mu 2,b} + \delta T_b \Gamma_{T 2,b}) ,$$

$$\frac{\tilde{c}_4}{3} (\delta T'_{bg} + v_w T \delta v'_{bg}) = N_t T \delta v_t \Gamma_{v,t} + \sum_{\text{bosons}} N_b T \delta v_b \Gamma_{v,b} , \quad \mu_{bg} = 0 ,$$



Phase transition Dynamics

Collision terms

$$\Gamma_{\mu 1,t} \simeq (5.0 \times 10^{-4} g_s^4 + 5.8 \times 10^{-4} g_s^2 y_t^2) T ,$$

$$\Gamma_{T1,t} \simeq \Gamma_{\mu 2,t} \simeq (1.1 \times 10^{-3} g_s^4 + 1.3 \times 10^{-3} g_s^2 y_t^2) T ,$$

$$\Gamma_{T2,t} \simeq (1.1 \times 10^{-2} g_s^4 + 4.0 \times 10^{-3} g_s^2 y_t^2) T ,$$

$$\Gamma_{v,t} \simeq (2.0 \times 10^{-2} g_s^4 + 1.8 \times 10^{-3} g_s^2 y_t^2) T ,$$

$$\Gamma_{\mu 1,W} \simeq (2.3 \times 10^{-3} g_s^2 g_w^2 + 2.0 \times 10^{-3} g_w^4) T ,$$

$$\Gamma_{T1,W} \simeq \Gamma_{\mu 2,W} \simeq (4.7 \times 10^{-3} g_s^2 g_w^2 + 4.1 \times 10^{-3} g_w^4) T$$

$$\Gamma_{T2,W} \simeq (1.5 \times 10^{-2} g_s^2 g_w^2 + 1.5 \times 10^{-2} g_w^4) T ,$$

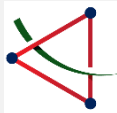
$$\Gamma_{v,W} \simeq (5.7 \times 10^{-2} g_s^2 g_w^2 + 1.5 \times 10^{-2} g_w^4) T ,$$

$$\Gamma_{\mu 1,A} \simeq 1.0 \times 10^{-2} \lambda_3^4 T ,$$

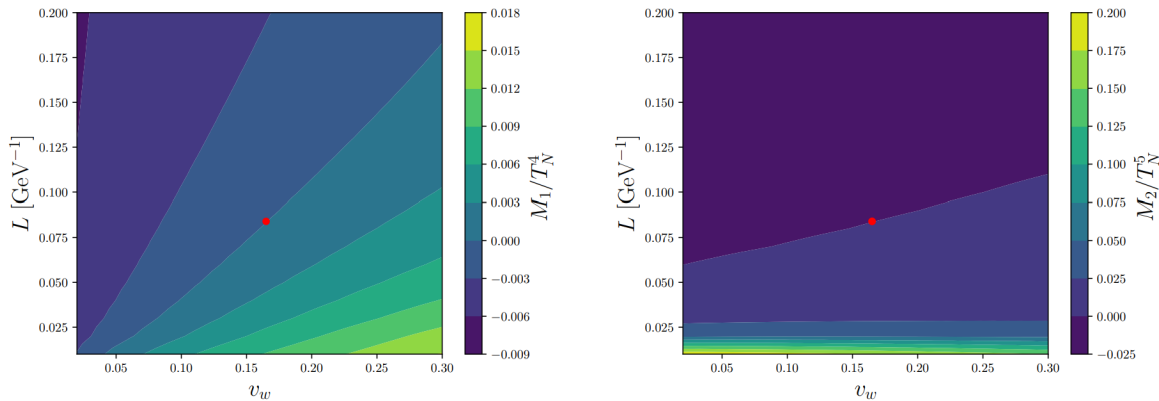
$$\Gamma_{T1,A} \simeq \Gamma_{\mu 2,A} \simeq 4.9 \times 10^{-3} \lambda_3^4 T ,$$

$$\Gamma_{T2,A} \simeq 5.1 \times 10^{-3} \lambda_3^4 T ,$$

$$\Gamma_{v,A} \simeq 1.8 \times 10^{-3} \lambda_3^4 T .$$



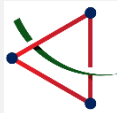
Phase transition Dynamics



| | T_c [GeV] | T_N [GeV] | v_w | L [GeV $^{-1}$] |
|-------------|-------------|-------------|-------|--------------------|
| Benchmark A | 118.3 | 117.1 | 0.165 | 0.084 |
| Benchmark B | 118.6 | 117.5 | 0.164 | 0.085 |
| Benchmark C | 119.4 | 118.4 | 0.164 | 0.088 |

In the allowed parameter spaces of IDM, the bubble wall velocity varies slightly around 0.165.

Siyu Jiang, FPH, Xiao Wang, Phys.Rev.D 107 (2023) no.9, 095005



Phase transition Dynamics

Precise predictions on the phase transition dynamics and its GW signals

Finite-temperature effective potential

$$V_{eff}(\phi, T)$$

- (1). Daisy resummation problem: Pawani scheme vs. Arnold scheme
- (2). Gauge dependence problem: see Michael J. Ramsey-Musolf's works
- (3). No perturbative calculations: lattice calculations
and dim-reduction method: by D. Weir, Michael J. Ramsey-Musolf et.al

Bubble wall velocity

$$v_b$$

Energy budget

$$\kappa$$

S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang, arXiv:2007.10343,
 Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith,
 arXiv:2009.14295v2
 Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903
 Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005

F. Giese, T. Konstandin, K. Schmitz and J. van de Vis, arXiv:2010.09744
 Xiao Wang, **FPH** and Xinmin Zhang, Phys.Rev.D 103 (2021) 10, 103520
 Xiao Wang, Chi Tian, **FPH**, arXiv: 2301.12328

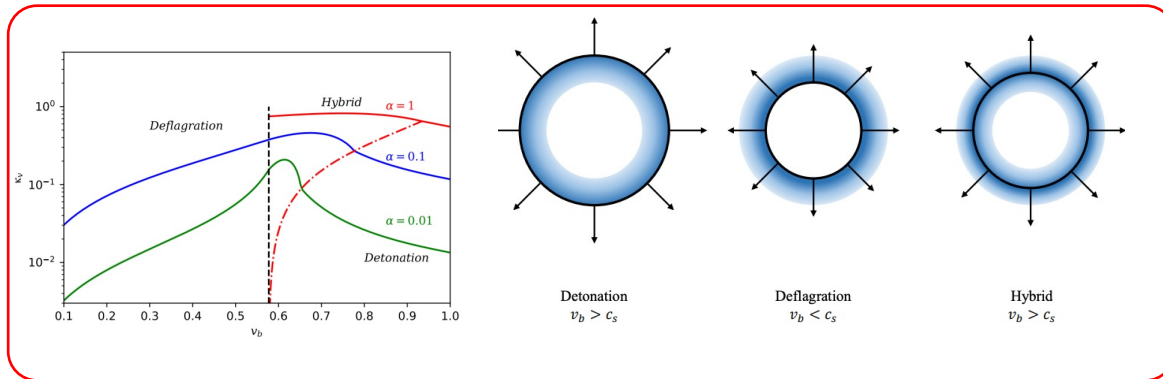


Energy budget

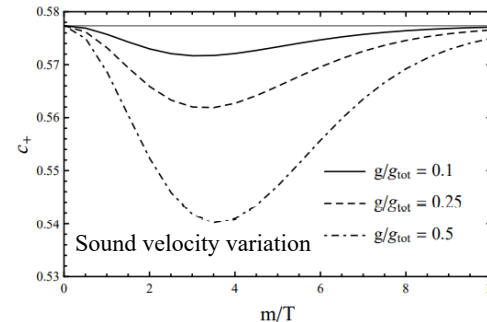
Energy budget (to measure the efficiency of the energy released by a SFOPT converting to the kinetic energy of sounding plasma)

LISA Cosmology working group *JCAP* 03 (2020) 024

$$h^2 \Omega_{\text{sw}}(f) \simeq 1.64 \times 10^{-6} (H_* \tau_{\text{sw}}) (H_* R_*) \left(\frac{\kappa_V \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{q_*} \right)^{1/3} (f/f_{\text{sw}})^3 \left(\frac{7}{4 + 3(f/f_{\text{sw}})^2} \right)^{7/2}$$



J. R. Espinosa, T. Konstandin, J. M. No, and G. Servant, *JCAP* 06 (2010) 028.



L. Leitao and A. Megevand, *Nucl. Phys. B* **891**, 159-199 (2015)

But for a realistic SFOPT, particle can obtain the mass, hence, the sound velocity can deviate from pure radiation phase. The energy budget/efficiency parameter beyond the bag model are study in our recent work:

Xiao Wang, **FPH** and Xinmin Zhang, *PRD* 103 (2021) 10, 103520

Xiao Wang, **FPH** and Xinmin Zhang, *JCAP* 05, 045 (2020)

Xiao Wang, Chi Tian, **FPH**, [arXiv: 2301.12328](https://arxiv.org/abs/2301.12328)



Phase transition Dynamics

Energy budget for phase transition GW

- Matching condition

$$\begin{aligned} w_- v_-^2 \gamma_-^2 + p_- &= w_+ v_+^2 \gamma_+^2 + p_+, \\ w_- v_- \gamma_-^2 &= w_+ v_+ \gamma_+^2. \end{aligned} \quad \longrightarrow \quad v_+ v_- = \frac{p_+ - p_-}{e_+ - e_-}, \quad v_+ = \frac{e_- + p_+}{e_+ + p_-}.$$

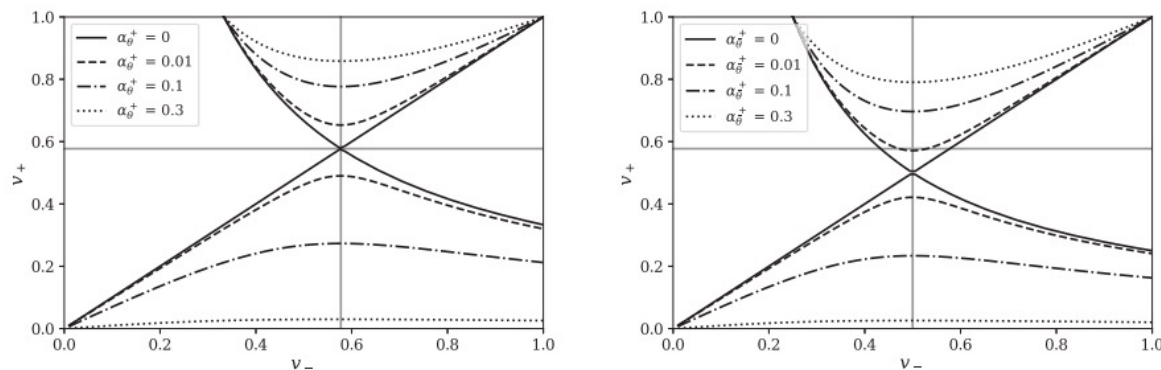


FIG. 1. The fluid velocities v_+ and v_- in the reference frame of bubble wall for different definitions and values of phase transition strength parameter. The horizontal and vertical gray lines indicate the sound velocities of symmetric and broken phase. Left panel: the bag model. Right panel: the DSVM with $c_+^2 = 1/3$ and $c_-^2 = 0.25$.



Phase transition Dynamics

Energy budget for phase transition GW

GW spectrum and SNR for different EoS with different parameter combination:

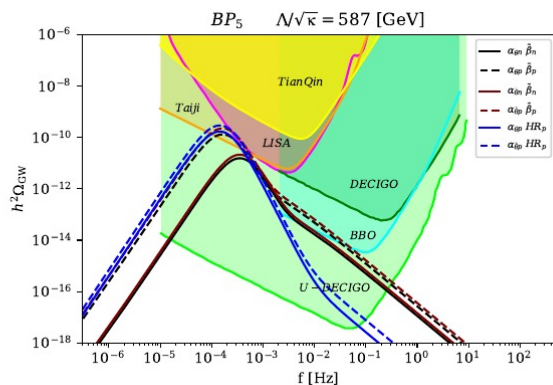


TABLE II. The SNR of BP_5 for different experiment configurations with different combinations of phase transition parameters and models of EOS.

| | $\alpha_{0n} \tilde{\beta}_n$ | $\alpha_{0p} \tilde{\beta}_p$ | $\alpha_{0n} \tilde{\beta}_n$ | $\alpha_{0p} \tilde{\beta}_p$ | $\alpha_{0p} HR_p$ | $\alpha_{0n} HR_p$ |
|--------------------------|-------------------------------|-------------------------------|-------------------------------|-------------------------------|--------------------|--------------------|
| SNR _(LISA) | 7.949 | 16.930 | 10.913 | 28.836 | 16.009 | 27.468 |
| SNR _(Taiji) | 14.760 | 58.607 | 20.271 | 100.343 | 66.216 | 113.609 |
| SNR _(TianQin) | 0.452 | 1.506 | 0.620 | 2.576 | 1.629 | 2.794 |

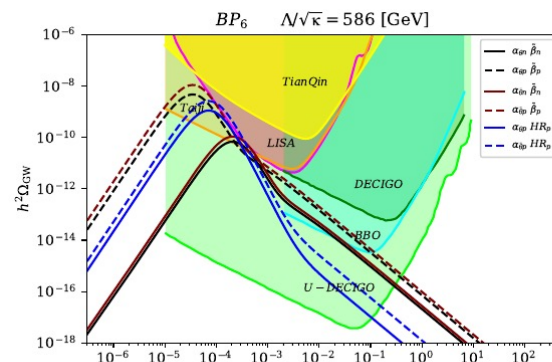
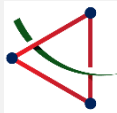


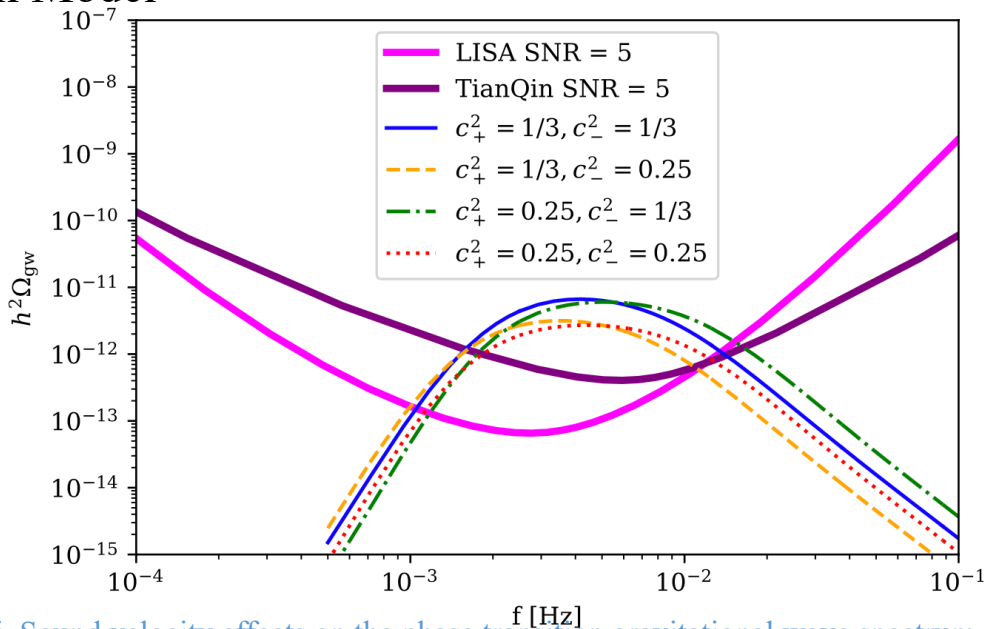
TABLE III. The SNR of BP_6 for different experiment configurations with different combinations of phase transition parameters and models of EOS.

| | $\alpha_{0n} \tilde{\beta}_n$ | $\alpha_{0p} \tilde{\beta}_p$ | $\alpha_{0n} \tilde{\beta}_n$ | $\alpha_{0p} \tilde{\beta}_p$ | $\alpha_{0p} HR_p$ | $\alpha_{0n} HR_p$ |
|--------------------------|-------------------------------|-------------------------------|-------------------------------|-------------------------------|--------------------|--------------------|
| SNR _(LISA) | 14.230 | 15.368 | 22.470 | 26.382 | 17.367 | 40.816 |
| SNR _(Taiji) | 38.666 | 427.813 | 61.208 | 1000.501 | 213.123 | 500.668 |
| SNR _(TianQin) | 1.060 | 5.569 | 1.678 | 12.934 | 3.973 | 9.333 |



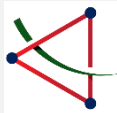
Phase transition Dynamics

Sound velocity effects on the phase transition gravitational wave spectrum in the Sound Shell Model

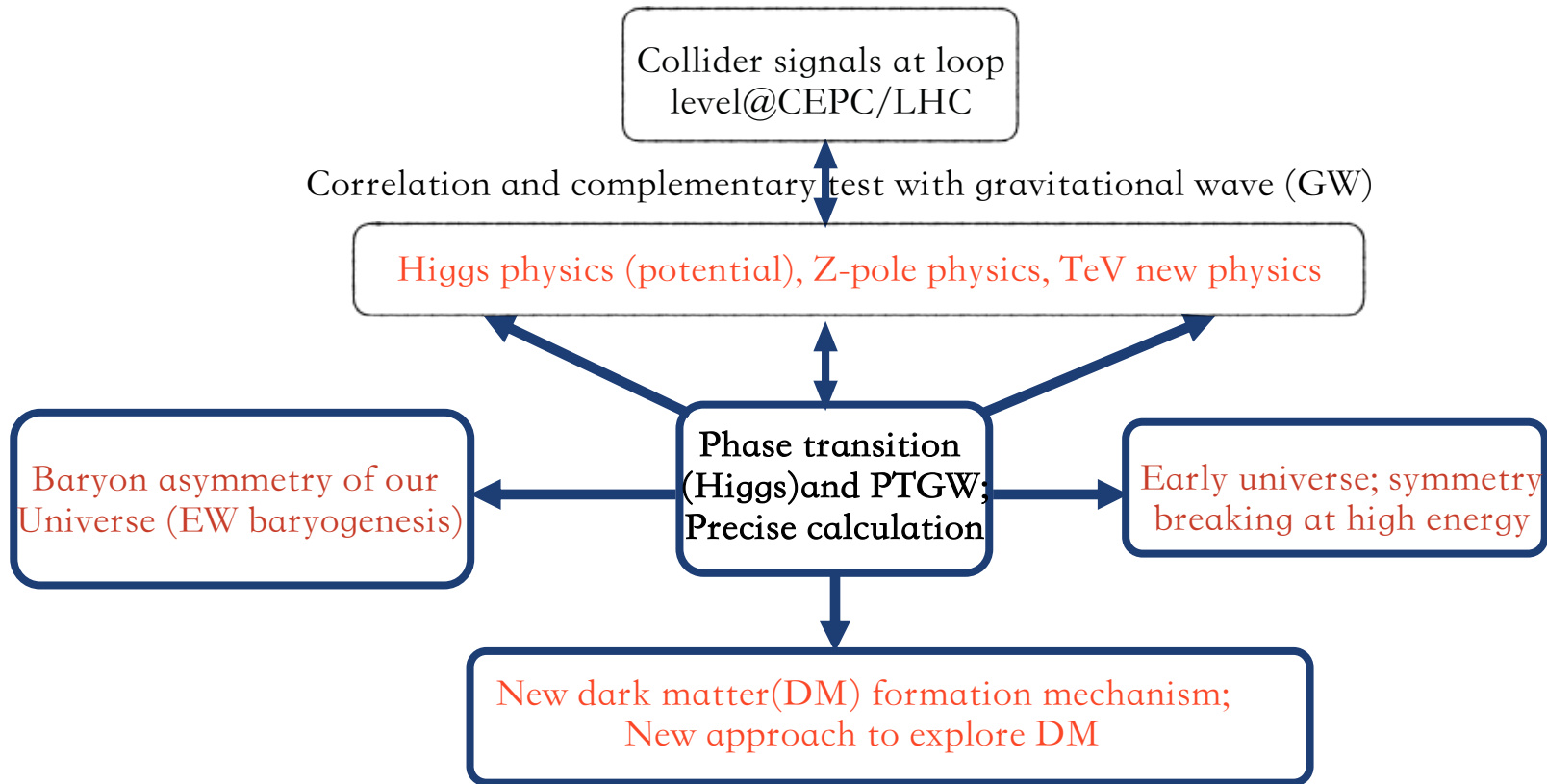


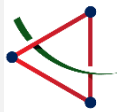
Xiao Wang, **FPH**, Yongping Li, Sound velocity effects on the phase transition gravitational wave spectrum in the Sound Shell Model, arXiv:2112.14650

Xiao Wang, **FPH**, Xinmin Zhang, Energy budget and the gravitational wave spectra beyond the bag model, Phys.Rev.D 103 (2021) 10, 103520



Summary





Summary and Outlook

- **EW first-order phase transition has abundant collider and cosmological effects in baryogenesis, dark matter, GW...**
- **The correlation between GW and collider signals at CEPC can make complementary test on the Higgs nature, baryogenesis, dark matter and the cosmic evolution history at 100 GeV.**
- **More precise study are needed: resummation, non-perturbation, bubble dynamics (wall velocity, energy budget).**

Thanks! **Comments and collaborations are welcome!**

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