

Spotlight talks

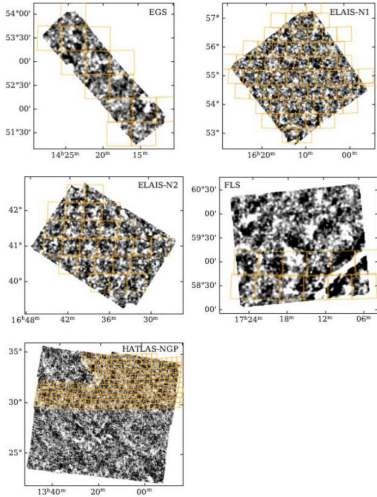
4:30-5 pm Monday 3 June

Images (even with all sources removed) contain LSS information!

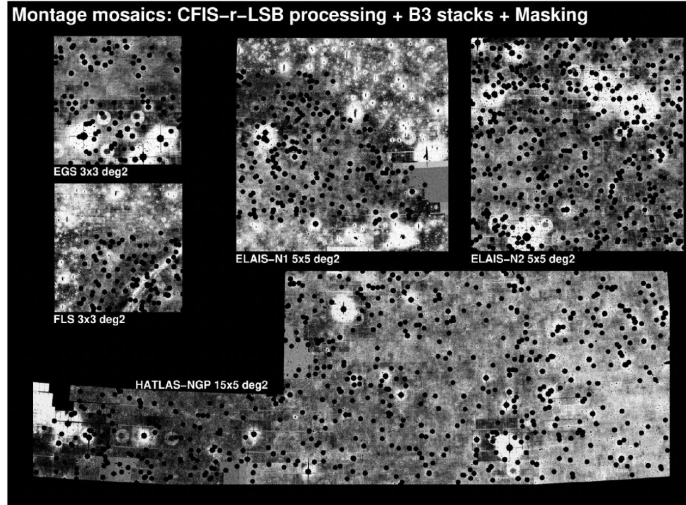
Douglas Scott (UBC) – see Lim et al. 2022, arXiv:2203.16545

Pilot study: Herschel-SPIRE X CFHT

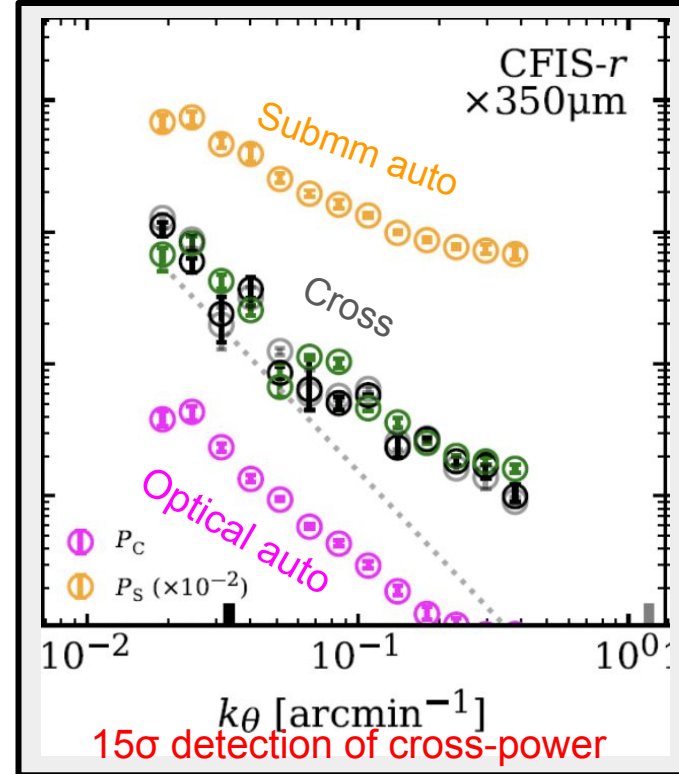
Overlap of **5 fields, $\sim 90\text{sq.deg.}$** , are used for CIB-optical cross-correlation



FIR/submm



Optical (*r* band)



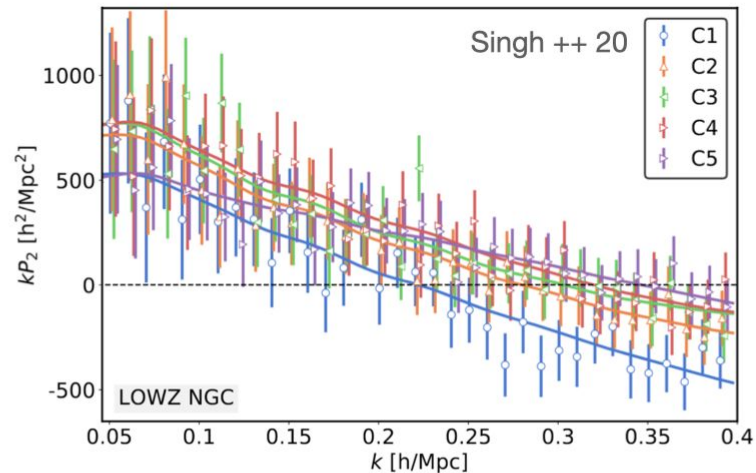
Will be dramatically better using Euclid images correlated with CCAT etc.



Can we select galaxy samples with fewer Fingers of God?

with U. Seljak

- FoG are sourced by a small fraction of galaxies with large virialized velocities
- They are **non-perturbative** on rather large scales
- The **zero-crossing of the quadrupole** is a marker of the prevalence of FoG — fewer FoG bring it to higher k
- We can define **density-independent selections** that push the zero-crossing to higher k !

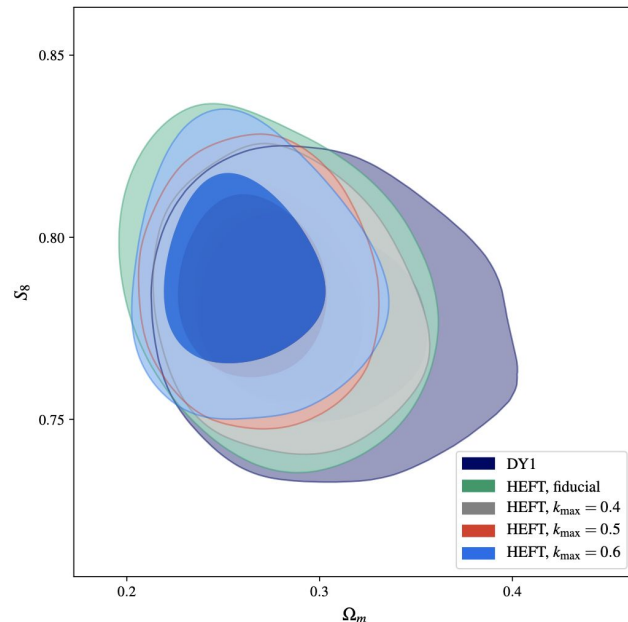


For example, one could discard the 25% reddest galaxies of BOSS LOWZ

Jahmour Givans (Flatiron Institute and Princeton)

How should we model HSC clustering on small scales?

- HSC is a testbed for future analyses with Rubin
- Access to increased information by moving beyond linear regime
- Exploring methods for combining EFT frameworks with insights from simulations
 - See Tuesday talks for more discussion
- Previously applied to data from DES Y1 by Hadzhiyska+ 2021
 - Improved constraints on (S_8, Ω_m) by (10%,35%), see figure



Finn Roper (PhD student at the IfA)

Supervised by Yan-Chuan Cai and John Peacock



Directly detecting the kinematic

Sunyaev–Zel’dovich effect in the CMB.

$$\frac{\delta T_{\text{kSZ}}(\hat{\mathbf{r}})}{T_0} = - \int n_e \sigma_T \left(\frac{\mathbf{v}}{c} \cdot \hat{\mathbf{r}} \right) dl$$

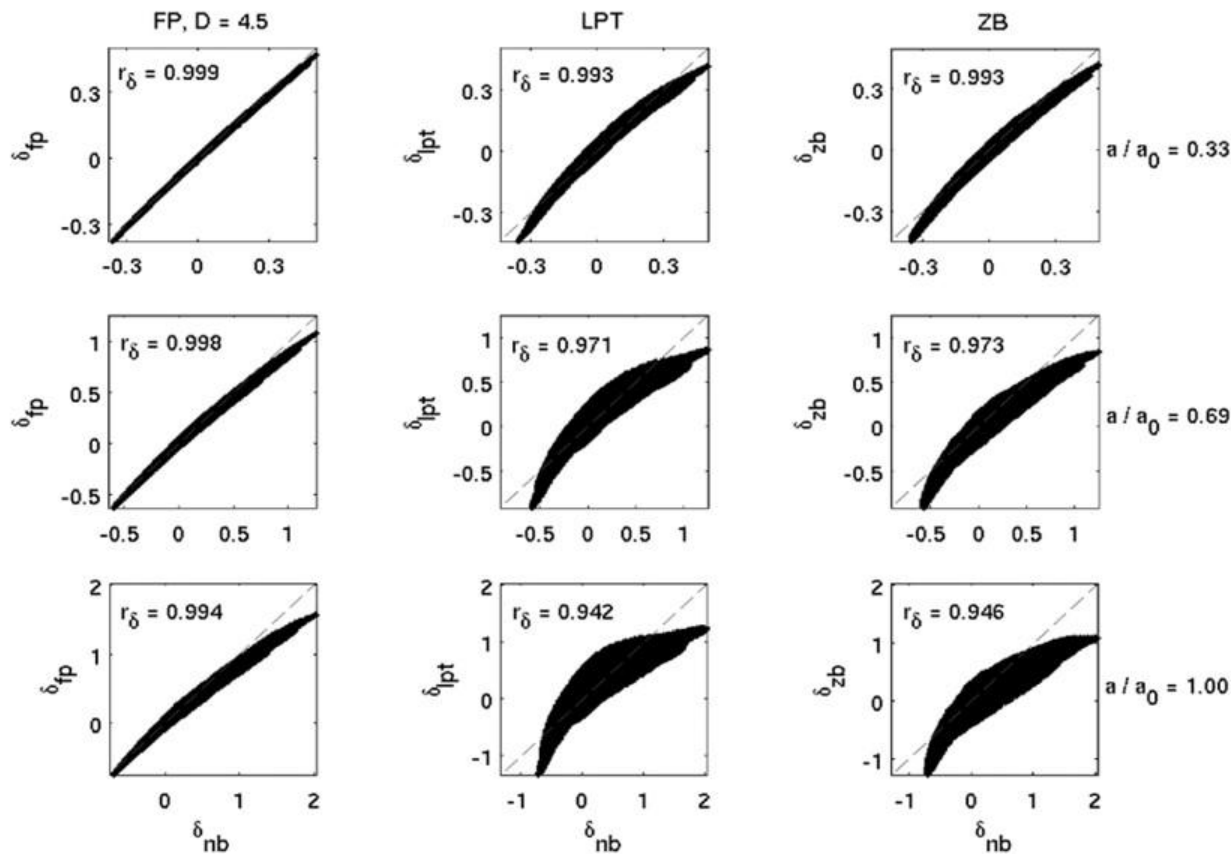
- **kSZ is a tiny effect.** Each galaxy group contributes $|\delta T| \approx 1 \mu\text{K}$. Overall, this accounts for $\sim 10^{-5}$ of the observed CMB power.
- Given group masses and positions, we can **reconstruct the density field and then the peculiar velocity field** in the linear regime.
- With this knowledge, we can **stack this signal** over millions of groups to obtain a significant direct detection.
- This allows us to, for example, obtain the average gas mass fraction in haloes.

Reconstructing with wave mechanics

Aoibhinn Gallagher supervised by Peter Coles



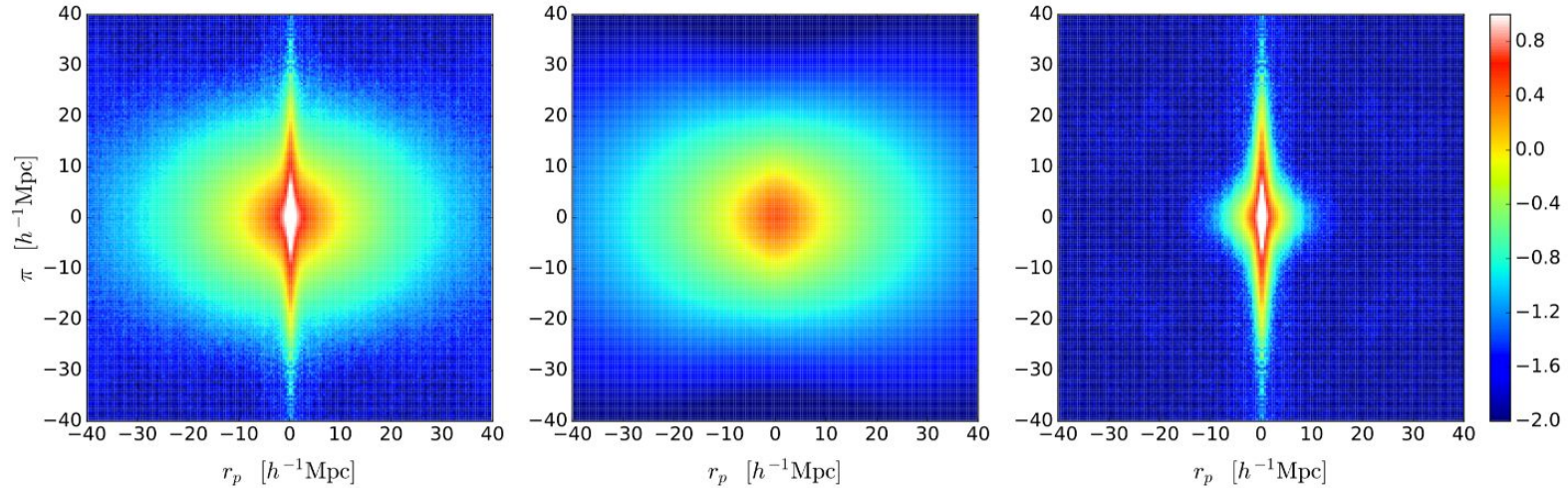
Maynooth University
National University of Ireland Maynooth



Point by point comparison

$$r_\delta = \frac{\langle \delta_{nb} \delta \rangle}{\langle \delta_{nb}^2 \rangle^{\frac{1}{2}} \langle \delta^2 \rangle^{\frac{1}{2}}}$$

John Peacock: other things



Halo streaming: use simulations to extract 1-halo template (2206.05065)

CMB lensing tomography: low Ω_m as a solution to S_8 tension (2010.00466)

Anthropic galaxy formation: simulate SFR(t) for large Λ (2209.08783)

Tidal anisotropy assembly bias: detected in GAMA survey (2305.01266)



Tripolar Spherical Harmonic Bispectrum

MikeSWang ✉ mikeshengbo.wang@ed.ac.uk



arXiv 2304.03643 JOSS 10.21105/joss.05571

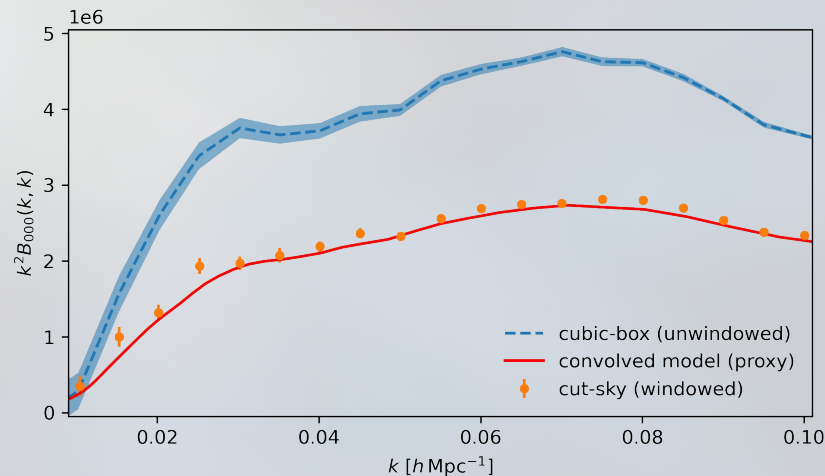
Tripolar spherical-harmonic decomposition (Sugiyama+ 2019)

$$B(k_1, k_2, \hat{n}) = \sum_{\ell_1 \ell_2 L} H_{\ell_1 \ell_2 L}^{-1} B_{\ell_1 \ell_2 L}(k_1, k_2) \sum_{m_1 m_2 M} \begin{pmatrix} \ell_1 & \ell_2 & L \\ m_1 & m_2 & M \end{pmatrix} y_{\ell_1}^{m_1}(\hat{k}_1) y_{\ell_2}^{m_2}(\hat{k}_2) y_L^M(\hat{n})$$

$$B_{\ell_1 \ell_2 L}(k_1, k_2) \propto \sum_M \int d \cos \theta y_{\ell_2}^M \cdots B_{LM}(k_1, k_2, \theta)$$

Comparison with Scoccimarro multipoles:

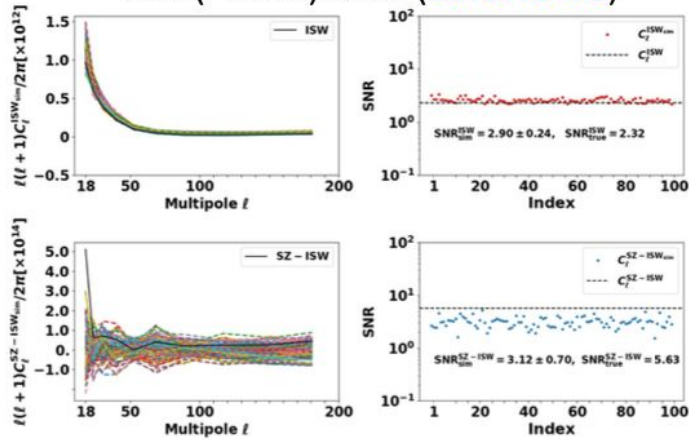
- Lower data vector dimensions (2D vs 3D binning);
- No explicit triangles (handled through an inverse FFT);
- Window convolution analogous to power spectrum.



Yin-Zhe Ma: Other observations

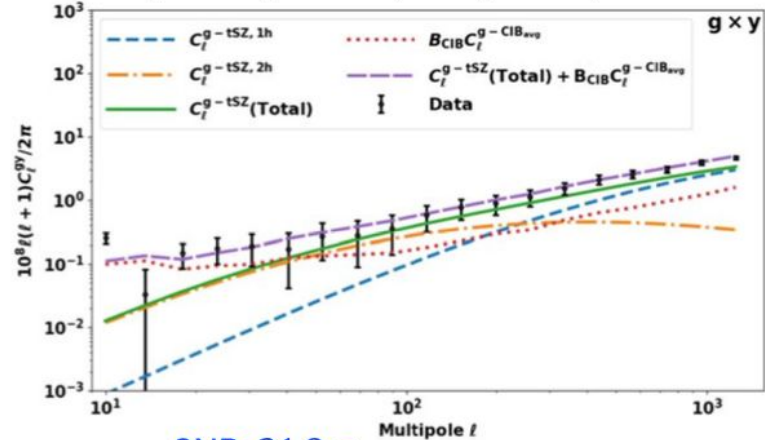
Head of Astrophysics Group, Stellenbosch University (South Africa)

ISW (Planck) \times tSZ (2310.18478)



SNR: $(5.63 - 3.12)/0.7 \approx 3.6$

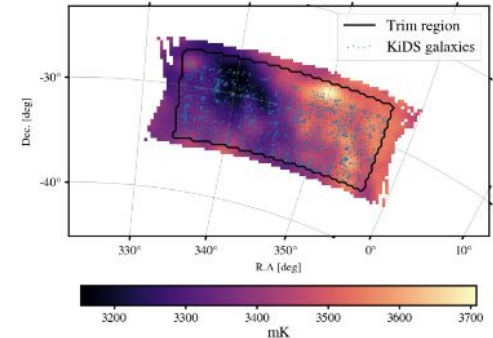
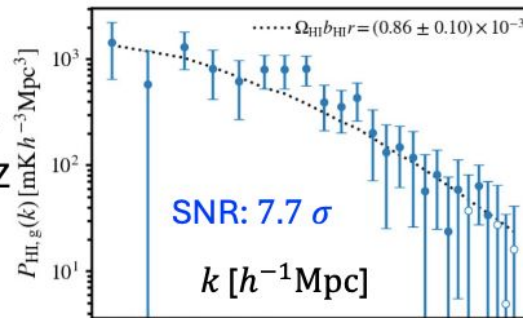
Projected galaxies (WISE) \times tSZ (2310.18478)

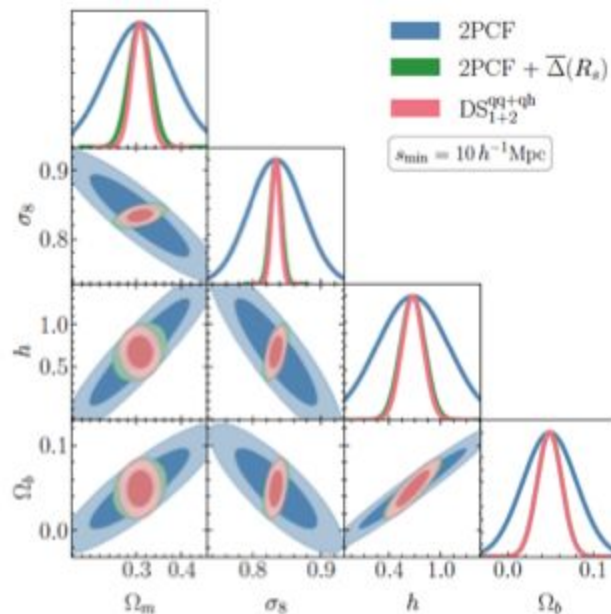
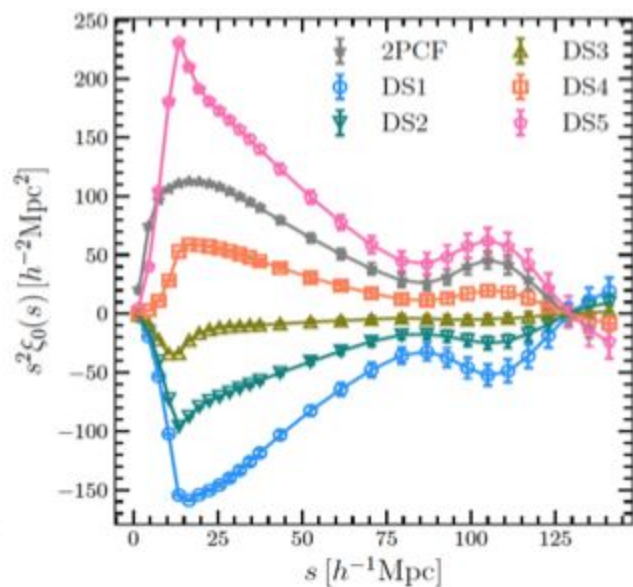
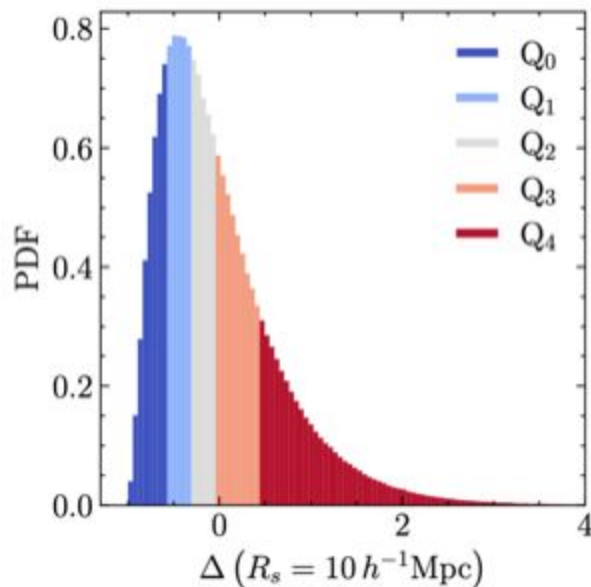


SNR: 21.8σ



MeerKAT
X WiggleZ





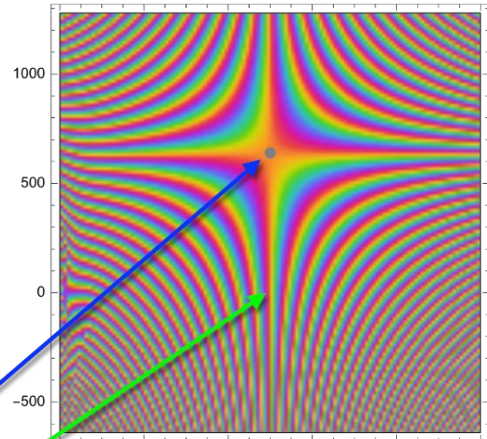
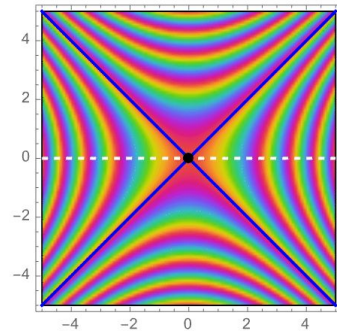
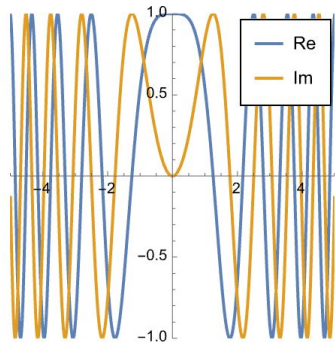
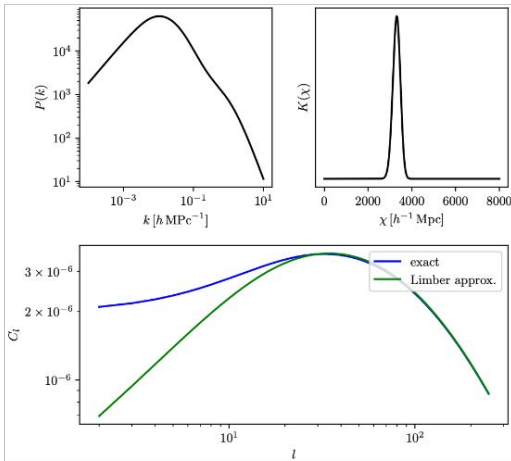
- E. Paillas, YC, N. Padilla, A. Sánchez, 2021, MNRAS.505.5731P
 E. Paillas, C. Cuesta-Lazaro, P. Zarrouk, YC... 2023, MNRAS.522.606P
 Giblin B., YC & Hornois-Deraps 2023, MNRAS.520.1721G
 E. Paillas...YC et al. 2024, MNRAS.531.898P
 C. Cuesta-Lazaro...YC et al. arXiv:2309.16539

See also:

- D. Gruen et al. 2018PhRvD..98b3507G
 Neyrinck, M. et al, 2018, MNRAS.478.2495N
 Abbas & Sheth 2007, MNRAS.378..641A
 Repp & Szapudi, 2022, MNRAS.509..586R

Complex Evaluation of Angular Power Spectra

Benjamin Hertzsch, Job Feldbrugge



Angular power spectrum: three-dimensional integral

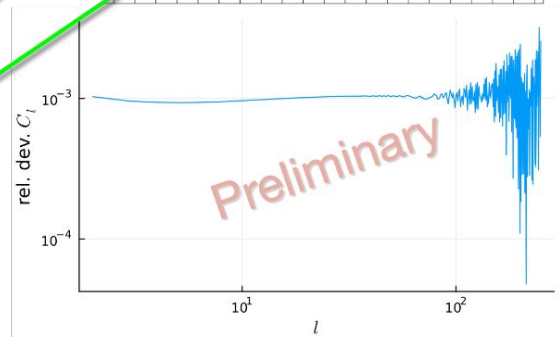
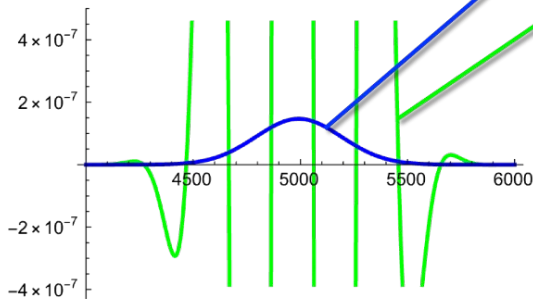
$$C_l = \frac{2}{\pi} \int \int \int dk d\chi_1 d\chi_2 k^2 P(k, \chi_1, \chi_2) K_A(\chi) j_l(k\chi) K_B(\chi) j_l(k\chi)$$

Problem: Oscillatory integral (cf. Fresnel integral)

$$F_l(k) = \int_0^\infty d\chi K(\chi) j_l(k\chi) \rightarrow \text{complex SPA [Feldbrugge (2023)]}$$

- spherical Hankel integrand
- SPA for all relevant parameter space

→ C_l one-dimensional integral, just like in Limber approximation



→ accurately evaluate $F_l(k)$ without integral

Question: Can we do integral in complex k ...?