

Dynamical Edge Modes and Entanglement in Maxwell Theory

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Based on 2403.14542 with Albert Law and Gabriel Wong

Outline

- Context: What are edge modes?

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 - Standard approach and its problems

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- A shrinkable boundary condition (BC) for Maxwell
- Bulk-edge split
- Resolving discrepancy

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- Hilbert space factorization? $\mathcal{H}_\Sigma = \mathcal{H}_A \otimes \mathcal{H}_{\bar{A}}$?

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 - Constraints generically violated
 - Extended objects (e.g. Wilson lines) get cut open

A unified framework

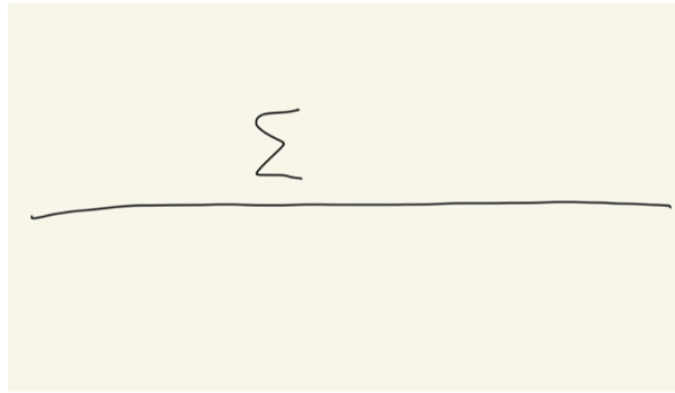
- Can address UV and IR obstructions together through **shrinkability**

A unified framework

- Can address UV and IR obstructions together through **shrinkability**
- Need to review path integral approach to EE first

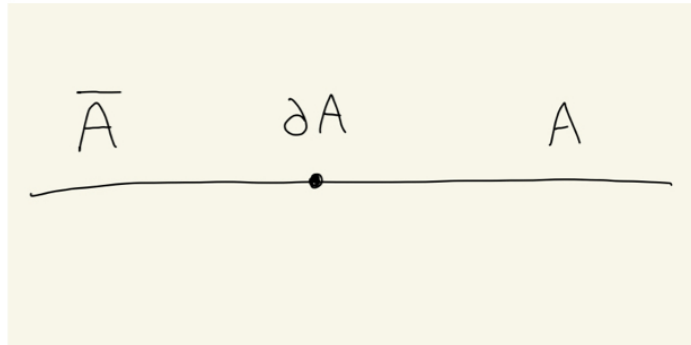
EE from path integral

For concreteness, let Σ
be a slice of Minkowski



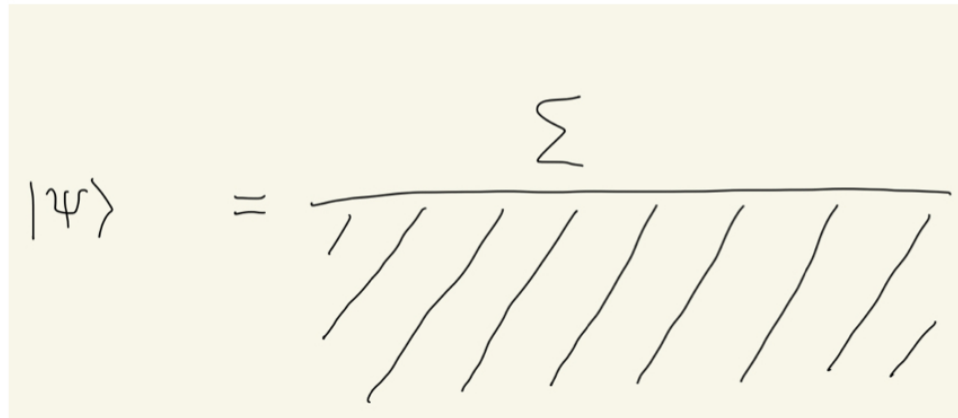
EE from path integral

Subregions and
entangling surface



EE from path integral

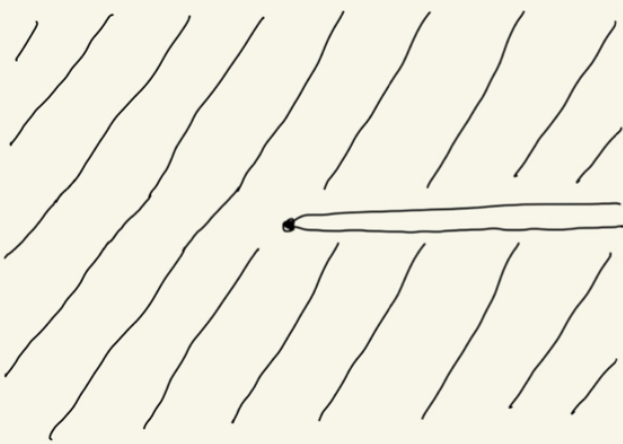
Prepare vacuum with
Euclidean half-plane



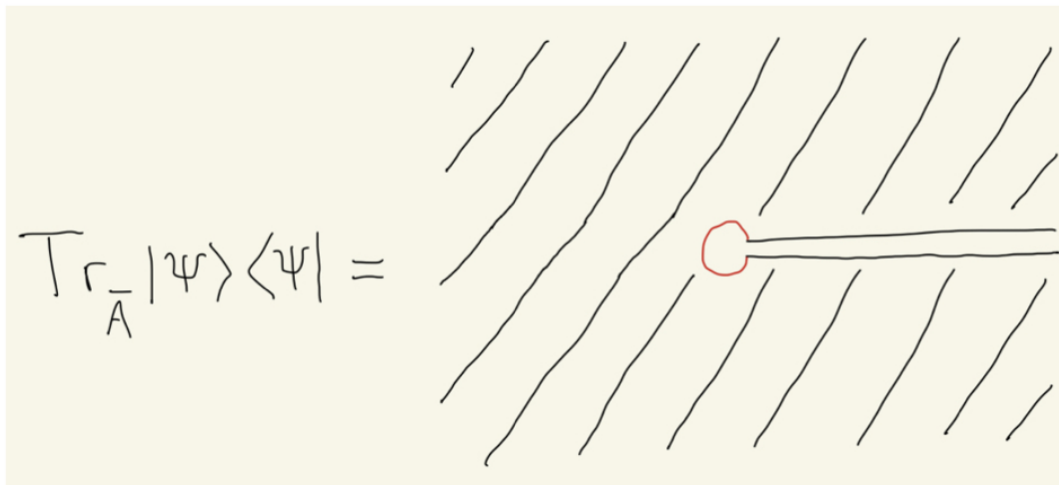
EE from path integral

$$|\psi\rangle\langle\psi| = \frac{\text{[Diagram of a path integral with two horizontal lines and diagonal hatching above and below]}{\text{[Diagram of a path integral with two horizontal lines and diagonal hatching above and below]}}$$

EE from path integral

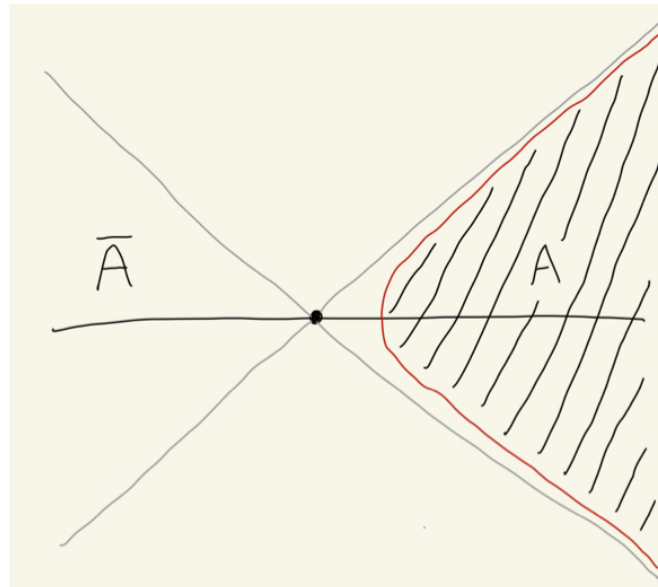
$$\text{Tr}_{\bar{A}} |\psi\rangle\langle\psi| =$$


EE from path integral



Cut out disk and choose a local BC.
Allows Hilbert space on radial slices

Brick wall



Lorentzian picture

EE from path integral

$$\text{Tr}_{\bar{A}} |\psi\rangle\langle\psi| = \text{[Diagram]} = \exp(-2\pi H_R)$$

The diagram shows a horizontal line with a red circle on the left side, representing a path integral. The area to the left of the red circle is filled with diagonal lines, representing a region of integration or a specific boundary condition.

EE from path integral

$$\text{Tr}_{\bar{A}} |\psi\rangle\langle\psi| = \text{[diagram of a path integral with a red circle on a horizontal line and diagonal hatching on either side]} = \exp(-2\pi H_R)$$

$$\rho_A \equiv \frac{\exp(-2\pi H_R)}{\text{Tr}_A[\exp(-2\pi H_R)]}$$

$$S_{\text{vN}} = -\text{Tr}_A[\rho_A \log \rho_A]$$

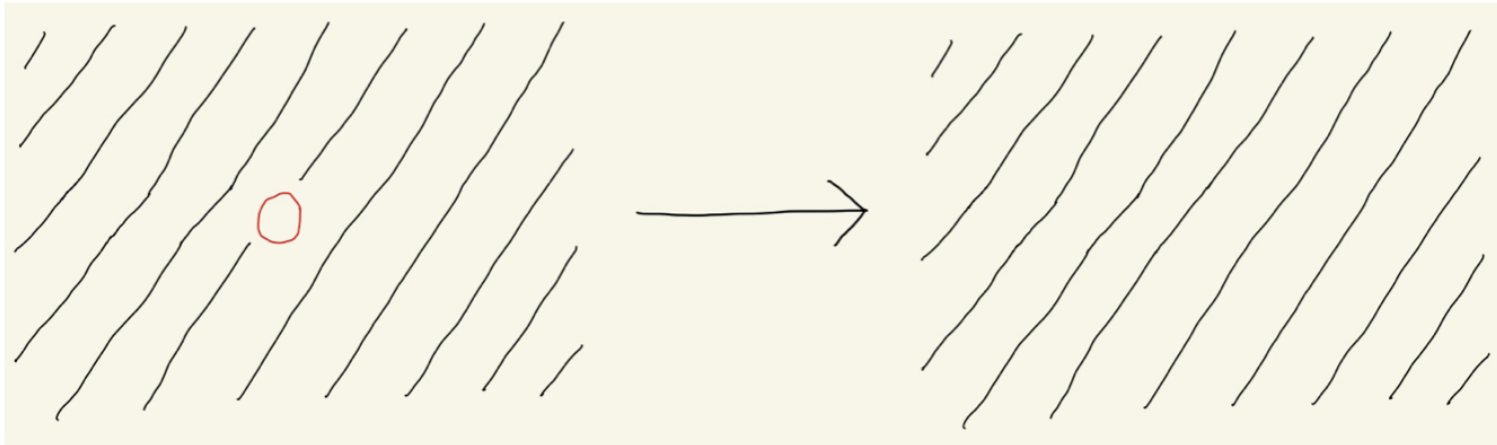
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$$S_{\text{vN}} = S_{\text{EE}}?$$

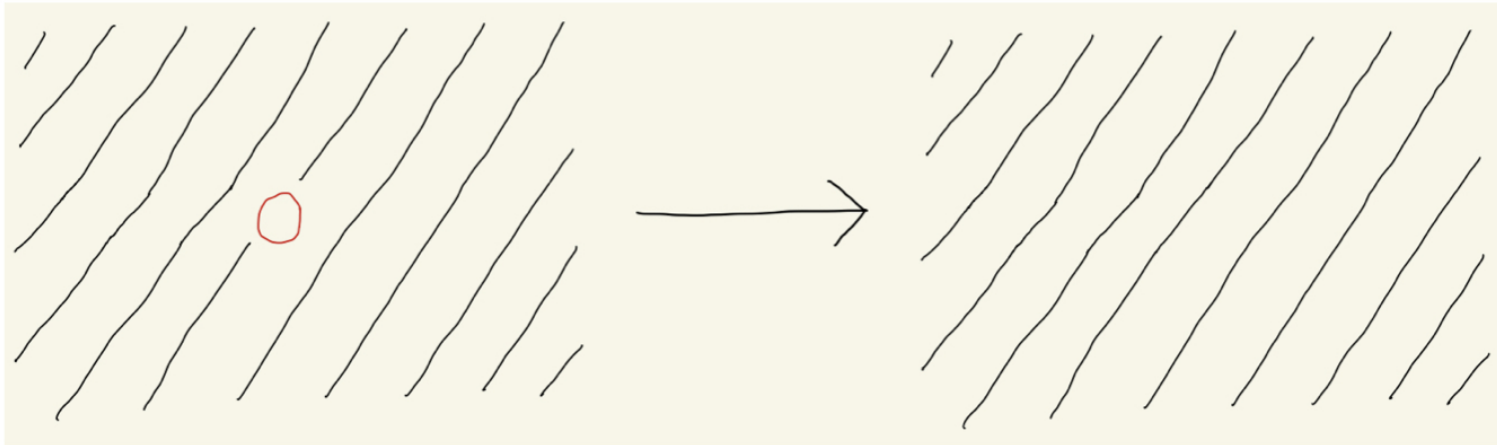
Shrinkability

A BC is shrinkable if it recovers no hole [Donnelly, Wong '18]



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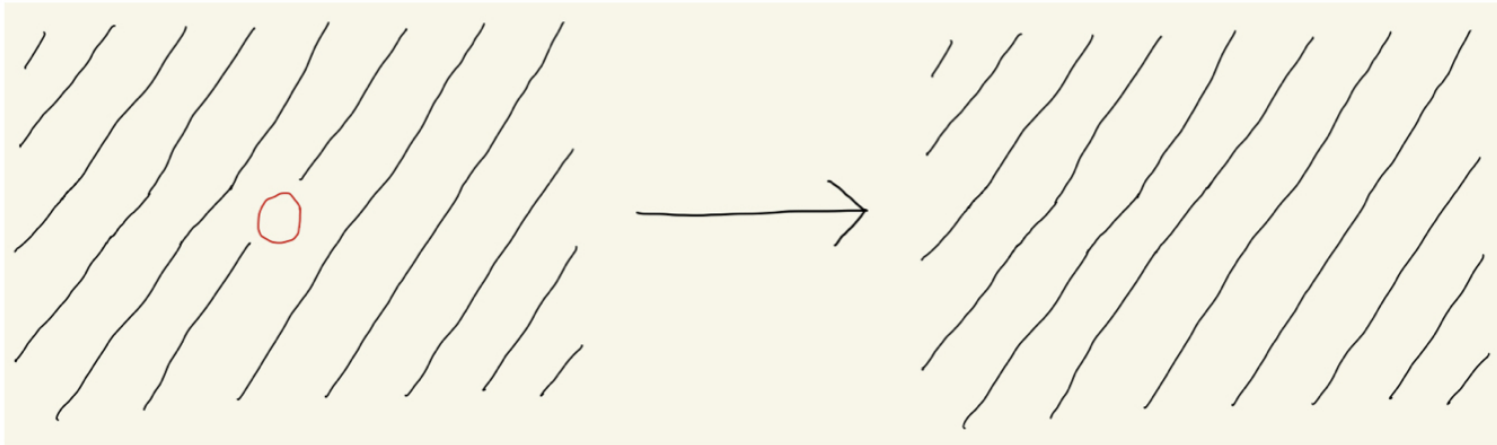
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Shrinkable BC's should give correct EE

Shrinkability

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Spoiler: including edge modes yields shrinkability

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- 4D: $S_{\text{EE,bulk}} \sim -\frac{16}{45} \log \frac{r}{\delta}$ [Dowker '10]
- 2D: $S_{\text{EE,scalar}} \sim \frac{1}{3} \log \frac{r}{\delta}$

General $Z(S^D)$

- $Z(S^D)$ and $Z_{\text{bulk}}^{(D)}$ of static dS computed for all D in [Anninos et al. '20]
 - Found systematic discrepancy $1/Z_{\text{scalar}}(S^{D-2})$

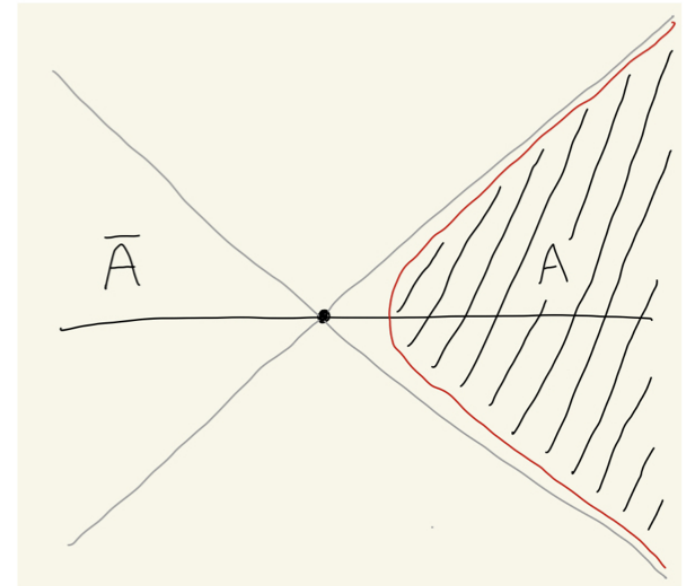
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To resolve, need $Z_{\text{edge}}^{(D)} = 1/Z_{\text{scalar}}(S^{D-2})$

BC's for Maxwell

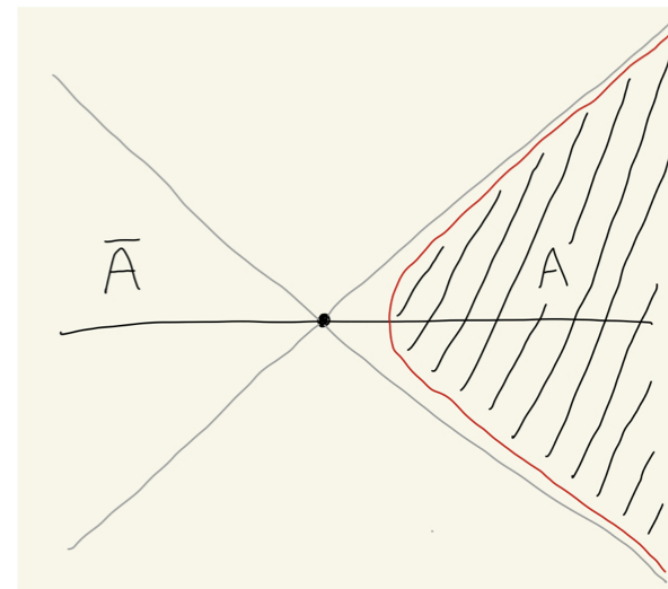
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$$\Omega = \int_A \delta A \wedge \star \delta F = \int_A \delta A^i \delta E_i$$



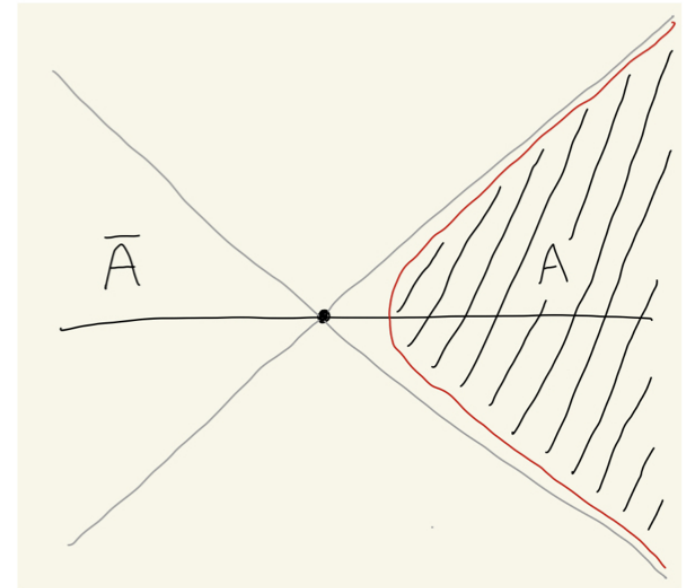
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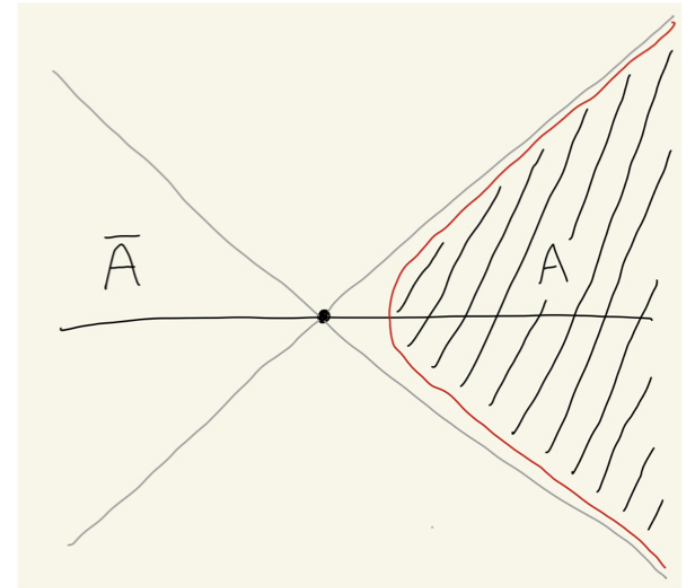
$$\Omega = \int_A \delta A \wedge \star \delta F = \int_A \delta A^i \delta E_i$$

- Plug $\delta A = d\lambda$, so $\Omega = \int_A \nabla^i \lambda \delta E_i = \int_{\partial A} \lambda \delta E_n$



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- Plug $\delta A = d\lambda$, so $\Omega = \int_A \nabla^i \lambda \delta E_i = \int_{\partial A} \lambda \delta E_n$
- Want a BC allowing both $\lambda|_{\partial A}$ and $E_n|_{\partial A}$



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 - Forces $\lambda|_{\partial M} = \text{const.}$
- Magnetically conducting BC: $F_{\mu n}|_{\partial M} = 0$
 - Forces $E_n|_{\partial M} = 0$
- Dynamical edge mode (DEM) boundary condition: $F_{in}|_{\partial M} = 0 = A_t|_{\partial M}$
 - Allows $\lambda|_{\partial M}$ and $E_n|_{\partial M}$
 - Keeps all edge modes!

Dynamical Edge Modes

- With DEM BC, can parametrize data on A as
 - $A_i = \tilde{A}_i + \nabla_i \alpha$ where $\nabla^i \tilde{A}_i = 0 = \tilde{A}_n|_{\partial A}$
 - $E_i = \tilde{E}_i + \nabla_i \beta$ where $\nabla^i \tilde{E}_i = 0 = \tilde{E}_n|_{\partial A}$

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- Can show $\Omega = \int_A \delta \tilde{A}^i \delta \tilde{E}_i + \int_{\partial A} \delta \alpha \delta E_n$ where we used $\nabla_n \beta = E_n$

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- Phase space factorizes,
 - $\Gamma_{\text{DEM}} = \Gamma_{\text{bulk}} \times \Gamma_{\text{edge}}$ also note $\Gamma_{\text{bulk}} = \Gamma_{\text{MC}}$

Dynamical Edge Modes

- Also have $H = \int_A \left(\frac{1}{2} \tilde{E}^i \tilde{E}_i + \frac{1}{4} \tilde{F}^{ij} \tilde{F}_{ij} \right) + \int_{\partial A} E_n \frac{1}{K} E_n$

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 - Can write as integral kernel in terms of harmonic Green's function
 - In static horizon limit, simplifies to $K \leftrightarrow \log(\varepsilon^{-1}) \Delta_{\partial A}$
 - Here ε is spatial distance from horizon to brick wall

Edge partition function

$$\begin{aligned} Z_{\text{DEM}}(\beta) &= \text{Tr} e^{-\beta H} \\ &= \text{Tr}_{\text{bulk}} e^{-\beta H_{\text{bulk}}} \text{Tr}_{\text{edge}} e^{-\beta H_{\text{edge}}} \\ &= Z_{\text{bulk}}(\beta) Z_{\text{edge}}(\beta) \end{aligned}$$

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- Z_{bulk} is magnetically conducting case. Has been computed in examples
- Z_{edge} is our object of interest
 - $Z_{\text{edge}}(2\pi) = \text{Tr}_{\text{edge}} e^{-2\pi H_{\text{edge}}} \sim \det(K)^{+1/2} \sim \det(\Delta_{\partial A})^{+1/2} \sim 1/Z_{\text{scalar}}(\partial A)$

Edge partition function

$$Z_{\text{edge}} = \text{Tr}_{\text{edge}} e^{-2\pi H_{\text{edge}}} \sim 1/Z_{\text{scalar}}(\partial A)$$

Found the codimension-two scalar
and wrote as trace!
Resolves discrepancies in literature.

$$4\text{D: } -\frac{31}{45} = -\frac{16}{45} - \frac{1}{3}$$

$$\text{Any D: } Z(S^D) = Z_{\text{bulk}} + Z_{\text{edge}}$$

Thank you!