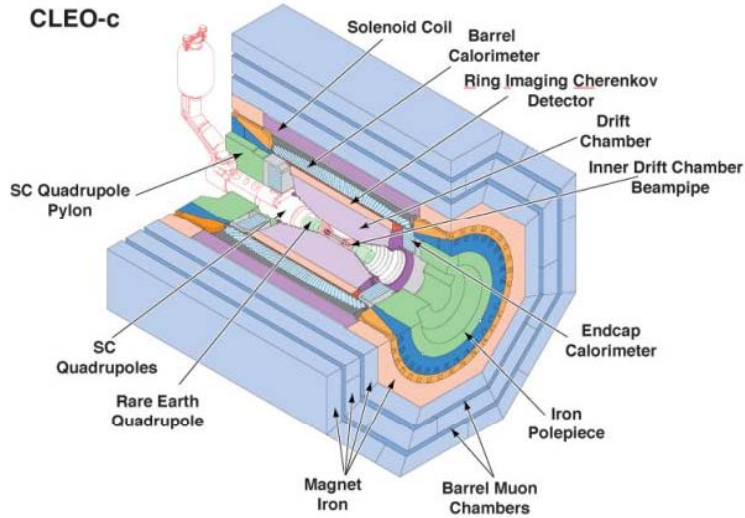


Belle Status and Prospects

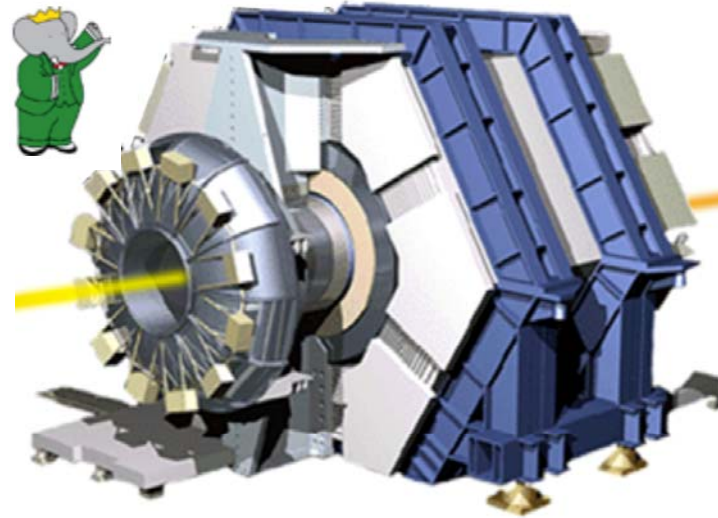
Alexey Garmash
Budker INP & Novosibirsk State U.
(on behalf of the BELLE collaboration)

- Introduction
- XYZ States @ Belle
- Prospects with Belle II

Main Suppliers of Exotics

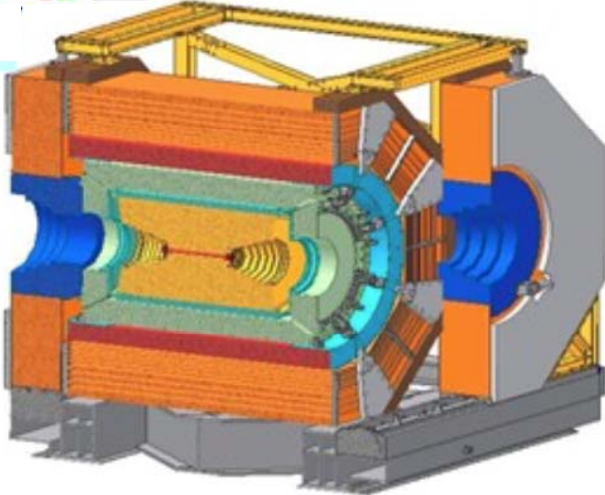


$$E_{\text{cms}} = 2.0 \dots 4.6 \text{ GeV}$$

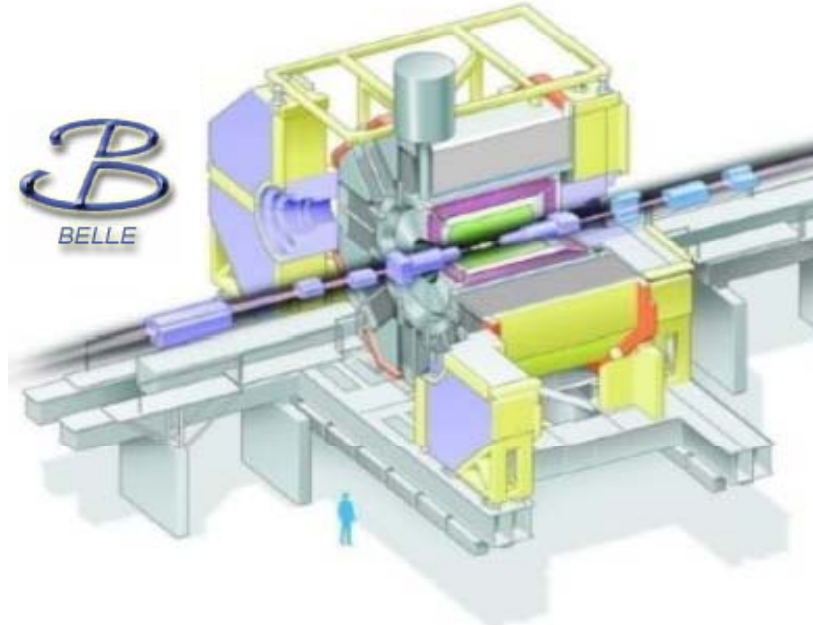


$$E_{\text{cms}} = 9.46 \dots 11.02 \text{ GeV}$$

BES III



BELLE

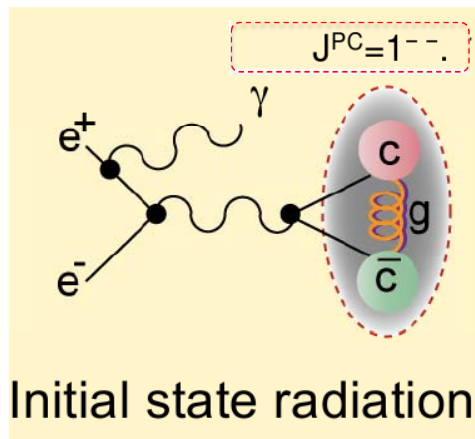


Introduction

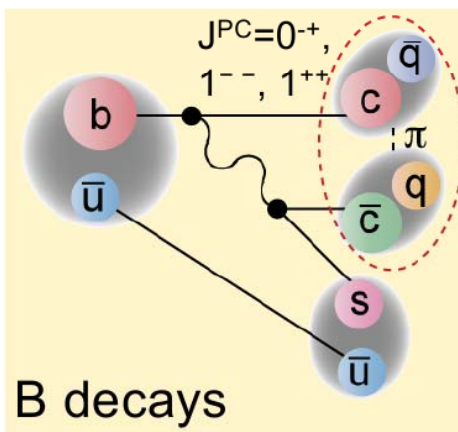
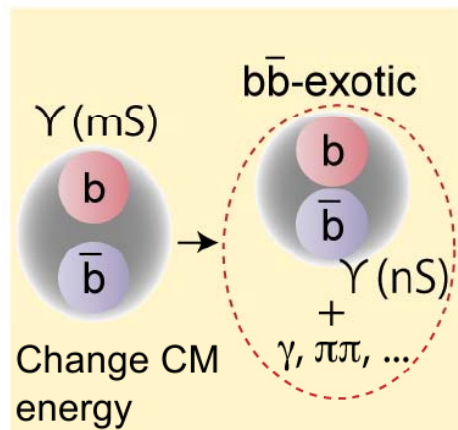
Recently, PDG switched to a notation scheme common for all new states: $X(NNNN)$, where $NNNN$ is a resonance mass in MeV.

In this talk I'll stick to an original scheme as it (to some extent) reflects production mechanism for a particular resonance.

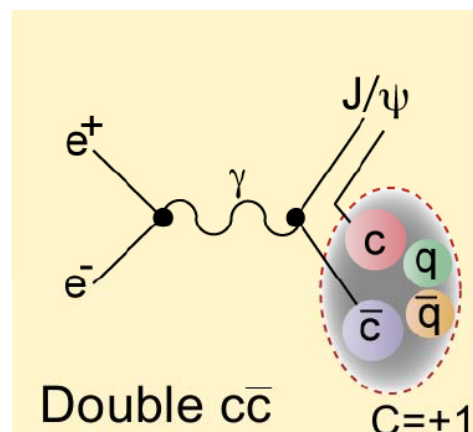
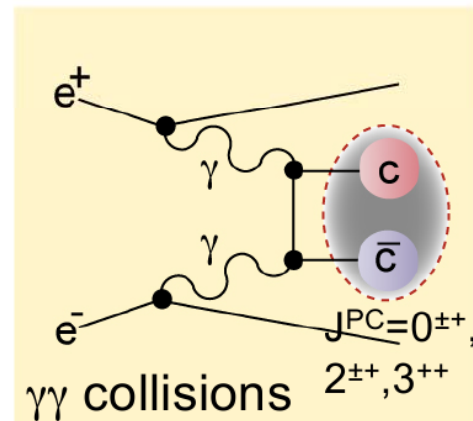
Y states ($J^{PC}=1^{--}$)



Z states (charged)



X states (all the rest)



X States

X(3872), X(3915) ...

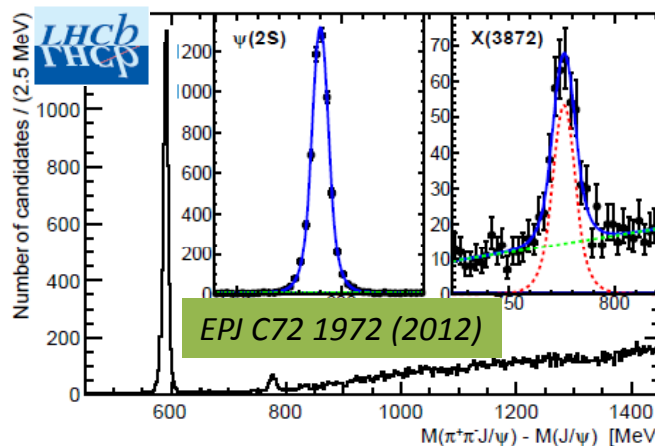
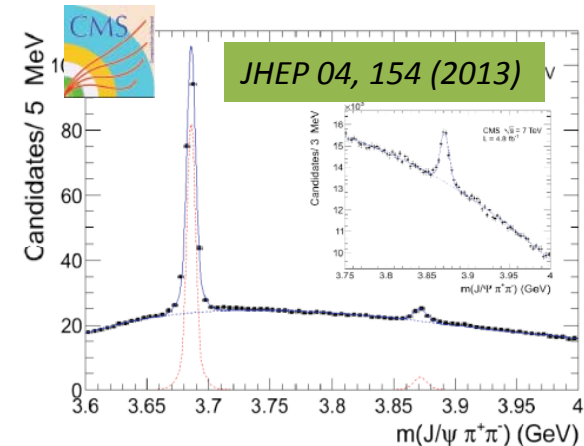
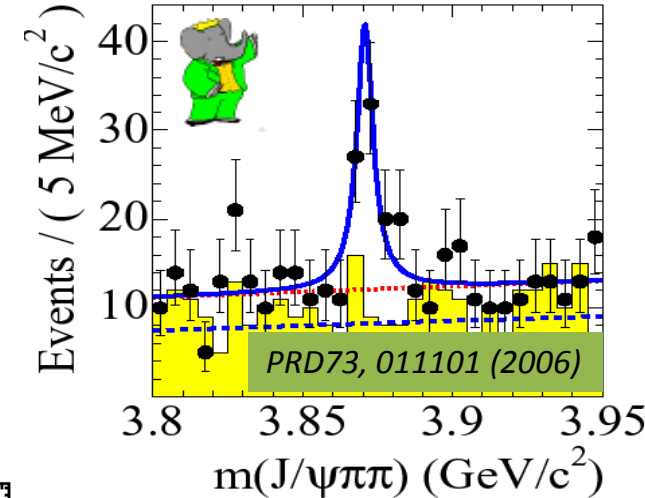
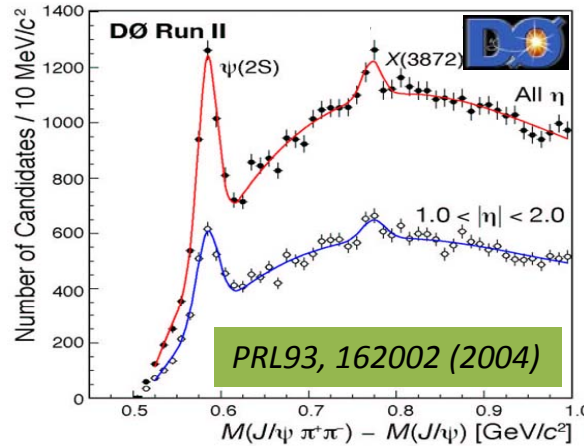
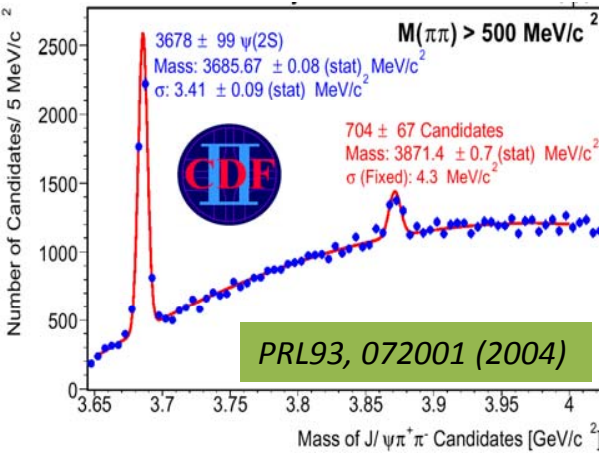
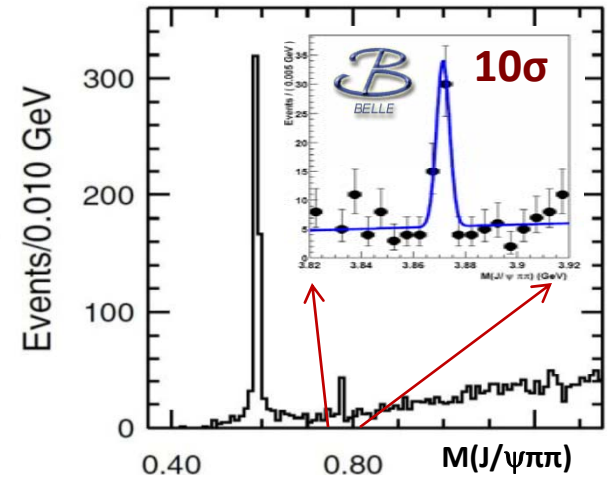
Observed in B decays and (suggestively) in decays of Y type states

$X(3872) \rightarrow J/\psi \pi^+ \pi^-$

Belle's most cited paper: ~1300

first observed by Belle in $B \rightarrow K J/\psi \pi^+ \pi^-$ *PRL91, 262001 (2003)*

- ◆ M_X close to $D^0 D^{*0}$ threshold $M = 3871.68 \pm 0.17$ MeV (not clear below or above: $\Delta m = -0.16 \pm 0.32$ MeV)
- ◆ surprisingly narrow: $\Gamma_{\text{tot}} < 1.2$ MeV at 90% CL



$X(3872) \rightarrow J/\psi \gamma$: **C-even**

Angular analysis:

Belle 2006: $J^{PC} = 1^{++}$ or ≥ 2

CDF 2008: $J^{PC} = 1^{++}$ or 2^{-+}

Belle 2011: $J^{PC} = 1^{++}$ or 2^{-+}

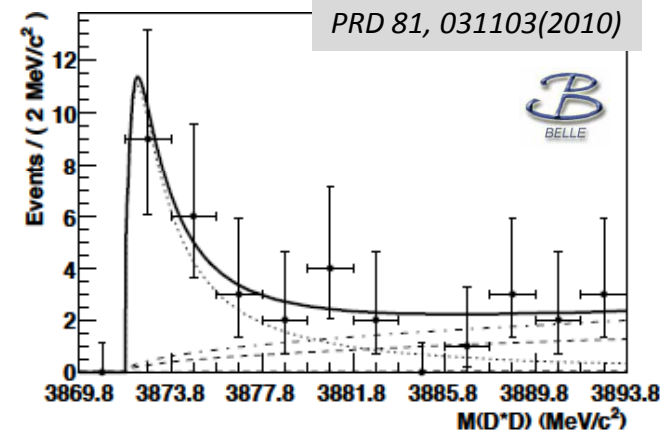
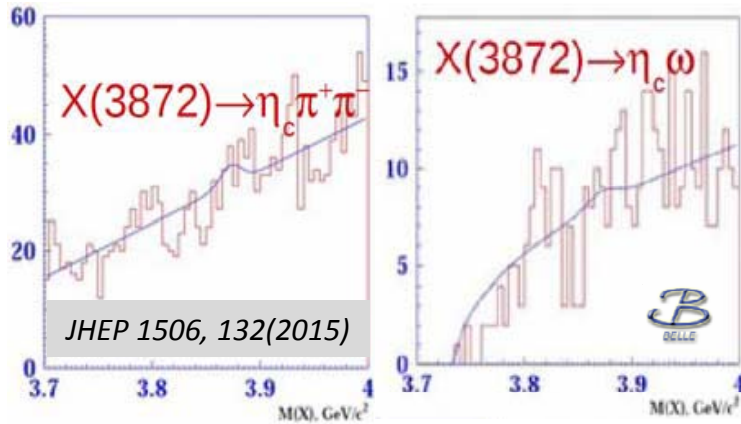
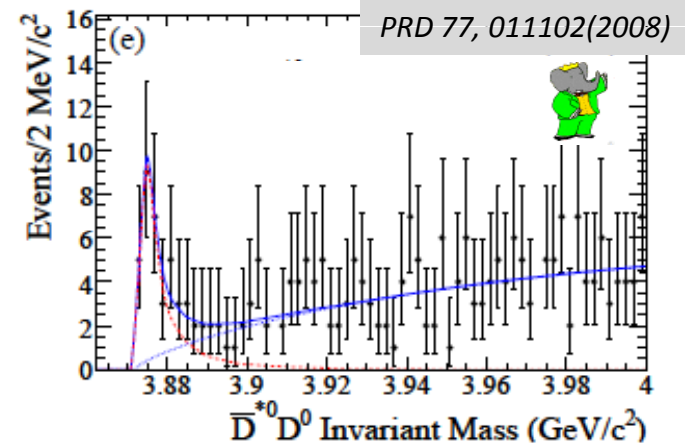
***LHCb 2013:* $J^{PC} = 1^{++}$**

X(3872): Other Decay Modes

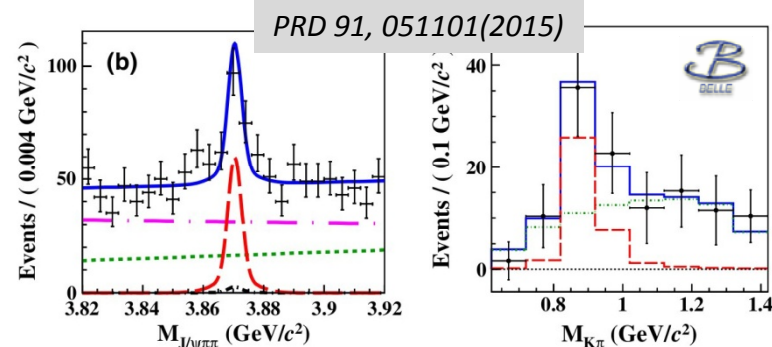
- ◆ $X(3872) \rightarrow J/\psi \rho$ & $X(3872) \rightarrow J/\psi \omega$ is seen: isospin violation
 $B(X(3872) \rightarrow J/\psi \omega) / B(X(3872) \rightarrow J/\psi \pi \pi) = 0.8 \pm 0.3$

- ◆ Radiative decays: Belle & Babar good agreement for $X \rightarrow J/\psi \gamma$;
 not consistent for $X \rightarrow \psi(2S) \gamma$.
 LHCb confirms BaBar's not vanishing $X \rightarrow \psi(2S) \gamma$.

- ◆ $X(3872) \rightarrow D \bar{D}^*$ - dominant mode:
 $B(X(3872) \rightarrow D^0 \bar{D}^{*0}) / B(X(3872) \rightarrow J/\psi \pi \pi) = 8.92 \pm 2.42$

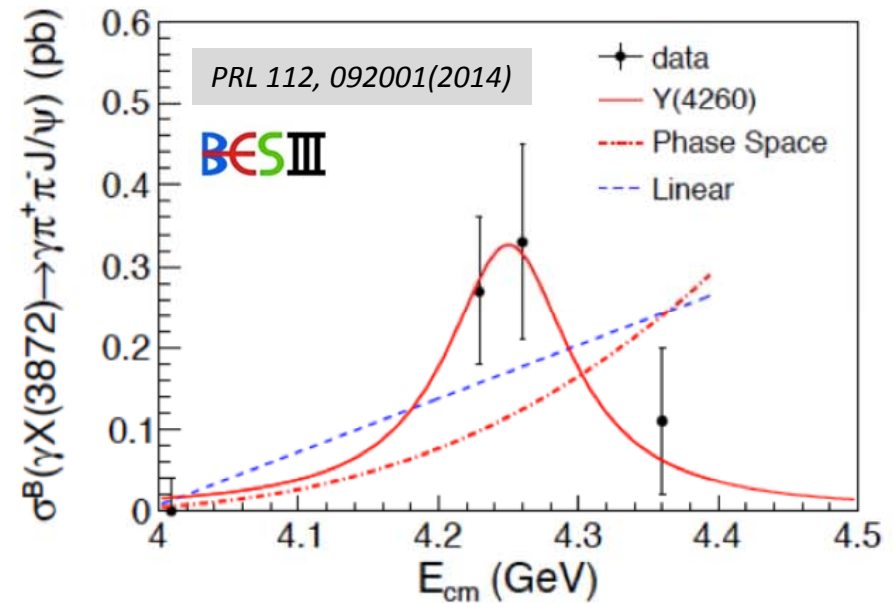
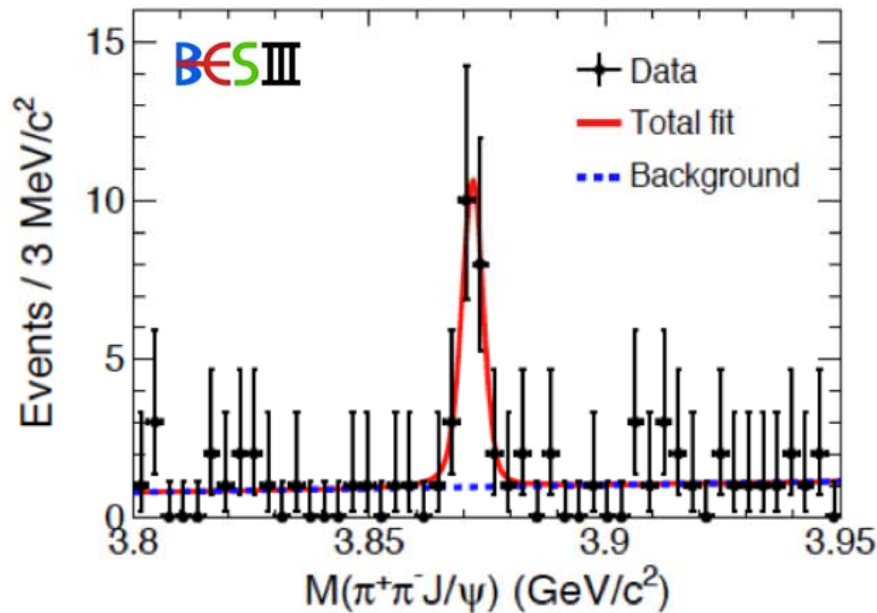


- ◆ $X(3872) \rightarrow \eta_c \pi \pi / \omega$ no signal found
- ◆ $B \rightarrow X(3872) K \pi$ non-resonant $K \pi$ dominates!



X(3872): Other Production Mechanisms

BESIII measures the X(3872) yield as a function of E_{cms} :



➤ The first observation of $e^+e^- \rightarrow \gamma X(3872) \rightarrow \gamma \pi^+ \pi^- J/\psi$

➤ $M = 3871.9 \pm 0.7 \pm 0.2 \text{ MeV}/c^2$, $\Gamma < 2.4 \text{ MeV}$, consistent with Belle's result

➤ Suggestive of $Y(4260) \rightarrow \gamma X(3872)$

➤ If $B(X(3872) \rightarrow \pi^+ \pi^- J/\psi) = 5\%$, $\mathcal{R} = \frac{\text{Br}(e^+e^- \rightarrow \gamma X(3872))}{\text{Br}(e^+e^- \rightarrow \pi^+ \pi^- J/\psi)} = 0.1$

X(3872) & Alike

Are there other X-like particles near D(*)D(*) thresholds?

Candidate	Combination	Quantum number J^{PC}	Decay modes
$X_1(3872)$	$D^0 \bar{D}^{*0} - \bar{D}^0 D^{*0}$	1^{+-}	$\eta_c \omega, \eta_c \rho$
$X(3730)$	$D^0 \bar{D}^0 + \bar{D}^0 D^0$	0^{++}	$\eta_c \eta, \eta_c \pi^0$
$X(4014)$	$D^{*0} \bar{D}^{*0} + \bar{D}^{*0} D^{*0}$	0^{++}	$\eta_c \eta, \eta_c \pi^0$

Upper limits of $\mathcal{B}(B^\pm \rightarrow K^\pm X(\rightarrow \eta_c h)) (\times 10^{-5})$

	Decay mode	Yield	UL
$X_1(3872)$	$\eta_c \pi^+ \pi^-$	17.9 ± 16.5	3.0
	$\eta_c \omega$	6.0 ± 12.5	6.9
$X(3730)$	$\eta_c \eta(\gamma\gamma)$	13.8 ± 9.9	4.6
	$\eta_c \eta(\pi^+ \pi^- \pi^0)$	1.4 ± 1.0	
$X(3730)$	$\eta_c \pi^0$	-25.6 ± 10.4	5.7
$X(4014)$	$\eta_c \eta(\gamma\gamma)$	8.9 ± 11.0	3.9
	$\eta_c \eta(\pi^+ \pi^- \pi^0)$	1.3 ± 1.6	
$X(4014)$	$\eta_c \pi^0$	-8.1 ± 13.2	1.2

Upper limits of

$\mathcal{B}(B^\pm \rightarrow K^\pm + \eta_c h) (\times 10^{-5})$ at 90% C.L.

Mode	Yield	UL
$\eta_c \pi^+ \pi^-$	155 ± 72	3.9
$\eta_c \omega$	-41 ± 27	5.3
$\eta_c \eta(\gamma\gamma)$	-14.1 ± 26.1	2.2
$\eta_c \eta(3\pi)$	-1.8 ± 3.4	
$\eta_c \pi^0$	-1.9 ± 12.1	6.2

JHEP 1506, 132 (2015)



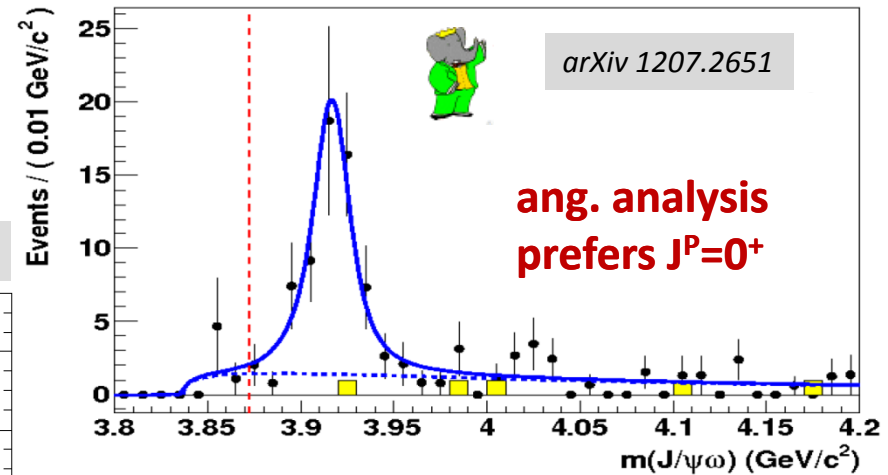
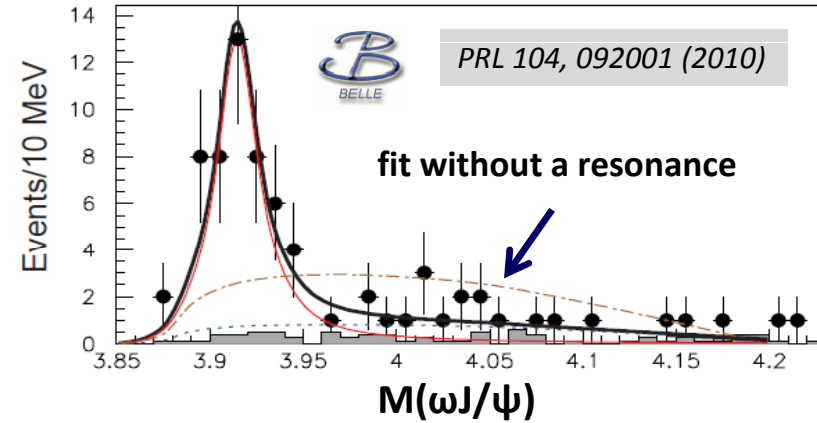
No signal is observed in any of the studied decay channels. The upper limits on their production are set at 90% C.L.

X(3915) [aka Y(3940)]

Observed in B decays and in $\gamma\gamma$

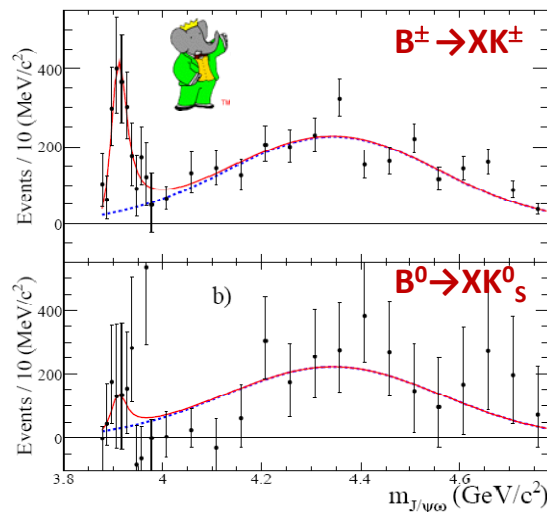
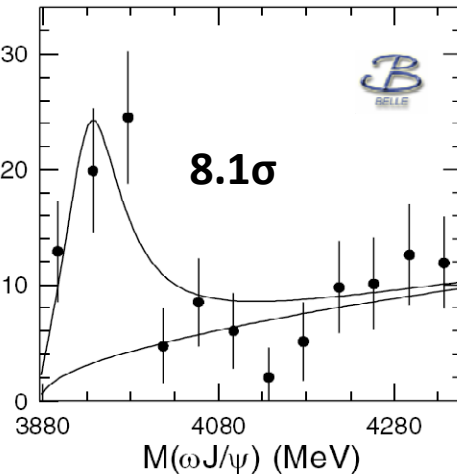
	Mass, MeV	Width, MeV
$B \rightarrow YK$ (Belle)	$3943 \pm 11 \pm 13$	$87 \pm 22 \pm 26$
$B \rightarrow YK$ (BaBar)	$3914.6 \pm 2 \pm 1.9$	$34^{+12}_{-8} \pm 6$
$\gamma\gamma \rightarrow Y$ (Belle)	$3915 \pm 3 \pm 2$	$17 \pm 10 \pm 3$
$\gamma\gamma \rightarrow Y$ (BaBar)	$3919.6 \pm 2.2 \pm 1.6$	$13 \pm 6 \pm 3$

	$\Gamma_{\gamma\gamma}$ Br (J=0)	$\Gamma_{\gamma\gamma}$ Br (J=2)
Belle	$61 \pm 17 \pm 8$	$18 \pm 5 \pm 2$
BaBar	$52 \pm 10 \pm 3$	$10.5 \pm 1.9 \pm 0.6$



PRL 94, 182002 (2005)

PRL 101, 082001 (2008)



Not observed in DD^* final state

Particle Data Group
 $Y(3940) = X(3915) = \chi_{c0}(2P)$

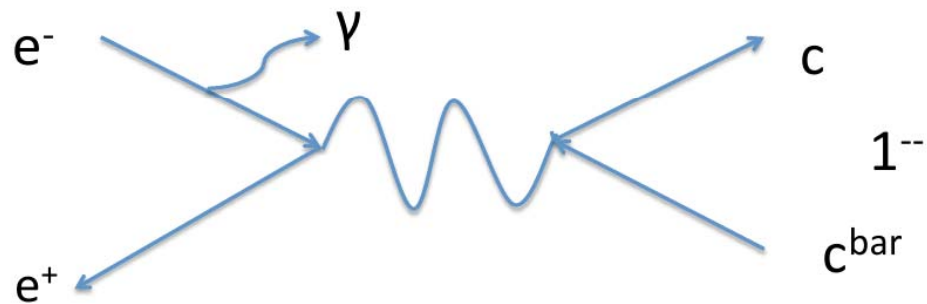
Theory

- $\chi_{c0}(2P)$ production in two body B decays is suppressed
- $\chi_{c0}(2P) \rightarrow DD$ should be dominant

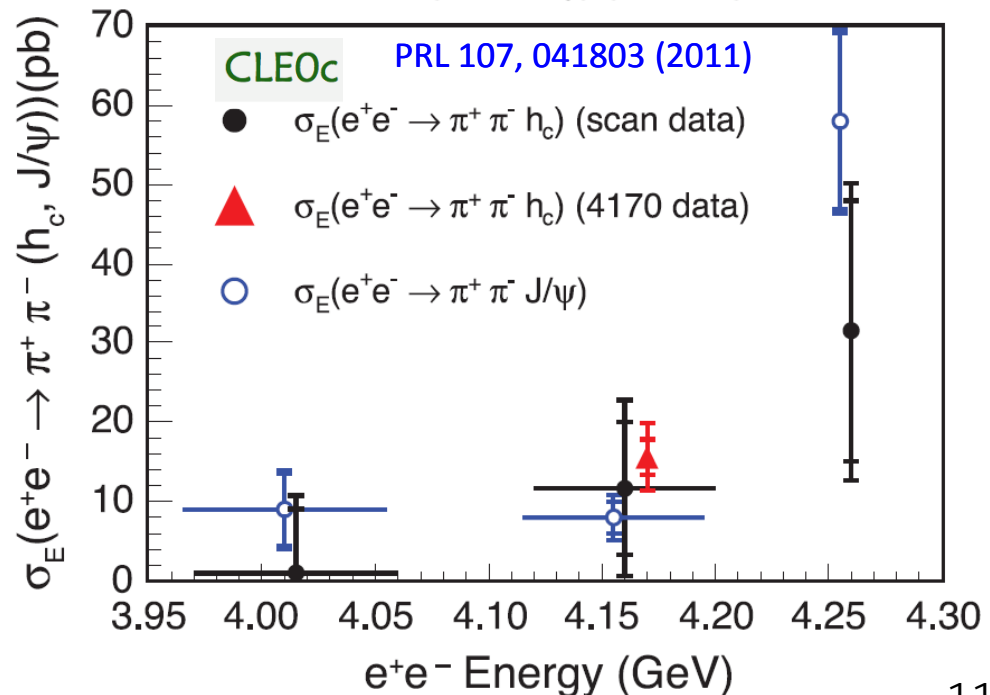
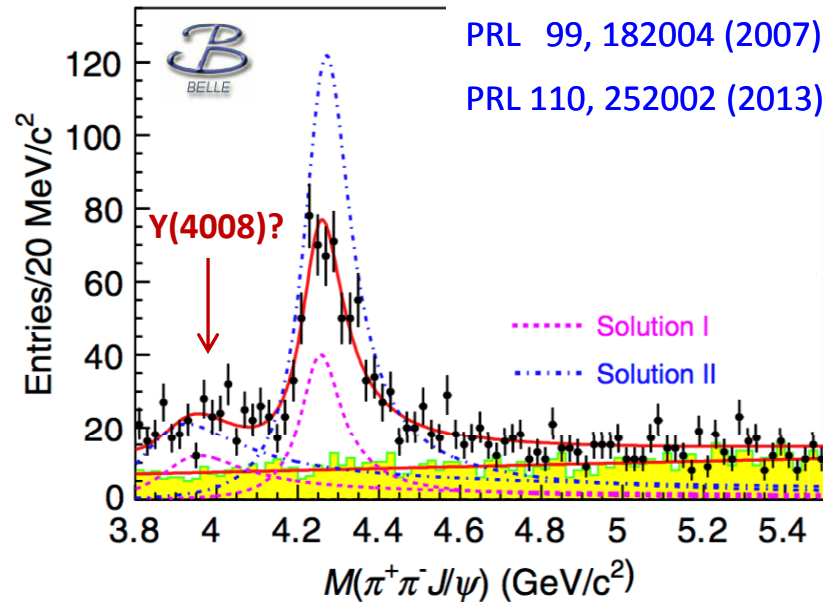
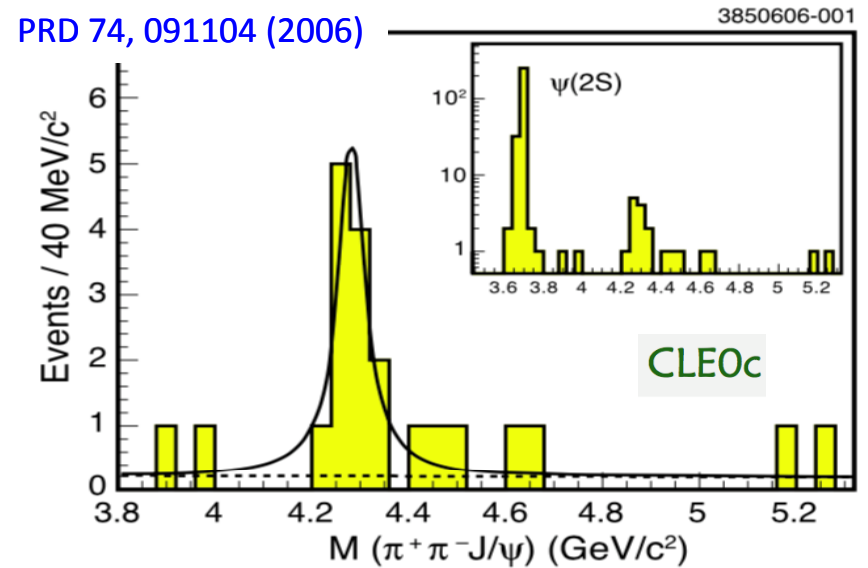
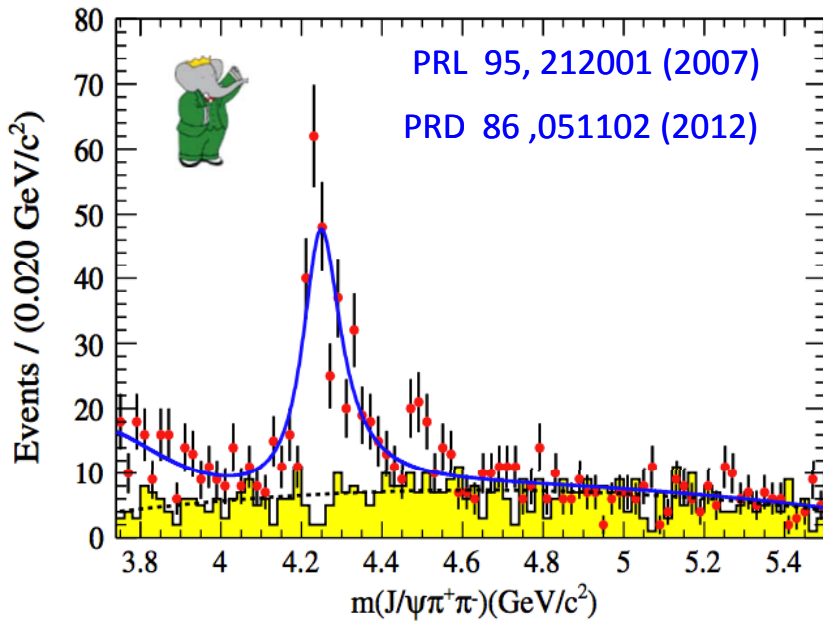
Y States

Y(4260), Y(4360), Y(4660)...

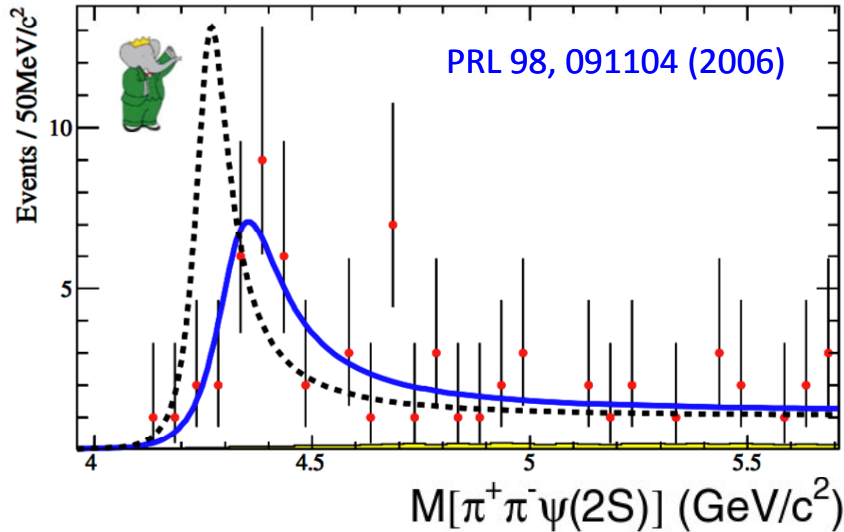
Observed in the e^+e^- annihilation (with ISR)



Studies with ISR: $\Upsilon(4260)$

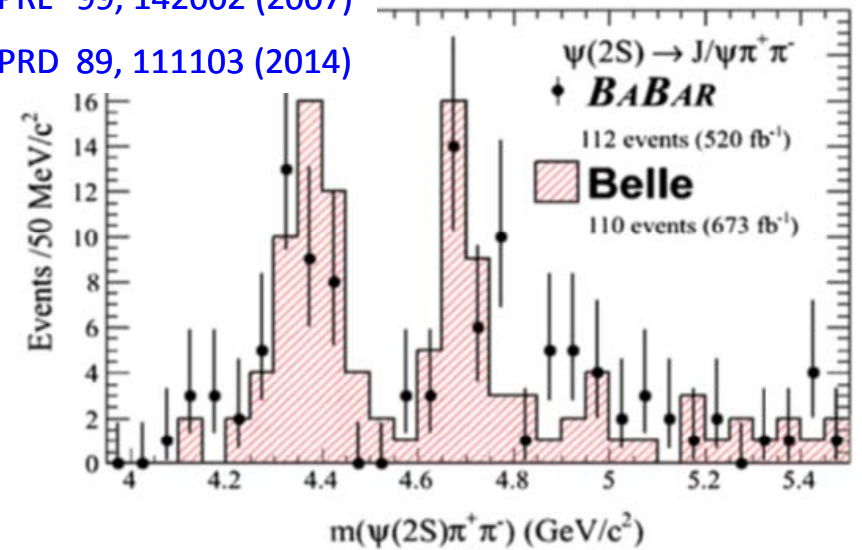


Studies with ISR: $Y(4360)$ & $Y(4660)$



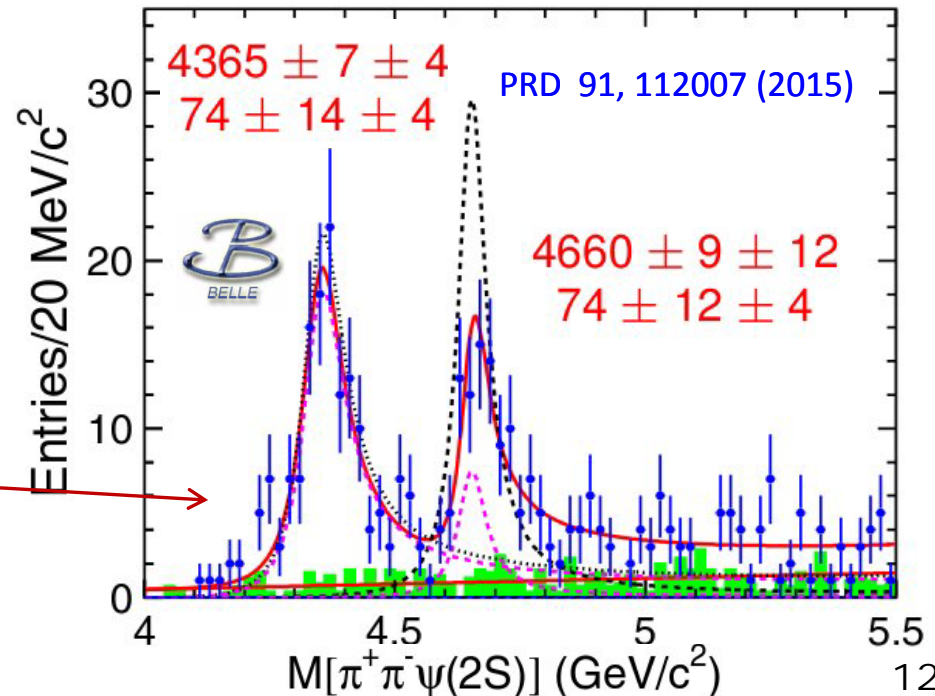
PRL 99, 142002 (2007)

PRD 89, 111103 (2014)

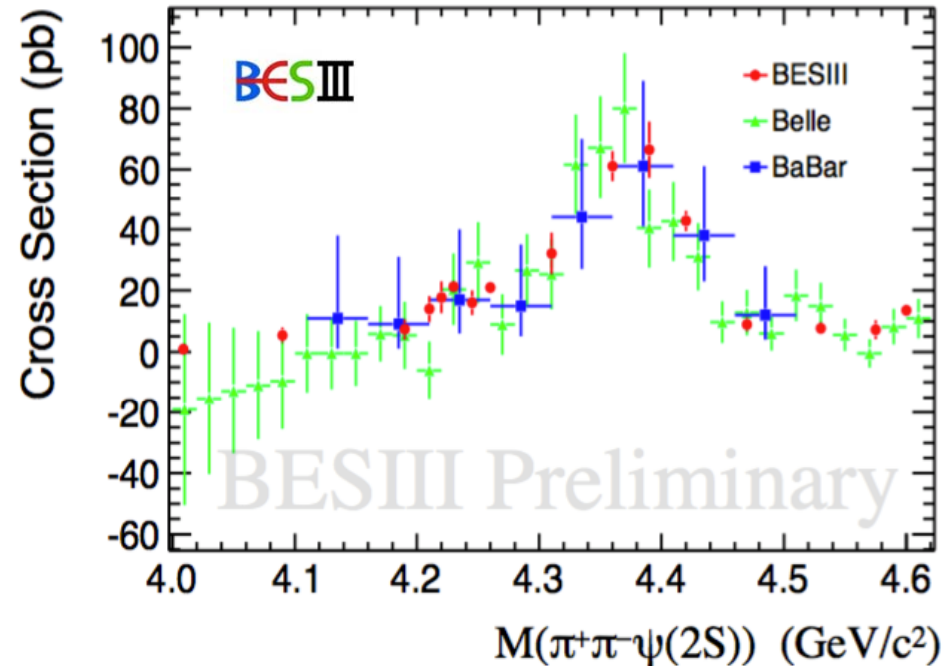
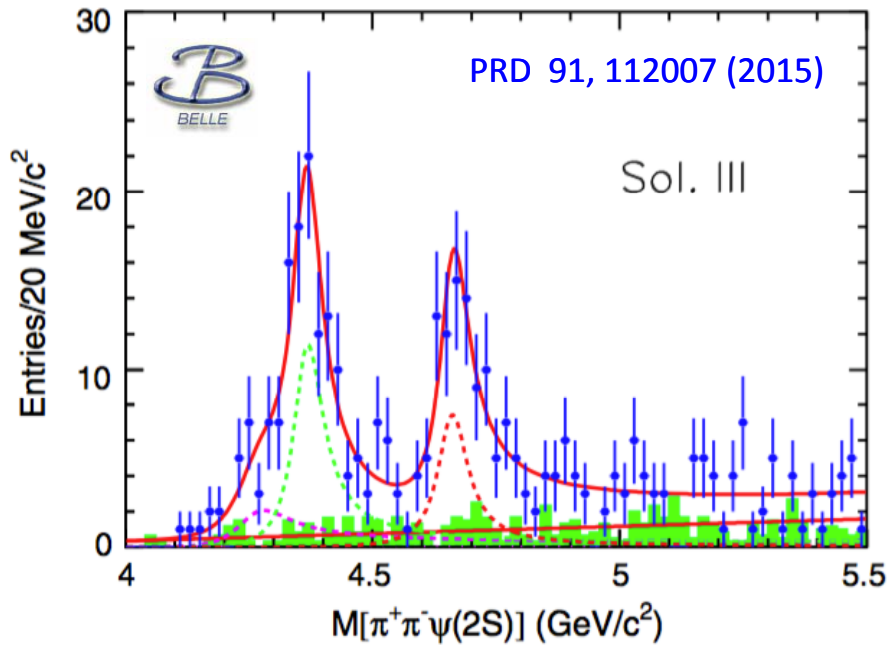


Belle update of the 2007 analysis with extra Lumi & much improved efficiency.

Some excess of signal events is observed at lower mass side. $Y(4260)$?



Studies with ISR: Y(4360) & Y(4660)



- In $\pi^+\pi^-\psi(2S)$, there are clear indications of the Y(4360) and Y(4660).
- Significance of the Y(4260) is $< 3\sigma$.

- BESIII confirms the lineshape for the Y(4360).
- More data will be taken soon to thoroughly study the region between 4.2 and 4.3 GeV.
- An analysis of the $\pi^\pm\psi(2S)$ substructure will be released soon.

None of the Y(4260), Y(4360) or Y(4660) is observed in the $D^{(*)}D^{(*)}$ final state

Z States

Z_c(3900), Z_c(4020), Z_c(4430), Z_b(10610), Z_b(10650), ...

Observed in B decays and in decays of Y states

Z_c(4430)

Z(4430)⁺ the first charged charmonium-like state

Charged – cannot be conventional charmonium or hybrid

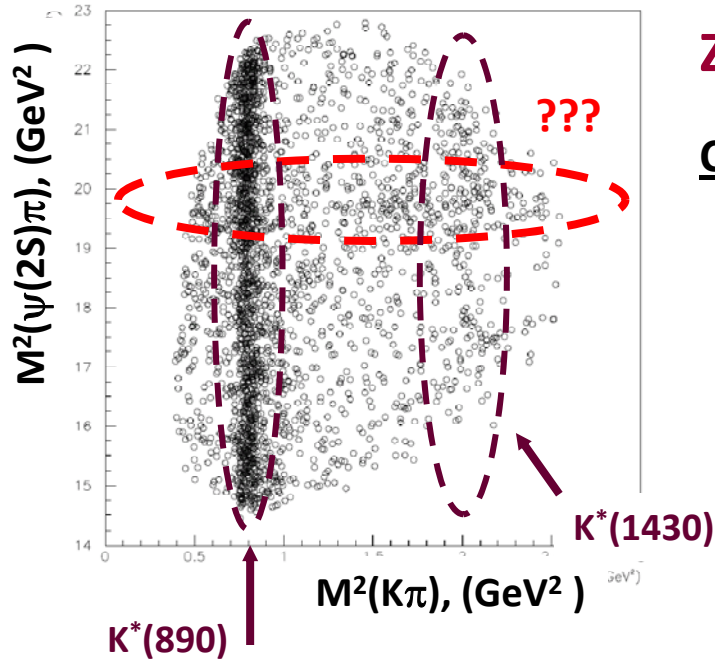
$B \rightarrow KZ, Z(4430)^+ \rightarrow \pi^+\psi(2S)$

$K=K^-, K^0_s ; \psi(2S) \rightarrow e^+e^-, \pi^+\pi^-J/\psi$

Fit: S-wave BW + phase space like function

$M = (4433 \pm 4 \pm 2) \text{ MeV}$

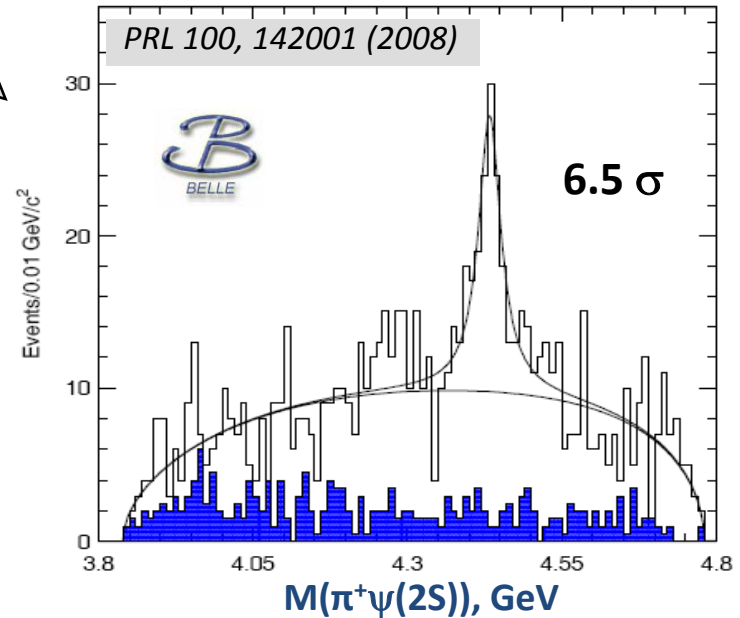
$\Gamma = (45^{+18}_{-13} \text{ } ^{+30}_{-13}) \text{ MeV}$



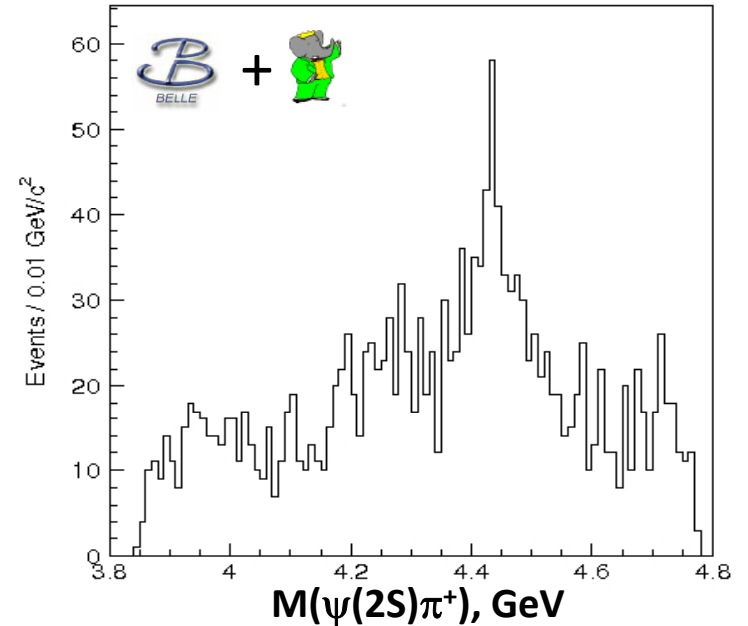
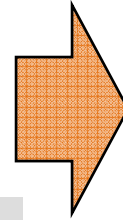
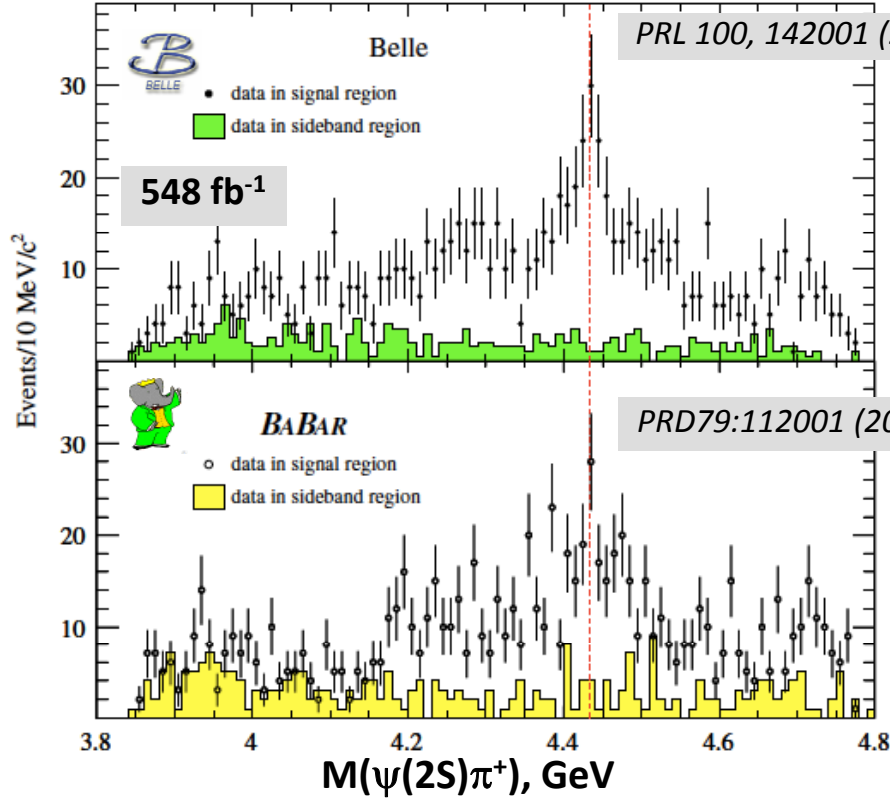
after the K* veto

$Br(B \rightarrow KZ) \times Br(Z \rightarrow \psi(2S)\pi) = (4.1 \pm 1.0 \pm 1.3) \times 10^{-5}$

- Could the Z(4430) be due to a reflection from the Kπ channel?
 - S- P- & D-waves cannot make a peak (+ nothing else)



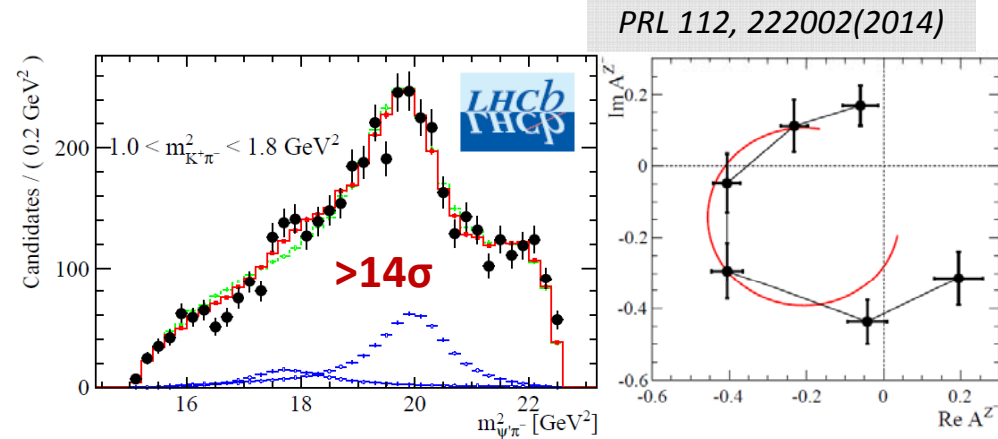
$Z_c(4430)$



4D-fit by Belle: confirms $Z_c(4430)$ & $J^P=1^+$

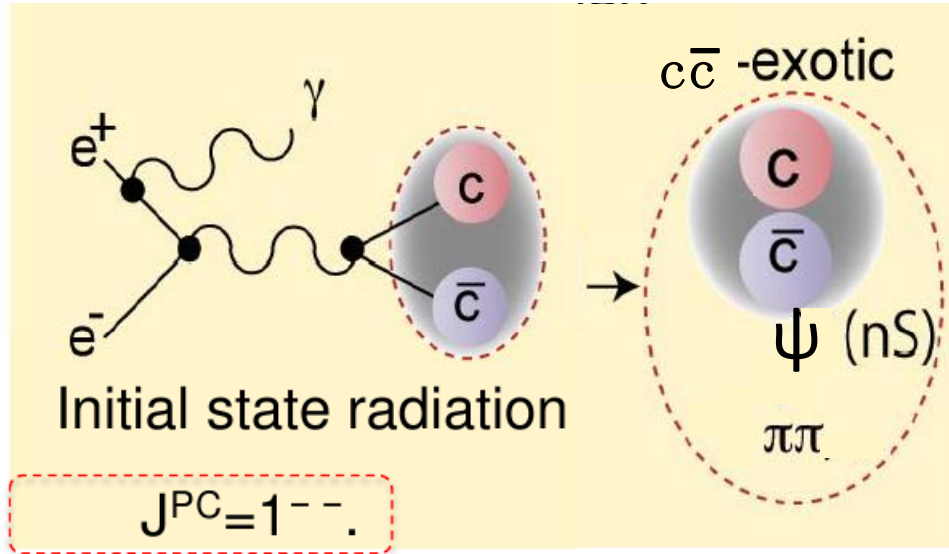
Belle and BaBar data look very similar; conclusions are different:

- **Belle:** observation of new resonance
- **BaBar:** after taking into account many $K\pi$ waves, the peak is not significant

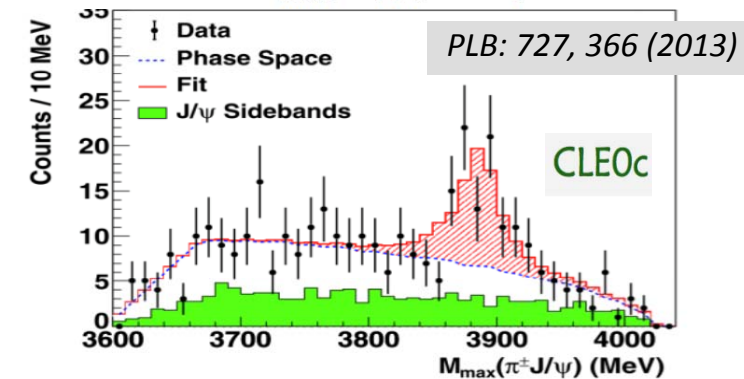
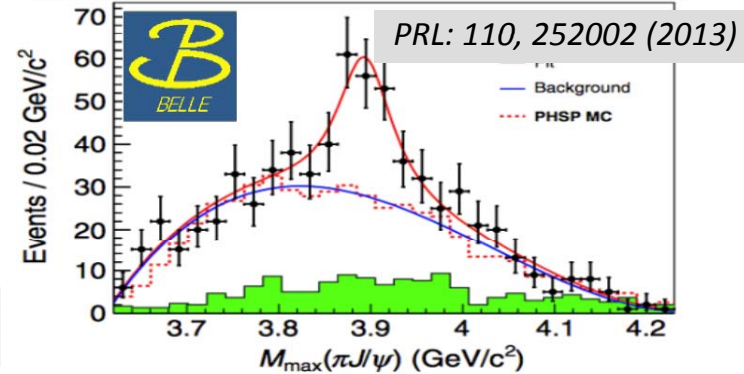
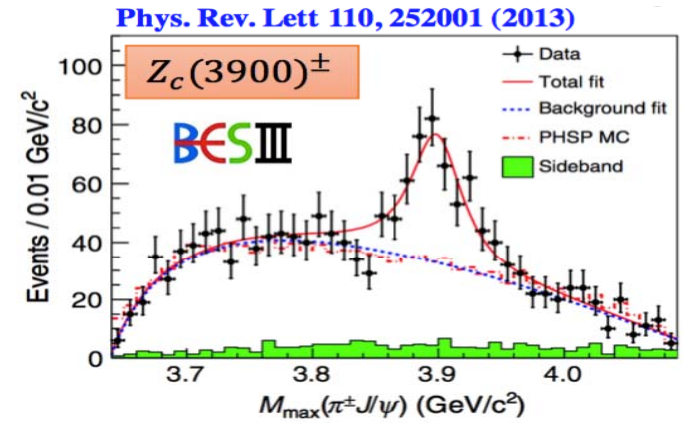


$Z_c(3900)$ State

Observed in $e^+e^- \rightarrow (\gamma)Y(4260) \rightarrow J/\psi\pi^+\pi^-$



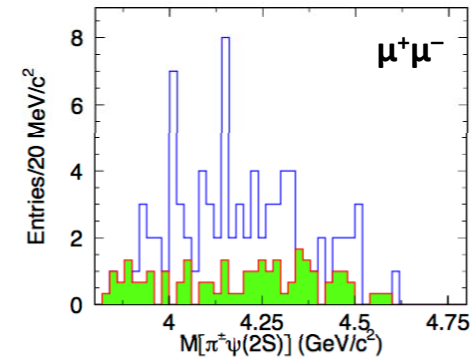
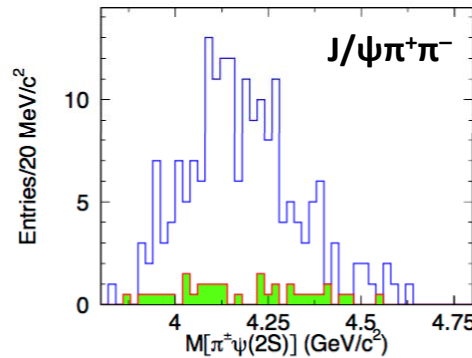
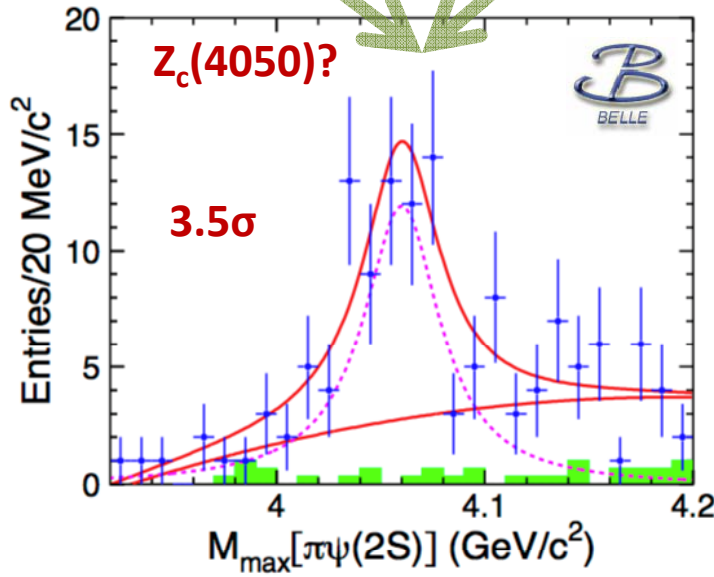
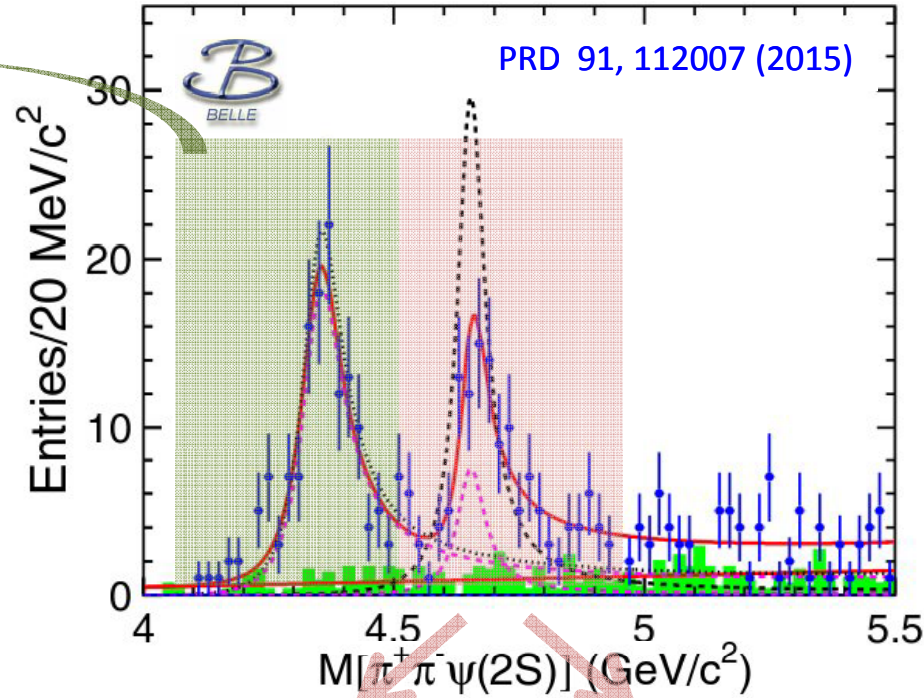
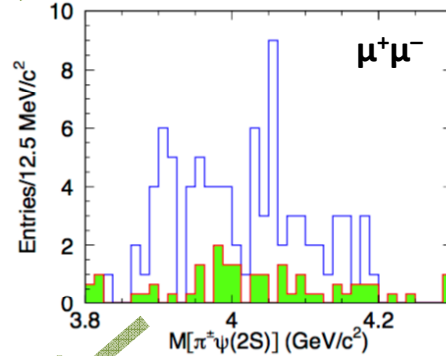
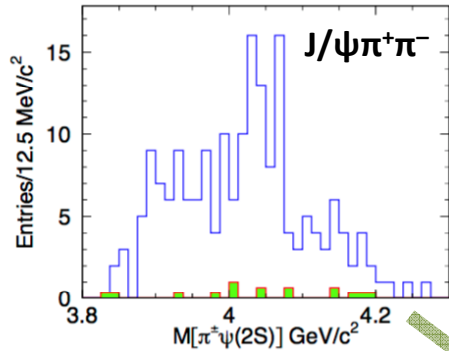
- Charged charmonium-like structure ($>10 \sigma$)
- Decay to J/ψ ($c\bar{c}$) and electric charge ($u\bar{d}$ or $d\bar{u}$)
- $M = 3899.0 \pm 3.6 \pm 4.9 \text{ MeV}/c^2$, $\Gamma = 46 \pm 10 \pm 20 \text{ MeV}$
- $\sigma(e^+e^- \rightarrow \pi^+\pi^- J/\psi) = 62.9 \pm 1.9 \pm 3.7 \text{ pb}$ at 4.26 GeV
- $\frac{\sigma(e^+e^- \rightarrow \pi^\mp Z_c(3900)^\pm \rightarrow \pi^+\pi^- J/\psi)}{\sigma(e^+e^- \rightarrow \pi^+\pi^- J/\psi)} = 21.5 \pm 3.3 \pm 7.5 \%$
- The first Z_c state observed by more than one experiment (Belle and CLEO-c)!



$Z_c(4050)$

Observed by Belle in ISR

$e^+e^- \rightarrow \gamma Y(4360) \rightarrow \psi(2S)\pi^+\pi^-$



No clear signal found in the $Y(4660)$ region

$M = 4054 \pm 3 \pm 1 \text{ MeV}/c^2$
 $\Gamma = 45 \pm 11 \pm 6 \text{ MeV}$

Another Z_c state? Need confirmation

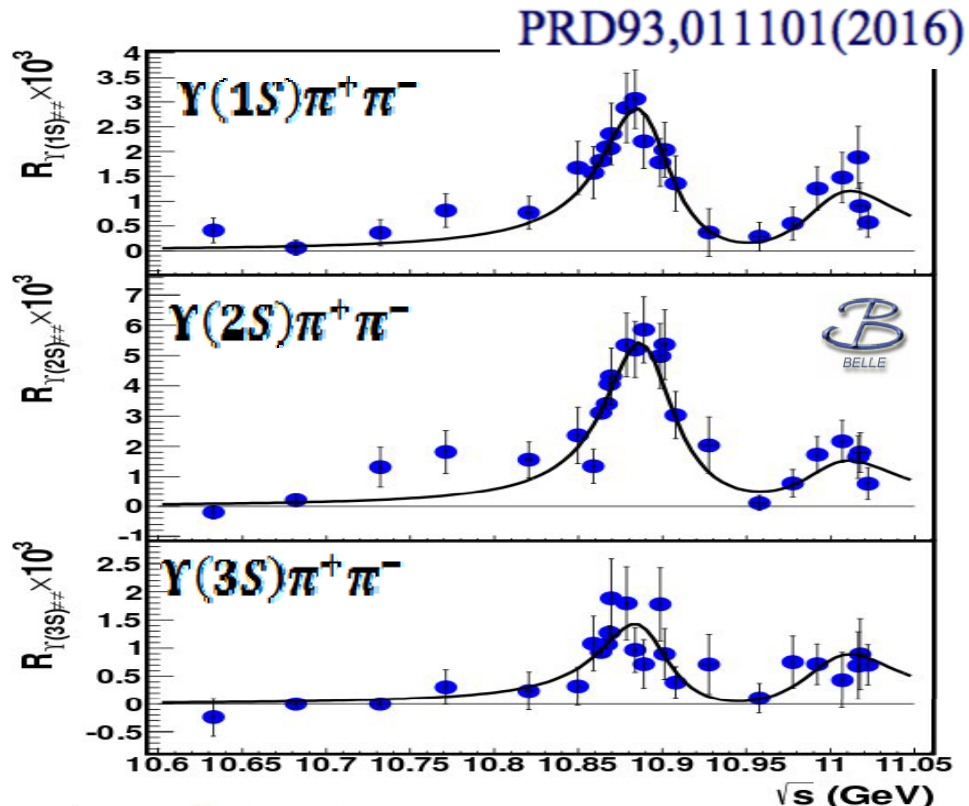
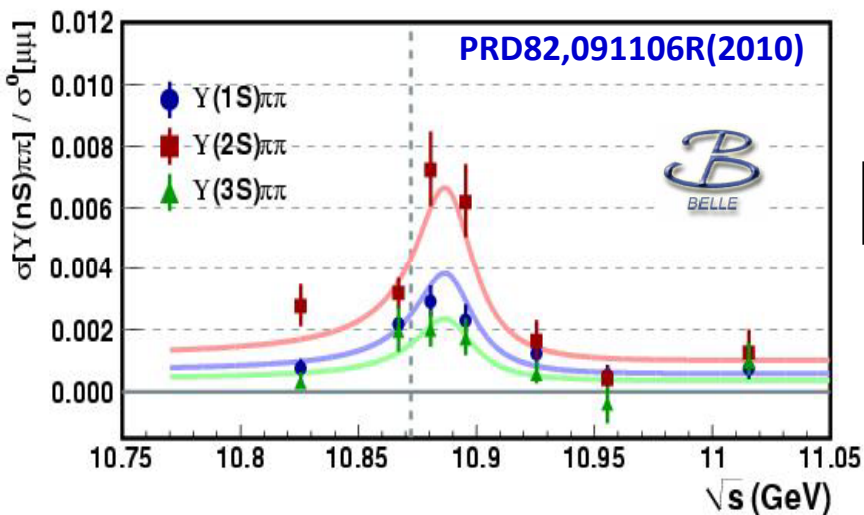
Exotics in Bottomonium Sector

- ◆ No X_b states found so far
- ◆ Mass splitting between $\Upsilon(4S)$ and $\Upsilon(5S)$ states is in strong disagreement with potential model expectation
- ◆ Anomalously (\sim factor 100) large coupling to $\Upsilon\pi^+\pi^-$
- ◆ No $h_b\pi^+\pi^-$ suppression

(2) An exotic resonance Y_b near $\Upsilon(5S)$ - analogue of the $Y(4260)$ resonance with anomalous $\Gamma(J/\psi\pi^+\pi^-)$

(3) Tetraquarks

Anomalous $\Upsilon(5S) \rightarrow \Upsilon(nS)\pi^+\pi^-$ Rates



$\Upsilon(10860)$:

$\Upsilon(11020)$:

$$10891.1 \pm 3.2^{+0.6}_{-1.7}$$

$$10987.5^{+6.4+9.0}_{-2.5-2.1}$$

$$53.7^{+7.1+1.3}_{-5.6-5.4}$$

$$61^{+9+2}_{-19-20}$$

$$\phi_{6S} - \phi_{5S}(\delta) \text{ (rad)} -1.0 \pm 0.4^{+1.4}_{-0.1}$$

$$10884.7^{+3.6+8.9}_{-3.4-1.0}$$

$$10999.0^{+7.3+16.9}_{-7.8-1.0}$$

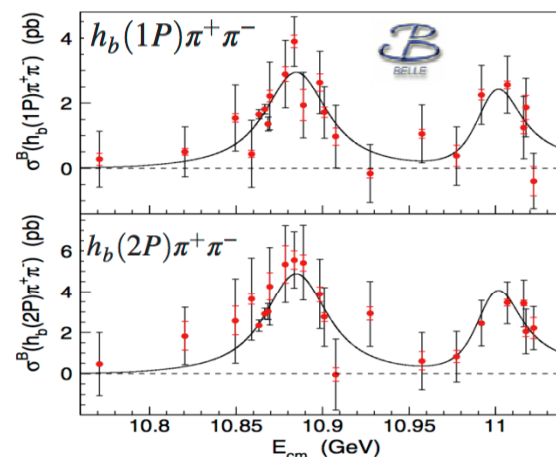
$$40.6^{+12.7+1.1}_{-8.0-19.1}$$

$$27^{+27+5}_{-11-12}$$

$$\phi = (0.1^{+0.4}_{-0.8} \pm 0.3)\pi$$



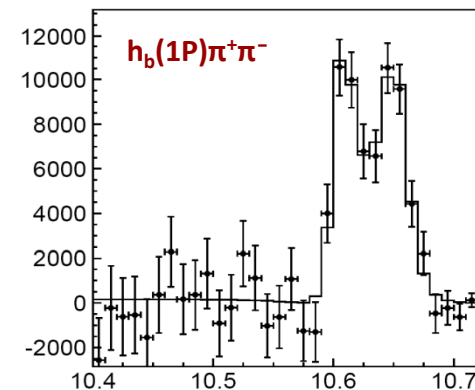
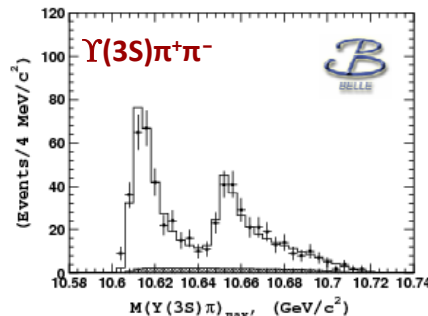
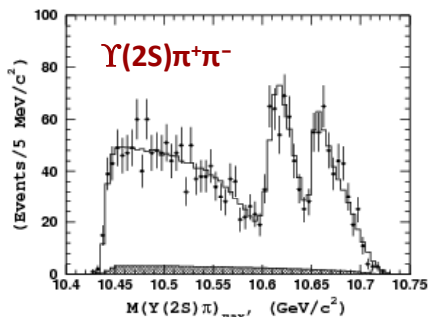
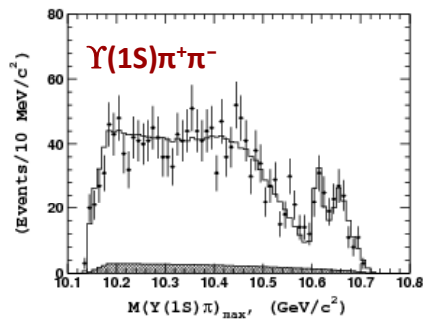
PRL117,142001(2016)



Summary on $Z_b \rightarrow$ Hidden Beauty

$\Upsilon(5S)$ Belle Data: 121.4 fb⁻¹

PRL 108, 122001(2012)



$Z_b(10610)$

$Z_b(10650)$

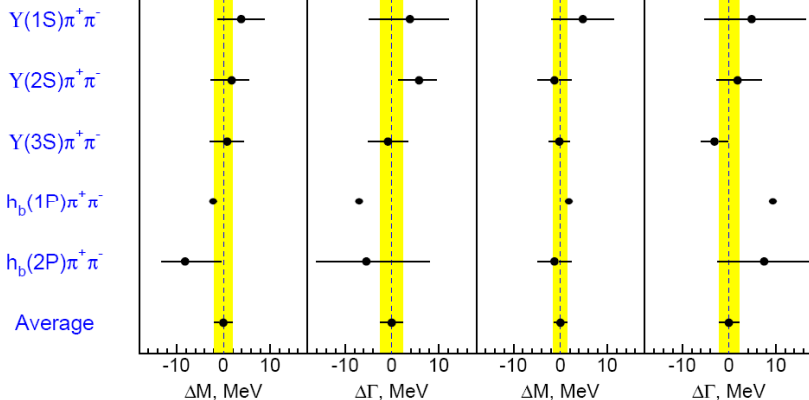
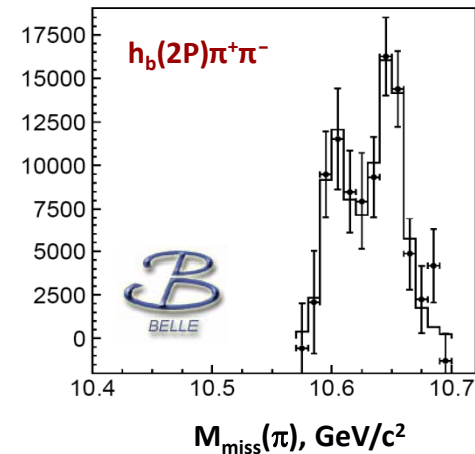
Average for Z_b^\pm :

$$\langle M_1 \rangle = 10607.2 \pm 2.0 \text{ MeV}$$

$$\langle \Gamma_1 \rangle = 18.4 \pm 2.4 \text{ MeV}$$

$$\langle M_2 \rangle = 10652.2 \pm 1.5 \text{ MeV}$$

$$\langle \Gamma_2 \rangle = 11.5 \pm 2.2 \text{ MeV}$$

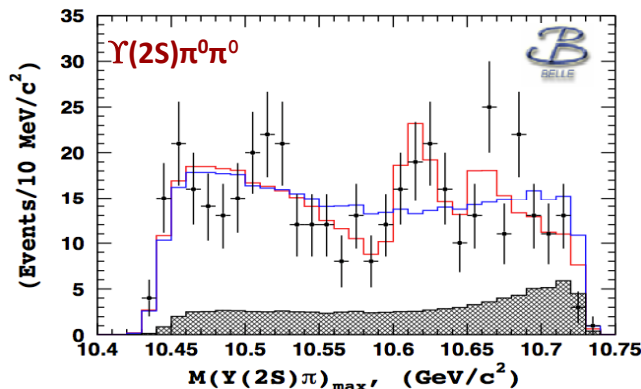


Z_b^0 Results:

$$\langle M_1 \rangle = 10609 \pm 7 \pm 6 \text{ MeV}$$

Consistent with Z_b^\pm

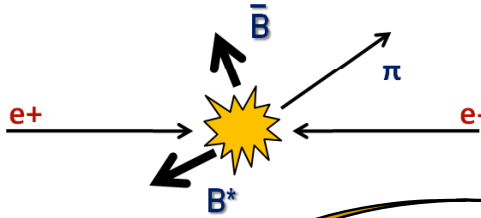
PRD88,052016 (2013)



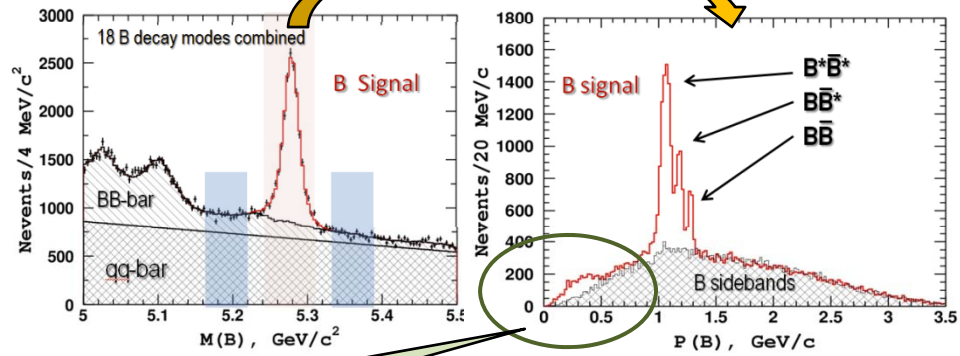
Angular analysis strongly favors $J^P=1^+$ assignment

$Z_b^\pm \rightarrow \text{Open Beauty}$

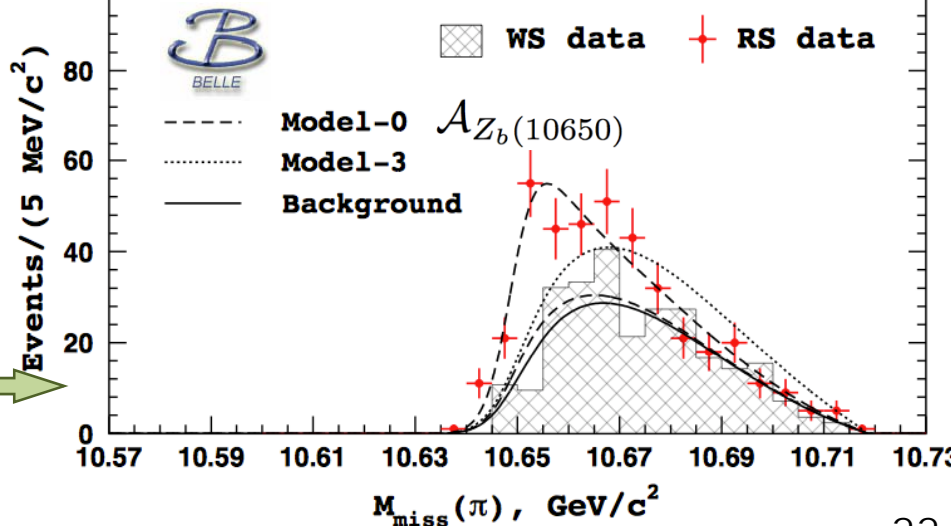
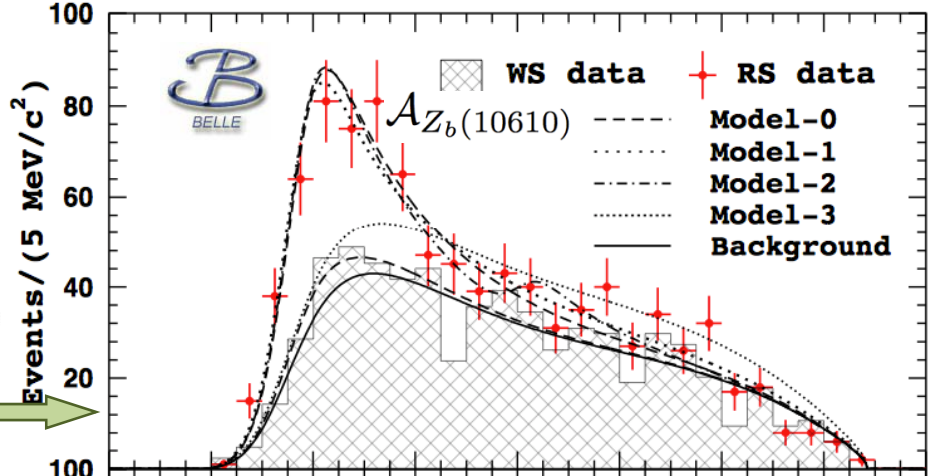
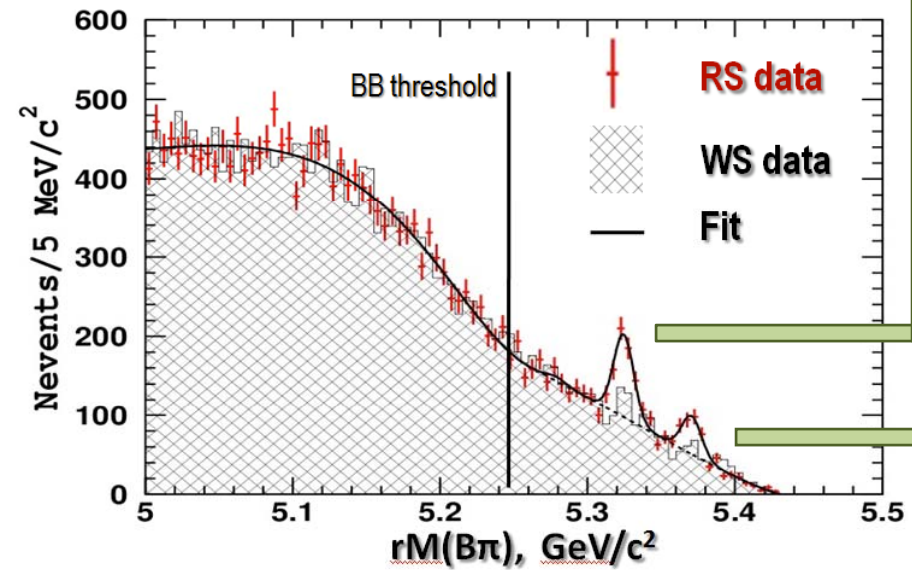
PRL116,212001(2016)



$$S(m) = |\mathcal{A}_{Z_b(10610)} + \mathcal{A}_{Z_b(10650)} + \mathcal{A}_{\text{nr}}|^2$$



$B^{(*)}B^{(*)}\pi + BB\gamma$



Summary on Z_b Decays by Belle

Assuming that Z_b decays are saturated by the $\Upsilon(nS)\pi$, $h_b(mP)\pi$ and $B^{(*)}B^*$ channels, one can calculate a table of relative branching fractions:

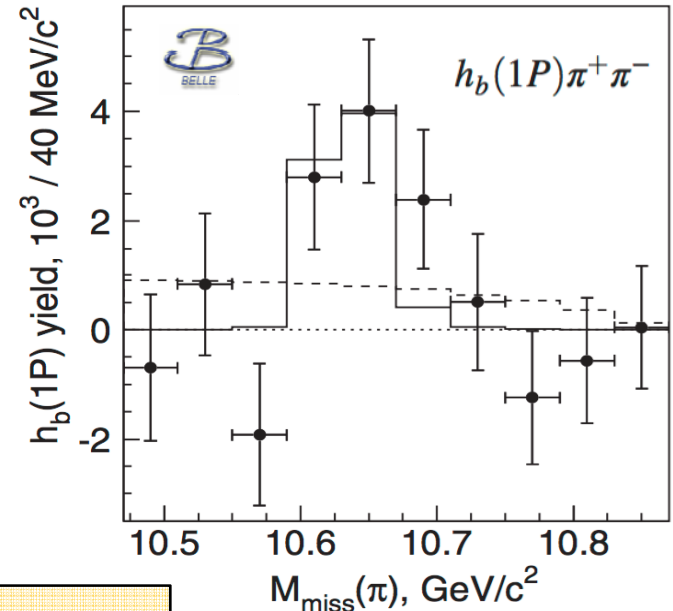
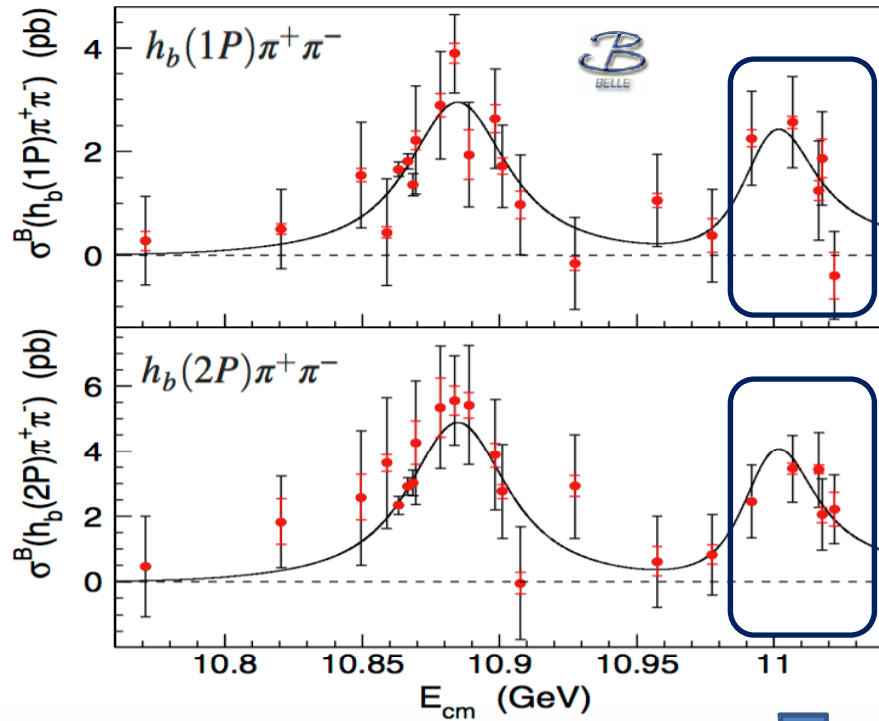
Channel	Fraction, %	
	$Z_b(10610)$	$Z_b(10650)$
$\Upsilon(1S)\pi^+$	$0.60 \pm 0.17 \pm 0.07$	$0.17 \pm 0.06 \pm 0.02$
$\Upsilon(2S)\pi^+$	$4.05 \pm 0.81 \pm 0.58$	$1.38 \pm 0.45 \pm 0.21$
$\Upsilon(3S)\pi^+$	$2.40 \pm 0.58 \pm 0.36$	$1.62 \pm 0.50 \pm 0.24$
$h_b(1P)\pi^+$	$4.26 \pm 1.28 \pm 1.10$	$9.23 \pm 2.88 \pm 2.28$
$h_b(2P)\pi^+$	$6.08 \pm 2.15 \pm 1.63$	$17.0 \pm 3.74 \pm 4.1$
$B^+ \bar{B}^{*0} + \bar{B}^0 B^{*+}$	$82.6 \pm 2.9 \pm 2.3$	—
$B^{*+} \bar{B}^{*0}$	—	$70.6 \pm 4.9 \pm 4.4$

$$\frac{\text{Br}(Z_b(10610)^+ \rightarrow \bar{B}B^*)}{\text{Br}(Z_b(10610)^+ \rightarrow b\bar{b})} = 5.93 + 0.99 / -0.59 + 1.01 / -0.73$$

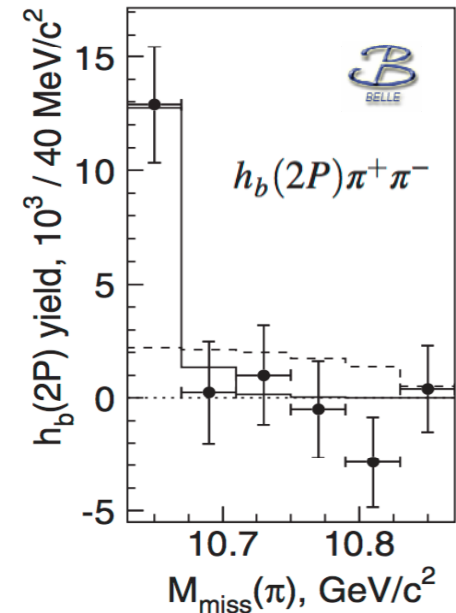
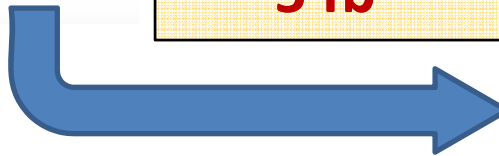
$$\frac{\text{Br}(Z_b(10650)^+ \rightarrow B^* \bar{B}^*)}{\text{Br}(Z_b(10650)^+ \rightarrow b\bar{b})} = 2.80 + 0.69 / -0.40 + 0.54 / -0.36$$

$B^{(*)}B^*$ channels dominate the Z_b decays

Z_b^\pm Production at $\Upsilon(6S)$



**$\Upsilon(6S)$ Data:
5 fb⁻¹**



PRL117,142001(2016)



The two Z_b states can not be separated with current statistics

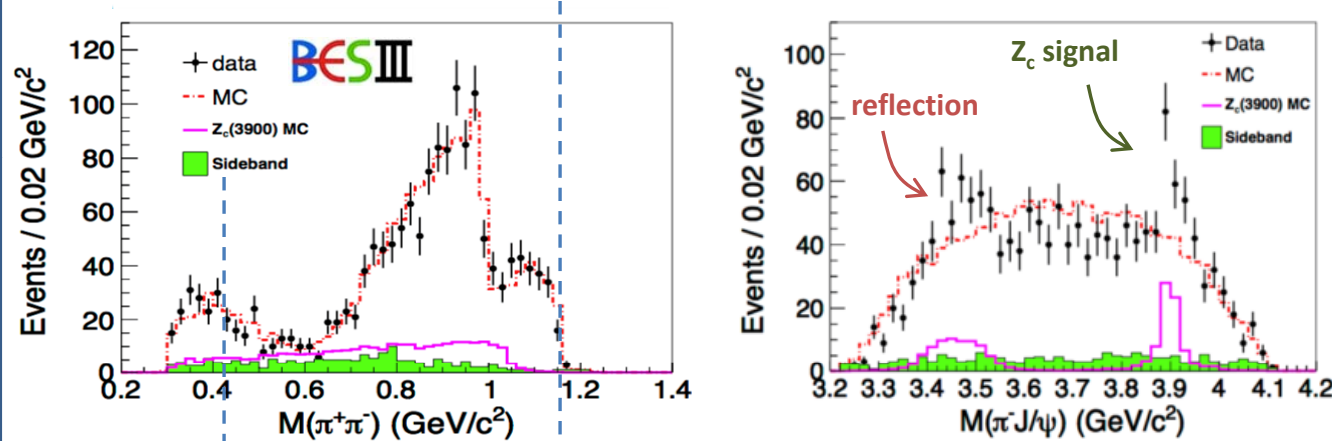
$\Upsilon(6S) \rightarrow h_b(mP)\pi^+\pi^-$ transition is dominated (saturated) by the intermediate Z_b^\pm production

Charm vs. Beauty: I

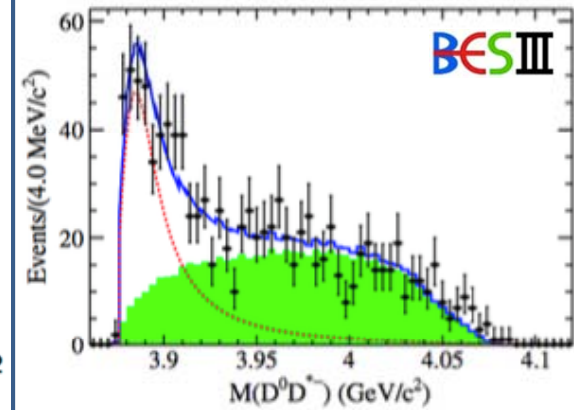
What is in common?

- $\Upsilon(10860)/\Upsilon(4260)$ demonstrates anomalously large coupling to $\Upsilon\pi^+\pi^-/\psi\pi^+\pi^-$ and $h_b\pi^+\pi^-/h_c\pi^+\pi^-$ final states.

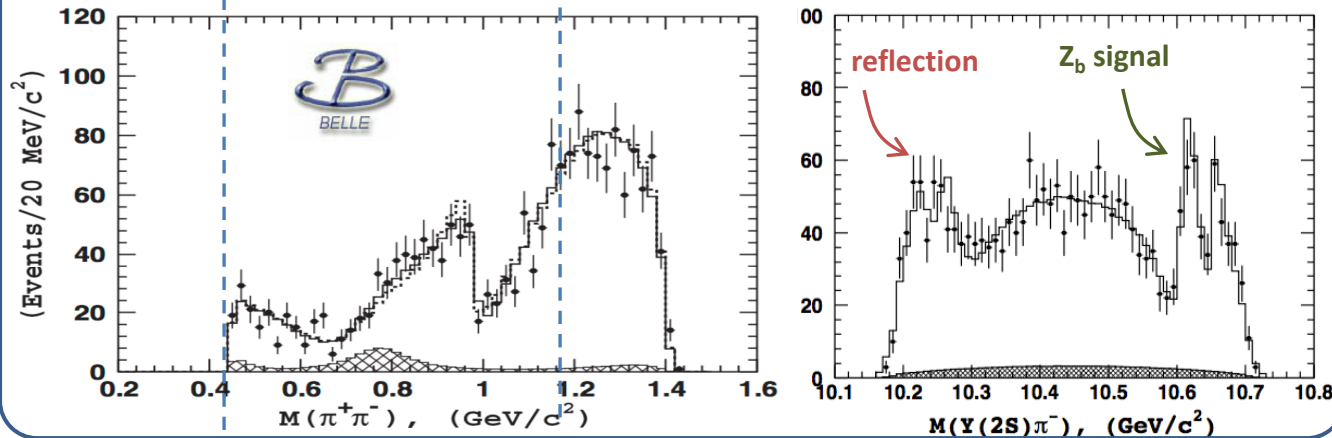
$Z_c(3900)$ is produced in $\Upsilon(4260) \rightarrow Z_c\pi \rightarrow \psi\pi^+\pi^-$



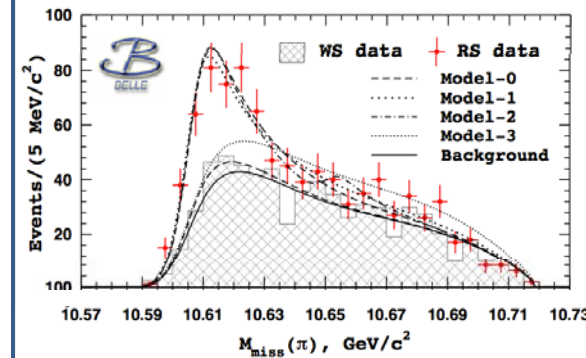
$\rightarrow Z_c\pi \rightarrow DD^*$



$Z_b(10610)$ is produced in $\Upsilon(10860) \rightarrow Z_b\pi \rightarrow \Upsilon(nS)\pi^+\pi^-$



$\rightarrow Z_b\pi \rightarrow BB^*$



Charm vs. Beauty: II

What is different?

- Both $Z_b(10610)$ and $Z_b(10650)$ isotriplets are observed in the $\Upsilon(nS)\pi$, ($n=1,2,3$) and $h_b\pi$ final states.

- $\Upsilon(10860) \rightarrow h_b\pi^+\pi^-$ is saturated by the intermediate two-body $Z_b\pi$ production.

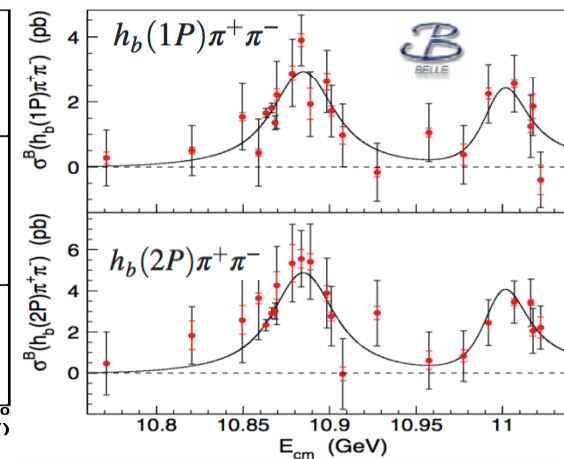
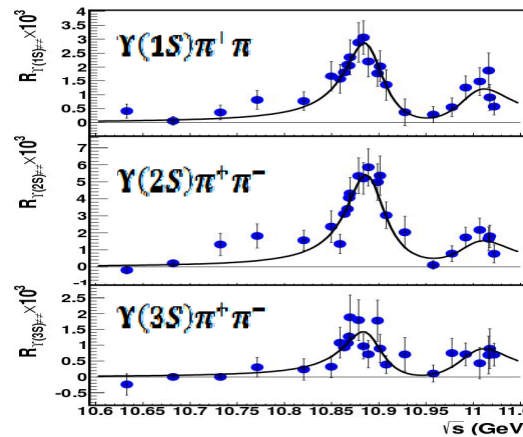
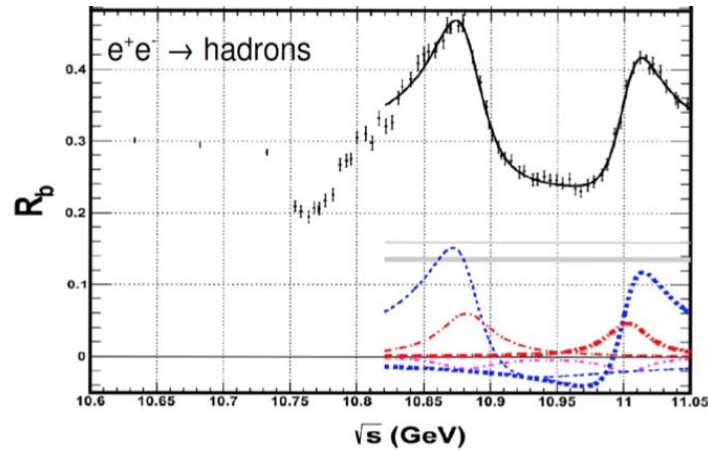
- Only $Z_c(3900)$ is observed in the $J/\psi\pi$ while both $Z_c(3900)$ and $Z_c(4020)$ are observed in the $h_c\pi$ final state. None of them is observed in the $\psi(2S)\pi$ final state (instead, another $Z_c(4430)$ is found).

- Large non $Z_c\pi$ component is observed in the $\Upsilon(4260) \rightarrow h_c\pi^+\pi^-$ amplitude.

Belle II Prospects

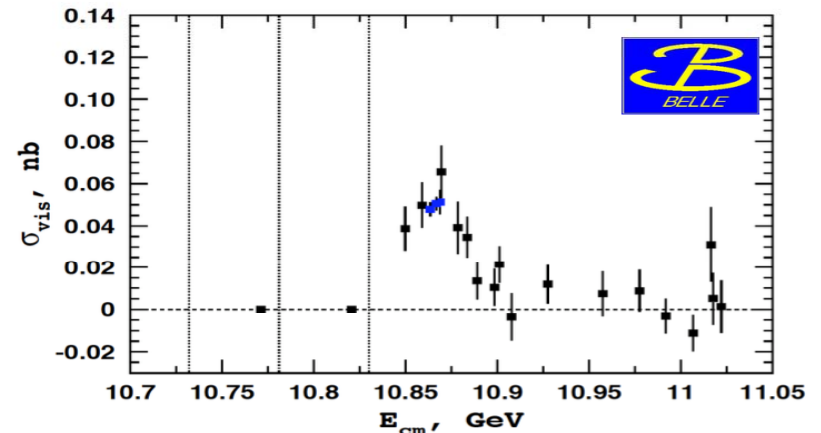
Other Bottomonium-like States @ Belle II

The energy dependence of the inclusive bb cross section in the $\Upsilon(4S) - \Upsilon(6S)$ exhibits very different behavior compared to exclusive channels measured so far



One needs to measure energy dependence for other exclusive cross sections such as $B^{(*)}B^{(*)}$, $B^{(*)}B^{(*)}\pi$, $B_S^{(*)}B_S^{(*)}$ to see how the inclusive cross section is being filled

Belle released preliminary results on the $e^+e^- \rightarrow B_S^{(*)}B_S^{(*)}$, still no significant signal outside the resonance regions...



Other Bottomonium-like States @ Belle II

- Up to 40 1/fb of data might be collected by Belle II during Phase-2 (early 2018) at various c.m. energies including energy range of $\Upsilon(6S)$ and above
- Up to 300 1/fb of data might be collected by Belle II during Phase-3 (late 2018 – ...) at various c.m. energies including energy range of $\Upsilon(6S)$ and above
- With 10 (50) 1/ab of data at (4S) Belle II will be able to perform precision measurements (with an equivalent luminosity of ~ 0.5 (1.5) 1/fb per a 10MeV bin) of Υ states in charm sector via ISR

Other Bottomonium-like States @ Belle II

Molecular states are naturally located near corresponding thresholds:

Particles	Threshold, GeV/ c^2	
$B^{(*)} \bar{B}^{**}$	11.00 – 11.07	arxiv:1610.01102
$B_s^{(*)} \bar{B}_s^{**}$	11.13 – 11.26	
$\Lambda_b \bar{\Lambda}_b$	11.24	
$B^{**} \bar{B}^{**}$	11.44 – 11.49	
$B_s^{**} \bar{B}_s^{**}$	11.48 – 11.68	
$\Lambda_b \bar{\Lambda}_b^{**}$	11.53 – 11.54	
$\Sigma_b^{(*)} \bar{\Sigma}_b^{(*)}$	11.62 – 11.67	
$\Lambda_b^{**} \bar{\Lambda}_b^{**}$	11.82 – 11.84	

- Belle-II maximal energy of **11.24 GeV** covers $B_s^{(*)} \bar{B}_s^{**}$ threshold region.
- Increase to **11.35 GeV** will give information about $\Lambda_b \bar{\Lambda}_b$ threshold region.
- Increase to **11.5 GeV** is crucial to search for partners of Z_b states.

Summary

- ❏ A whole new field of exotic physics discovered in last decade; new information is still coming from both completed (Belle & BaBar) and currently ongoing (BESIII) e+e- experiments
- ❏ Much more data required for a better understanding (BelleII, BESIII)
- ❏ Belle extends searches for exotics to $\Upsilon(6S)$ energy. Start of the Belle II is approaching.
- ❏ Common features between charmonium and bottomonium sectors is gradually emerging. No direct identity though.
- ❏ No clear understanding of the nature of new states yet. Substantial progress in theory achieved recently; more efforts are required from both experimental and theoretical sides.
- ❏ Input from Belle II is crucial to push exotic studies into bottomonium sector. Energy scan up to 11.24 GeV is readily accessible by the SuperKEKB; higher energies are under discussion.

Bottomonium-like States @ Belle II

11.5GeV \longrightarrow

PRD84, 031502 (2011)

$$|Z'_b\rangle = \frac{1}{\sqrt{2}} 0_{bb}^- \otimes 1_{Qq}^- - \frac{1}{\sqrt{2}} 1_{bb}^- \otimes 0_{Qq}^-$$

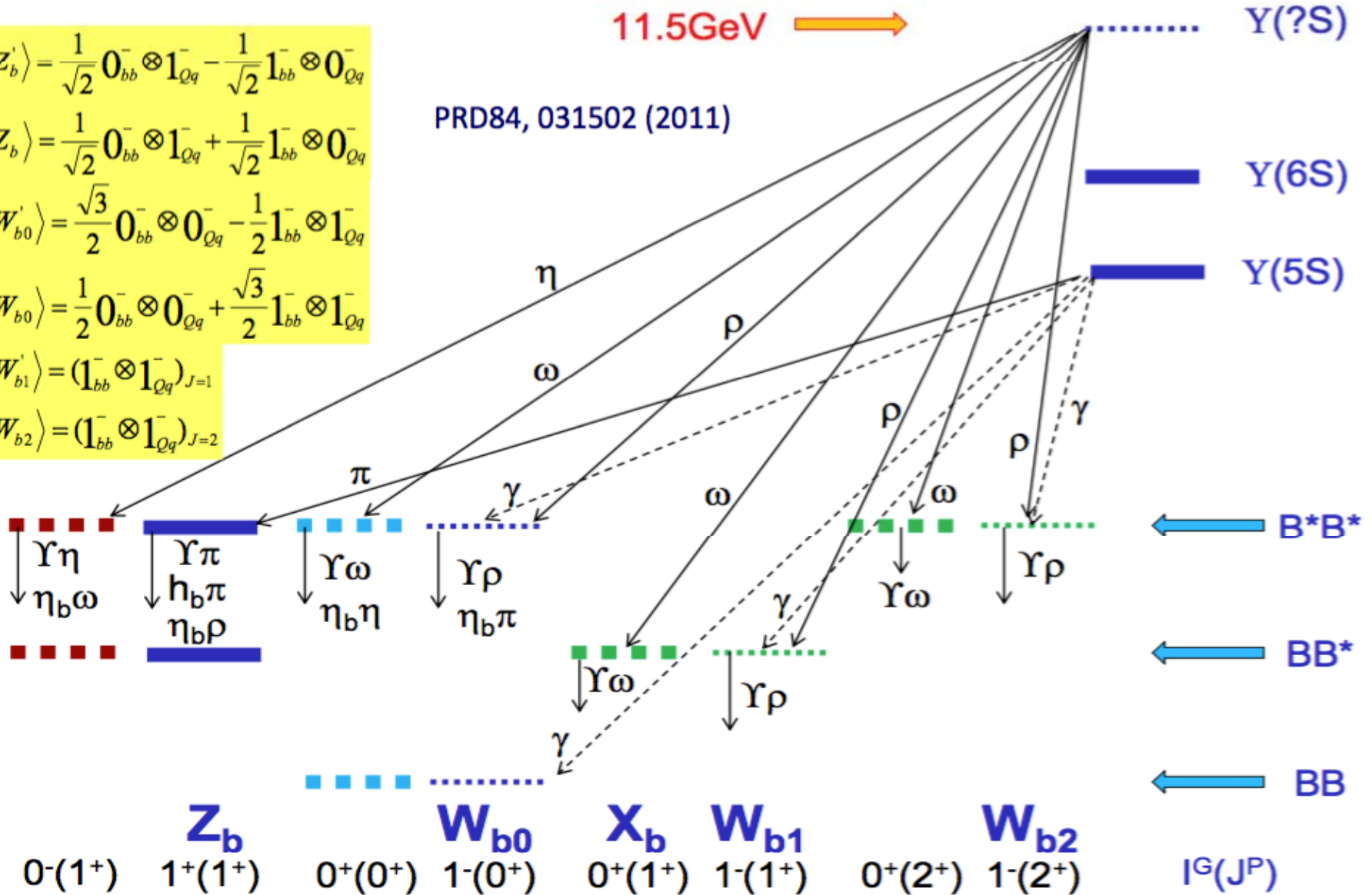
$$|Z_b\rangle = \frac{1}{\sqrt{2}} 0_{bb}^- \otimes 1_{Qq}^- + \frac{1}{\sqrt{2}} 1_{bb}^- \otimes 0_{Qq}^-$$

$$|W'_{b0}\rangle = \frac{\sqrt{3}}{2} 0_{bb}^- \otimes 0_{Qq}^- - \frac{1}{2} 1_{bb}^- \otimes 1_{Qq}^-$$

$$|W_{b0}\rangle = \frac{1}{2} 0_{bb}^- \otimes 0_{Qq}^- + \frac{\sqrt{3}}{2} 1_{bb}^- \otimes 1_{Qq}^-$$

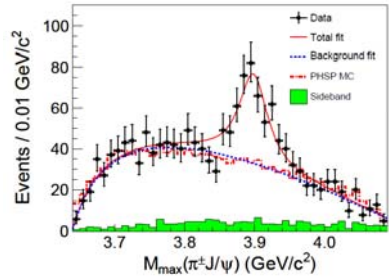
$$|W'_{b1}\rangle = (1_{bb}^- \otimes 1_{Qq}^-)_{J=1}$$

$$|W_{b2}\rangle = (1_{bb}^- \otimes 1_{Qq}^-)_{J=2}$$

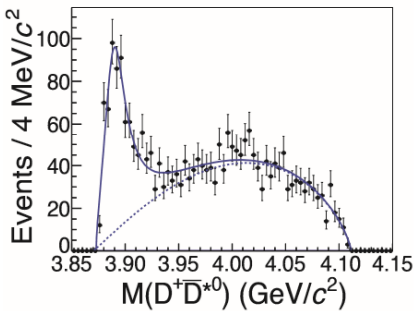


Input from the Belle II is critical!

Summary on Z_c States by BESIII

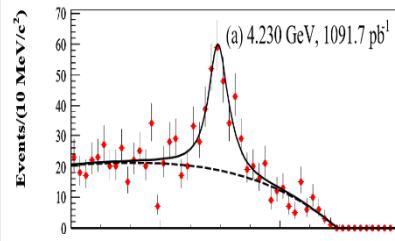


$$e^+e^- \rightarrow \pi^+ \pi^- J/\psi$$

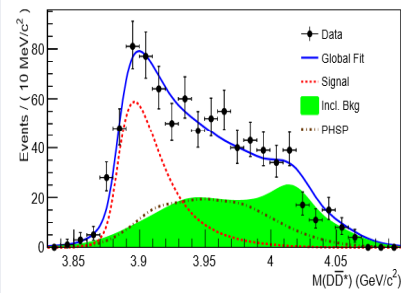


$$e^+e^- \rightarrow \pi^+ (D\bar{D}^*)^-$$

$$\underline{Z_c(3900)^{\pm?}}$$

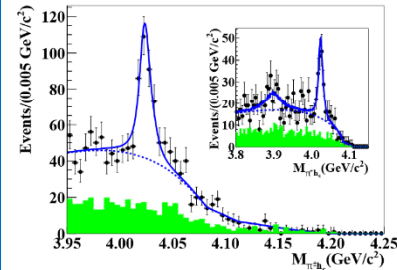


$$e^+e^- \rightarrow \pi^0 \pi^0 J/\psi$$

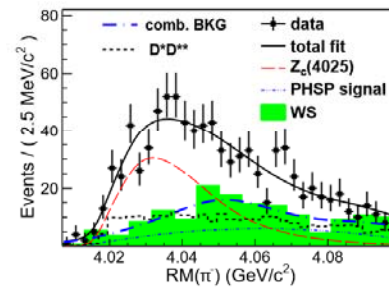


$$e^+e^- \rightarrow \pi^0 (D\bar{D}^*)^0$$

$$\underline{Z_c(3900)^0?}$$

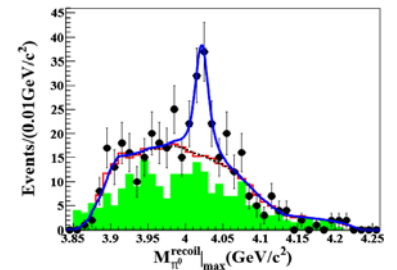


$$e^+e^- \rightarrow \pi^+ \pi^- h_c$$

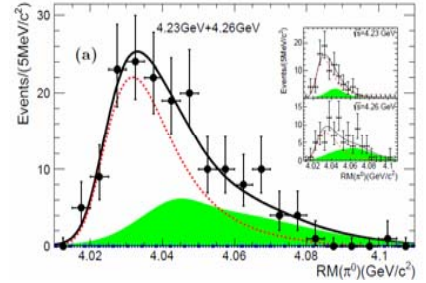


$$e^+e^- \rightarrow \pi^+ (D^* \bar{D}^*)^-$$

$$\underline{Z_c(4020)^{\pm?}}$$



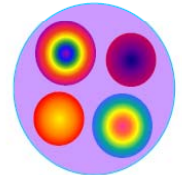
$$e^+e^- \rightarrow \pi^0 \pi^0 h_c$$



$$e^+e^- \rightarrow \pi^0 (D^* \bar{D}^*)^0$$

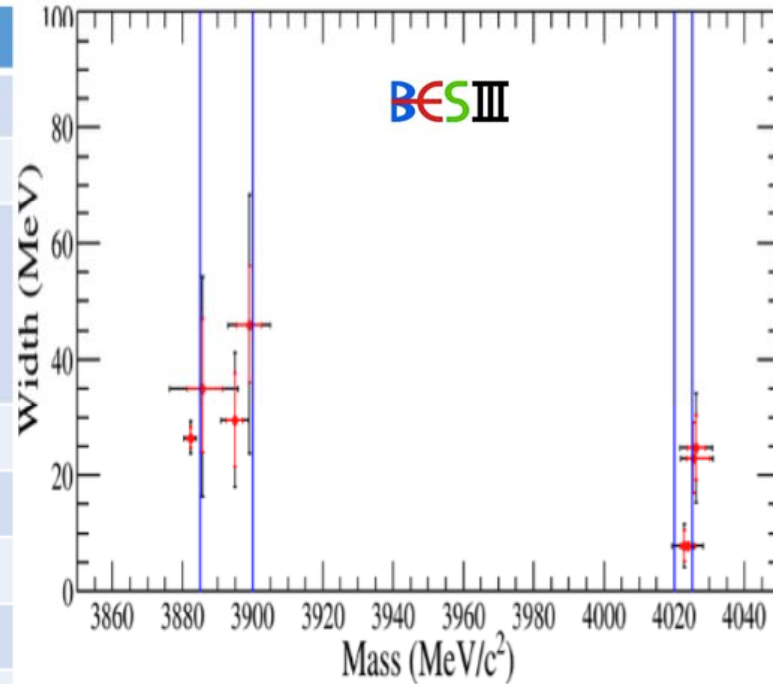
$$\underline{Z_c(4020)^0?}$$

- If these structures are real QCD states, charged Z_c decays into $\pi^{\pm} J/\psi$ ($\pi^{\pm} h_c$) \rightarrow at least four valence quarks to satisfy charge= ± 1 and strong couplings to $c\bar{c}$ components.



Summary on Z_c States by BESIII

State	Mass (MeV/c ²)	Width (MeV)	Process
$Z_c(3900)^\pm$	$3899.0 \pm 3.6 \pm 4.9$	$46 \pm 10 \pm 20$	$e^+e^- \rightarrow \pi^+\pi^- J/\psi$
$Z_c(3900)^0$	$3894.8 \pm 2.3 \pm 3.2$	$29.6 \pm 8.2 \pm 8.2$	$e^+e^- \rightarrow \pi^0\pi^0 J/\psi$
$Z_c(3885)^\pm$ (ST)	$3883.9 \pm 1.5 \pm 4.2$	$24.8 \pm 3.3 \pm 11.0$	$e^+e^- \rightarrow \pi^\mp(D\bar{D}^*)^\pm$
$Z_c(3885)^\pm$ (DT)	$3881.7 \pm 1.6 \pm 1.6$	$26.6 \pm 2.0 \pm 2.1$	
Weighted average	$3882.2 \pm 1.1 \pm 1.5$	$26.5 \pm 1.7 \pm 2.1$	
$Z_c(3885)^0$ (DT)	$3885.7^{+4.3}_{-5.7} \pm 8.4$	$35^{+11}_{-12} \pm 15$	$e^+e^- \rightarrow \pi^0(D\bar{D}^*)^0$
$Z_c(4020)^\pm$	$4022.9 \pm 0.8 \pm 2.7$	$7.9 \pm 2.7 \pm 2.6$	$e^+e^- \rightarrow \pi^+\pi^- h_c$
$Z_c(4020)^0$	$4023.9 \pm 2.2 \pm 3.8$	7.9 (fixed)	$e^+e^- \rightarrow \pi^0\pi^0 h_c$
$Z_c(4025)^\pm$	$4026.3 \pm 2.6 \pm 3.7$	$24.8 \pm 5.6 \pm 7.7$	$e^+e^- \rightarrow \pi^\mp(D^*\bar{D}^*)^\pm$
$Z_c(4025)^0$	$4025.5^{+2.0}_{-4.7} \pm 3.1$	$23.0 \pm 6.0 \pm 1.0$	$e^+e^- \rightarrow \pi^0(D^*\bar{D}^*)^0$



➤ $Z_c(3885)^\pm$ mass is about 2.6σ lower and the width 1.5σ lower than $Z_c(3900)^\pm$ value. If $Z_c(3885) = Z_c(3900)$,

$$\frac{\Gamma(Z_c(3885)^\pm \rightarrow (D\bar{D}^*)^\pm)}{\Gamma(Z_c(3900)^\pm \rightarrow \pi^\pm J/\psi)} = 6.2 \pm 1.1 \pm 2.7, \text{ coupling to } D\bar{D}^* \text{ is larger than to } \pi J/\psi;$$

➤ $Z_c(4020)^\pm$ and $Z_c(4025)^\pm$ mass and width are consistent within 1.5σ . If $Z_c(4020) = Z_c(4025)$,

$$\frac{\Gamma(Z_c(4025)^\pm \rightarrow (D^*\bar{D}^*)^\pm)}{\Gamma(Z_c(4020)^\pm \rightarrow \pi^\pm h_c)} = 12 \pm 5, \text{ coupling to } D^*\bar{D}^* \text{ is larger than to } \pi h_c.$$

Z_b^\pm Spin & Parity

Belle: Full (6D) amplitude analysis of the
 $\Upsilon(5S) \rightarrow Z_b^\pm \pi^\mp \rightarrow \Upsilon(nS) \pi^+ \pi^-$

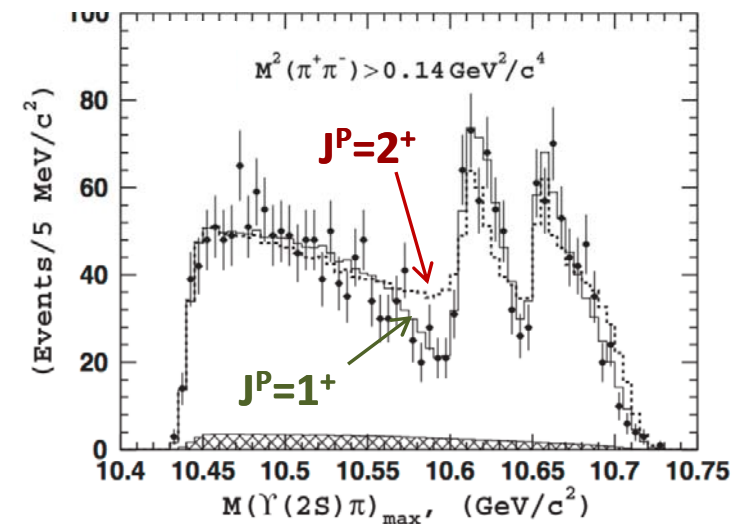
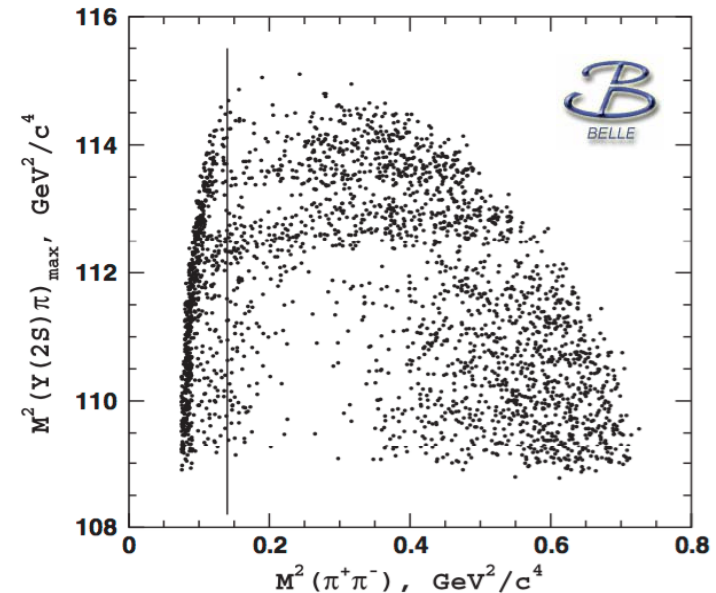
PRD91,072003 (2015)

The $\Upsilon(5S) \rightarrow \Upsilon(nS) \pi^+ \pi^-$ amplitude is parameterized as

$$\mathcal{M}_{\Upsilon\pi\pi} = \mathcal{A}_{Z_1\pi} + \mathcal{A}_{Z_2\pi} + \mathcal{A}_{\Upsilon\sigma} + \mathcal{A}_{\Upsilon f_0} + \mathcal{A}_{\Upsilon f_2} + \mathcal{A}_{\text{NR}}$$

TABLE III. Results of the fit to $\Upsilon(2S) \pi^+ \pi^-$ [$\Upsilon(3S) \pi^+ \pi^-$] events with different J^P values assigned to the $Z_b(10610)$ and $Z_b(10650)$ states. Shown in the table is the difference in \mathcal{L} values for fits to an alternative model and the nominal one.

$Z_b(10610)$	$Z_b(10650)$			
	1^+	1^-	2^+	2^-
1^+	0(0)	60(33)	42(33)	77(63)
1^-	226(47)	264(73)	224(68)	277(106)
2^+	205(33)	235(104)	207(87)	223(128)
2^-	289(99)	319(111)	321(110)	304(125)



Angular analysis strongly favors $J^P=1^+$ assignment

Spin Structure of Z_b

Molecules are not eigenstates of the total $b\bar{b}$ spin

PRD84,054010(2011)

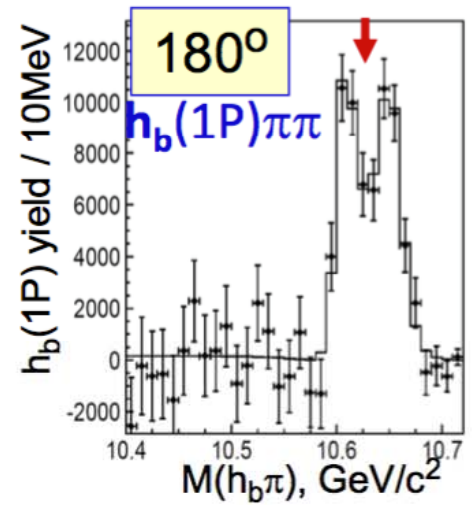
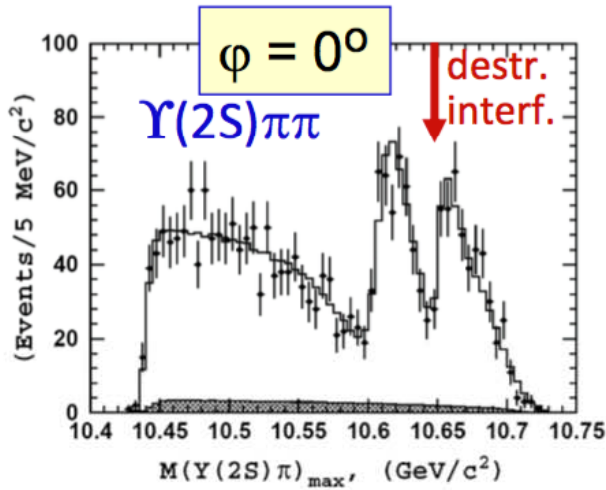
Decomposition \Rightarrow

$$|Z'_b\rangle = (0_{b\bar{b}}^- \otimes 1_{q\bar{q}}^- - 1_{b\bar{b}}^- \otimes 0_{q\bar{q}}^-) / \sqrt{2}$$

$$|Z_b\rangle = (0_{b\bar{b}}^- \otimes 1_{q\bar{q}}^- + 1_{b\bar{b}}^- \otimes 0_{q\bar{q}}^-) / \sqrt{2}$$

\swarrow $h_b(mP)\pi$ \searrow $\Upsilon(nS)\pi$
↓ relative phase

$$BW(s, M_1, \Gamma_1) + ae^{i\phi} BW(s, M_2, \Gamma_2)$$



Assumption of molecular structure allows to explain all properties of Z_b states.