# Numerical conformal bootstrap for the 3d Ising CFT

arXiv:2411.15300 with Chang, Dommes, Erramilli, Homrich, Liu, Mitchell, Poland, Simmons-Duffin

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- Introduction: 3d Ising CFT & numerical bootstrap
- Setting up  $T,\sigma,\epsilon$  in 3d Ising
- Results

3d Ising CFT describes phase transitions associated with breaking of  $\mathbb{Z}_2$  symmetry, e.g. in the 3d Ising model

$$H[\sigma] = J \sum_{\langle ij \rangle} (1 - \sigma_i \sigma_j) + h \sum_i \sigma_i, \quad Z = \sum_{\sigma_i = \pm 1} e^{-H[\sigma]/T}$$

Second-order phase transition at  $T = T_c$  and h = 0. Conformal invariance believed to emerge at large distances.



Why 3d Ising CFT?

- A very common universality class (from  $H_2O$  to QCD?).
- The simplest interacting quantum field theory in d > 2?
- Conformal bootstrap is very effective.
- Good testing ground for new methods.

Conformal bootstrap works directly in the scaling limit.

3d Ising a unitary CFT with the symmetries:

- 1. connected conformal group  $\widetilde{\mathrm{SO}}(2,3)$
- 2. space parity
- 3.  $\mathbb{Z}_2$  global symmetry

Local operators are labelled by scaling dimension  $\Delta$  and  $[j^{\pm},\pm]$ . Only two relevant scalars:

1. 
$$\sigma \in [0^+, -], \quad \Delta_{\sigma} = \frac{1}{2} + \frac{\eta}{2} \approx 0.52$$
  
2.  $\epsilon \in [0^+, +], \quad \Delta_{\epsilon} = 3 - \frac{1}{\nu} \approx 1.41$ 

3d Ising model/CFT has been studied using many methods: Monte-Carlo, High-T expansions, ERG,  $\epsilon$ -expansion, skeleton expansions, fuzzy sphere...

The most precise values of  $\Delta_{\sigma}, \Delta_{\epsilon}$  from non-bootstrap methods come from Monte-Carlo [Hasenbusch'21]

 $\Delta_{\sigma} = 0.518142(20), \quad \Delta_{\epsilon} = 1.41265(13)$ 

Recent fuzzy-sphere regularization seems to offer some advantages [Zhu, Han, Huffman, Hofman, He'22]

Conformal bootstrap constrains properties of CFTs based on crossing symmetry of four-point functions. [Rattazzi, Rychkov, Tonni, Vichi'08]

For  $\langle \sigma \sigma \sigma \sigma \rangle$  one finds



[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi'12]

# 3d Ising & conformal bootstrap

Conformal invariance implies 
$$\left(z\bar{z} = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}, (1-z)(1-\bar{z}) = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}\right)$$

$$\langle \sigma(x_1)\sigma(x_2)\sigma(x_3)\sigma(x_4)\rangle = \frac{g(z,\bar{z})}{x_{12}^{2\Delta\sigma}x_{34}^{2\Delta\sigma}}.$$

Using operator product expansion  $\sigma \times \sigma = \sum_{\mathcal{O}} \lambda_{\sigma\sigma\mathcal{O}} \mathcal{O}$ ,

$$g(z,\bar{z}) = (z\bar{z})^{\Delta_{\sigma}} \sum_{\mathcal{O}} \lambda^2_{\sigma\sigma\mathcal{O}} g_{\Delta_{\mathcal{O}},j_{\mathcal{O}}}(z,\bar{z}).$$

The operators  $\mathcal{O}$  have quantum numbers  $[j^+, +], j \in 2\mathbb{Z}$ .

 $g_{\Delta,j}(z,\bar{z}) = \text{conformal block (computable function)}$ 

# 3d Ising & conformal bootstrap

Permutation invariance under  $x_1 \leftrightarrow x_3$  implies

$$g(z, \bar{z}) = g(1 - z, 1 - \bar{z}).$$

Together with  $j \in 2\mathbb{Z}$  this implies full  $S_4$  permutation symmetry.

$$\sum_{\mathcal{O}} \lambda_{\sigma\sigma\mathcal{O}}^2 g_{\Delta_{\mathcal{O}},j_{\mathcal{O}}}(z,\bar{z}) = \sum_{\mathcal{O}} \lambda_{\sigma\sigma\mathcal{O}}^2 g_{\Delta_{\mathcal{O}},j_{\mathcal{O}}}(1-z,1-\bar{z})$$
$$\sum_{\mathcal{O}} \lambda_{\sigma\sigma\mathcal{O}}^2 F_{\Delta_{\mathcal{O}},j_{\mathcal{O}}}(z,\bar{z}) = 0.$$

At  $z = \overline{z} = \frac{1}{2}$  and odd m + n, using  $\lambda_{\sigma\sigma 1} = 1$  $\partial^m \overline{\partial}^n F_{0,0} + \sum_{\mathcal{O} \neq 1} \lambda_{\sigma\sigma\mathcal{O}}^2 \partial^m \overline{\partial}^n F_{\Delta_{\mathcal{O}},j_{\mathcal{O}}} = 0.$ 

Crucially,  $\lambda^2_{\sigma\sigma\mathcal{O}} \geq 0$  due to unitarity.

$$\partial^m \bar{\partial}^n F_{0,0} + \sum_{\mathcal{O} \neq 1} \lambda^2_{\sigma\sigma\mathcal{O}} \partial^m \bar{\partial}^n F_{\Delta_{\mathcal{O}}, j_{\mathcal{O}}} = 0.$$

Toy model: for odd n,

$$(0-1)^n + \sum_{\Delta} \lambda_{\Delta}^2 (\Delta - 1)^n = 0.$$

Suppose all  $\Delta \geq 3$  in the sum. We prove it is impossible:

1. Take  $-4 \times (n = 1 \text{ equation})$  plus n = 3 equation.

2. Obtain

$$P(0) + \sum_{\Delta} \lambda_{\Delta}^2 P(\Delta) = 0,$$
  
$$P(\Delta) = -4(\Delta - 1) + (\Delta - 1)^3 = (\Delta + 1)(\Delta - 1)(\Delta - 3)$$

3. Contradiction from P(0) = 3 and  $P(\Delta) \ge 0$  for  $\Delta \ge 3$ .

In the real problem

$$\partial^m \bar{\partial}^n F_{0,0} + \sum_{\mathcal{O} \neq 1} \lambda^2_{\sigma\sigma\mathcal{O}} \partial^m \bar{\partial}^n F_{\Delta_{\mathcal{O}}, j_{\mathcal{O}}} = 0.$$

[Rattazzi, Rychkov, Tonni, Vichi'08]

- 1. Take linear combinations with coefficients  $\alpha_{mn}$ .
- 2. Cutoff at  $m + n \leq \Lambda$ . Typical  $\Lambda \lesssim 50$ .
- 3. Search for coefficients  $\alpha$  numerically so as to exclude potential solutions.

# 3d Ising & conformal bootstrap

Studying four-point functions of  $\sigma$  and  $\epsilon$ ,

 $\langle \sigma \sigma \sigma \sigma \rangle, \quad \langle \epsilon \epsilon \epsilon \epsilon \rangle, \quad \langle \sigma \epsilon \sigma \epsilon \rangle,$ 

one finds ( $\Lambda = 19$ )



#### [Kos,Poland,Simmons-Duffin'14]

# 3d Ising & conformal bootstrap

[Simmons-Duffin'15] introduced bootstrap-tuned solver SDPB and [Kos,Poland,Simmons-Duffin,Vichi'16] compute at  $\Lambda=43$ 



State of the art prior to this work.

What's new: add T to  $\sigma,\epsilon$ 

 $\begin{array}{ll} \langle \sigma\sigma\sigma\sigma\rangle, & \langle\epsilon\epsilon\epsilon\epsilon\rangle, & \langle\sigma\epsilon\sigma\epsilon\rangle \\ & \langle T\sigma\sigma\epsilon\rangle, & \langle T\epsilon\epsilon\epsilon\rangle \\ & \langle TT\sigma\sigma\rangle, & \langle TT\epsilon\epsilon\rangle \\ & & \langle TTTT\rangle \end{array}$ 



 $\{T,\sigma,\epsilon\}$  advantages:

- 1. Access to almost all symmetry sectors.  $_{(\text{except } [0^-, -], \ [1^-, -], \ [1^-, +], \ [1^+, +])}$
- 2. Access to more operators, OPE coefficients.
- 3. Improved precision.

 $\{T,\sigma,\epsilon\}$  challenges:

- 1. Many essentially different conformal blocks  $g_{\Delta,j}$  to compute.
- 2. Many more crossing equations to impose.
- 3. A large jump in computational complexity.

We use a fully algorithmic approach based on Zamolodchikov-like recursion relations, schematically

$$g_{\Delta,j}(z,\bar{z}) = g_{\infty,j}(z,\bar{z}) + \sum_{i} \frac{1}{\Delta - \Delta_i} r^{n_i} g_{\Delta_i + n_i, j_i}(z,\bar{z})$$

[Zamolodchikov'87] [PK'17] [Erramilli, Iliesiu, PK'19]

Implemented in blocks\_3d software. Can compute completely general 3d conformal blocks. [Erramilli,Iliesiu,PK,Landry,Poland,Simmons-Duffin'20] The full  $\{T, \sigma, \epsilon\}$  involves many more crossing equations.

Low-level manual implementation is error-prone and not practical.

We use the hyperion family of Haskell packages. Designed to handle very general systems of crossing equations given only a high-level description.

[Simmons-Duffin, PK, Erramilli, Liu, ...]

hyperion also handles the generation and bookkeeping of the conformal blocks, other intermediate data, and the final results.

	fresh build	using pre-built $\langle TTTT\rangle$ blocks	during OPE scan
Block3d	629	79	0
CompositeBlock	10538	4576	0
PartChunk	7926	7926	53
Part	363	363	1
total time	4.8 hours	1.4 hours	13 seconds

build times using 2048 compute cores

The amount of data that needs to be generated required us to develop a custom scheduling algorithm to distribute tasks to allocated compute cores.

This work used SDPB 3.0.0 [Simmons-Duffin, Landry, Dommes, ...]

- 1. Arbitrary-precision semidefinite program solver
- 2. Highly-optimized for numerical bootstrap applications
- 3. Scales well to > 10 compute nodes using MPI

- DiRAC Memory Intensive service Cosma8 at Durham University
- Expanse cluster at the San Diego Supercomputing Center (SDSC)
- CREATE High Performance Cluster at King's College London.
- Resnick High Performance Computing Center, Caltech
- Yale Grace computing cluster

O(10) MCPUhrs

# **Results**: $\Delta_{\sigma}, \Delta_{\epsilon}$

Method	$\Delta_{\sigma}$	$\Delta_{\epsilon}$
Monte-Carlo	0.518142(20)	1.41265(13)
$\{\sigma, e\}@\Lambda = 43$	0.51814 <mark>89(10)</mark>	1.412625(10)
$\{T, \sigma, \epsilon\}@\Lambda = 51$	0.5181488 <mark>04(13)</mark>	1.41262527(16)

 $<sup>\</sup>Lambda=51$  is preliminary



Gap assumptions:  $\Delta_{\text{other scalar}} \geq 3$ ,  $\Delta_{T'} \geq 4$ ,  $\tau_{\text{gap}} \geq 10^{-6}$ .

# **Results: OPE coefficients**

3-point functions between  $\sigma, \epsilon, T$  enter into the crossing equations and can be bounded (not as rigorous as  $\Delta_{\sigma}, \Delta_{\epsilon}$ )

	$T\sigma\epsilon$ (this work)	prev. bootstrap	monte carlo	fuzzy sphere
$\lambda_{\sigma\sigma\epsilon}$	1.051853 <mark>73(11)</mark>	1.05185 <mark>37(41)</mark>	1.05 <mark>1(1)</mark>	1.0539(18)
$\lambda_{\epsilon\epsilon\epsilon}$	1.532443 <mark>04(58)</mark>	1.5324 <mark>35(19)</mark>	1.53 <mark>3(5)</mark>	1.5441(23)
$c_T/c_B$	0.9465386 <mark>75(42)</mark>	0.9465 <mark>34(11)</mark>	0.952(29)	0.9 <mark>55(21)</mark>
$\lambda_{TT\epsilon}$	0.953315 <mark>13(42)</mark>	0.95 <mark>8(7)</mark>		0.9162(73)
$n_B$	0.9334445 <mark>59(75)</mark>	0.93 <mark>3(4)</mark>		
$n_F$	0.013094116(33)	0.014(4)		

$$\langle TT \rangle = c_T \times (\text{standard})$$
$$\langle TTT \rangle = n_B \langle TTT \rangle_B + n_F \langle TTT \rangle_F$$
$$c_T / c_B = n_B + n_F$$

Affine transformation of the  $\Lambda = 34, 43$  islands







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Summary:

- We have been able to study  $\{T,\sigma,\epsilon\}$  system at very high  $\Lambda.$
- High-precision determinations of  $\Delta_{\sigma}, \Delta_{\epsilon}$  and  $\{T, \sigma, \epsilon\}$  three-point functions.
- Analysis of extremal spectra in progress
- Possible thanks to substantially improved numerical algorithms for SDPs, conformal blocks, and software infrastructure.
- Requires O(10 MCPUhrs)

Future:

- Natural extensions to  ${\cal O}(N)$  models and other CFTs in 3d
- Bootstrap of 4d CFTs

# Zoom-out plot



# Zoom-in plot

