

# Lattice QCD calculations for Muon $g-2$

Vera Gülpers

School of Physics and Astronomy  
The University of Edinburgh

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THE UNIVERSITY  
*of* EDINBURGH

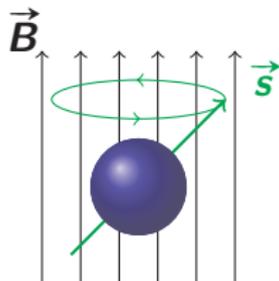
## Magnetic Moment of the Muon

- ▶ magnetic moment  $\vec{\mu}$  of the muon due to its spin  $\vec{s}$  and electric charge  $e$

$$\vec{\mu} = g \frac{e}{2m} \vec{s}$$

torque  $\vec{\tau} = \vec{\mu} \times \vec{B}$

- ▶ gyromagnetic-factor ( $g$ -factor) of the muon

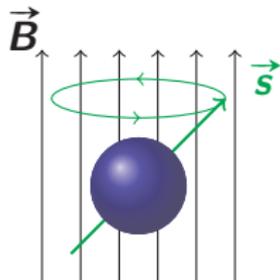


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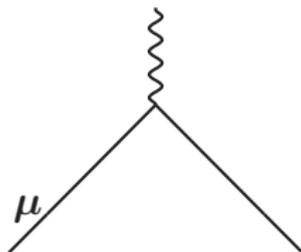
$$\vec{\mu} = g \frac{e}{2m} \vec{s}$$

torque  $\vec{\tau} = \vec{\mu} \times \vec{B}$



- ▶ gyromagnetic-factor ( $g$ -factor) of the muon without quantum effects:

$$g = 2$$

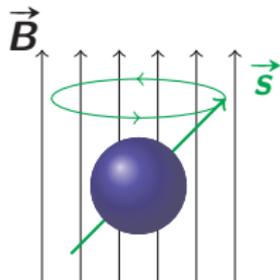


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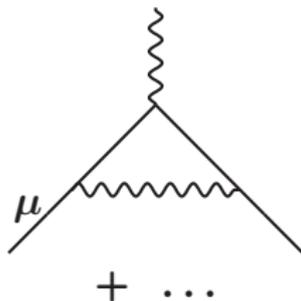


- ▶ gyromagnetic-factor ( $g$ -factor) of the muon with quantum effects:

$$g = 2.00233 \dots$$

anomalous magnetic moment of the muon  
“Muon  $g-2$ ”

$$a_{\mu} = \frac{g - 2}{2}$$



## Muon g-2: Experimental measurement

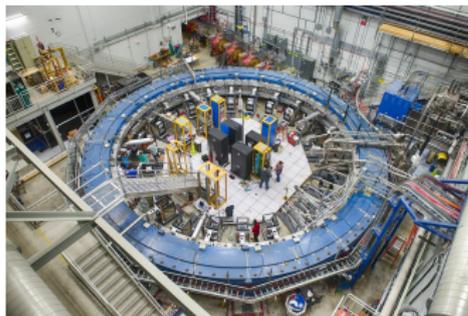
Previous: Muon g-2 @ BNL (2006)

[Phys.Rev. D73, 072003 (2006)]

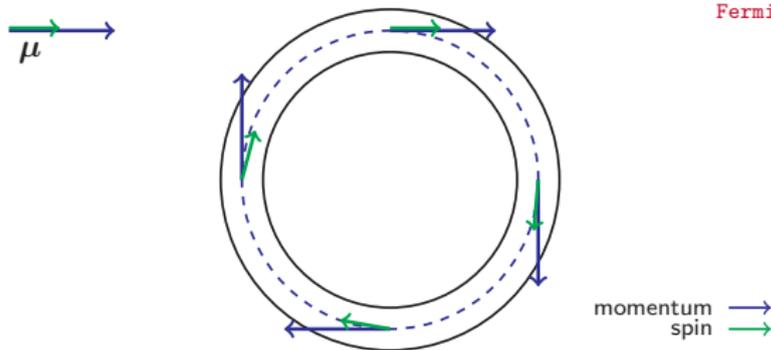
New: Muon g-2 @ FNAL (2021, 2023)

[PhysRevLett.126.141801 (2021)]

measure precession frequency of muons in magnetic field:



[[https://commons.wikimedia.org/wiki/File:Fermilab\\_g-2\\_\(E989\)\\_ring.jpg](https://commons.wikimedia.org/wiki/File:Fermilab_g-2_(E989)_ring.jpg)]



$$\omega_a = a_\mu \frac{eB}{m_\mu}$$

$$a_\mu(\text{exp}) = 11659205.9(2.2) \times 10^{-10}$$

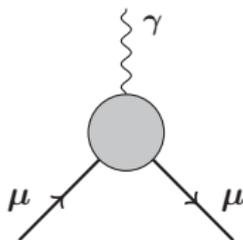
[Phys. Rev. Lett. 131, 161802 (2023)]

# Muon $g-2$ : Standard Model Prediction

## White Paper (2020) of the Muon $g-2$ Theory initiative

[Phys.Rept. 887 (2020) 1-166]

[<https://muon-gm2-theory.illinois.edu/>]

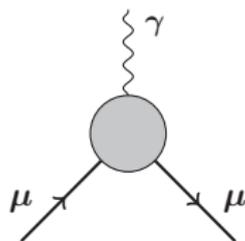


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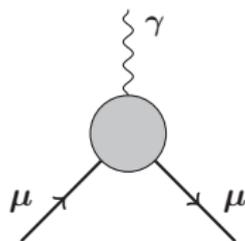


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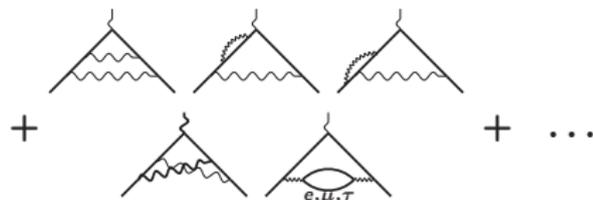
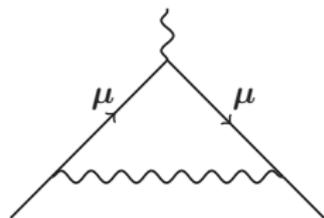
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electro-magnetism

$$11658471.8931(104) \times 10^{-10}$$



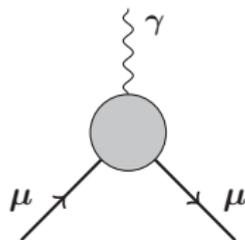
$O(10^4)$  diagrams  
at  $O(\alpha^5)$

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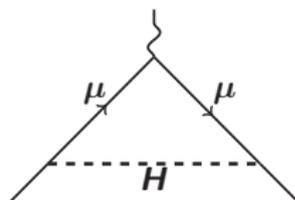
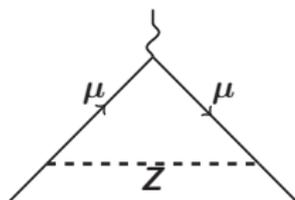
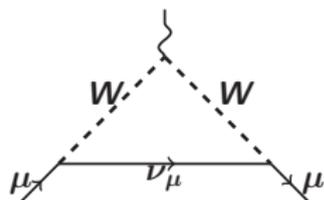


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$$15.36(10) \times 10^{-10}$$

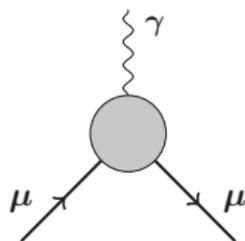


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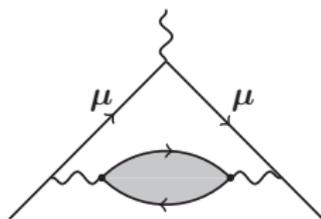
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Hadronic Vacuum Polarisation (HVP)

$$693.1(4.0) \times 10^{-10}$$

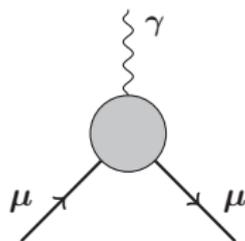


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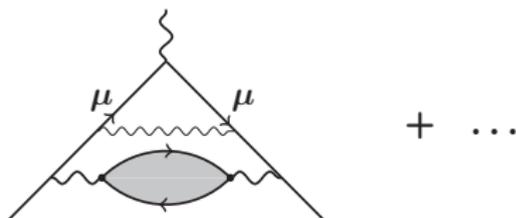
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HVP( $\alpha^3, \alpha^4$ )

$$-8.59(7) \times 10^{-10}$$

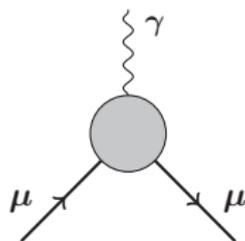


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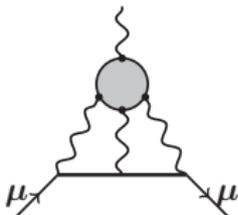
$$693.1(4.0) \times 10^{-10}$$

HVP( $\alpha^3, \alpha^4$ )

$$-8.59(7) \times 10^{-10}$$

Hadronic light-by-light scattering

$$9.2(1.8) \times 10^{-10}$$



## Experiment vs Standard Model prediction

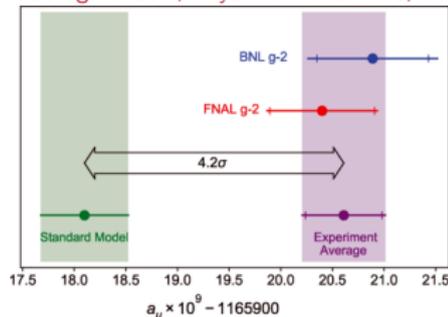
Exp ('23):  $a_\mu = 0.00116592059(22)$

Exp ('21):  $a_\mu = 0.00116592061(41)$

SM:  $a_\mu = 0.00116591810(43)$

► new physics?

Muon  $g-2$  Coll., Phys. Rev. Lett. 126, 141801



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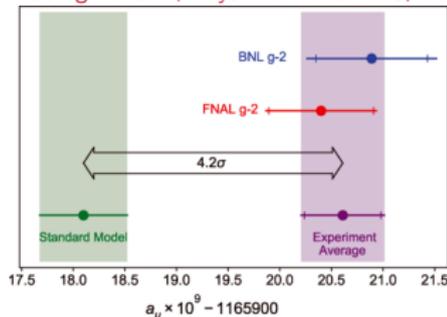
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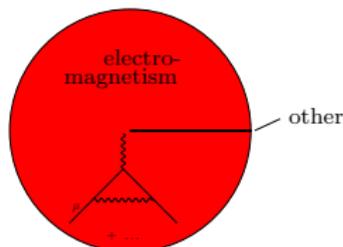
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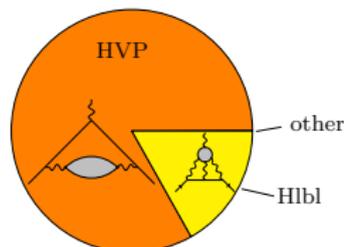
► FNAL to publish final result in 2025 (0.14 ppm), new upcoming experiment @JPARC

► Breakdown of Standard Model Prediction

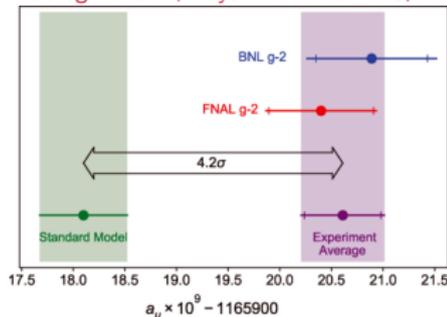
contribution to  $a_\mu$



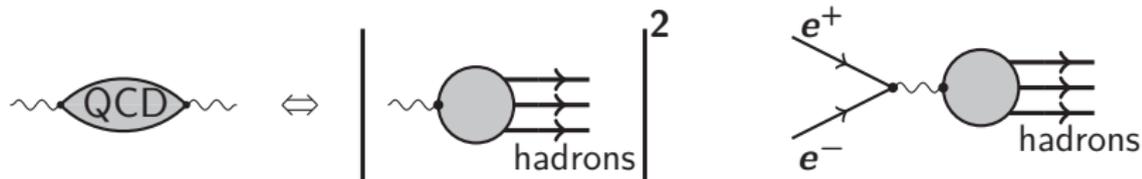
contribution to variance  $\Delta^2 a_\mu$



Muon  $g-2$  Coll., Phys. Rev. Lett. 126, 141801

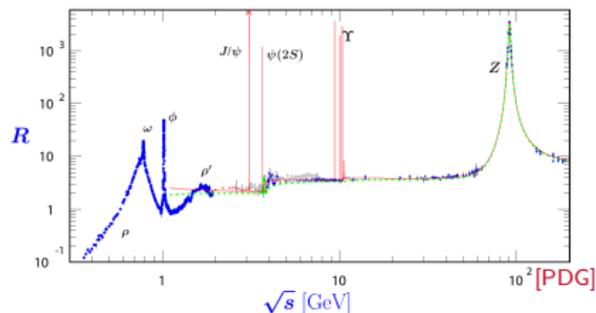


## The HVP from R-ratio



$$R(s) = \frac{\sigma(e^+ e^- \rightarrow \text{hadrons}, s)}{\sigma(e^+ e^- \rightarrow \mu^+ \mu^-, s)}$$

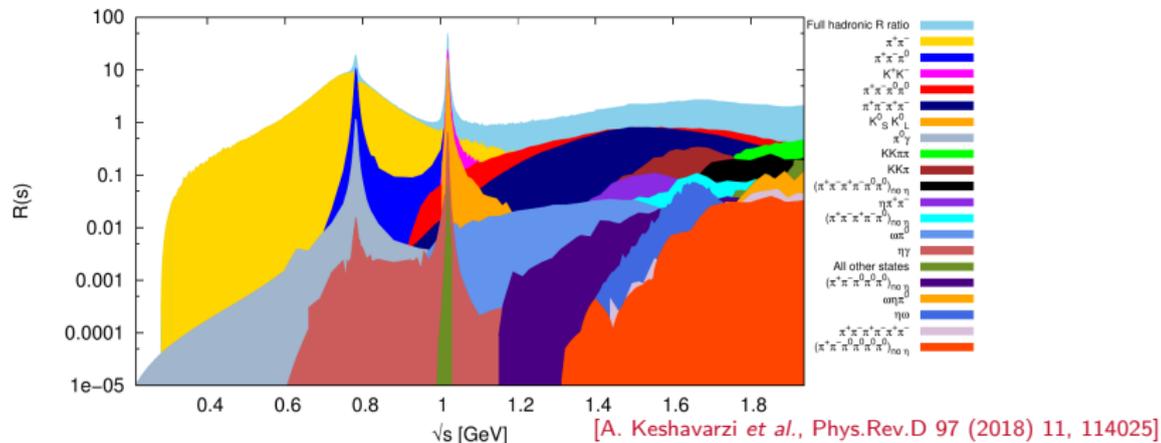
$$a_\mu^{\text{HVP}} = \left( \frac{\alpha m_\mu}{3\pi} \right)^2 \int_{m_\pi^2}^{\infty} ds \frac{R(s)K(s)}{s^2}$$



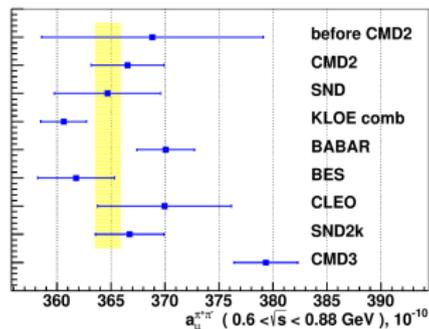
$a_\mu^{\text{HVP}} = 689.46(3.25)$	[Jegerlehner 18]
$a_\mu^{\text{HVP}} = 693.9(4.0)$	[DHMZ 19]
$a_\mu^{\text{HVP}} = 693.37(2.46)$	[KNT 18]
$a_\mu^{\text{HVP}} = 693.1(4.0)$	[white paper]

## Tensions in R-ratio

- various channels contributing to  $R$ -ratio



- 2023 results for  $e^+e^- \rightarrow \pi^+\pi^-(\gamma)$  from CMD-3 disagree with KLOE and BABAR [CMD-3, arXiv:2302.08834]



## Hadronic Vacuum Polarisation (HVP) from the lattice

- ▶ calculate **hadronic** part on the lattice



- ▶ vector two-point function

$$C_{\mu\nu}(t) = \sum_{\vec{x}} \langle J_{\mu}(t, \vec{x}) J_{\nu}(0) \rangle$$

- ▶ electromagnetic current

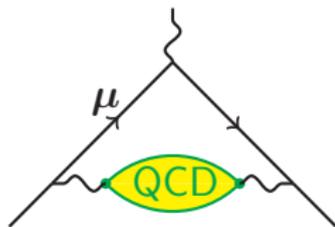
$$J_{\mu} = \frac{2}{3} \bar{u} \gamma_{\mu} u - \frac{1}{3} \bar{d} \gamma_{\mu} d - \frac{1}{3} \bar{s} \gamma_{\mu} s + \dots$$

- ▶  $a_{\mu}$  from  $C(t)$  [T. Blum, Phys.Rev.Lett.**91**, 052001 (2003); Bernecker and Meyer, Eur.Phys.J.**A47**, 148 (2011)]

$$a_{\mu}^{\text{HVP}} = \sum_t w_t C_{ii}(t) \quad \text{with kernel function } w_t$$

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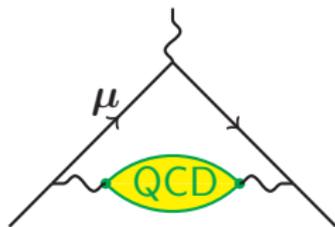
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- ▶ flavour decomposition (isospin symmetric QCD  $u = d = \ell$ )

$$C(t) = \frac{5}{9} C^{\ell}(t) + \frac{1}{9} C^s(t) + \frac{4}{9} C^c(t)$$

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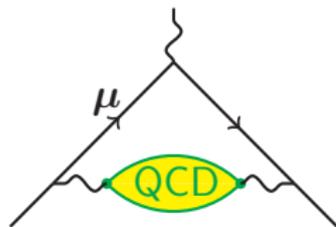
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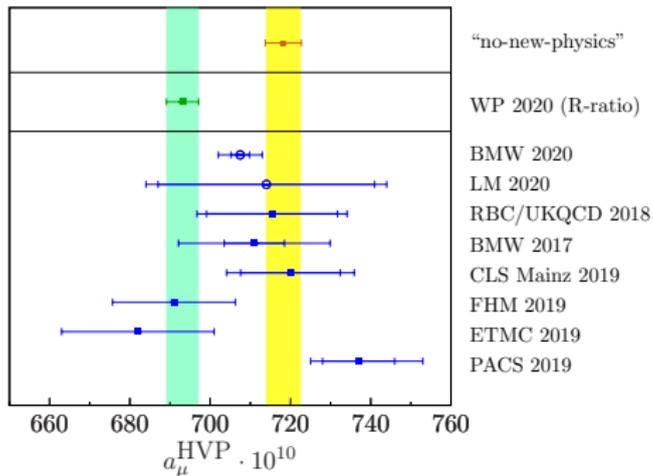
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## Lattice Calculations of HVP (2020 WP status)

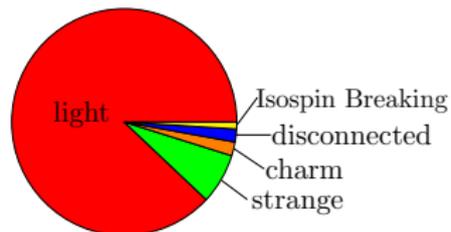
- ▶ White Paper lattice average  $a_{\mu}^{\text{HVP}}(\text{lat}) = 711.6(18.4) \times 10^{-10}$



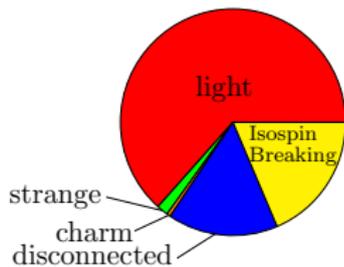
- ▶ BMW 2020

$$a_{\mu}^{\text{HVP}}(\text{BMW}) = 707.5(5.5) \times 10^{-10}$$

contributions to  $a^{\text{HVP}}$

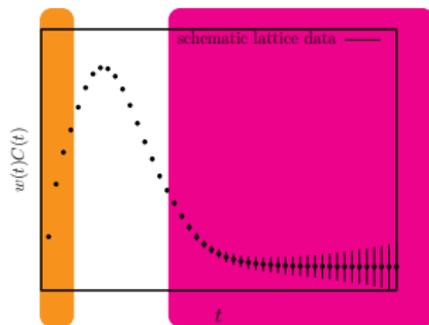


contributions to  $\Delta a_{\mu}^{\text{HVP}}$



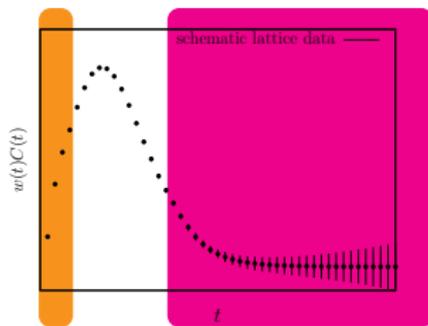
# The window method

- ▶ main challenges:
  - ▶ statistical noise at large  $t$
  - ▶ finite volume effects (largest at large  $t$ )
  - ▶ discretisation effects at small  $t$



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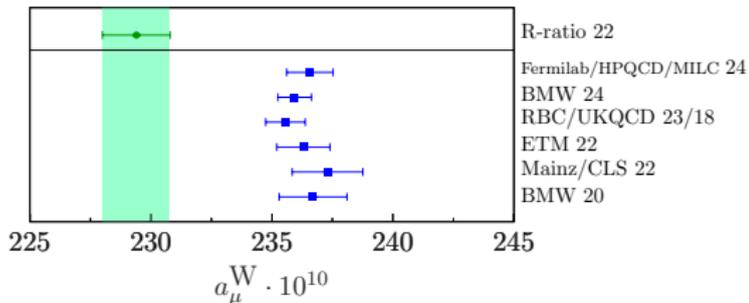
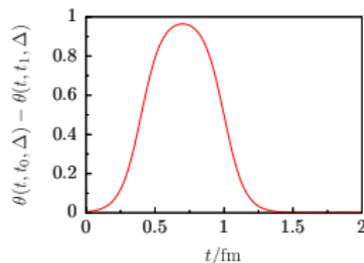
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- ▶ window method

$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

$$a_\mu^{\text{W}} = \sum_t w_t C(t) [\theta(t, t_0, \Delta) - \theta(t, t_1, \Delta)]$$



## The long distance contribution to the HVP

- ▶ tame **statistical noise** by using knowledge of the spectral representation of the vector correlator  $\mathbf{C}(t)$

$$\mathbf{C}(t) = \frac{1}{3} \sum_j \sum_x \langle J_j(t, \mathbf{x}) J_j(0) \rangle = \sum_i \mathbf{A}_i^2 e^{-E_i t} \quad \text{with} \quad \mathbf{A}_i^2 > 0$$

- ▶ lowest states: two pions with back-to-back momentum
- ▶ bounding method [S. Borsanyi *et al.*, (2017)], [T. Blum, VG, *et al.*, (2018)]

$$0 \leq \mathbf{C}(t_c) e^{-E_{t_c}(t-t_c)} \leq \mathbf{C}(t) \leq \mathbf{C}(t_c) e^{-E_0(t-t_c)}$$

- ▶ explicitly reconstruct the lowest  $N$  states
- ▶ improved bounding method [A. Meyer, *et al.*, (2019)]

$$\rightarrow \text{bounding method for } \mathbf{C}(t) - \sum_{i=1}^N \mathbf{A}_i^2 e^{-E_i t}$$

# RBC/UKQCD collaboration

## Boston University

Nobuyuki Matsumoto

## BNL and BNL/RBRC

Peter Boyle

Taku Izubuchi

Christopher Kelly

Shigemi Ohta (KEK)

Amarjit Soni

Masaaki Tomii

Xin-Yu Tuo Shuhei Yamamoto

## University of Cambridge

Nelson Lachini

## CERN

Matteo Di Carlo

Felix Erben

Andreas Jüttner (Southampton)

Tobias Tsang

## Columbia University

Norman Christ

Sarah Fields

Ceran Hu

Yikai Huo

Yong-Chull Jang

Joseph Karpie (JLab)

Erik Lundstrum

Bob Mawhinney

Bigeng Wang (Kentucky)

## University of Connecticut

Tom Blum

Jonas Hildebrand

Luchang Jin

Vaishakhi Moning

Anton Shcherbakov

Douglas Stewart

Joshua Swaim

## DESY Zeuthen

Raoul Hodgson

## Edinburgh University

Luigi Del Debbio

Vera Gülpers

Maxwell T. Hansen

Nils Hermansson-Truedsson

Ryan Hill

Antonin Portelli

Azusa Yamaguchi

## University of Liverpool

Nicolas Garron

## LLNL

Aaron Meyer

## Autonomous University of

### Madrid

Nikolai Husung

## Milano Bicocca

Mattia Bruno

## Nara Women's University

Hiroshi Ohki

## Peking University

Xu Feng

Tian Lin

## University of Regensburg

Andreas Hackl

Danile Knüttel

Christoph Lehner (BNL)

Sebastian Spiegel

## RIKEN CCS

Yasumichi Aoki

## University of Siegen

Matthew Black

Anastasia Boushmelev

Oliver Witzel

## University of Southampton

Bipasha Chakraborty

Ahmed Elgaziari

Jonathan Flynn

Joe McKeon

Rajnandini Mukherjee

Callum Radley-Scott

Chris Sachrajda

## Stony Brook University

Fangcheng He

Sergey Syritsyn (RBRC)

The long-distance window of the hadronic vacuum polarization for the muon  $g - 2$

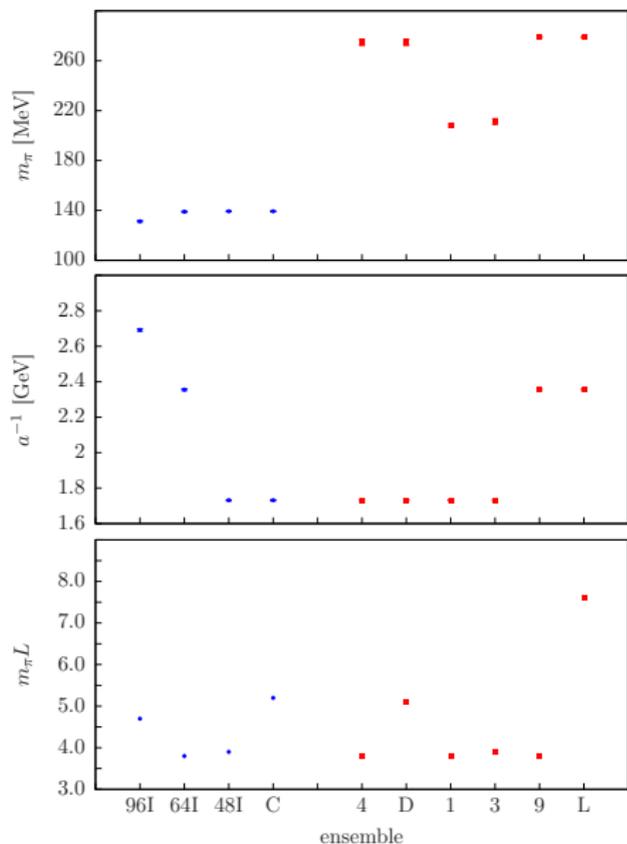
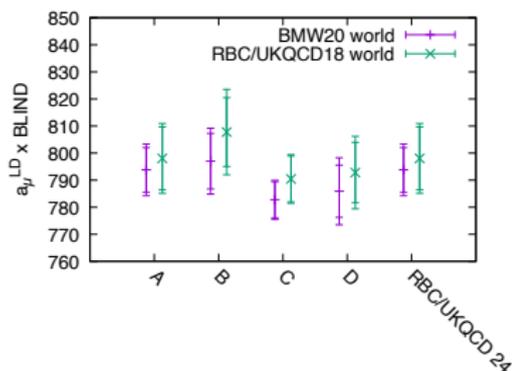
T. Blum,<sup>1</sup> P. A. Boyle,<sup>2,3</sup> M. Bruno,<sup>4,5</sup> B. Chakraborty,<sup>6</sup> F. Erben,<sup>7</sup> V. Gülpers,<sup>3</sup>  
A. Hackl,<sup>8</sup> N. Hermansson-Truedsson,<sup>3</sup> R. C. Hill,<sup>3</sup> T. Izubuchi,<sup>2,9</sup> L. Jin,<sup>1</sup> C. Jung,<sup>2</sup>  
C. Lehner,<sup>8</sup> J. McKeon,<sup>6</sup> A. S. Meyer,<sup>10</sup> M. Tomii,<sup>1,9</sup> J. T. Tsang,<sup>7</sup> and X.-Y. Tuo<sup>2</sup>

(RBC and UKQCD Collaborations)

arXiv:2410.20590

## RBC/UKQCD long distance result – computational details

- ▶  $N_f = 2 + 1$  Domain Wall Fermions
- ▶ four physical point ensembles
- ▶ three lattice spacings
- ▶ several volumes
- ▶ four groups
- ▶ blinding (overall factor)



## RBC/UKQCD long distance result – spectrum

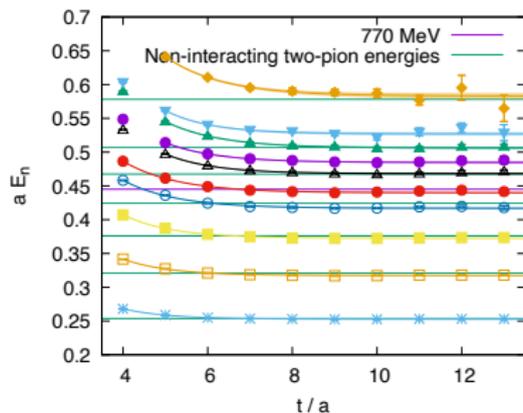
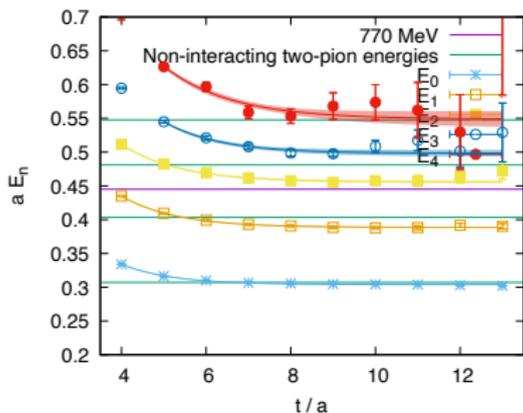
- ▶ spectrum/overlaps of vector correlator using GEVP

[M. Luscher and U. Wolff (1990); B. Blossier *et al* (2009)]

$$C^{ij}(t) = \langle O^i(t) O^j(t)^\dagger \rangle$$

- ▶ vector and two-pion operators with momenta up to  $(2, 0, 0)$ , or  $(2, 2, 0)$  (96l, C), using distillation [M. Peardon *et al*, Phys. Rev. D 80, 054506 (2009)]

- ▶ spectrum ensembles 48l (left) and C (right)

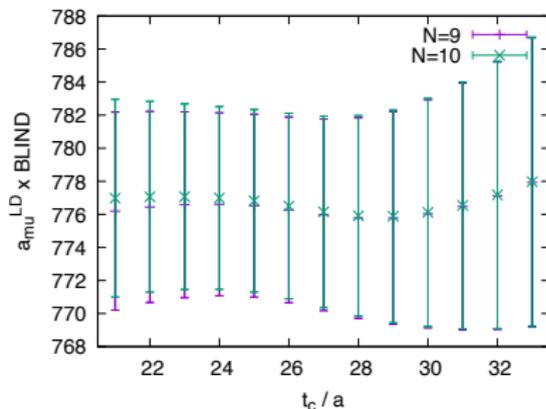
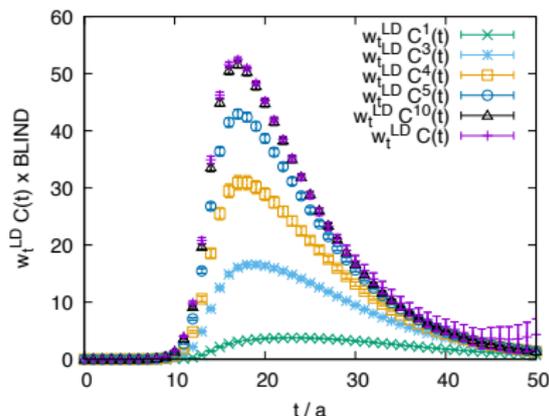


## RBC/UKQCD long distance result – reconstruction

- ▶ long-distance  $\mathbf{a}_\mu$  from vector correlator  $\mathbf{a}_\mu^{\text{LD}} = \sum_{\mathbf{t}} \mathbf{w}_{\mathbf{t}} \mathbf{C}(\mathbf{t}) \theta(\mathbf{t}, \mathbf{t}_1, \Delta)$
- ▶ for large times  $\mathbf{t} > \mathbf{t}_c$ : use reconstruction

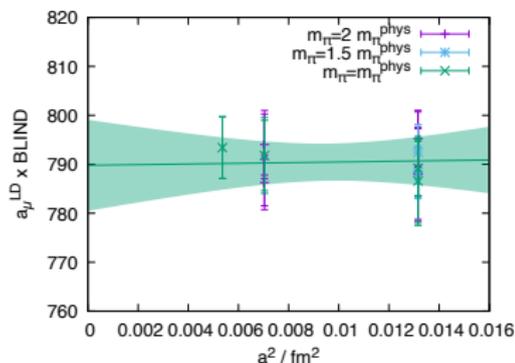
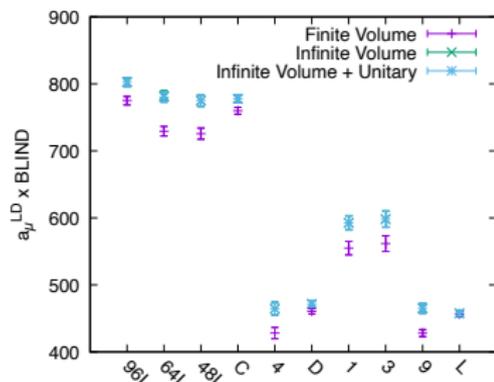
$$\mathbf{C}^N(\mathbf{t}) = \sum_{i=1}^N \mathbf{A}_i^2 e^{-E_i \mathbf{t}}$$

- ▶ Left: reconstruction for 96l
- ▶ Right: dependence of  $\mathbf{a}_\mu$  on  $\mathbf{t}_c$  and number of states  $N$  for 96l



## RBC/UKQCD long distance result – FV and continuum

- ▶ FV corrections using Hansen-Patella formalism [M. T. Hansen and A. Patella (2019), (2020)]
- ▶ ensembles that only differ by volume (48l&C, 4&D, 9&L) agree after FV corrections
- ▶ combined extrapolation: continuum and “physical” isospin-symmetric point
- ▶ continuum extrapolation:  $a^2$ -term
- ▶ two choices of  $Z_V$



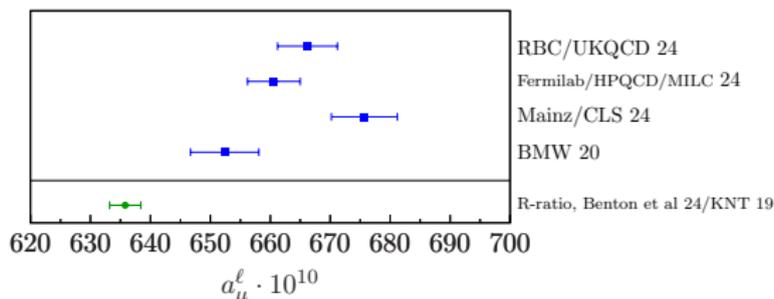
## Comparison

- ▶ Final results (BMW isospin symmetric scheme):

$$a_{\mu}^{\text{LD},\ell} = 411.4(4.3)(2.4) \times 10^{-10}$$

$$a_{\mu}^{\ell} = 666.2(4.3)(2.5) \times 10^{-10}$$

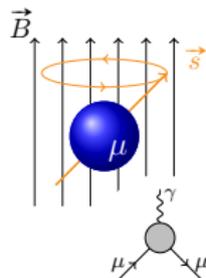
- ▶ systematic error dominated by choices of extrapolation ( $2.2 \times 10^{-10}$ )
- ▶ comparison with other (recent and precise) light-quark results



- ▶ good agreement between lattice calculations
- ▶ **Disclaimer:** not all results shifted to same IB symmetric point

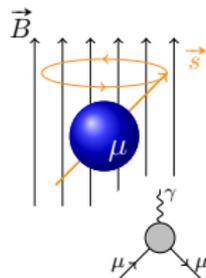
## Summary

- ▶ Muon  $g - 2$  promising quantity for finding new physics  
→ currently:  $4.2\sigma$  tension between experiment and theory
- ▶ Theory error dominated by Hadronic Vacuum Polarisation  
→ R-ratio input: tensions between experiments
- ▶ precision of lattice calculations now competitive with  $R$ -ratio, further improvements needed to match precision of FNAL experiment
- ▶ recent lattice calculations closer to experiment, in tension with  $R$ -ratio
- ▶ Are we about to find new physics with  $a_\mu$ ?



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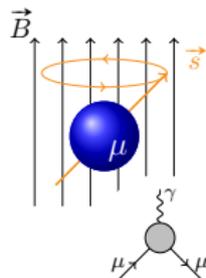


Probably not.



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Probably not.

# Thank you!



Backup

## Isospin symmetric schemes

- ▶ Decomposing into isospin symmetric and isospin breaking

$$\mathbf{X}^\phi = \bar{\mathbf{X}} + \mathbf{X}_{SU(2)} + \mathbf{X}_\gamma \quad \hat{\mathbf{X}} \equiv \bar{\mathbf{X}} + \mathbf{X}_{SU(2)}$$

observable at the physical point  
isospin symmetric  
strong isospin breaking  
QED isospin breaking

- ▶ observable at the physical point unambiguously defined
- ▶ separation into  $\bar{\mathbf{X}}$ ,  $\mathbf{X}_{SU(2)}$  and  $\mathbf{X}_\gamma$  requires separation scheme  
 → define arbitrary values  $\bar{\Pi}$  and  $\hat{\Pi}$  for a set of quantities  $\Pi$
- ▶ RBC/UKQCD long-distance calculation uses two schemes:
  - ▶ BMW20:  $m_\pi = 0.13497\text{GeV}$ ,  $m_{ss^*} = 0.6898\text{GeV}$ ,  $w_0 = 0.17236\text{fm}$
  - ▶ RBC/UKQCD18:  $m_\pi = 0.135\text{GeV}$ ,  $m_K = 0.4957\text{GeV}$ ,  
 $m_\Omega = 1.67225\text{GeV}$

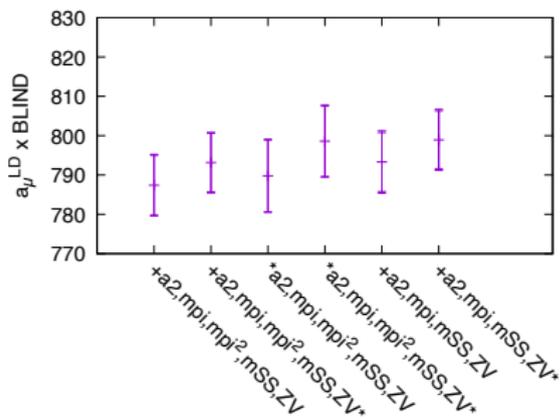
## $a_\mu^{\text{LD}}$ combined extrapolation

- ▶ two fit-ansetze (e.g., for BMW scheme):

$$f_+ = f_0 + f_1 a^2 + f_2 \left( w_0 m_\pi - (w_0 m_\pi)_{\text{phys}} \right) \\ + f_3 \left( w_0 m_\pi - (w_0 m_\pi)_{\text{phys}} \right)^2 + f_4 \left( w_0 m_{\text{SS}^*} - (w_0 m_{\text{SS}^*})_{\text{phys}} \right)$$

$$f_* = f_0 \left( 1 + f_1 a^2 \right) \left( 1 + f_2 \left( w_0 m_\pi - (w_0 m_\pi)_{\text{phys}} \right) \right) \\ + f_3 \left( w_0 m_\pi - (w_0 m_\pi)_{\text{phys}} \right)^2 + f_4 \left( w_0 m_{\text{SS}^*} - (w_0 m_{\text{SS}^*})_{\text{phys}} \right)$$

- ▶ either  $f_3 = 0$  or  $f_3 \neq 0$
- ▶ two different  $Z_V$
- ▶ total:  $\rightarrow$  8 fit choices  
 $\rightarrow$  model average



## RBC/UKQCD window result – ensembles

- ▶ Möbius Domain Wall Fermions
- ▶ three lattice spacings at the physical point with  $N_f = 2 + 1$

ID	$a^{-1}/\text{GeV}$	$L^3 \times T \times L_s$	$m_\pi/\text{MeV}$	$m_K/\text{MeV}$
48I	<b>1.7312(28)</b>	<b><math>48^3 \times 96 \times 24</math></b>	<b>139.32(30)</b>	<b>499.44(88)</b>
64I	<b>2.3549(49)</b>	<b><math>64^3 \times 128 \times 12</math></b>	<b>138.98(43)</b>	<b>507.5(1.5)</b>
96I	<b>2.6920(67)</b>	<b><math>96^3 \times 192 \times 12</math></b>	<b>131.29(66)</b>	<b>484.5(2.3)</b>

- ▶ additional “helper” ensembles at heavier masses

ID	$a^{-1}/\text{GeV}$	$N_f$	$L^3 \times T \times L_s$	$m_\pi/\text{MeV}$	$m_K/\text{MeV}$	$m_{D_s}/\text{GeV}$
1	<b>1.7310(35)</b>	2+1	<b><math>32^3 \times 64 \times 24</math></b>	<b>208.1(1.1)</b>	<b>514.0(1.8)</b>	–
2	<b>1.7257(74)</b>	2+1	<b><math>24^3 \times 48 \times 32</math></b>	<b>285.4(2.9)</b>	<b>537.8(4.6)</b>	–
3	<b>1.7306(46)</b>	2+1	<b><math>32^3 \times 64 \times 24</math></b>	<b>211.3(2.3)</b>	<b>603.8(6.1)</b>	–
4	<b>1.7400(73)</b>	2+1	<b><math>24^3 \times 48 \times 24</math></b>	<b>274.8(2.5)</b>	<b>530.1(3.1)</b>	–
5	<b>1.7498(73)</b>	2+1+1	<b><math>24^3 \times 48 \times 24</math></b>	<b>279.8(3.5)</b>	<b>539.1(5.3)</b>	<b>1.9902(69)</b>
7	<b>1.7566(81)</b>	2+1+1	<b><math>24^3 \times 48 \times 24</math></b>	<b>272.5(5.9)</b>	<b>523(10)</b>	<b>1.3882(57)</b>
A	<b>1.7556(83)</b>	2+1	<b><math>24^3 \times 48 \times 8</math></b>	<b>307.4(3.5)</b>	<b>557.3(5.7)</b>	–

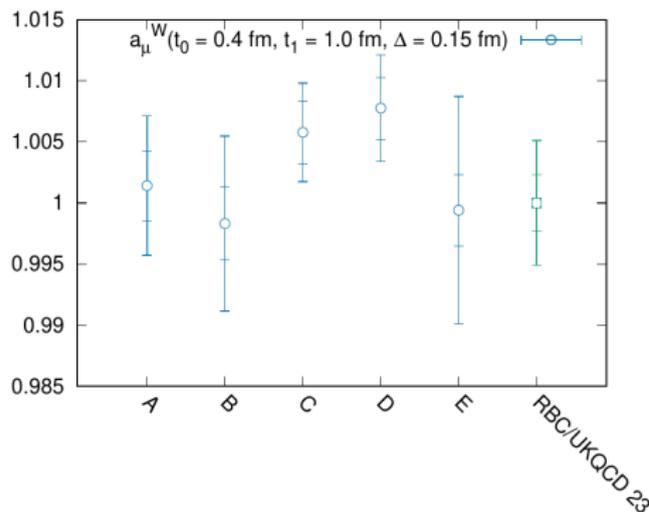
## RBC/UKQCD window result – blinding

- ▶ blinded analysis, to avoid bias towards other results
- ▶ five different analysis groups
- ▶ vector two-point function blinded

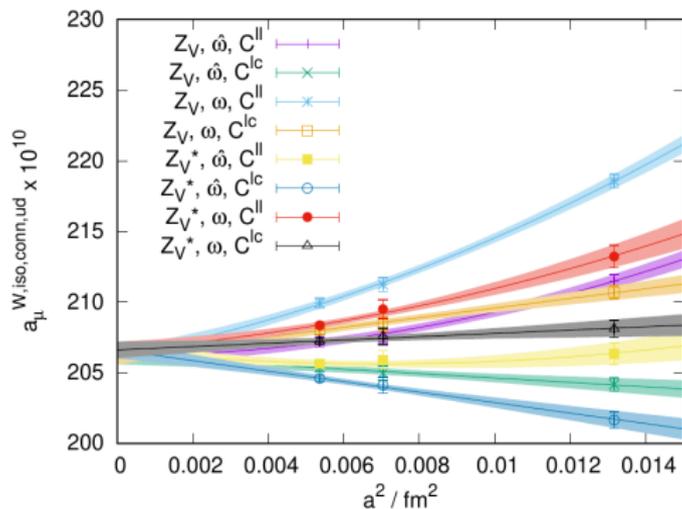
$$C_b(t) = (b_0 + b_1 a^2 + b_2 a^4) C_0(t)$$

with random coefficients  $b_0$ ,  $b_1$  and  $b_2$  different for each group

- ▶ relative unblinding



## RBC/UKQCD window result – continuum extrapolation



- ▶ isospin symmetric light-quark connected window

$$a_{\mu}^{W, \text{ iso, conn, ud}} = 206.36(44)(43) \times 10^{-10}$$

- ▶ total window (using RBC/UKQCD 2018 results for other flavours)

$$a_{\mu}^W = 235.56(65)(50) \times 10^{-10}$$