

## 1 Problem 1 - new physics bound from the neutrino masses

### 1.1 Problem

Consider the Weinberg operator

$$\frac{1}{\Lambda} \left( \epsilon^{\alpha\beta} H^i H^j L_\alpha^i L_\beta^j \right) \quad (1)$$

where  $H$  is the Higgs field,  $L$  is the left-handed lepton field,  $i$  and  $j$  are  $SU(2)$  indices and  $\alpha$  and  $\beta$  are spinor indices. Derive the bound on the scale of new physics  $\Lambda$  to ensure that the neutrino mass is below 0.1 eV.

### 1.2 Solution

Below electroweak symmetry breaking scale we can replace  $H \rightarrow v + h$ , where  $v \approx 246$  GeV is the vacuum expectation value of the Higgs and  $h$  is the Higgs boson field. The mass term will arise then both  $H$  are replaced by  $v$ . We can then read off the neutrino mass from the Lagrangian and we see that it will be proportional to

$$m_\mu \propto \frac{v^2}{\Lambda}. \quad (2)$$

Solving for  $\Lambda$  gives

$$\Lambda \gtrsim 10^{15} \text{ GeV}. \quad (3)$$

## 2 Problem 2 - proton decay

### 2.1 Problem

Consider a baryon number violating operator

$$\frac{1}{\Lambda^2} \epsilon^{IJK} \epsilon^{ij} (Q_{Ii}^T C \sigma^\mu u_J) (L_j^T C \sigma_\mu d_K) \quad (4)$$

Using the constraint on the proton lifetime  $\tau \gtrsim 10^{34}$  years, derive a bound on the new physics scale  $\Lambda$ .

### 2.2 Solution

We can see that the single insertion of the operator will induce proton decay by converting  $uud$  to  $\bar{d}d$  and  $e^+$ . The amplitude is then proportional to  $1/\Lambda^2$  and the decay width will be proportional to amplitude squared giving:

$$\Gamma(p \rightarrow e^+ \pi^0) \propto \frac{m_p^5}{\Lambda^4} \quad (5)$$

where the proton mass,  $m_p$ , which is the characteristic energy scale of this process, was inserted on dimensional grounds (decay width has dimension 1). Proton lifetime is then the inverse of its decay width:

$$\tau_p = \frac{1}{\Gamma(p \rightarrow X)} = \frac{1}{\Gamma(p \rightarrow e^+ \pi^0)} = \frac{\Lambda^4}{m_p^5} \quad (6)$$

The rest is a matter of substitution:

$$1 = \hbar = 6.582 \times 10^{-25} \text{ GeVs} \quad (7)$$

$$m_p \approx 1 \text{ GeV} \quad (8)$$

$$1 \text{ year} \approx \pi \times 10^7 \text{ s} \quad (9)$$

This results in  $\Lambda \gtrsim 10^{16} \text{ GeV}$ .

### 3 Problem 3 - electric dipole moment of the electron

#### 3.1 Problem

Consider the operator in the Lagrangian:

$$\frac{1}{\Lambda^2} \bar{L}^i Y_e \sigma^{\mu\nu} e_R (\sigma_{ij}^a) H^j W_{\mu\nu}^a \quad (10)$$

where the indices  $i$  and  $j$  are SU(2) indices in the fundamental representation and  $a$  - in the adjoint representation and  $Y_e$  is the Yukawa coupling matrix of the electron.

Using the experimental result for the electron's electric dipole moment

$$d_e \lesssim 10^{-28} \text{ ecm} \quad (11)$$

obtain the constraint on the scale  $\Lambda$ .

#### 3.2 Solution

The electric dipole moment is defined from the effective Lagrangian:

$$\mathcal{L}_{eff} = -\frac{i}{2} d_e \bar{\psi} \gamma^5 \sigma^{\mu\nu} F_{\mu\nu} \psi \quad (12)$$

and it can be read off from our effective operator after replacing the Higgs field with its vev. We can then see that

$$d_e \approx \frac{(Y_e)_{11} v}{\Lambda^2} \quad (13)$$

Recalling that  $(Y_e)_{11} \sim m_e/v$  and using the values from the previous problem together with

$$1 = c \approx 3 \times 10^{10} \text{cms}^{-1} \quad (14)$$

$$m_e \approx 0.5 \text{MeV} \quad (15)$$

we end up with

$$\Lambda \approx 10^5 \text{GeV} \quad (16)$$

## 4 Problem 4 - $K$ - $\bar{K}$ mixing

### 4.1 Problem

Find the dependence on the Yukawa couplings of  $K^0$ - $\bar{K}^0$  mixing amplitude.

### 4.2 Solution

The leading order diagram has 4 Yukawa insertions as shown in figure. This is necessary because:

- In flavour basis,  $W$  interactions are diagonal in flavour basis
- $W$  couples only to left-handed fields
- $H$  couples to one left-handed and one right-handed field

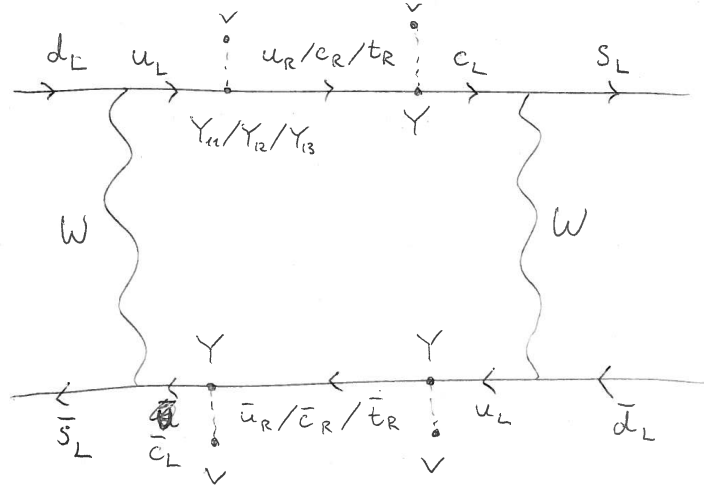


Figure 1: Leading Standard Model contribution to  $K$  $\bar{K}$  mixing

The remaining loop integral, taking large  $M_W$  limit so the  $W$ -propagators get contracted to a point, will have 6 fermion propagators and is IR divergent in the massless quark limit, behaving as:

$$\int \frac{d^4 p}{p^6} \sim \frac{1}{m_c^2}. \quad (17)$$

Furthermore, we choose the basis of the Yukawa couplings to be related to the diagonal (mass) basis by

$$Y_d = Y_d^{\text{diag}} \quad (18)$$

$$Y_u = V_{CKM}^\dagger Y_u^{\text{diag}} \quad (19)$$

Adding the  $1/(16\pi^2)$  loop factor,  $g^4$  for  $W$  vertices,  $1/m_W^4$ , replacing  $Y_u^{\text{diag}} = M/v$  and recalling that  $(V_{CKM})_{11} \approx (V_{CKM})_{22} \approx 1$  and  $(V_{CKM})_{12} \approx \lambda \approx 0.2$  we get the following order-of-magnitude estimate:

$$\sim \frac{g^4 \lambda^2 m_c^2}{16\pi^2 m_W^4} = \frac{G_F^2}{2\pi^2} m_c^2 \lambda^2 \quad (20)$$

Using kaon decay constant  $f_K$  to compensate for the dimensions we find

$$\frac{\Delta m_K}{m_K} \sim \frac{G_F^2}{2\pi^2} m_c^2 \lambda^2 f_K^2 \approx 10^{-15} \quad (21)$$