

Gradient-flow coupling in QED₃ and QED₄

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The gradient-flow coupling
of three- and four-dimensional QED

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Abstract

We evaluate the QED coupling in the gradient-flow scheme in three and four space-time dimensions. Our general result applies to any theory with a U(1) gauge field coupled to arbitrary other fields via arbitrary interactions. As an example, we consider QED with n_f flavors in three and four space-time dimensions and evaluate the corresponding β functions. In four dimensions, we find that the perturbative expansion of the β function behaves much better than the corresponding expression in QCD. In three dimensions, we recover both the ultraviolet as well as the infrared fixed points of the QED coupling in the large- n_f limit.

Gradient-flow scheme

- See GF as regularization scheme:

- ➔ No flowed gluon renormalization
- ➔ No operator mixing

[Lüscher '10; Lüscher, Weisz '11]

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- Current status:

perturbation theory

$$\alpha_{\text{GF}}(t) = \alpha_s \left[1 + c_1(t) \alpha_s + c_2(t) \alpha_s^2 + c_3(t) \alpha_s^3 + \dots \right]$$

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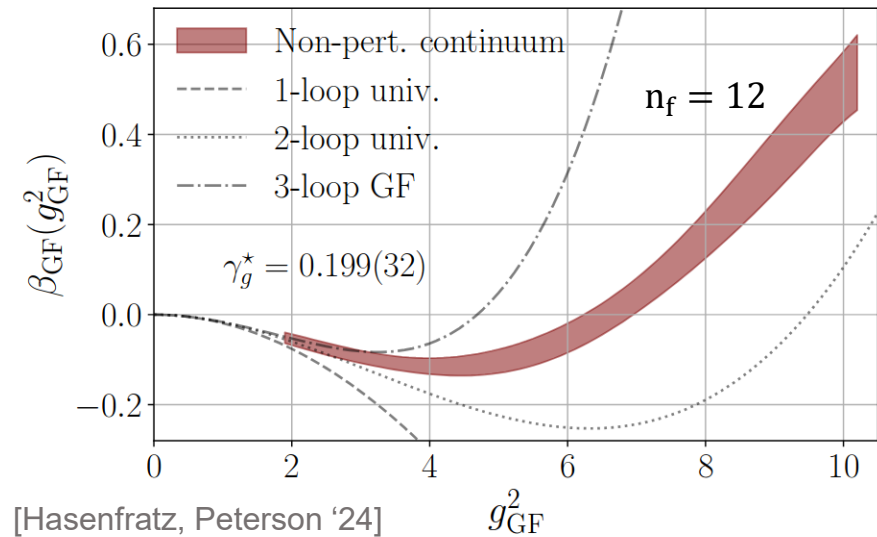
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[Lüscher '10]

[Harlander, Neumann '16]



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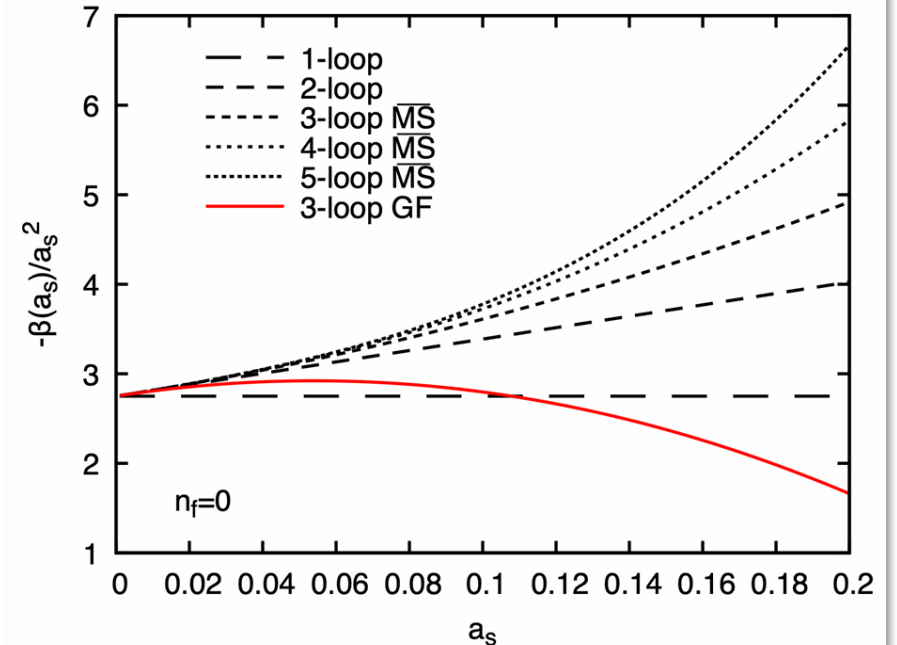
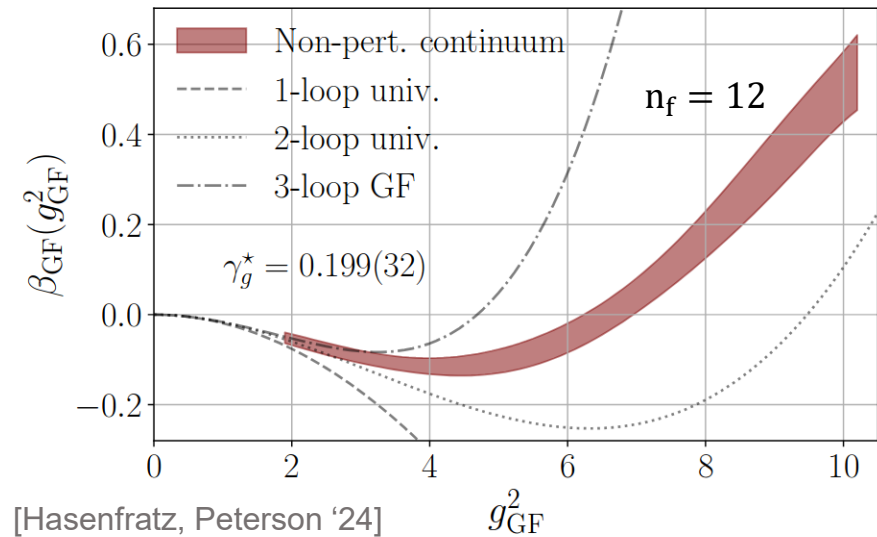
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[Lüscher '10]

[Harlander, Neumann '16]

Unknown



Perturbative gradient flow

[Lüscher '10; Lüscher, Weisz '11]

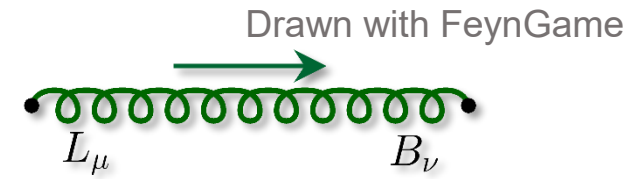
$$\mathcal{L}_{flow} = \mathcal{L}_{QCD} + \int dt \left[L_\mu \frac{\partial}{\partial t} B_\mu - L_\mu \mathcal{D}_\nu G_{\nu\mu} \right]$$

$$G_{\mu\nu}(t) = \partial_\mu B_\nu - \partial_\nu B_\mu + g_0 [B_\mu, B_\nu] \quad 3$$

Perturbative gradient flow

[Lüscher '10; Lüscher, Weisz '11]

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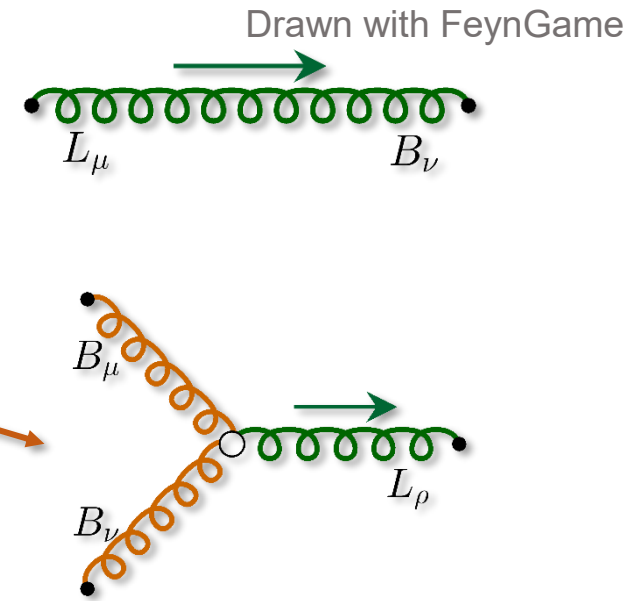


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Perturbative gradient flow

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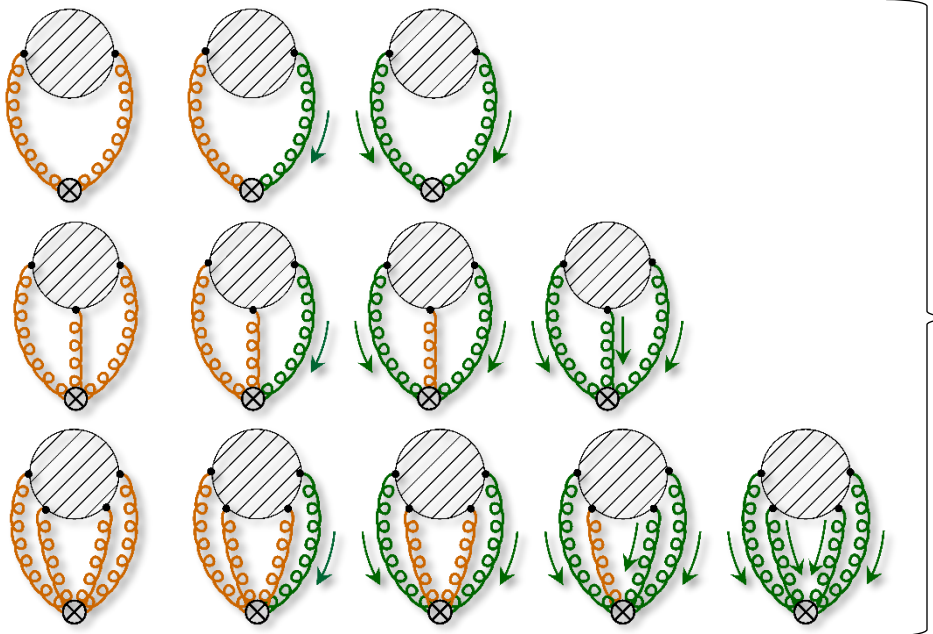
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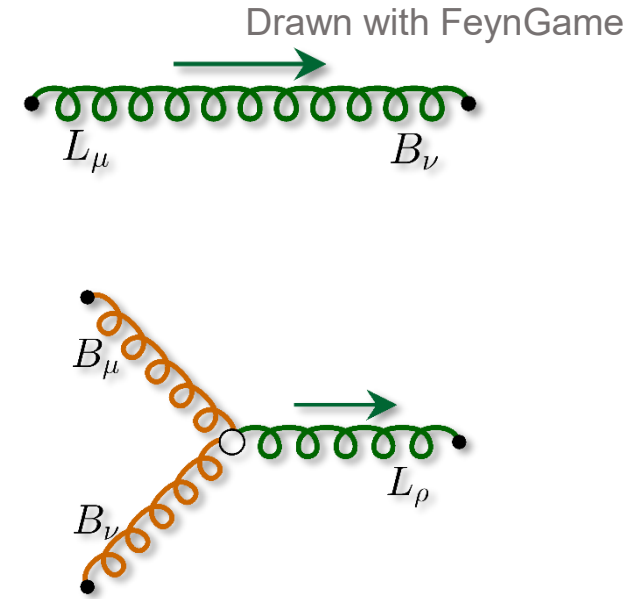
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Possible at 4-loop?

- ~ 5.500 Diagrams
- ~ 500.000 Integrals
- Computational heavy
- In progress



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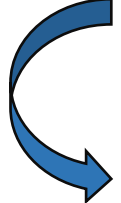
Digression: ftint

[Harlander, Nellopoulos, Olsson, Wesle '24]

$$I\left(\{c_1, \dots, c_N\}, \{a_1, \dots, a_9\}, \{b_1, \dots, b_9\}\right) = \left(\prod_{j=1}^N \int_0^1 du_j u_j^{c_j} \right) \int_{p_1, \dots, p_4} \frac{\exp\left[-t \sum_{i=1}^9 a_i(u) P_i^2\right]}{(P_1^2)^{b_1} \dots (P_9^2)^{b_9}}$$

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$$\frac{1}{(p^2)^b} \sim \int_0^\infty dx x^{b-1} e^{-xp^2} \quad \left(\xrightarrow{\text{map}} \int_0^1 dx \dots \right)$$
$$\sim \left(\prod_{j=1}^N \int_0^1 du_j u_j^{c_j} \right) \left(\prod_{i=1}^9 \int_0^1 dx_i x_i^{b_i-1} \right) \int_{p_1, \dots, p_4} \exp [\mathbf{p}^T \mathbf{A}(x, u) \mathbf{p}]$$

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 \end{aligned}$$

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$$\int_{\mathbf{p}} \exp [\mathbf{p}^T \mathbf{A} \mathbf{p}] \sim [\det \mathbf{A}]^{-D/2}$$

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Sector decomposition:

→ pySecDec

[Borowka et al. '17]

$$\sim \sum_{l=1}^{\alpha} \sum_{n=-r}^p \mathcal{I}_{l,n} \frac{1}{\epsilon^n} + \mathcal{O}(\epsilon^{r+1})$$

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Sector decomposition:
 → pySecDec
 [Borowka et al. '17]

Numerical integration:
 → pySecDec/Disteval
 [Borowka et al. '17]

Digression: ftint at 4-loop

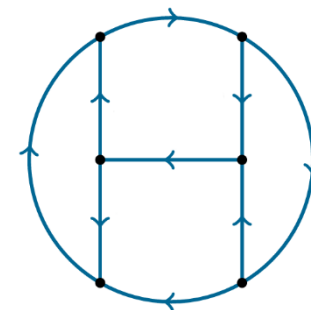
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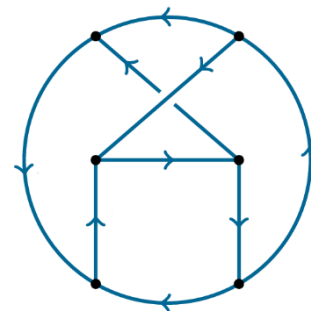
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(a) **H**



(b) **X**

Digression: ftint at 4-loop

[Harlander, Nellopoulos, Olsson, Wesle '24]

$$I(\{c_1, \dots, c_N\}, \{a_1, \dots, a_9\}, \{b_1, \dots, b_9\}) = \left(\prod_{j=1}^N \int_0^1 du_j u_j^{c_j} \right) \int_{p_1, \dots, p_4} \frac{\exp \left[-t \sum_{i=1}^9 a_i(u) P_i^2 \right]}{(P_1^2)^{b_1} \dots (P_9^2)^{b_9}}$$

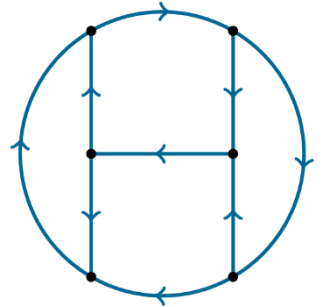
Example:

$$I(\{\}, \{2, 0, \dots, 0\}, \{2, 1, \dots, 1, -1, -1\}, H) \sim \left(\prod_{j=1}^N \int_0^1 du_j u_j^{c_j} \right) \left(\prod_{i=1}^9 \int_0^1 dx_i x_i^{b_i-1} \right) \int_{p_1, \dots, p_4} \exp [\mathbf{p}^T A(x, u) \mathbf{p}]$$

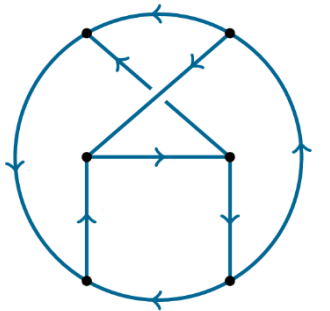
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Digression: ftint at 4-loop

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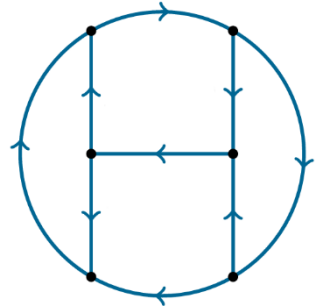
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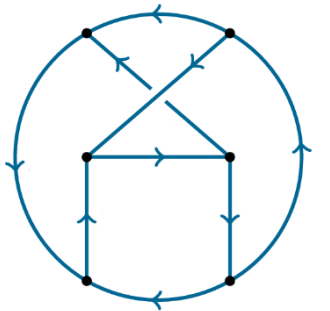
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Digression: ftint at 4-loop

[Harlander, Nellopoulos, Olsson, Wesle '24]

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< 1 min



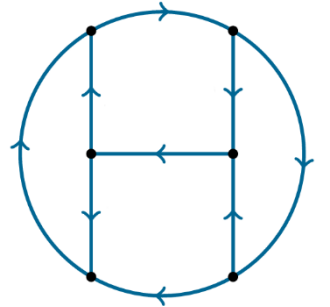
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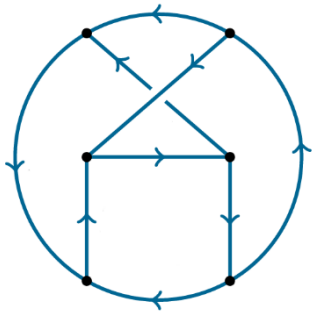
~ 2 min



$$= \sum_{n=-r}^p d_n \frac{1}{\epsilon^n} + \mathcal{O}(\epsilon^{r+1})$$



(a) **H**



(b) **X**

Digression: ftint at 4-loop

[Harlander, Nellopoulos, Olsson, Wesle '24]

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< 1 min



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~ 56 hours

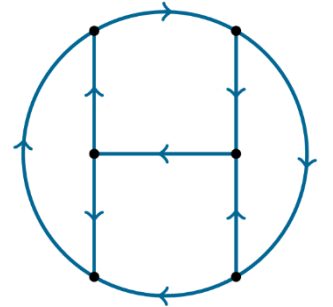


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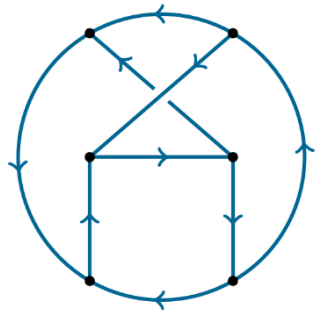
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$$= \sum_{n=-r}^p d_n \frac{1}{\epsilon^n} + \mathcal{O}(\epsilon^{r+1})$$



(a) **H**

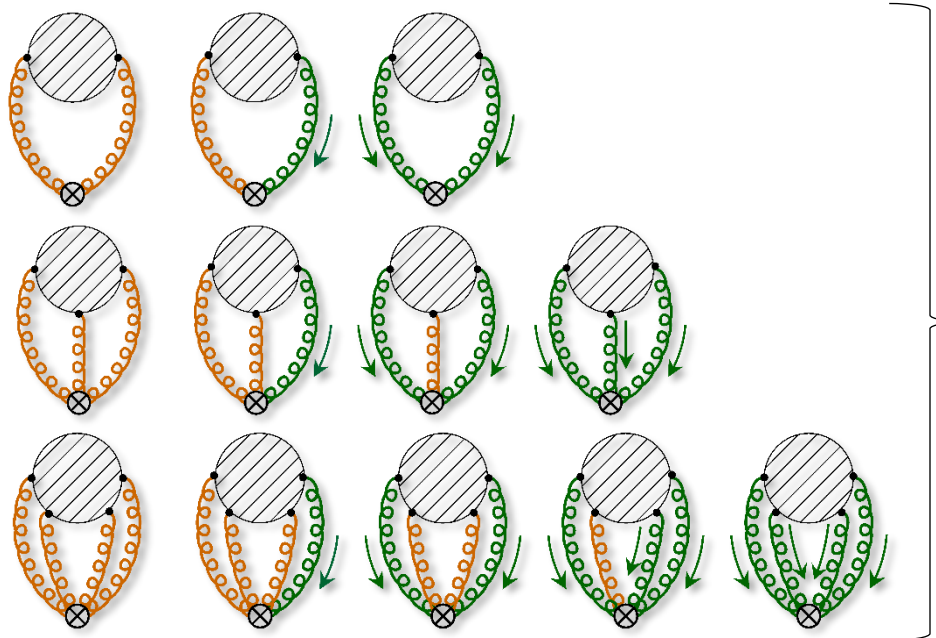


(b) **X**

Perturbative gradient flow

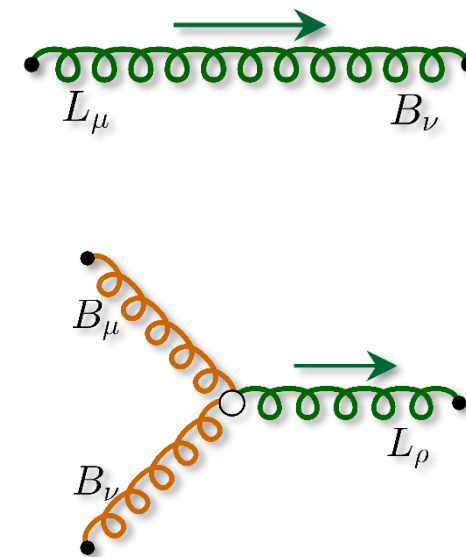
$$\mathcal{L}_{flow} = \mathcal{L}_{QCD} + \int dt \left[L_\mu \frac{\partial}{\partial t} B_\mu - L_\mu \mathcal{D}_\nu G_{\nu\mu} \right]$$

$$\alpha_{GF}(t) = \frac{4\pi t^2}{3} \langle G_{\mu\nu}(t) G^{\mu\nu}(t) \rangle$$



Possible at 4-loop?

- ~ 5.500 Diagrams
- ~ 500.000 Integrals
- Computational heavy
- In progress

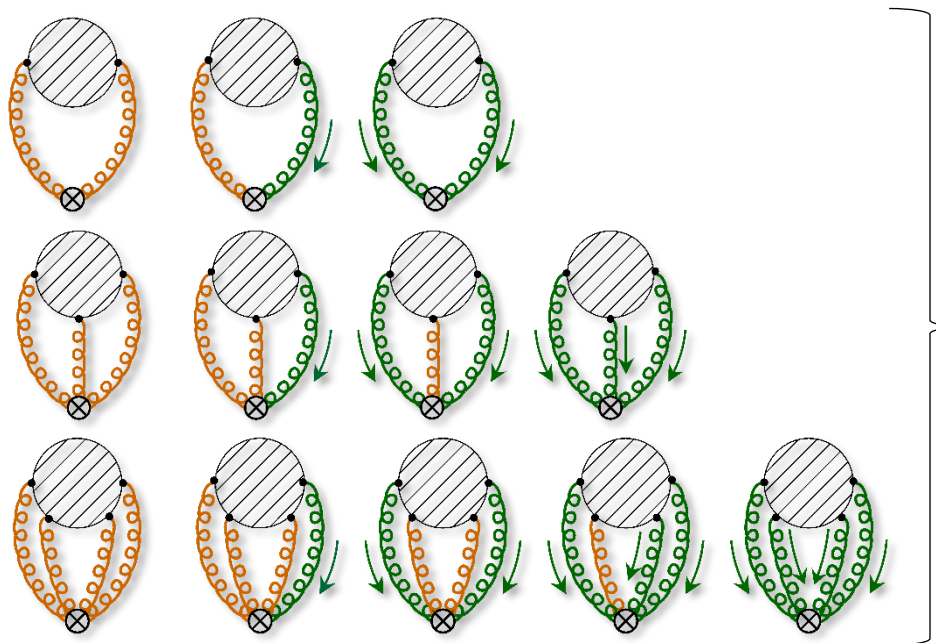


$$G_{\mu\nu}(t) = \partial_\mu B_\nu - \partial_\nu B_\mu + g_0 [B_\mu, B_\nu]$$

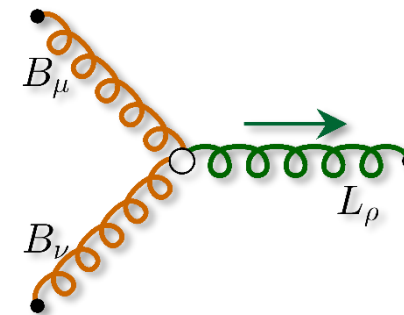
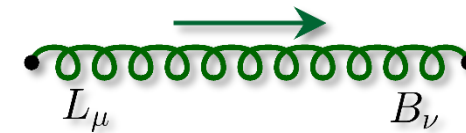
Perturbative gradient flow

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Simpler Coupling: QED

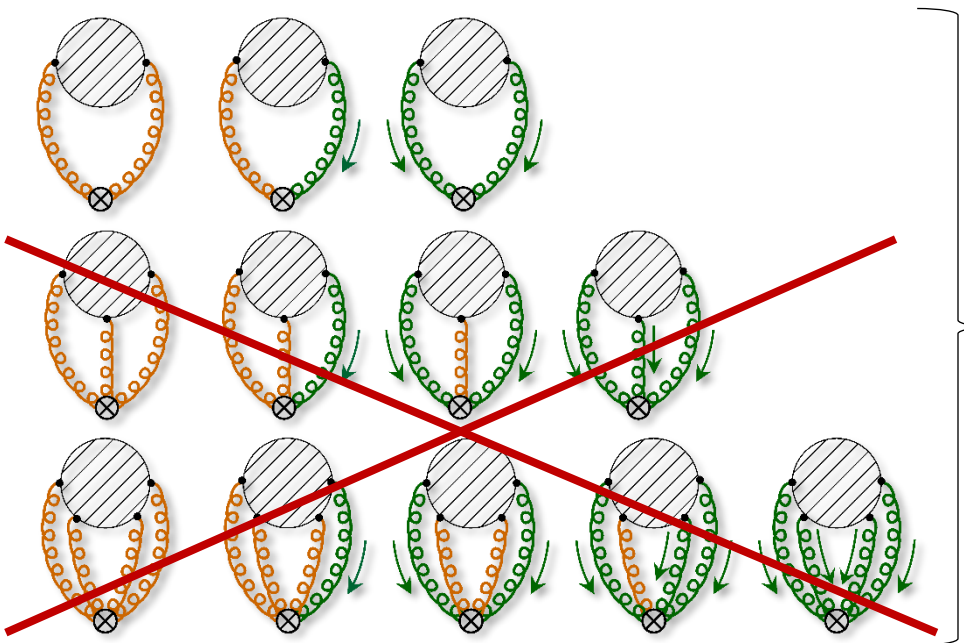


$$F_{\mu\nu}(t) = \partial_\mu B_\nu - \partial_\nu B_\mu + g_0 [\cancel{B_\mu}, \cancel{B_\nu}]_6$$

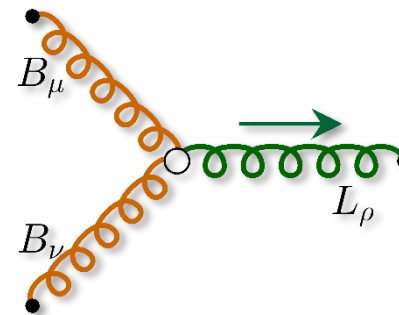
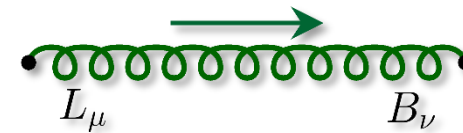
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Simpler Coupling: QED
 - $F_{\mu\nu}$ is linear in B_μ

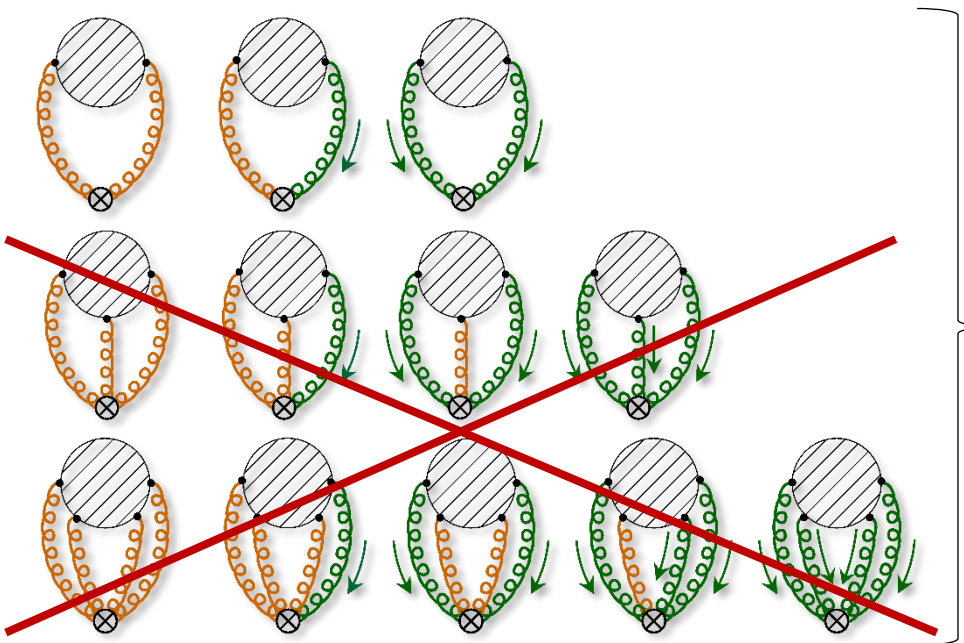


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Perturbative gradient flow

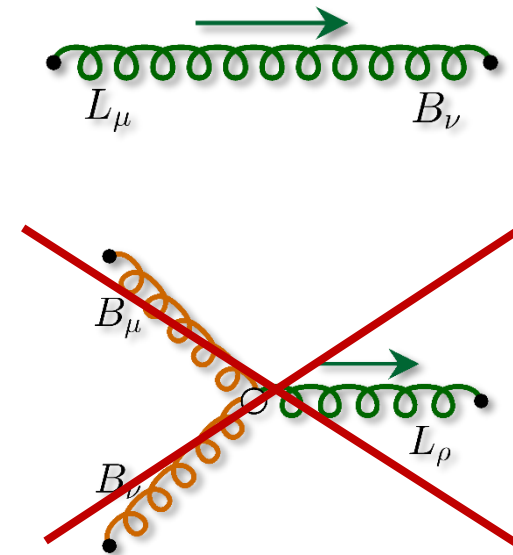
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Simpler Coupling: QED

- $F_{\mu\nu}$ is linear in B_μ
- No flow-vertices

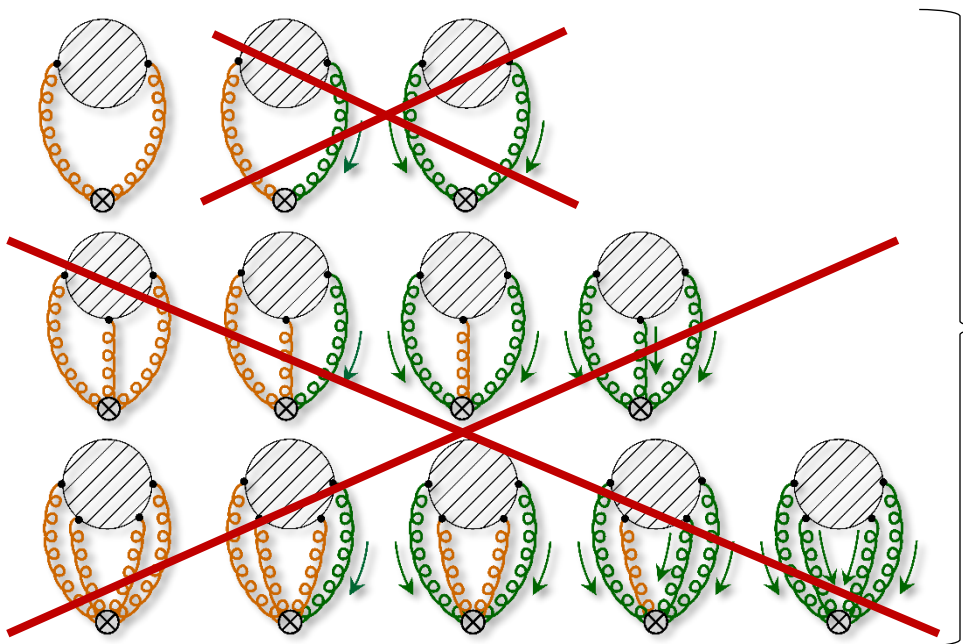


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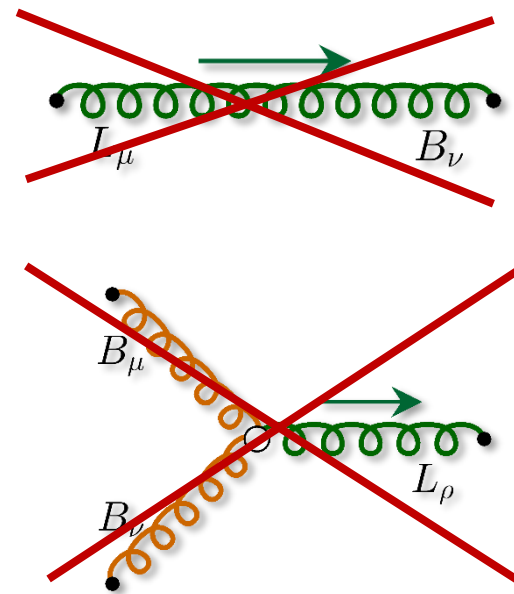
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- ➔ No photon flow-lines

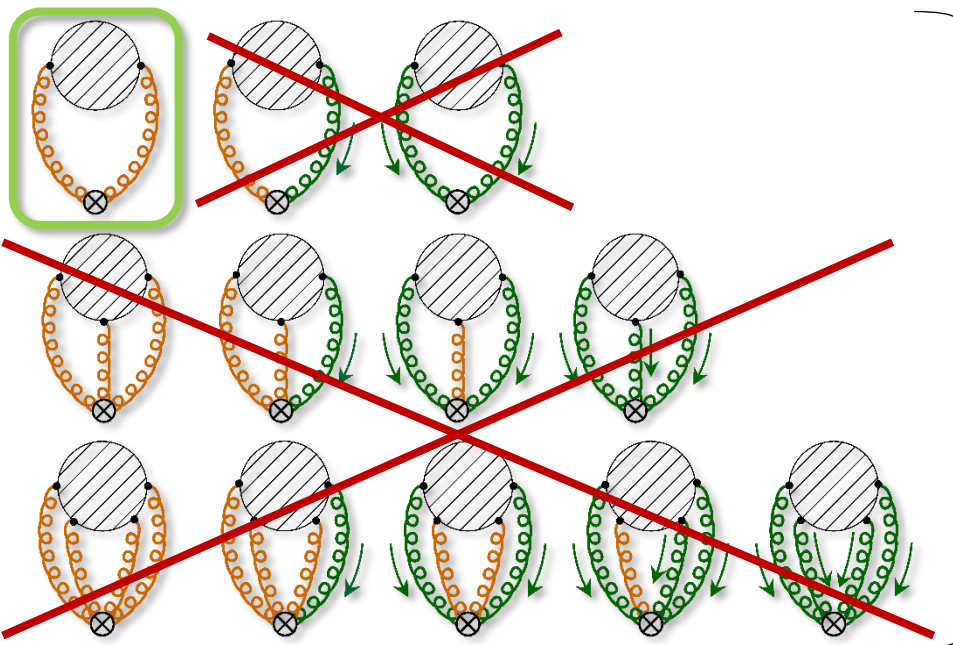


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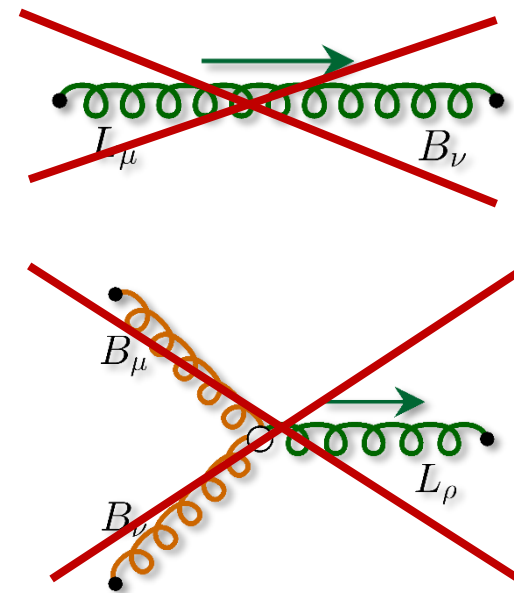
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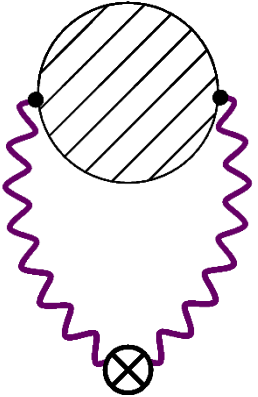
Simpler Coupling: QED

- $F_{\mu\nu}$ is linear in B_μ
- No flow-vertices
- ➔ No photon flow-lines
- Solution to flow-equation is exact

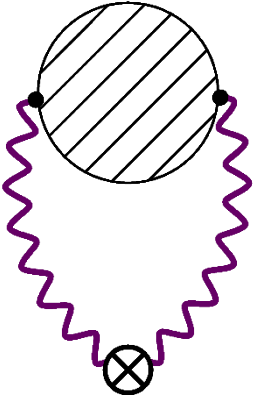


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QED coupling in the gradient flow

$$\alpha_e^{\text{GF}}(t) = \text{Diagram} = \alpha_e (8\pi t)^{D/2} \int \frac{d^D p}{(2\pi)^D} \frac{e^{-2tp^2}}{1 + \Pi_{\text{R}}(p)}$$


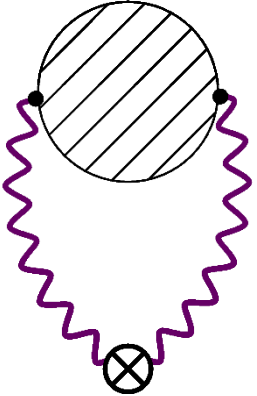
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Photon polarization
function of the unflowed
theory

Not restricted to QED
contributions
→ QCD contr.
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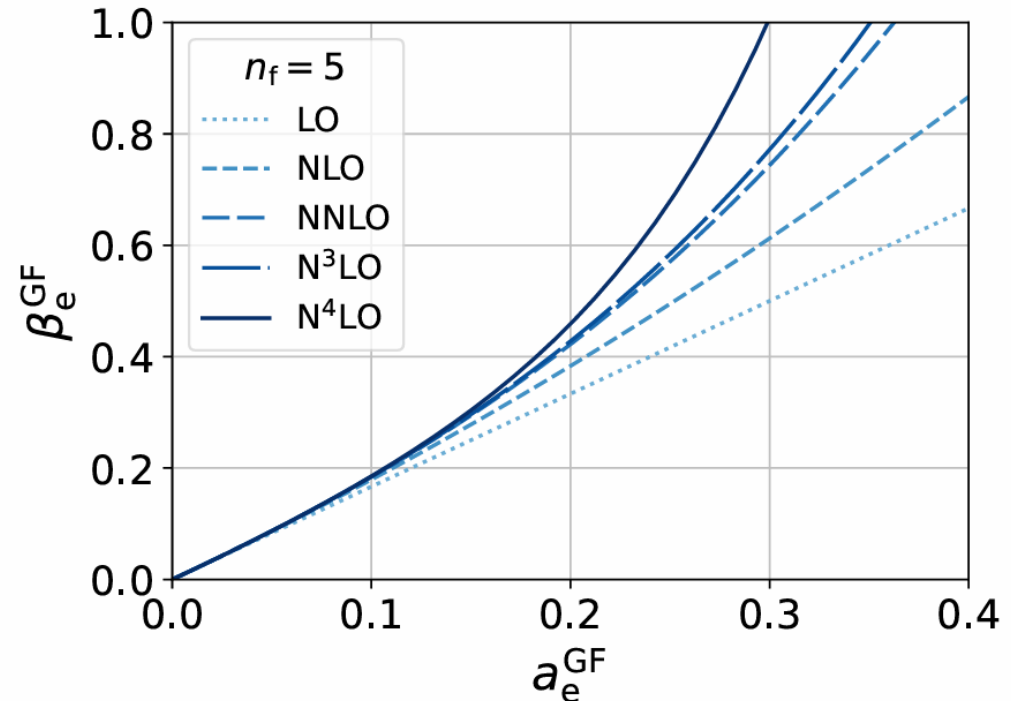
Not restricted to QED contributions
 → QCD contr.
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First example: QED₄ contr.

→ Photon polarization function known through N⁴LO

[Baikov, Chetyrkin, Kühn, Sturm '13]

- First 4-loop GF results
- Testing ground for QCD



QED in (2+1) dimensions

- Phenomenological relevance:

QED in (2+1) dimensions


- Phenomenological relevance:
 - Effective modeling of 2D materials

Two-loop fermion self-energy in reduced quantum electrodynamics and application to the ultrarelativistic limit of graphene #1

A.V. Kotikov (Dubna, JINR), S. Teber (Paris, LPTHE) (Dec 9, 2013)

Published in: *Phys.Rev.D* 89 (2014) 6, 065038 • e-Print: 1312.2430 [hep-ph]

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Regarding confinement and dynamical chiral symmetry breaking in QED3 #1

A. Bashir (IFM-UMSNH, Michoacan), A. Raya (IFM-UMSNH, Michoacan), I.C. Cloet (Argonne, PHY), C.D. Roberts (Argonne, PHY) (Jun, 2008)

Published in: *Phys.Rev.C* 78 (2008) 055201 • e-Print: 0806.3305 [hep-ph]

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QED in (2+1) dimensions

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- Lattice

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A Computer Simulation of Chiral Symmetry Breaking in (2+1)-Dimensional QED with N Flavors #3


Elbio Dagotto (Illinois U., Urbana), John B. Kogut (Illinois U., Urbana), Aleksandar Kocic (Arizona U.) (Aug, 1988)

Published in: *Phys.Rev.Lett.* 62 (1989) 1083

pdf DOI cite claim reference search 180 citations

GF-coupling in QED₃

- Ingredient: Photon polarization function $\Pi(p)$

$$\Pi_R(p) = c_1 \left(\frac{\alpha_e}{p} \right) + c_2 \left(\frac{\alpha_e}{p} \right)^2 + c_3 \left(\frac{\alpha_e}{p} \right)^3 + \dots$$


$$\alpha_e^{\text{GF}}(t) = \alpha_e (8\pi t)^{D/2} \int \frac{d^D p}{(2\pi)^D} \frac{e^{-2tp^2}}{1 + \Pi_R(p)}$$

- Highly unstable in the IR!!
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Large- n_f :

$$\Pi_R(p) = d_0 \left(\frac{\alpha_e}{p}\right) \left[1 + d_1 \frac{1}{n_f} + \dots \right]$$

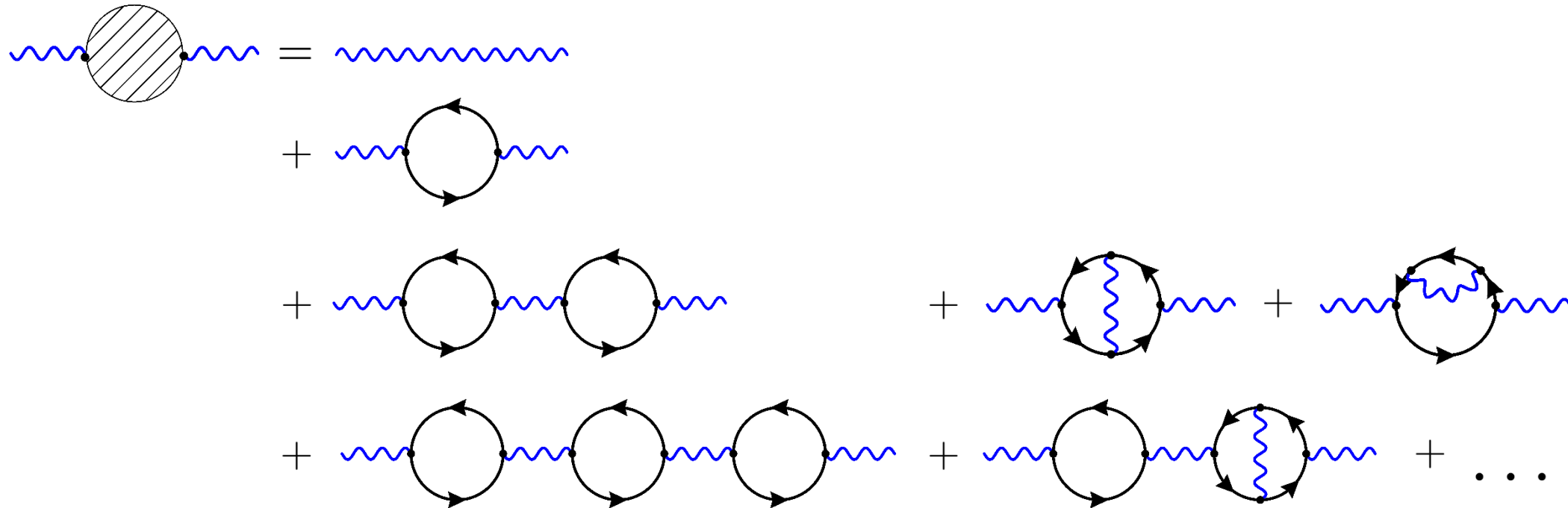
[Gusynin, Hams, Reenders '01]



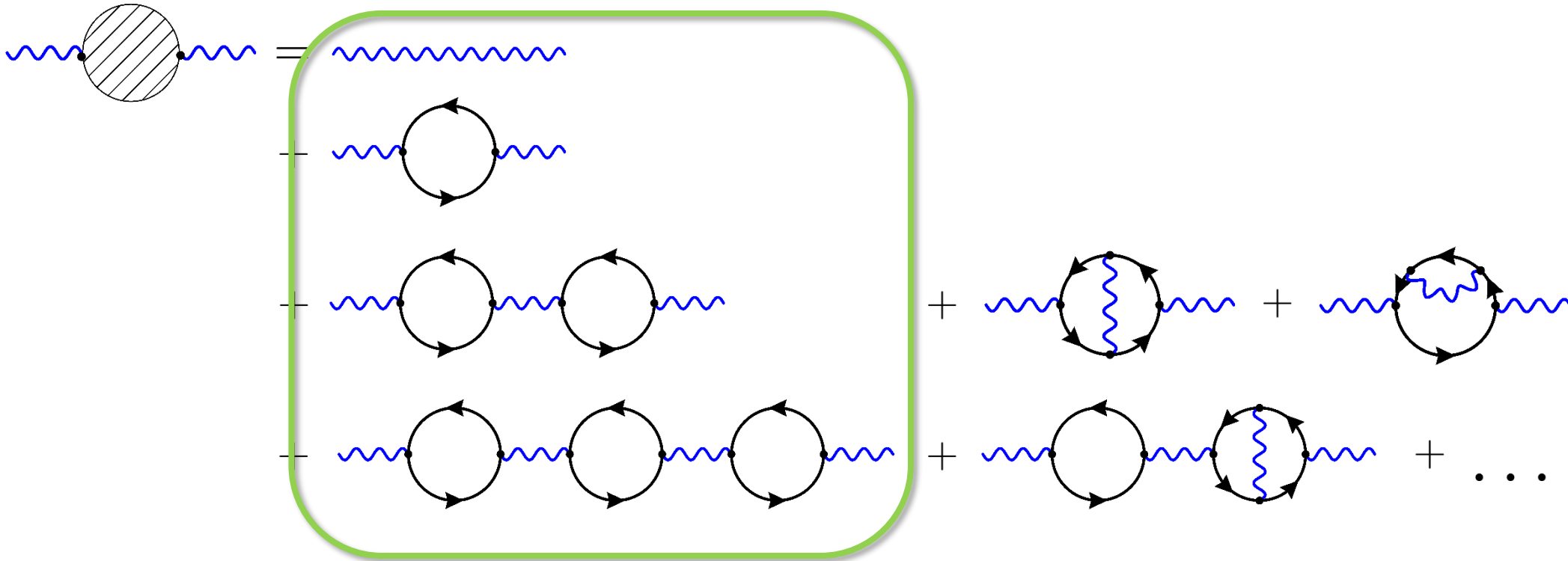
- IR-softened
- Non-trivial IR fixed point

[Appelquist, Bowick, Karabali, Wijewardhana '86]

Large- n_f expansion



Large- n_f expansion

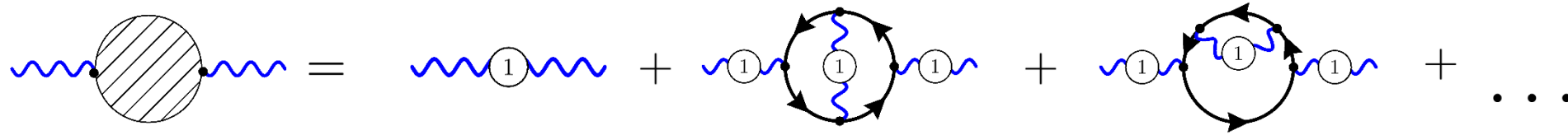


$$= \text{wavy line with } \textcircled{1} \sim \frac{1}{e^2 n_f p}$$

LO-softened photon propagator

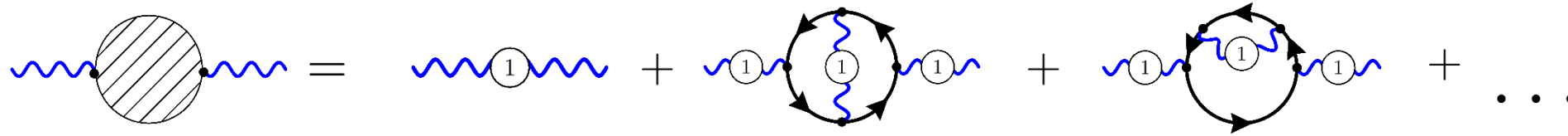
Large- n_f expansion

For more details: [Gracey '18] & [Metayer, Teber '23]



Large- n_f expansion

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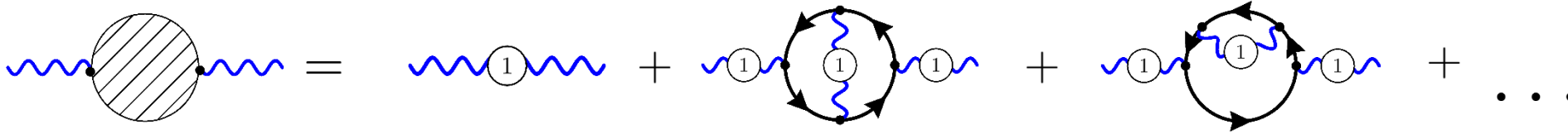


$$\Pi_R(p) = \left(\frac{\alpha_e}{p}\right) h_f \quad \text{with} \quad h_f = \frac{\pi n_f}{2} \left[1 + \left(\frac{184}{9\pi^2} - 2\right) \frac{1}{n_f} + \dots \right]$$

[Gusynin, Hams, Reenders '01]

Large- n_f expansion

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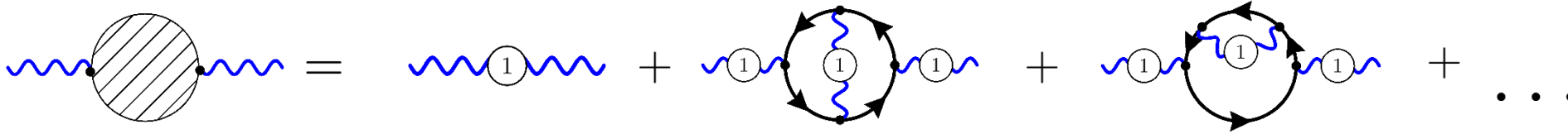
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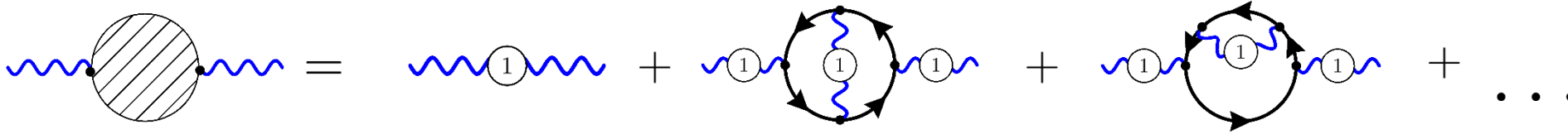
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[Gusynin, Hams, Reenders '01]

- Define rescaled dimensionless coupling: $\hat{\alpha}_e \sim h_f \alpha_e \sqrt{t}$

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- Define rescaled dimensionless coupling: $\hat{\alpha}_e \sim h_f \alpha_e \sqrt{t}$

- Determine $\hat{\alpha}_e^{\text{GF}}$ analytically

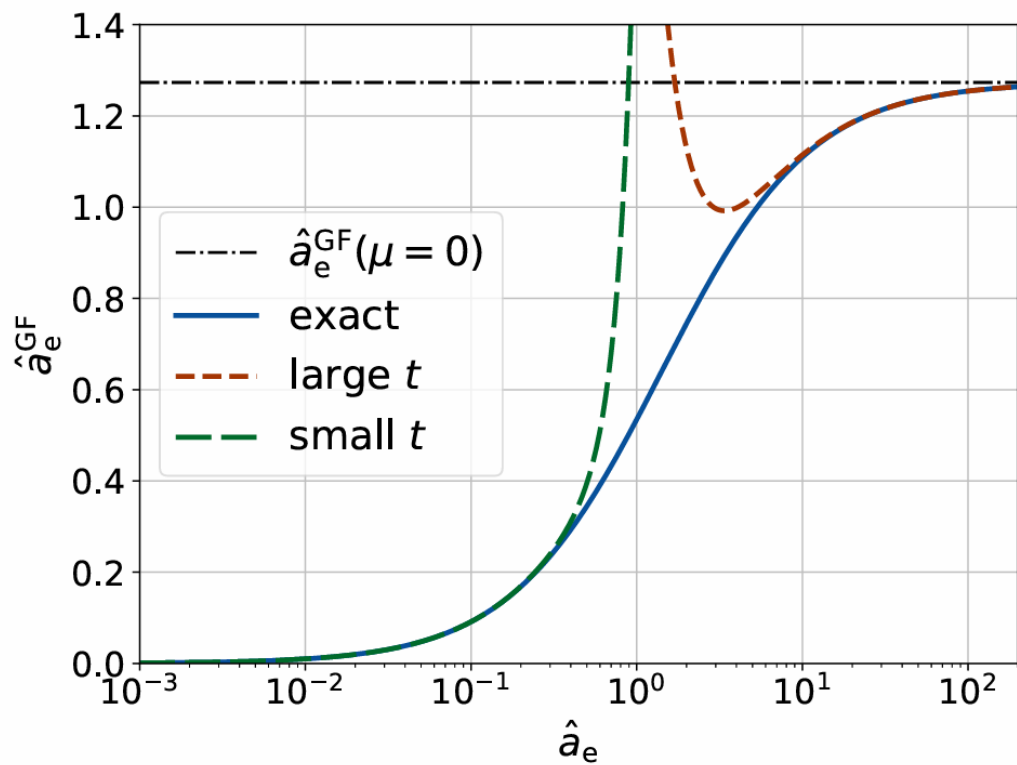
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Results for QED₃

$$\hat{\alpha}_e^{\text{GF}}(t) = \hat{\alpha}_e \left\{ 1 - \hat{\alpha}_e + \frac{\pi}{2} \hat{\alpha}_e^2 - \frac{\pi}{4} \hat{\alpha}_e^3 e^{-\frac{\pi}{4} \hat{\alpha}_e^2} \left[\Gamma \left(0, -\frac{\pi}{4} \hat{\alpha}_e^2 \right) + \sqrt{\pi} i \Gamma \left(\frac{1}{2}, -\frac{\pi}{4} \hat{\alpha}_e^2 \right) \right] \right\}$$

Results for QED₃

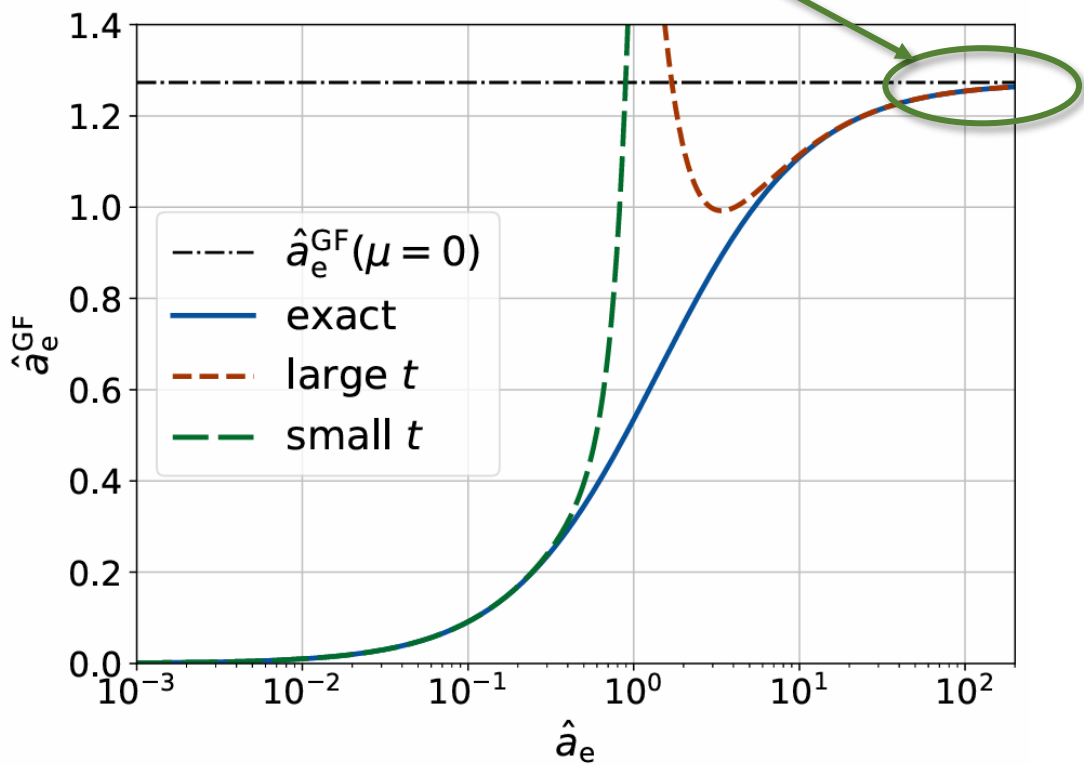
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$$\hat{\alpha}_{e,*}^{\text{GF}} = \frac{4}{\pi}$$

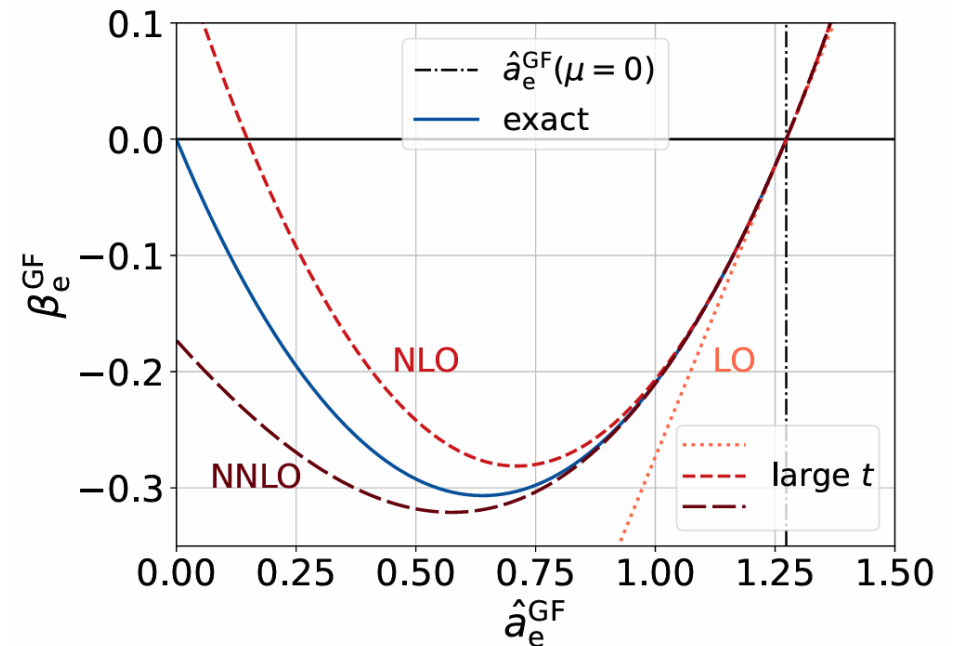
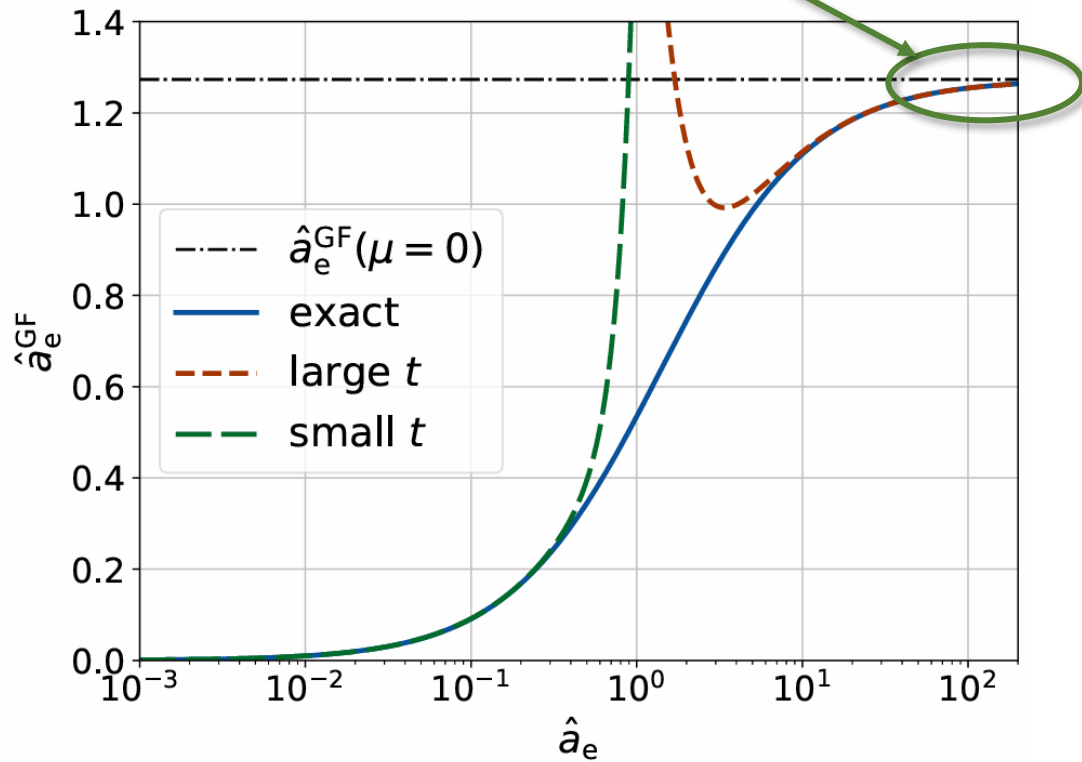


Results for QED₃

$$\hat{\alpha}_e^{\text{GF}}(t) = \hat{\alpha}_e \left\{ 1 - \hat{\alpha}_e + \frac{\pi}{2} \hat{\alpha}_e^2 - \frac{\pi}{4} \hat{\alpha}_e^3 e^{-\frac{\pi}{4} \hat{\alpha}_e^2} \left[\Gamma\left(0, -\frac{\pi}{4} \hat{\alpha}_e^2\right) + \sqrt{\pi} i \Gamma\left(\frac{1}{2}, -\frac{\pi}{4} \hat{\alpha}_e^2\right) \right] \right\}$$

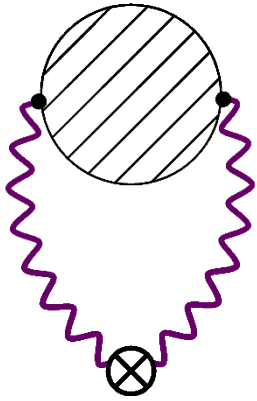
$$\hat{\alpha}_{e,*}^{\text{GF}} = \frac{4}{\pi}$$

- Recover Infra-red fixed point in the GF scheme predicted by [Appelquist et al. '86]
- Scheme independent critical exponent matches: $\beta'(\hat{\alpha}_{e,*}) = 1$



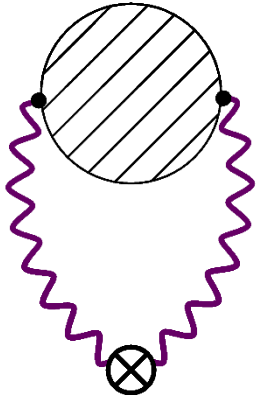
Summary & Conclusion

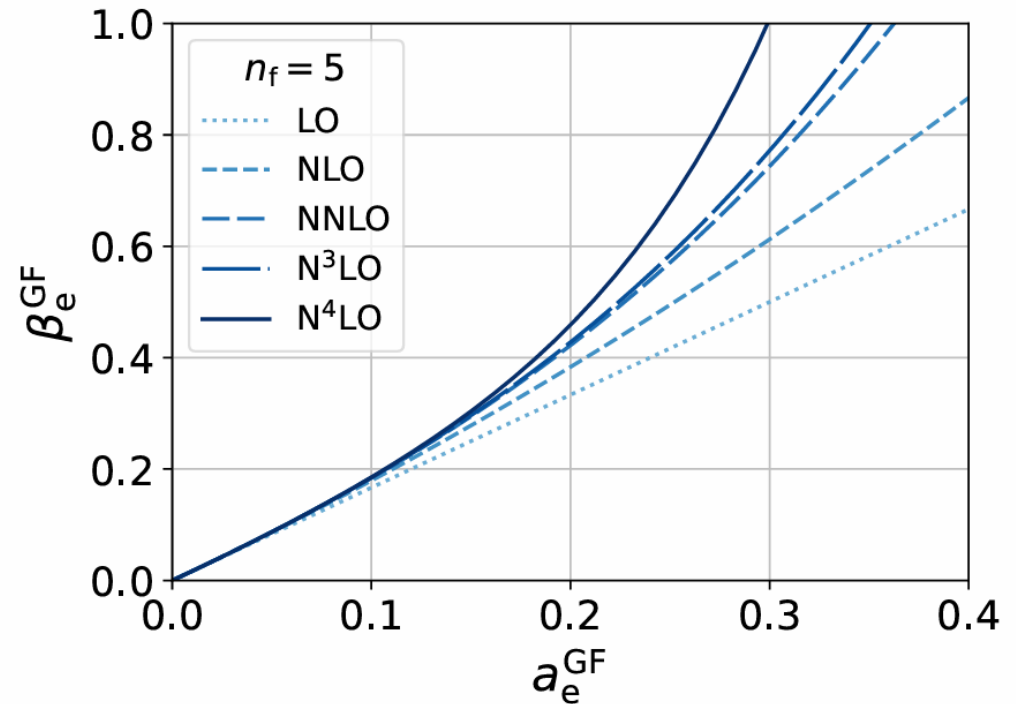
- QED gradient-flow coupling can be expressed in terms of the unflowed theory

$$\alpha_e^{\text{GF}}(t) = \text{Diagram} = \alpha_e (8\pi t)^{D/2} \int \frac{d^D p}{(2\pi)^D} \frac{e^{-2tp^2}}{1 + \Pi_R(p)}$$


Summary & Conclusion

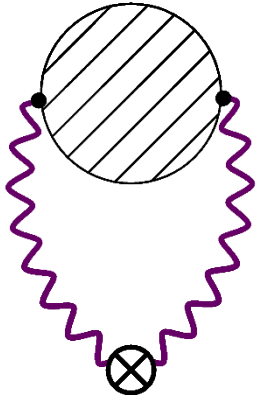
- QED gradient-flow coupling can be expressed in terms of the unflowed theory
- Obtained 4-loop gradient-flow results in analytical form

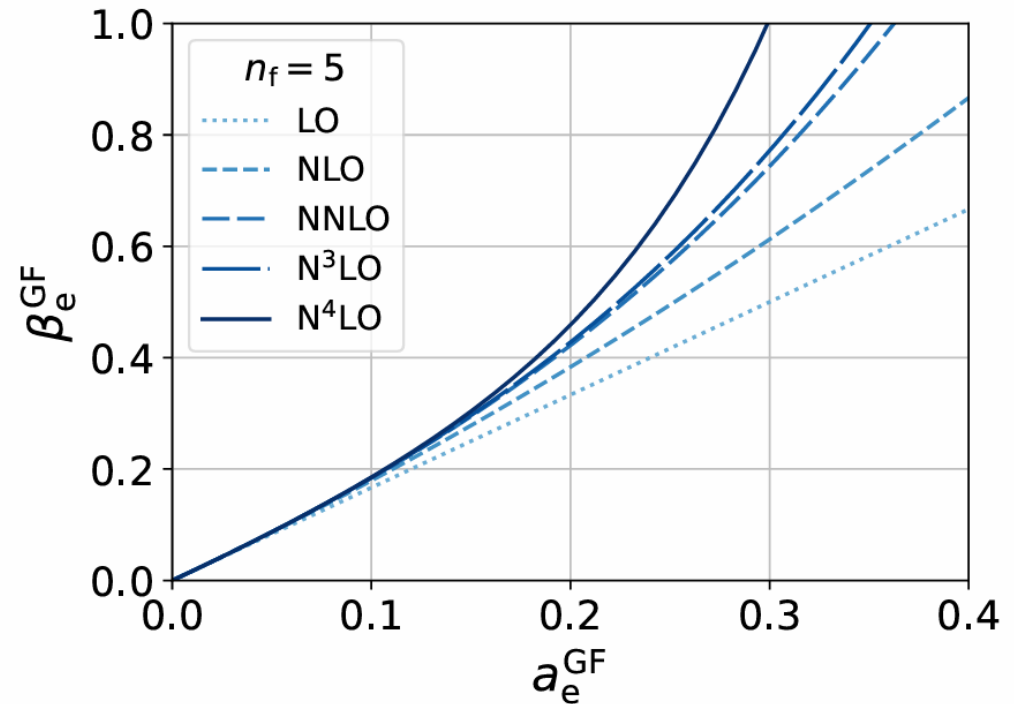
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Summary & Conclusion

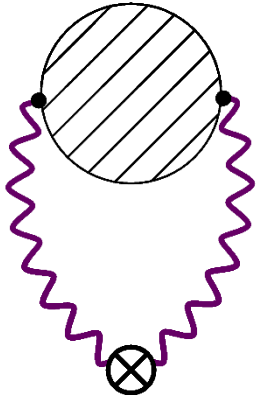
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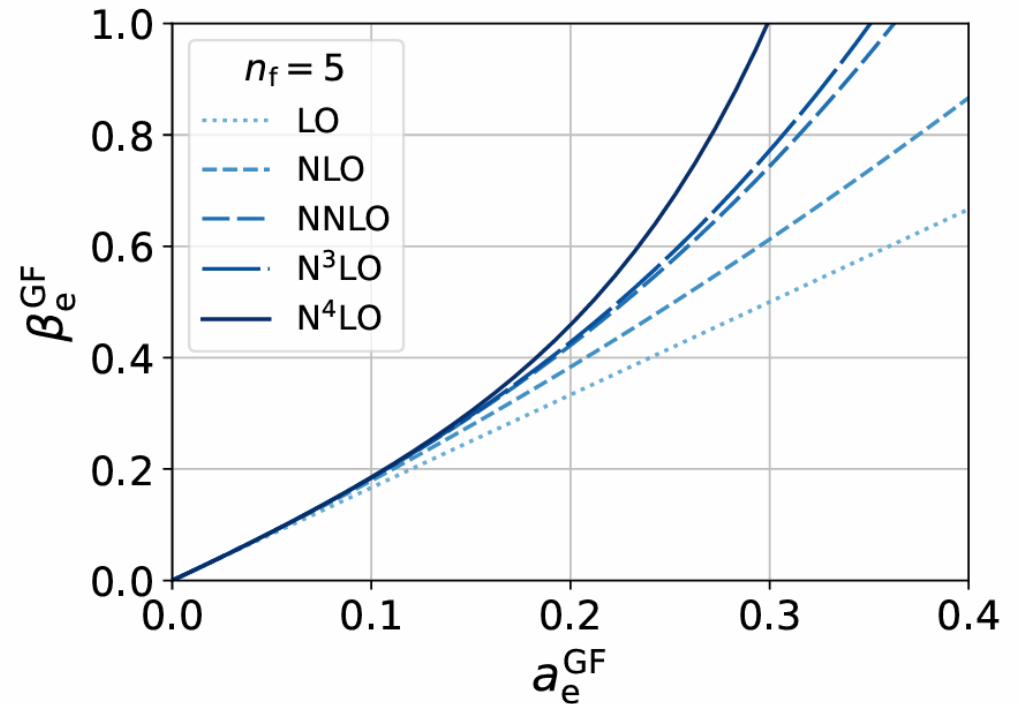
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Summary & Conclusion

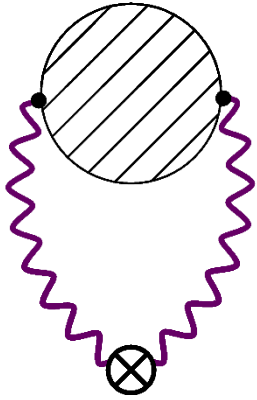
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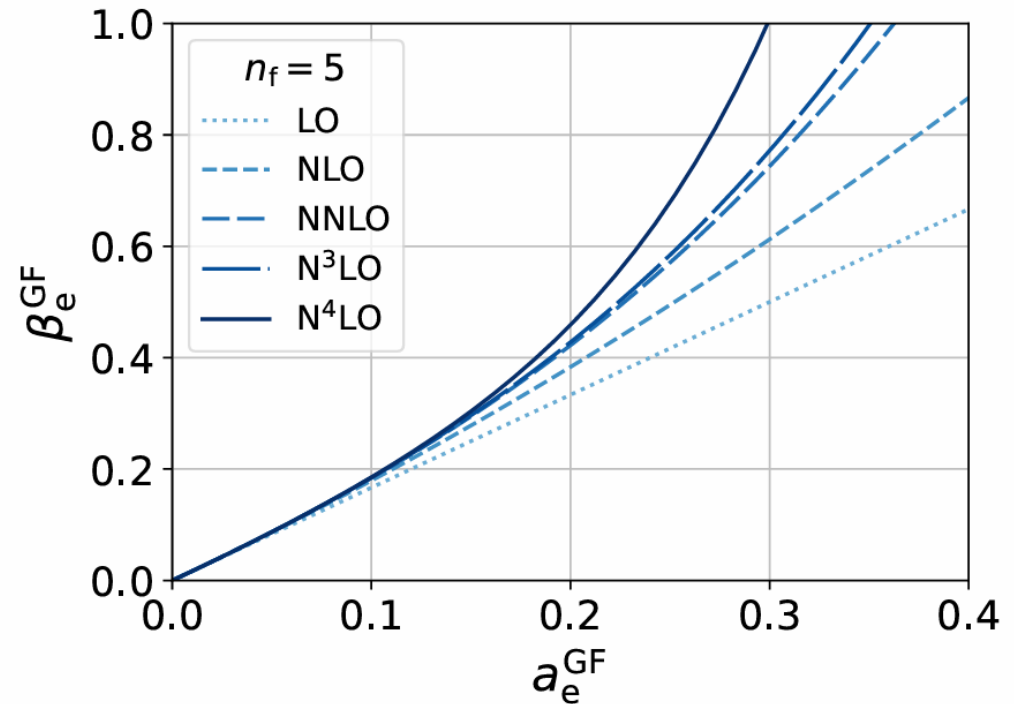
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Summary & Conclusion

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- Obtained 4-loop gradient-flow results in analytical form
- This provides first steps toward 4-loop gradient-flow coupling in QCD
- QED₃ as QCD-like test theory
- Studying QED₃ can offer insight into non-perturbative QCD-like dynamics
- Reproduced known IR fixed point in the gradient-flow scheme

$$\alpha_e^{\text{GF}}(t) = \text{[Diagram: A purple wavy line with a cross at the bottom and a shaded circle at the top]} = \alpha_e (8\pi t)^{D/2} \int \frac{d^D p}{(2\pi)^D} \frac{e^{-2tp^2}}{1 + \Pi_R(p)}$$

