



# Weak and Higgs physics from the lattice

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*In collaboration with Axel Maas, Georg Wieland and Patrick Jenny*

Nordic Lattice Meeting 2026  
19th of May 2026  
University of Edinburgh

# Introduction

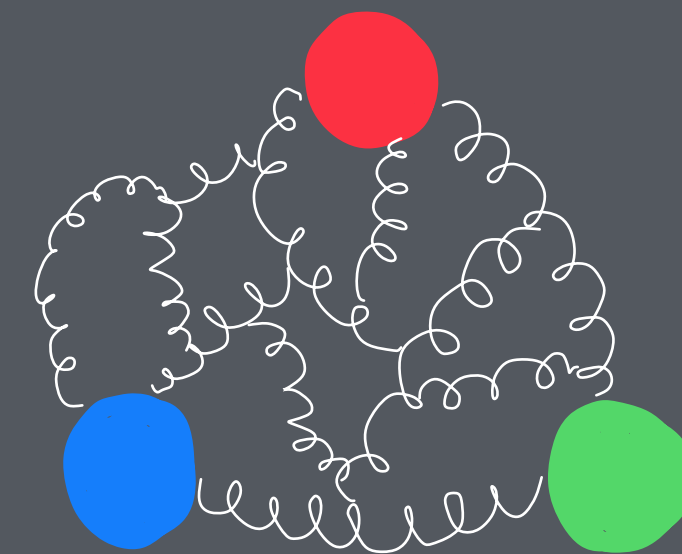
# Confinement in QCD (SU(3))

1. Observers are color blind



2. Physical objects are color singlets

3. Predictions require non-perturbative techniques



# Some claims about the weak interaction

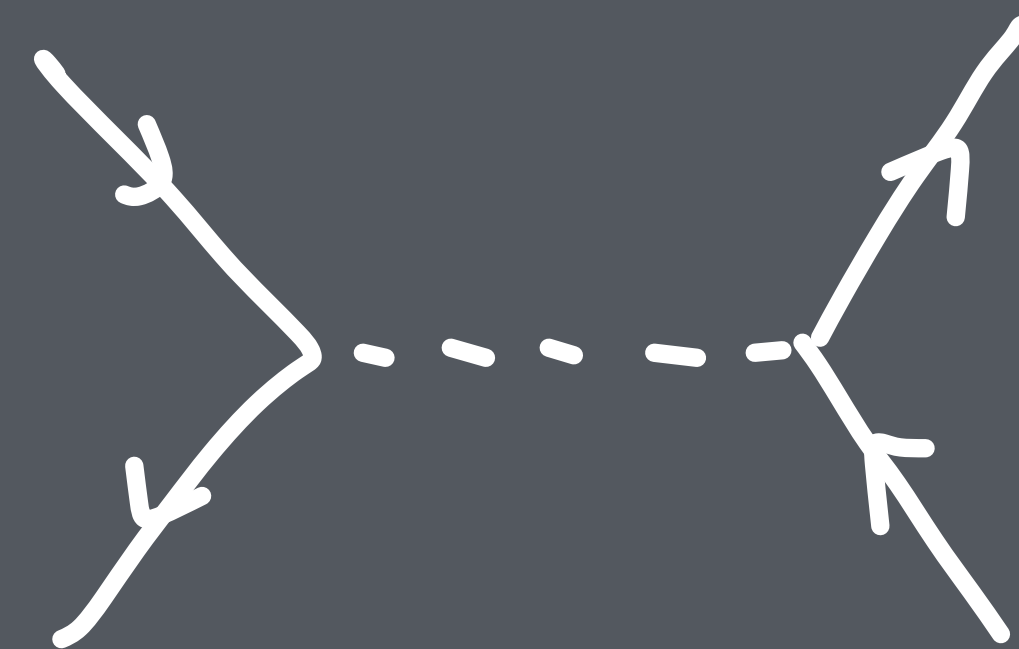
1. Observers are **not** weak isospin blind



3. Perturbative treatment possible using the BRST construction



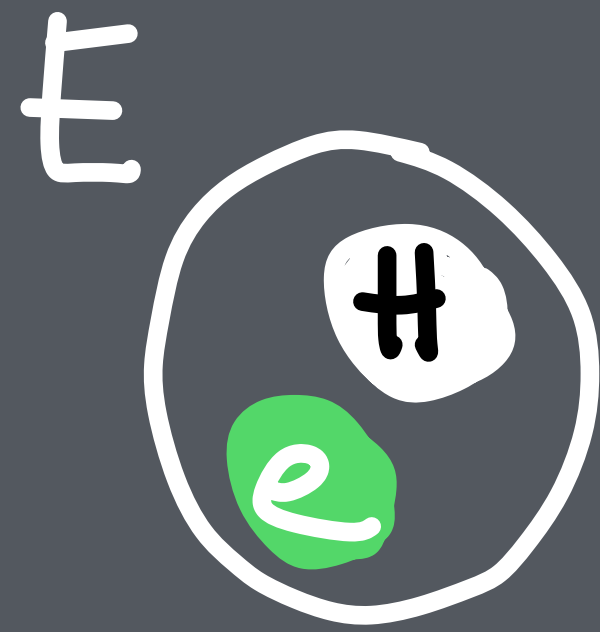
2. Leptons and neutrinos are individually observable



**Elitzur's theorem: Gauge symmetries cannot be spontaneously broken!**

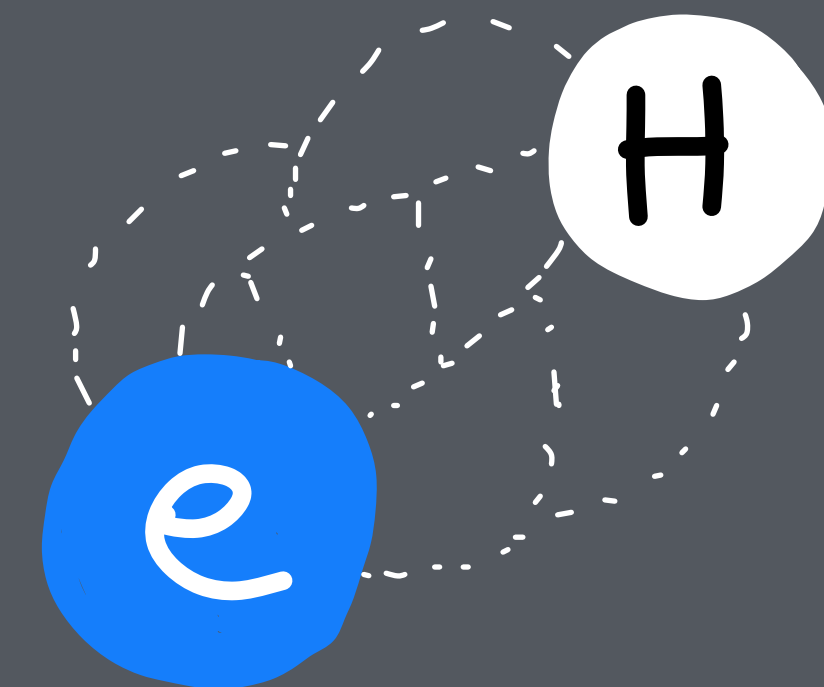
# What would make more sense

1. Observers are weak-isospin blind



2. Physical fermions are composite

3. Predictions require non-perturbative techniques

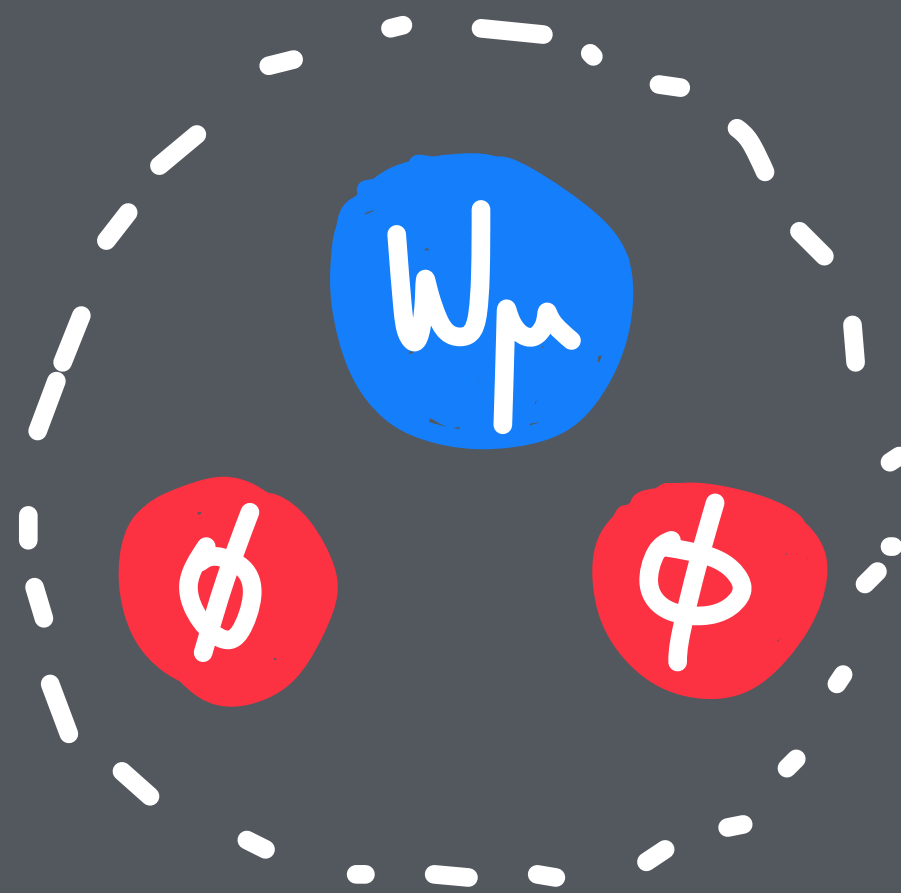


# More operators

SCALAR



VECTOR BOSON



LEFT-HANDED  
FERMION



[A. Maas, Prog. Part. Nucl. Phys., 2019, 1712.04721] and references therein

# Lattice Approach

Target theory

GAUGE

SCALAR

$$\mathcal{L} = -\frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} + \frac{1}{2} \text{tr}[(D_\mu X^\dagger)(D^\mu X)]$$

$$- \lambda \left( \frac{1}{2} \text{tr}(X^\dagger X) - f^2 \right)$$

+ LEFT-HANDED & RIGHT-HANDED  
FERMIONS

+ YUKAWA - COUPLINGS

LATTICE

PDFs:

GAUGE + SCALAR

SPECTROSCOPY:

GAUGE + SCALAR +

LEFT-HANDED FERMIONS

[A. Maas, Prog. Part. Nucl. Phys., 2019, 1712.04721] and references therein

# Augment established perturbation theory

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$$\langle \theta_{\phi}^{\dagger} \theta_{\phi} \rangle = v^2 h_i h_j \langle \theta_{\eta}^{\dagger} \theta_{\eta} \rangle + \mathcal{O}(v)$$

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$$\langle \theta_\phi^\dagger \theta_\phi \rangle = \underbrace{\nu^2 n_i n_j \langle \theta_\chi^\dagger \theta_\chi \rangle}_{\text{FMS mechanism}} + \mathcal{O}(\nu)$$

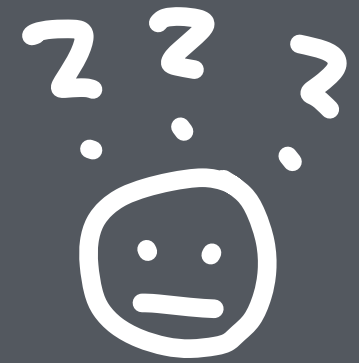
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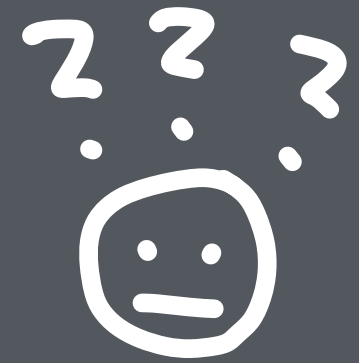
↑  
Actual Standard  
Model effects

# Why you should care

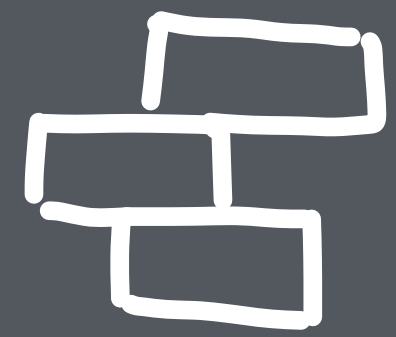


To not confuse a BSM with NP SM effects

# Why you should care

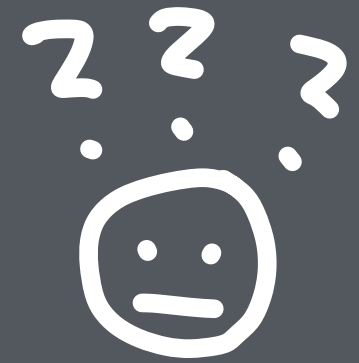


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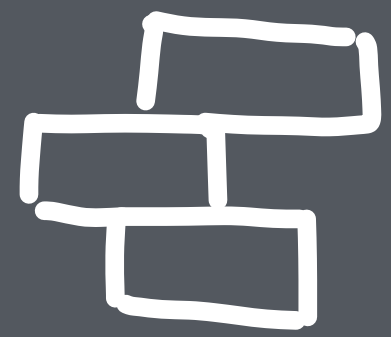


To correctly build extensions of the SM on top of the weak sector

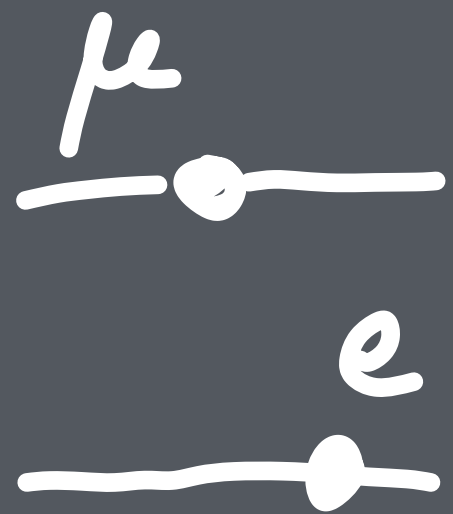
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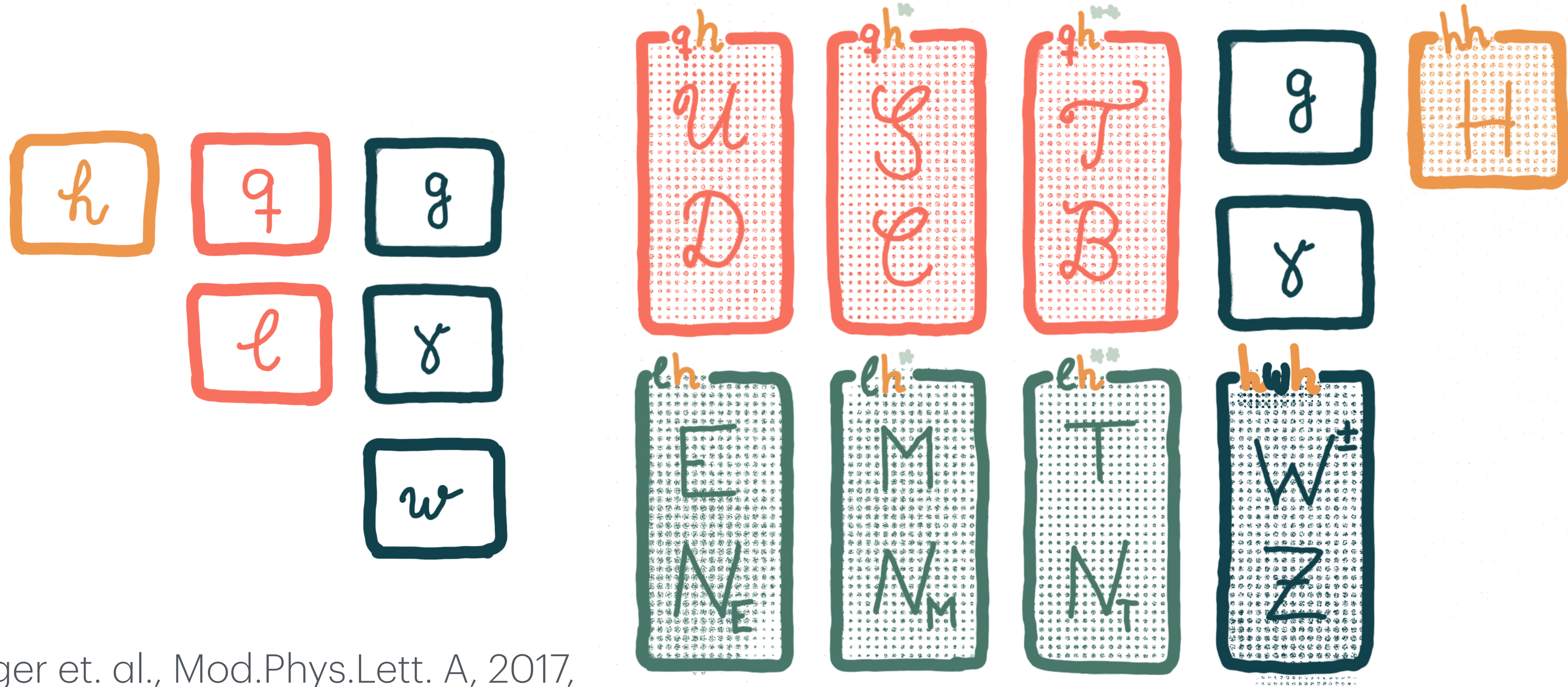
To correctly build extensions of the SM on top of the weak sector



It is consistent with the PDG that lepton & quark generations are levels of excited states -> compute CKM/PMNS matrix elements from the lattice

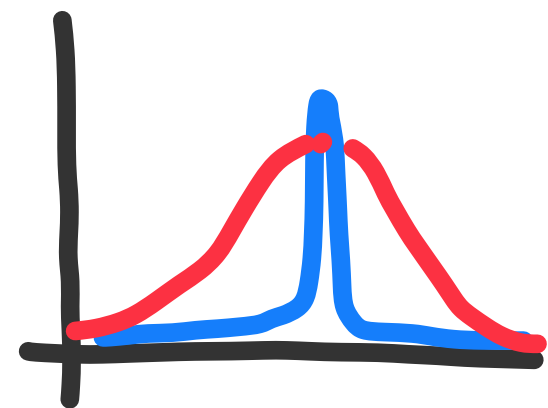
[Egger et. al., Mod.Phys.Lett. A, 2017, 1701.02881]

# The standard model but simpler



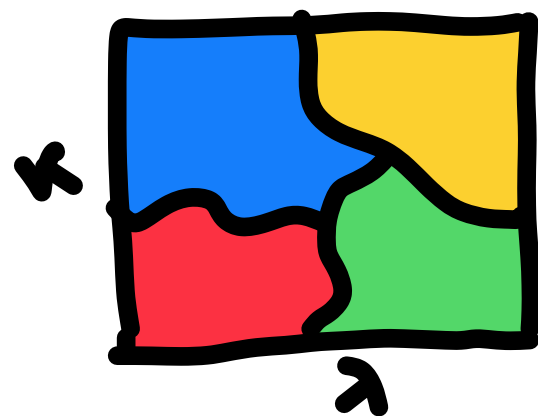
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# Outline



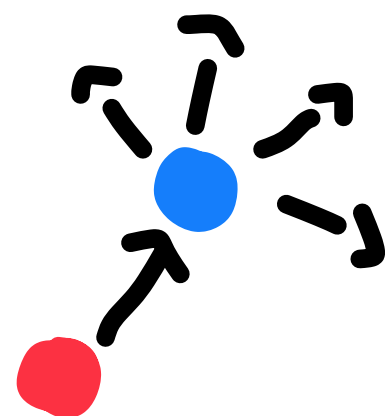
## 1. Quasi-PDFs

*How much internal structure do we see theoretically?*



## 2. Physical Phases and Mass Hierarchies in the system

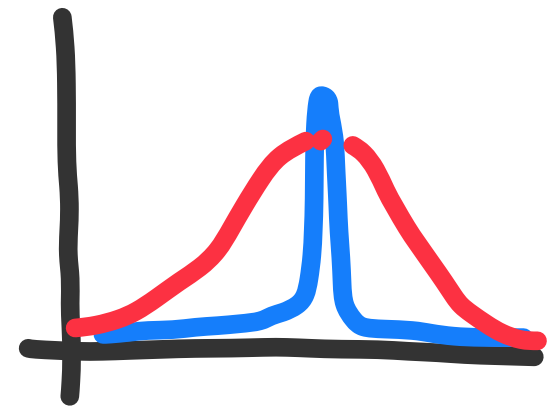
*How does the theory behave?*



## 3. Cross-sections obtained from spectral densities

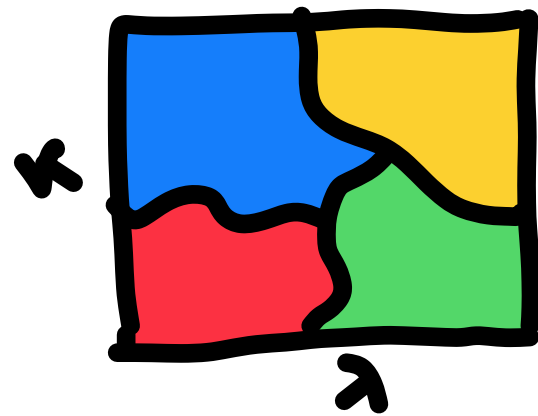
*Identify deviations from perturbation theory*

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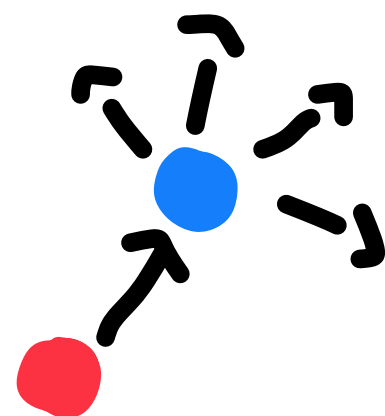
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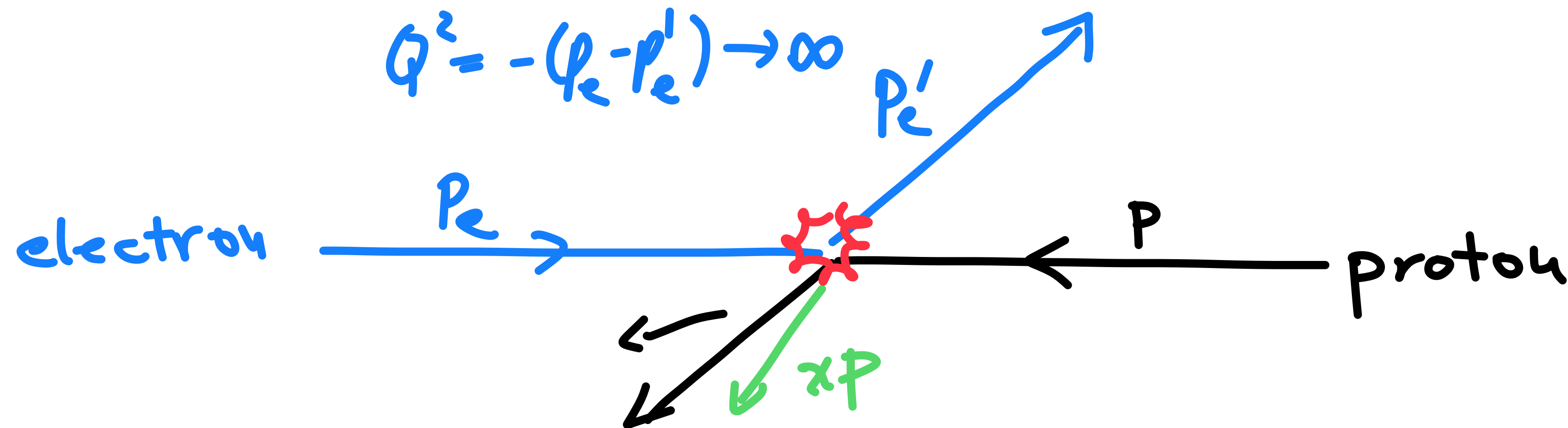
See excited states!

# Parton Distribution Functions for electroweak physics

# Parton Distribution Functions

Example: Deep-inelastic scattering of the proton

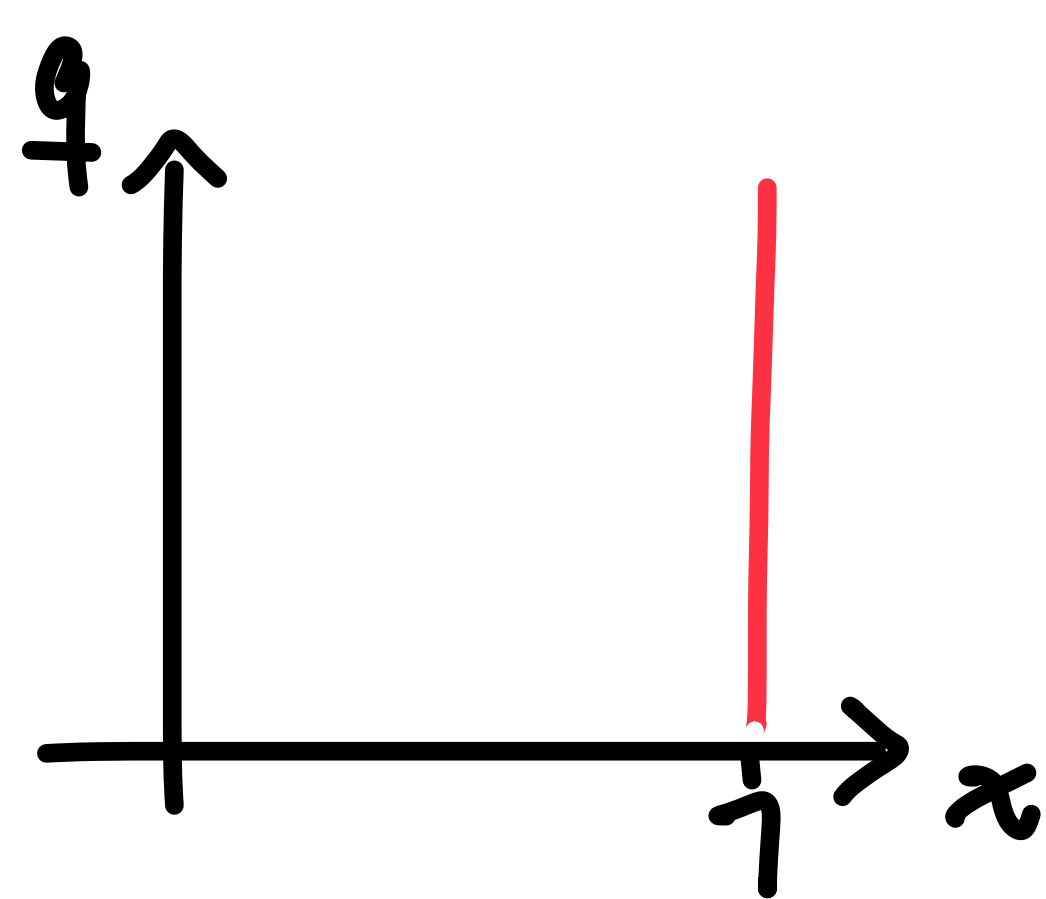
- Shows the distribution of a **single constituent momentum** relative to the total momentum
- Is computed in the Bjorken Limit: Infinite momentum transfer



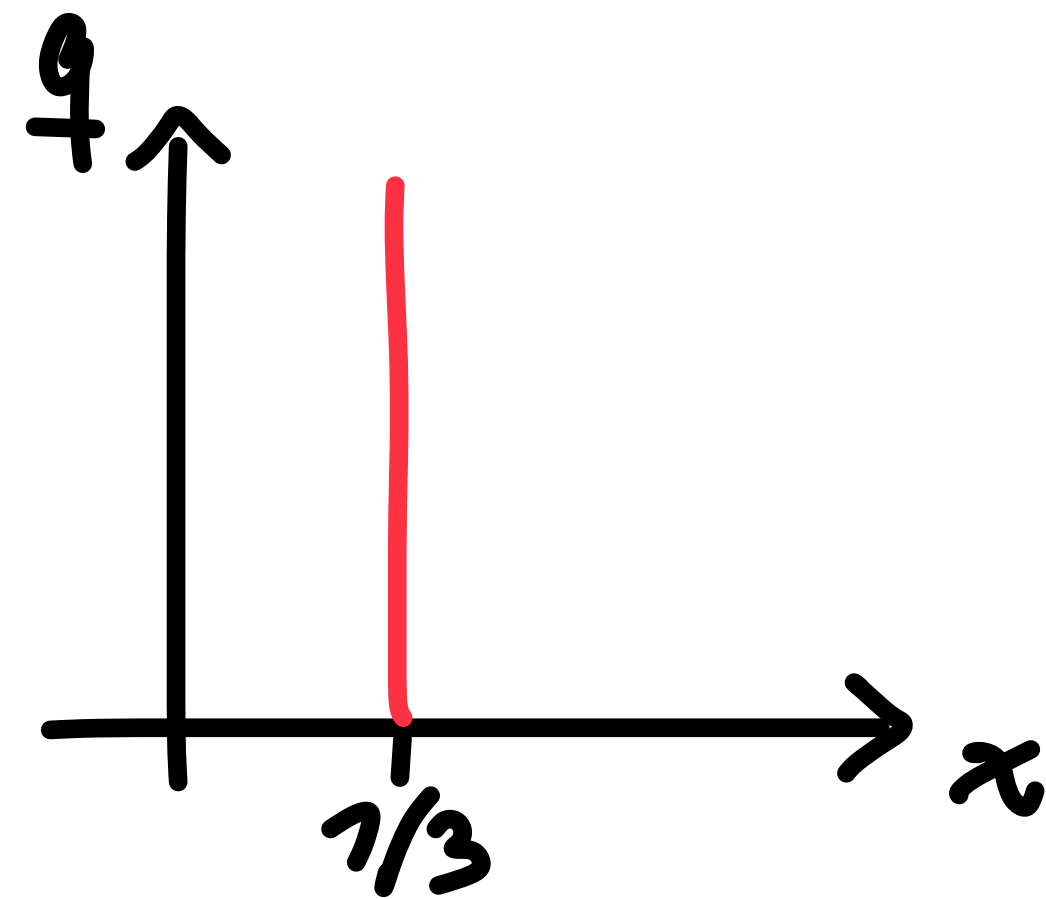
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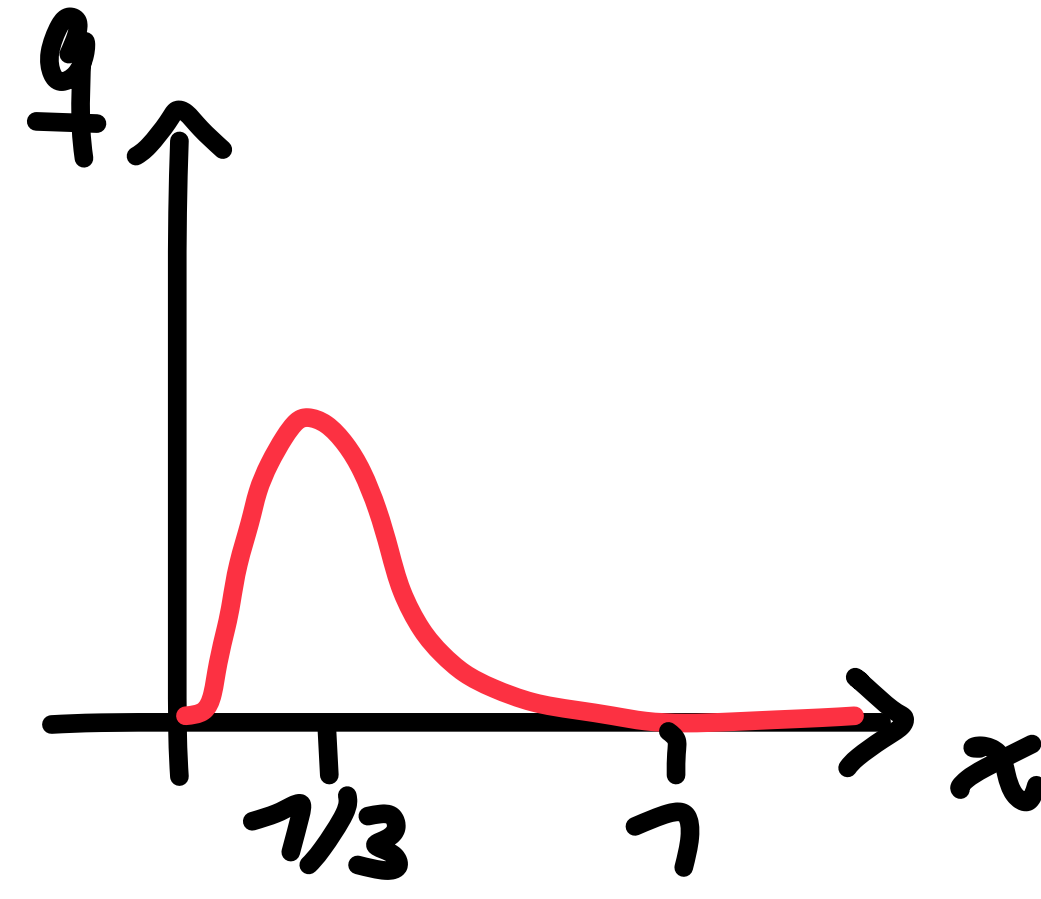
- Contains information about the **internal structure** of the proton



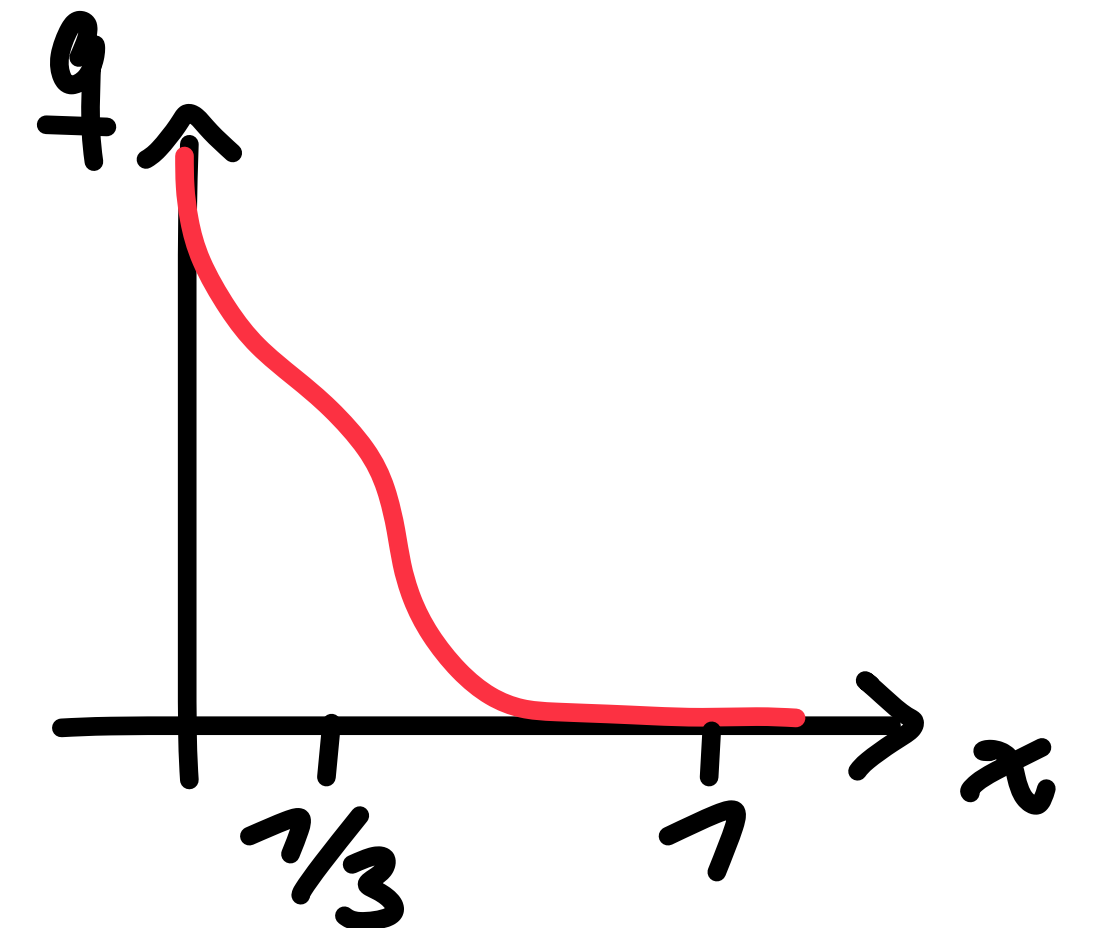
A fundamental proton



A proton consisting of 3 free quarks



A proton consisting of 3 bound quarks



A proton consisting of  
\* 3 bound quarks  
\* sea quarks  
\* self-interacting gluons

# Lattice PDFs: Quasi-PDFs

- We cannot achieve the Bjorken limit: Infinite momentum
- Instead we use large but **finite momenta**

PDF:  $q(x)$  from structure functions in Bjorken limit

$$\text{Quasi-PDF: } \tilde{q}(x, P_3) = \int_{-\infty}^{\infty} \frac{dz}{4\pi} e^{-izxP_3} \langle p | \sigma(z, 0) | p \rangle$$

Without proof:  $q(x) = \tilde{q}(x, P_3) + \mathcal{O}\left(\frac{\alpha}{2\pi}\right)$

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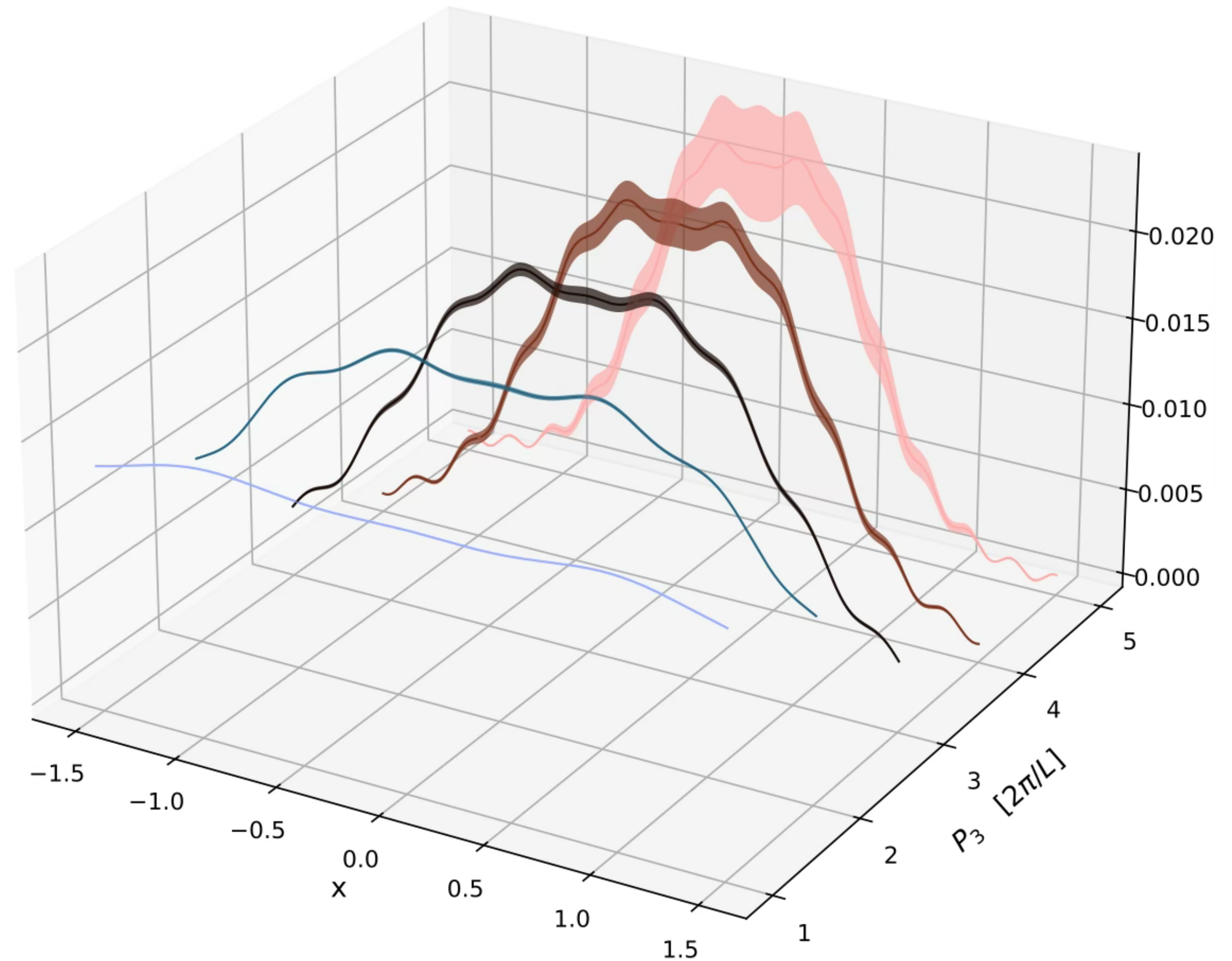
$$q(x) = \tilde{q}(x, P_3) + \mathcal{O}\left(\frac{\alpha}{2\pi}\right)$$

Strong interaction  
=  
Strong effect

Weak interaction  
=  
better control!

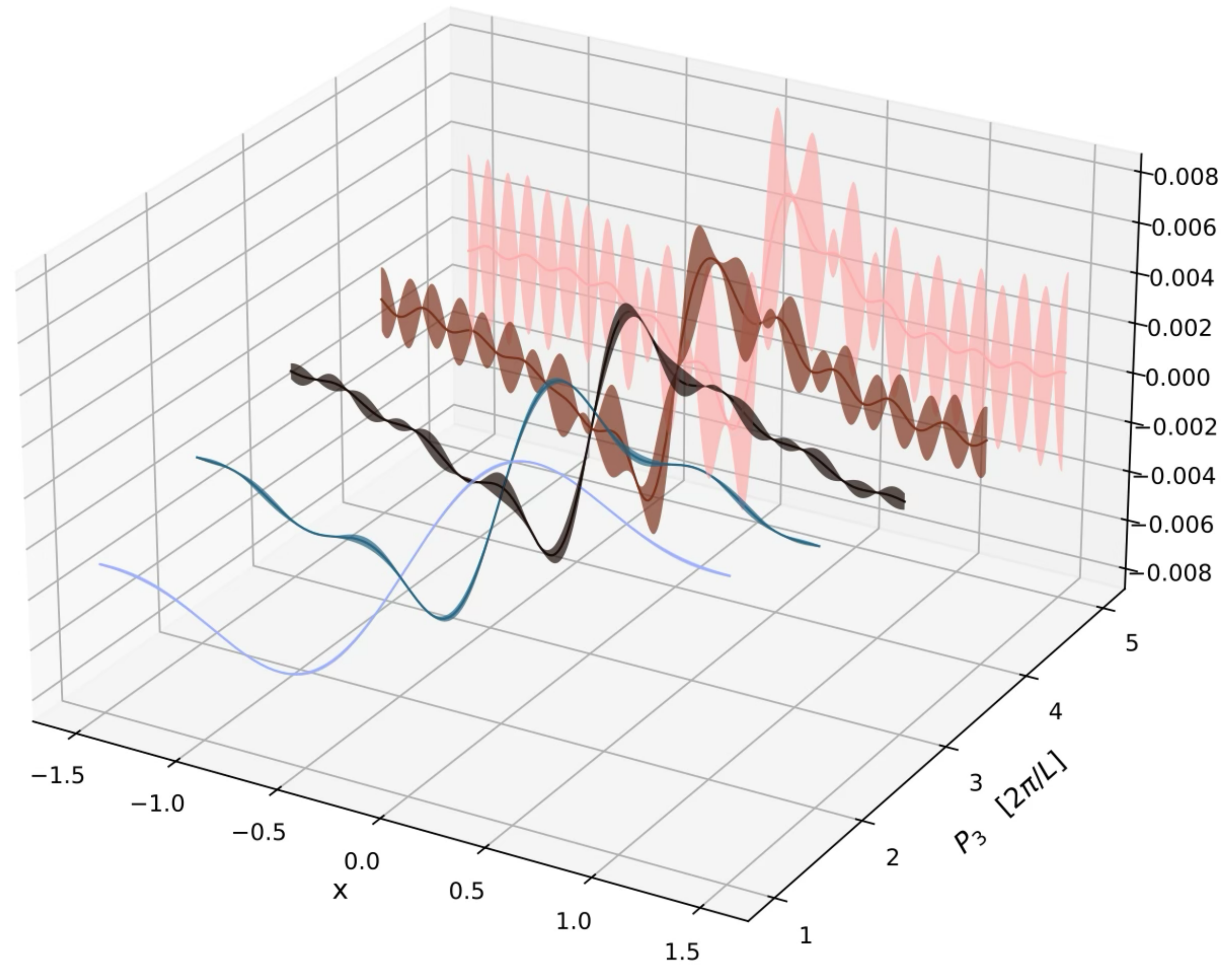
# FMS W-boson PDF

- $x\tilde{w}_{W_{\parallel}}^{(2)}(x; P_3)$  for  $L_T = 16$
- Shows the internal structure of physical W under the assumption of FMS



# FMS Higgs-boson PDF

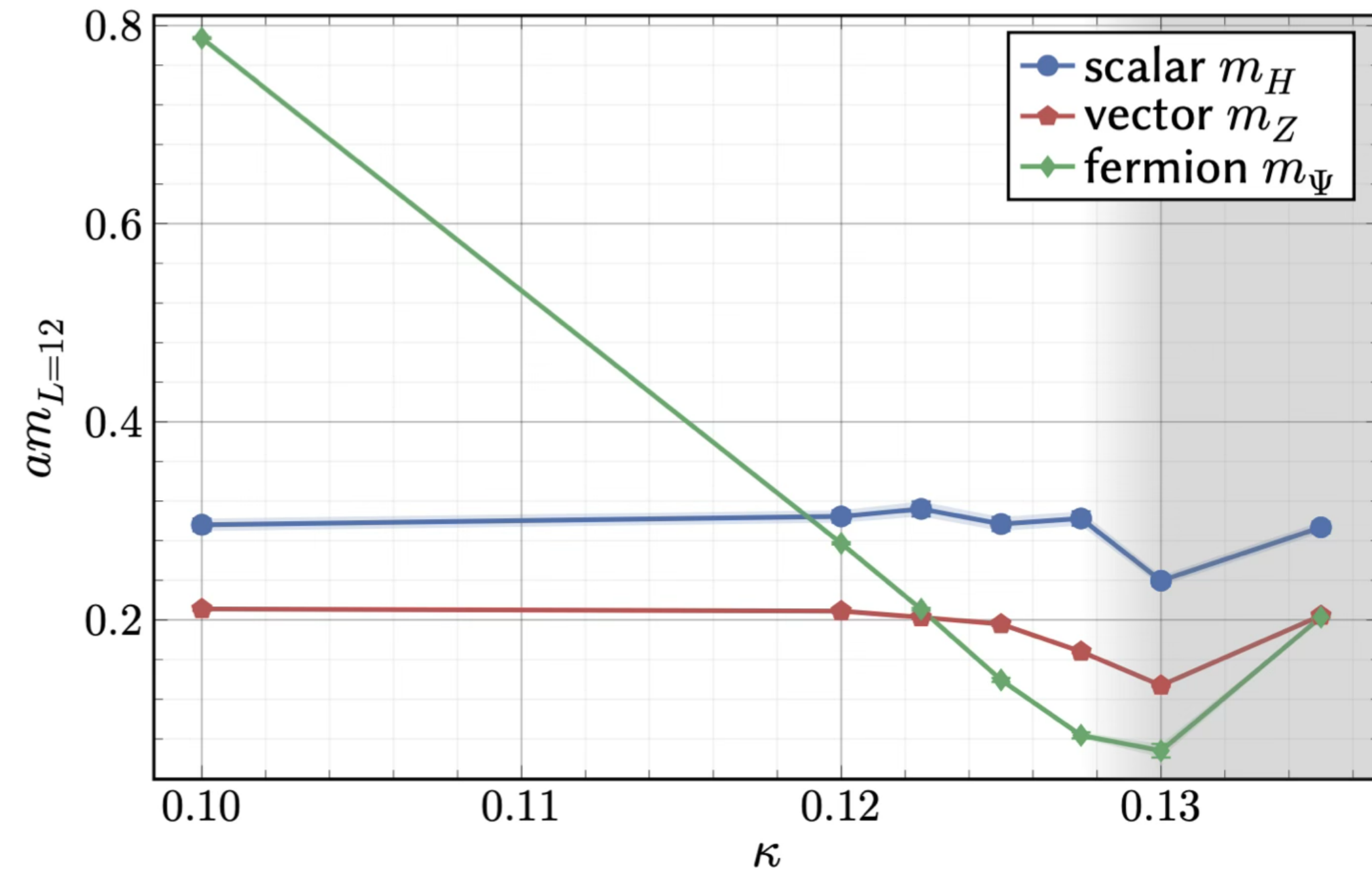
- $x\tilde{h}_H^{(1)}(x; P_3)$  for  $L_T = 16$
- We may be able to compare this to results from the High-Luminosity LHC



# Lattice Spectroscopy

# Parameter Regions

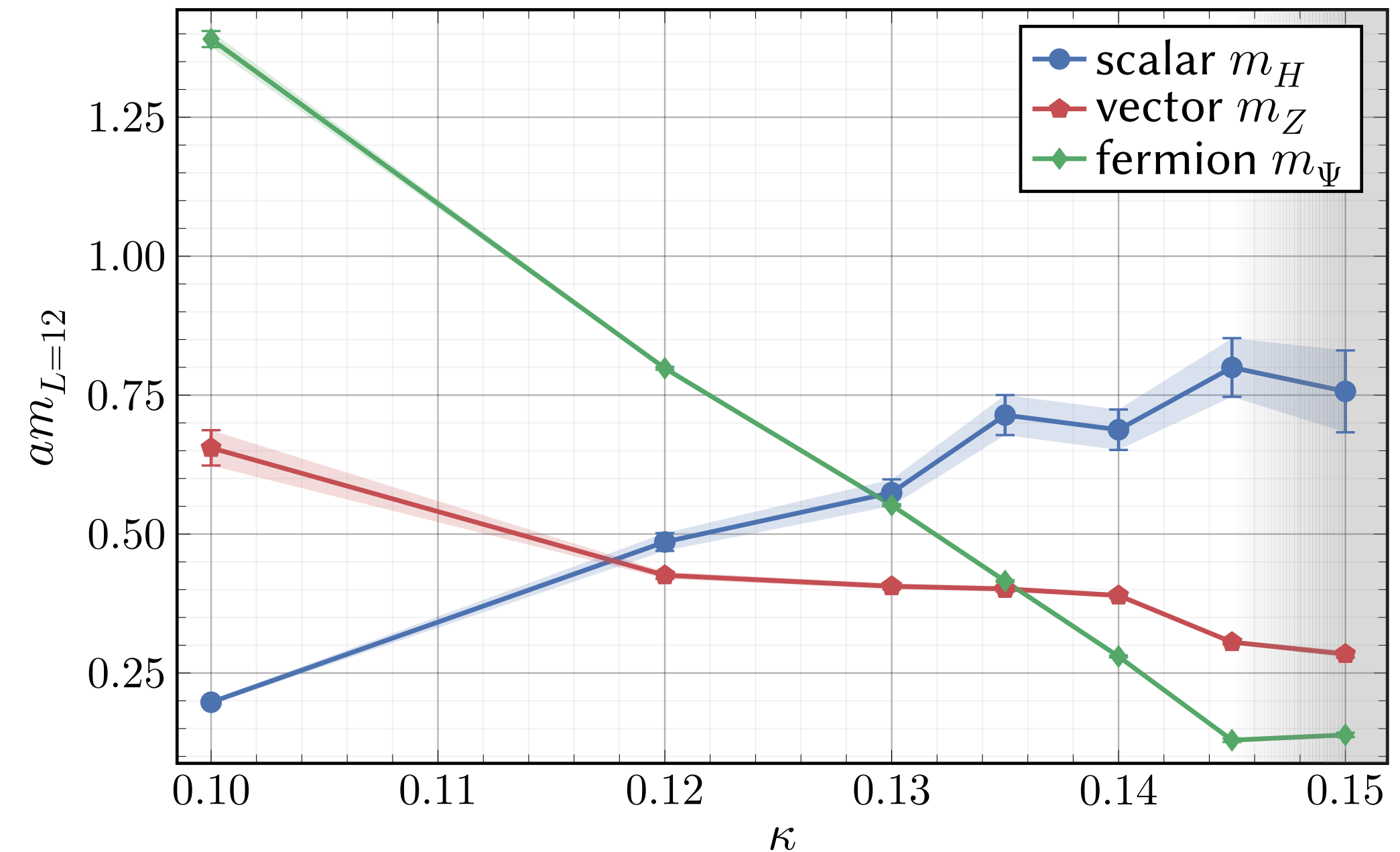
Phase space, without GEVP, future precision will be increased



**Higgs-like region**

stable Higgs

Scalar = Higgs much heavier than vector = W/Z



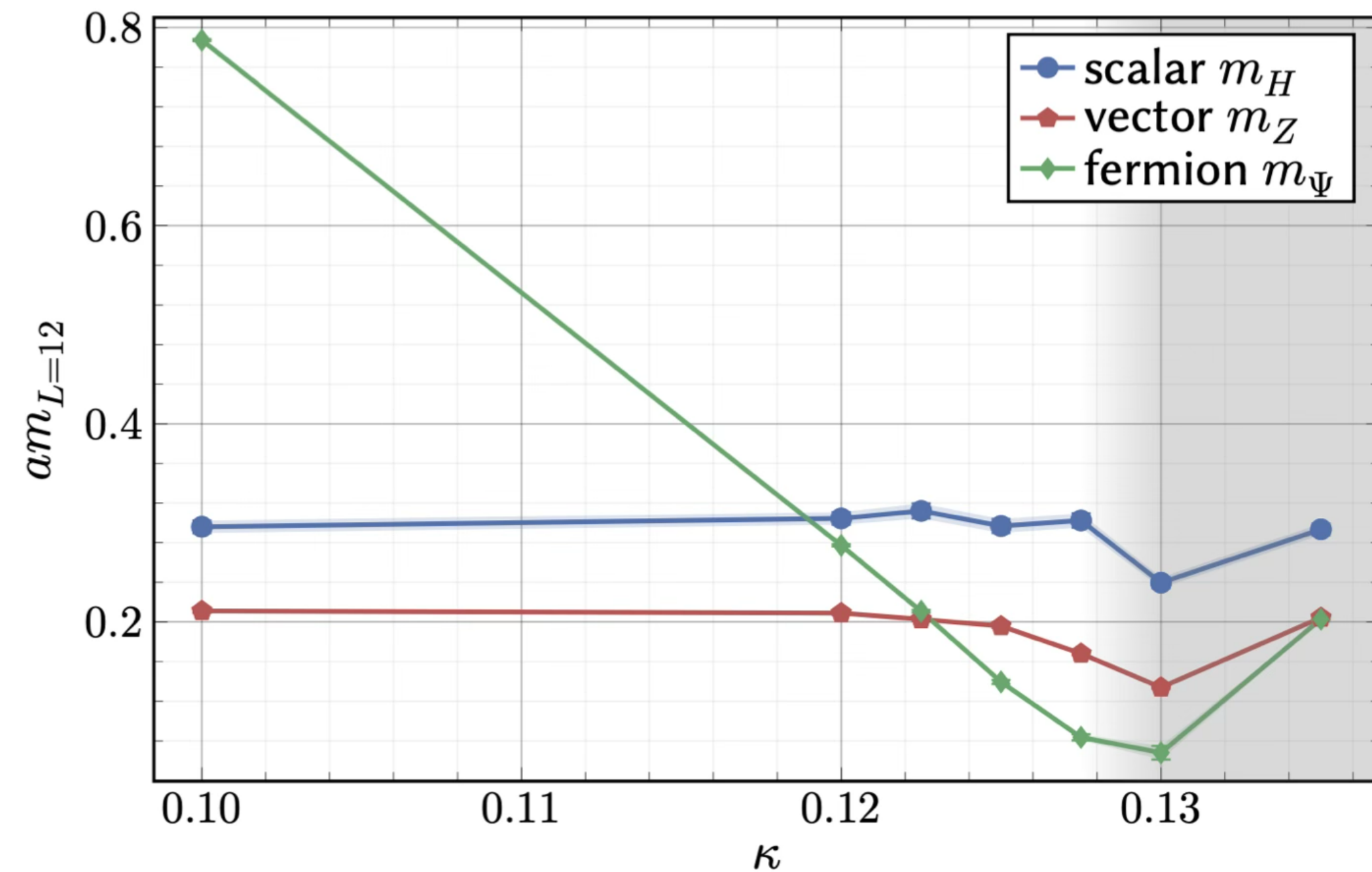
**QCD-like region**

inverted mass hierarchy with vector heavier than scalar

[Wieland, Maas, PoS, 2025]

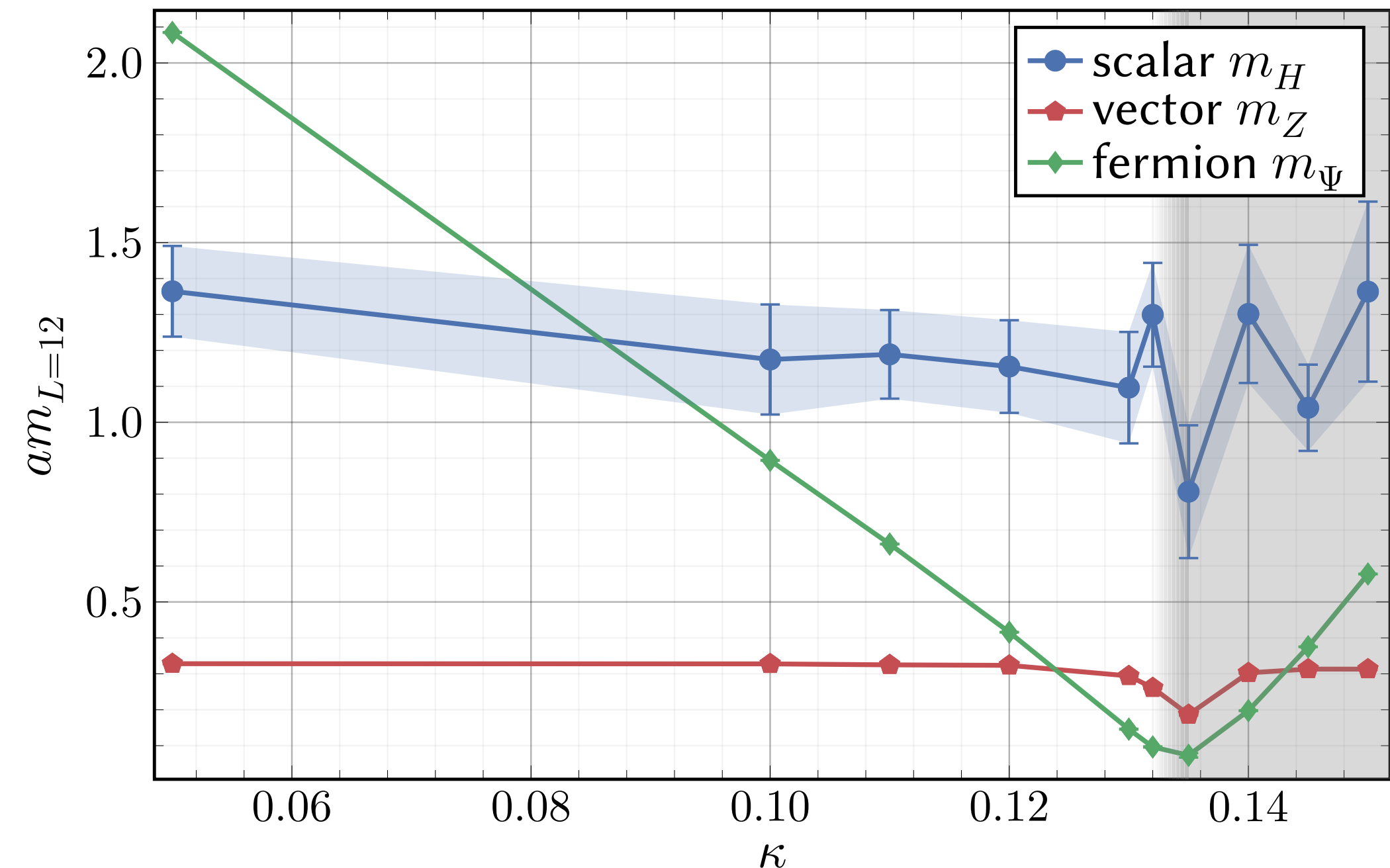
# Parameter Regions

Phase space, without GEVP, future precision will be increased



## Stable Higgs

Higgs mass lighter than two vector boson state



## Unstable Higgs

Higgs mass heavier than two vector boson state

[Wieland, Maas, PoS, 2025]

# Structure of asymptotic states

We expect the following states to be the physical Higgs, W/Z-bosons and leptons

$$\mathcal{O}_{0^+}(x) = \frac{1}{2} \text{tr} [\chi^\dagger(x) \chi(x)] = \phi^\dagger(x) \phi(x) \quad \text{HIGGS}$$

$$\mathcal{O}_{1^-}^{\alpha, \mu}(x) = -\frac{i}{2} \text{tr} [\tau^\alpha \chi^\dagger(x) \mathcal{U}^\mu(x) \chi(x + \hat{\mu})] \quad \text{W/Z}$$

$$\mathcal{O}_{\frac{1}{2}^-}^{i, \alpha}(x) = \chi_i^{\dagger i}(x) \psi^{i, \alpha}(x) \quad e, \nu_e$$

# Stable Higgs: structure of asymptotic states

A larger operator basis has largely no impact on ground state masses and unquenching effects are small

$$O_{\partial t}(p=0)$$

$$O_{\partial t}(p=0)O_{\partial t}(p=0)$$

$$O_{\partial_1-}(p=0)O_{\partial_1-}(p=0)$$

$$O_{\partial t}(p=(0,0,1))O_{\partial t}(p=(0,0,-1))$$

$$O_{\partial_1-}(p=(0,0,1))O_{\partial_1-}(p=(0,0,-1))$$

+ 2  $O_{\partial t}$  alternatives

	EVP Basis 1, Bootstrap	GEVP Basis 2, gamma method	Quenched, GEVP Basis 2, gamma
<b>Scalar mass</b>	0.296(10)	0.287(12)	0.299(14)
<b>Vector Mass</b>	0.1990(16)	0.1969(12)	0.1971(13)

Preliminary

See [Jenny, Maas, Riederer, Phys. Rev. D., 2022]

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Basis 1

$$O_{\partial t}(p=0)$$

$$O_{\partial t}(p=0)O_{\partial t}(p=0)$$

$$O_{1-}(p=0)O_{1-}(p=0)$$

$$O_{\partial t}(p=(0,0,1))O_{\partial t}(p=(0,0,-1))$$

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$$O_{\sigma^+}(p=0)$$

$$O_{\sigma^+}(p=0)O_{\sigma^+}(p=0)$$

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Basis 2

$$O_{\sigma^+}(p=(0,0,1))O_{\sigma^+}(p=(0,0,-1))$$

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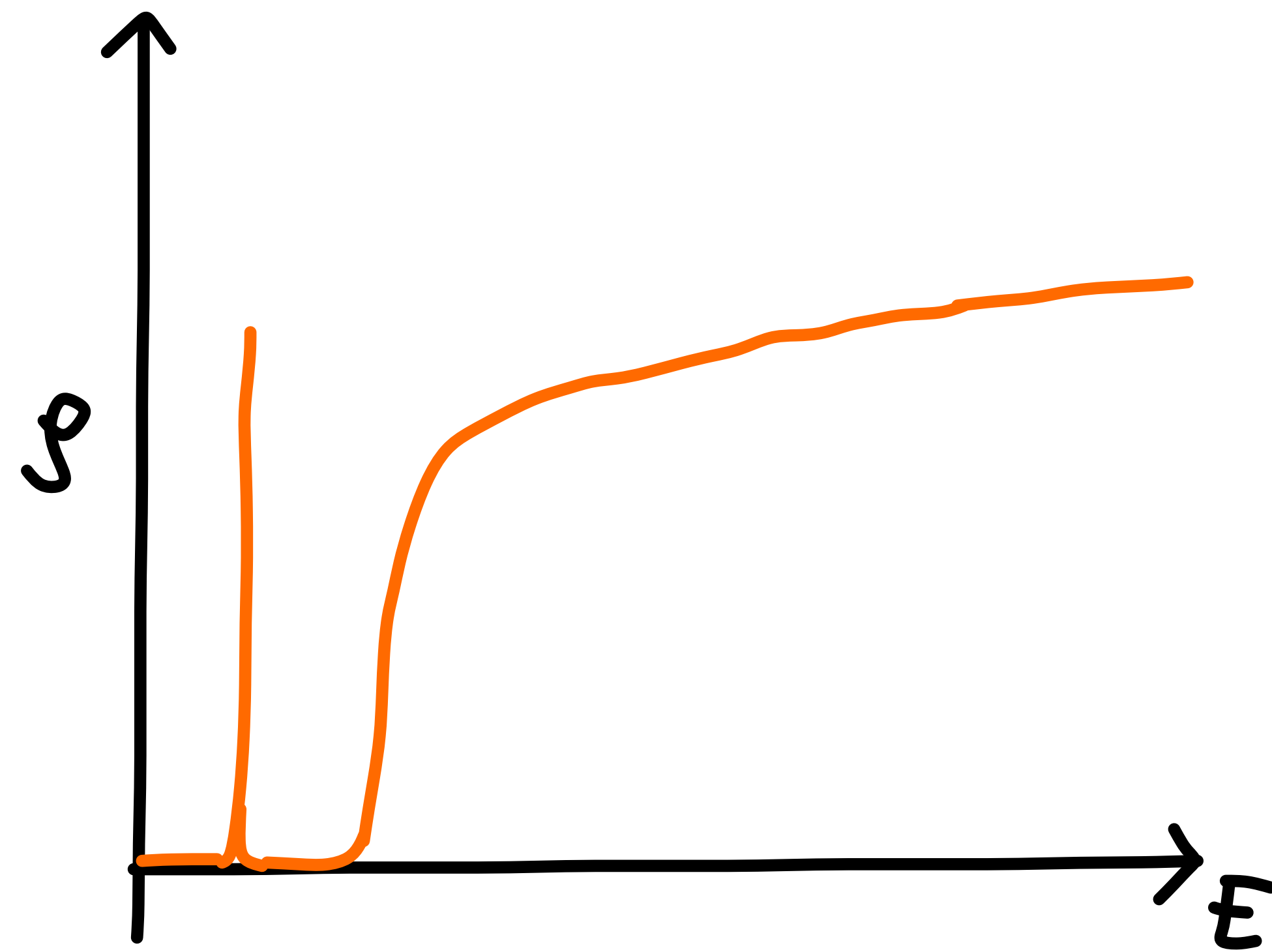
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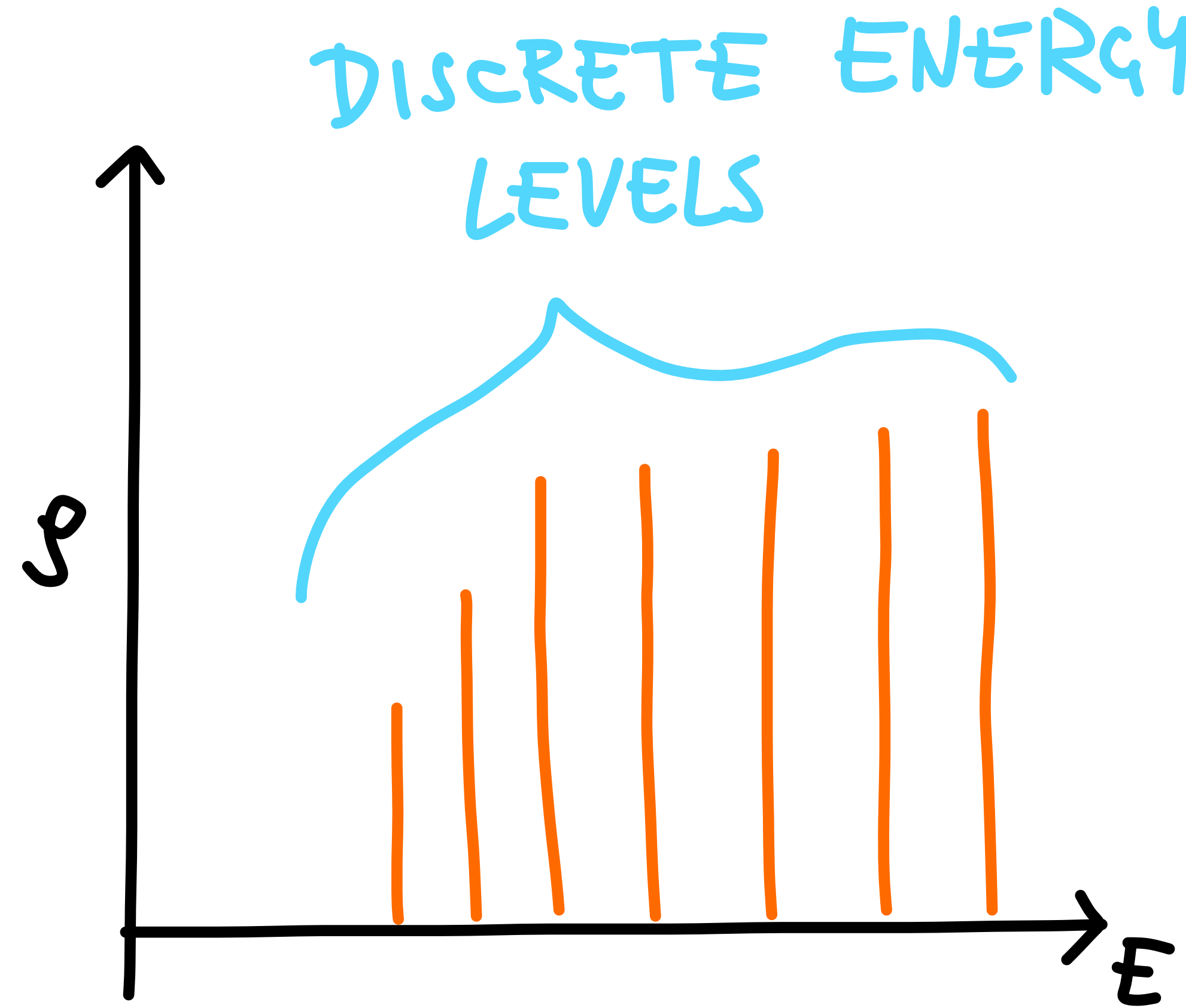
# Spectral densities

# Backus Gilbert Approach

[Backus, Gilbert, 1968 & 1970]



Continuum



Lattice / Finite Volume

# Hansen-Lupo-Tantalo Method

Expand kernel in functional basis preserving backpropagation

$$\Delta(E_x, E) = \sum_{t=0}^1 g_t(E_x) [e^{-tE} + e^{-(T-t)E}]$$

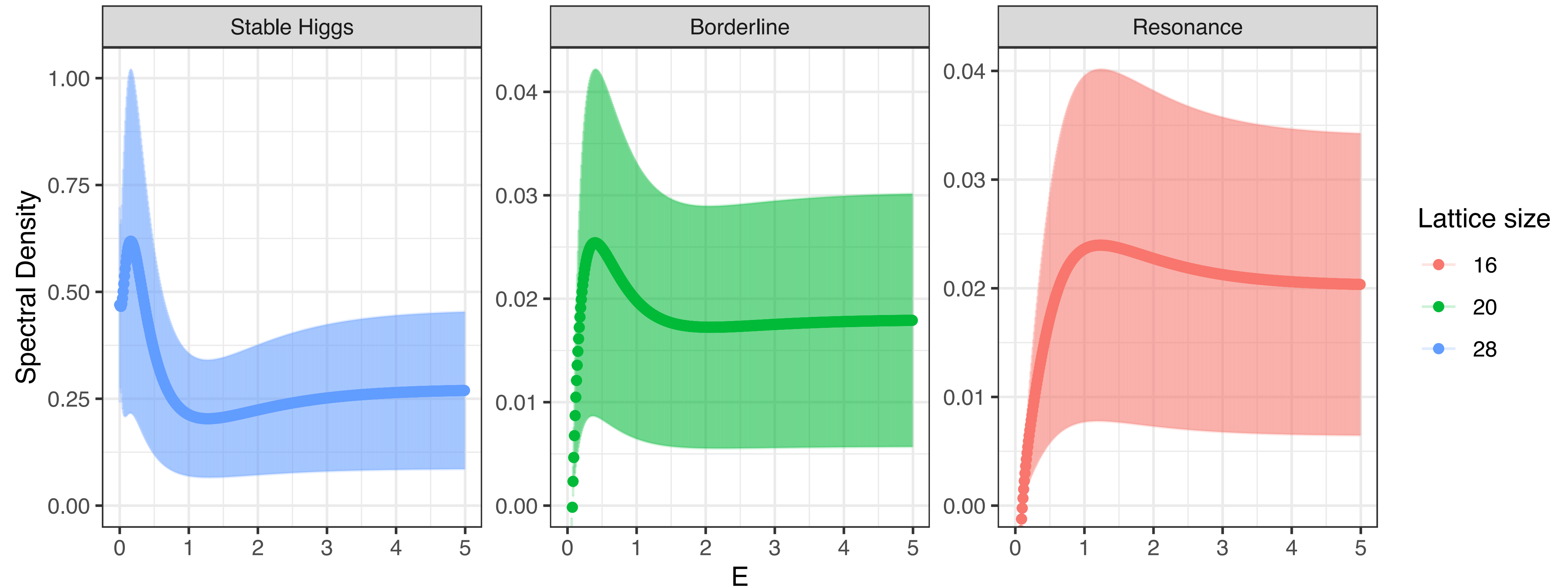
Balance **systematic** and **statistical** uncertainties

$$[g] = \int_{E_0}^{\infty} dE |\overline{\Delta}_\sigma(E_x, E) - \Delta_\sigma(E_x, E)|^2$$

$$\frac{\text{Var}(c(t))}{c^2(0)}$$

# Spectral Density Functions

Obtained using HLT, statistic and systematic errors

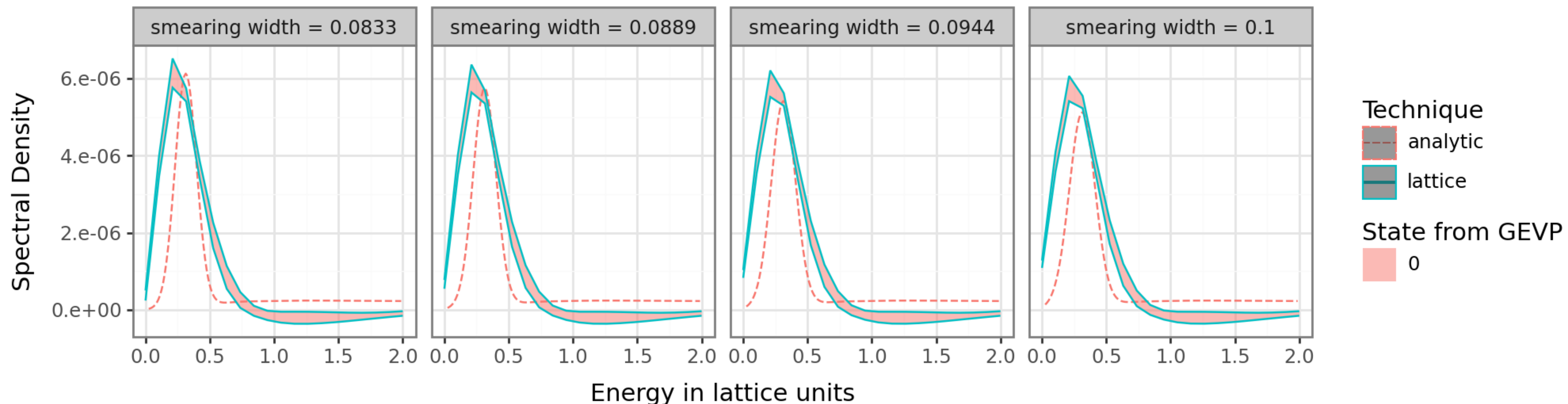


Technique: [Hansen, Lupo, Tantalò, Phys. Rev. D, 2019, 1903.06476]

Picture: [Martins et. al., submitted to PoS, 2026]

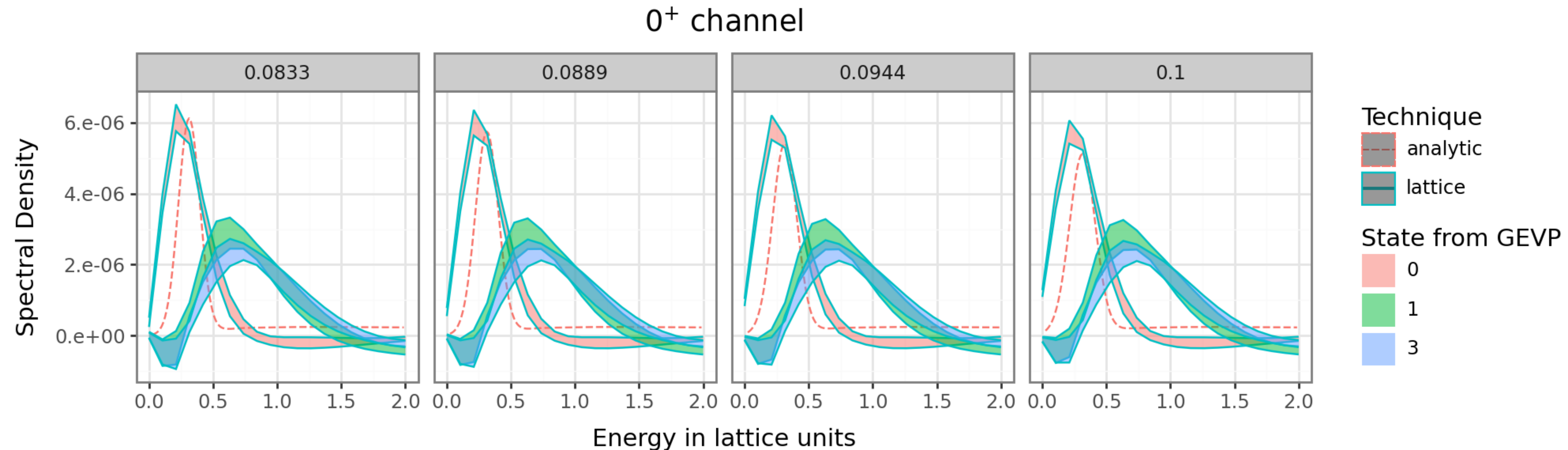
# Stable Higgs: Lattice vs. analytic

$0^+$  channel



Higgs spectral densities obtained from HLT with systematic effects suppressed compared with an analytic computation that was smeared with the same smearing parameter from [Maas, Sonderheimer, Phys. Rev. D., 2020]

# Stable Higgs: Lattice vs. analytic

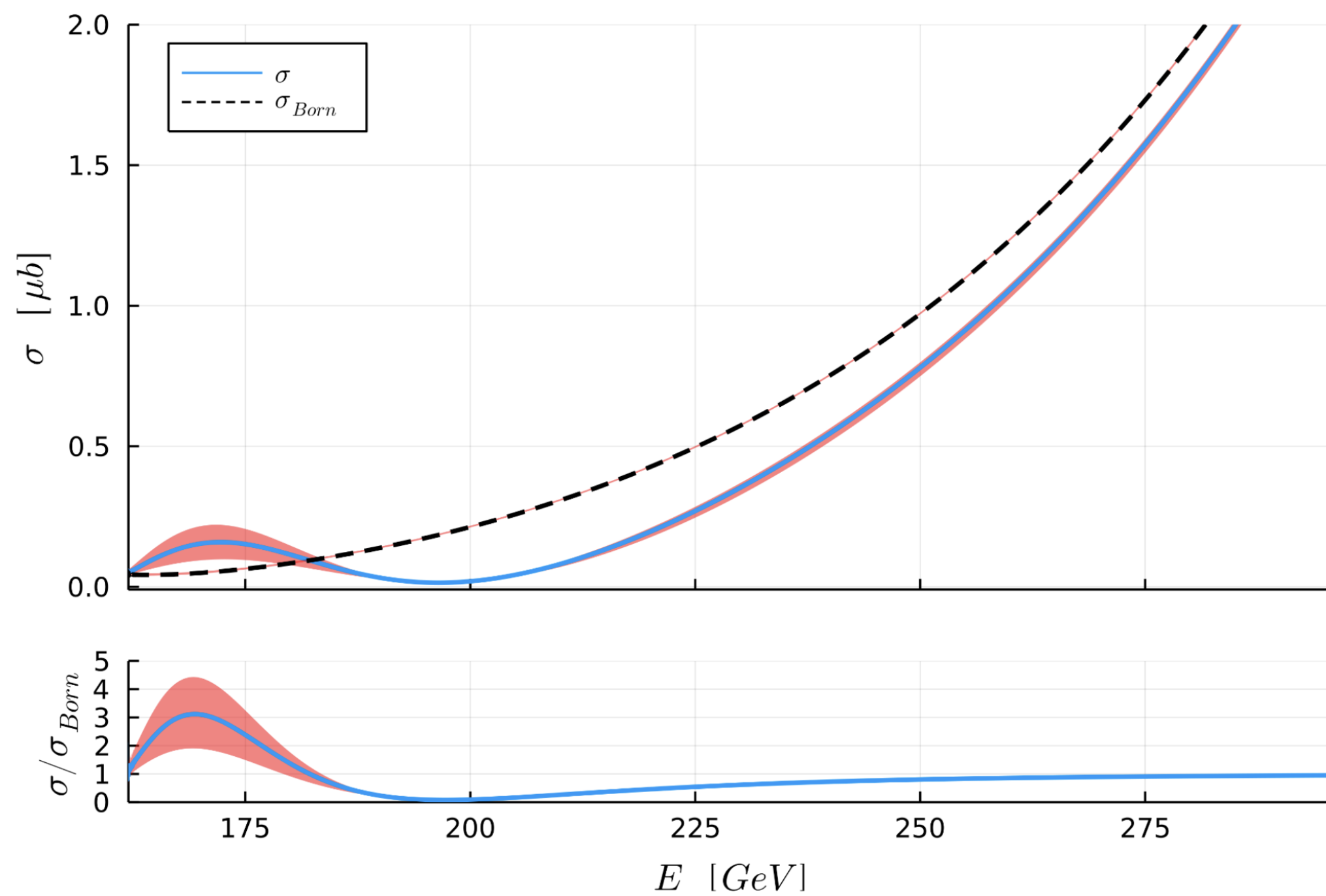


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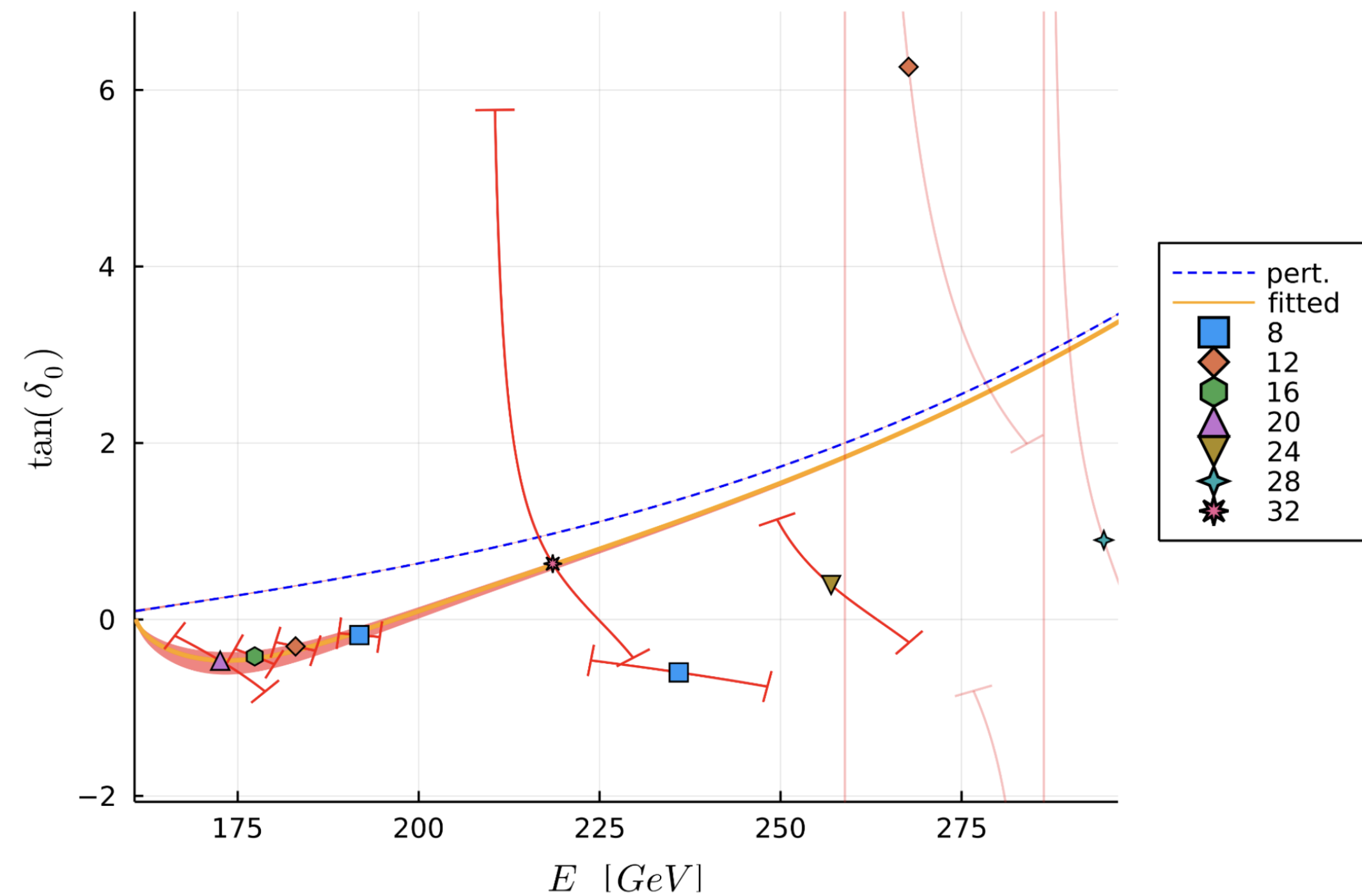
# Predictions for cross-sections of Higgs and Vector Boson Scattering

# Exclusive Decay Rates

Phenomenology predicts measurable differences in  $ZZ \rightarrow H \rightarrow ZZ$

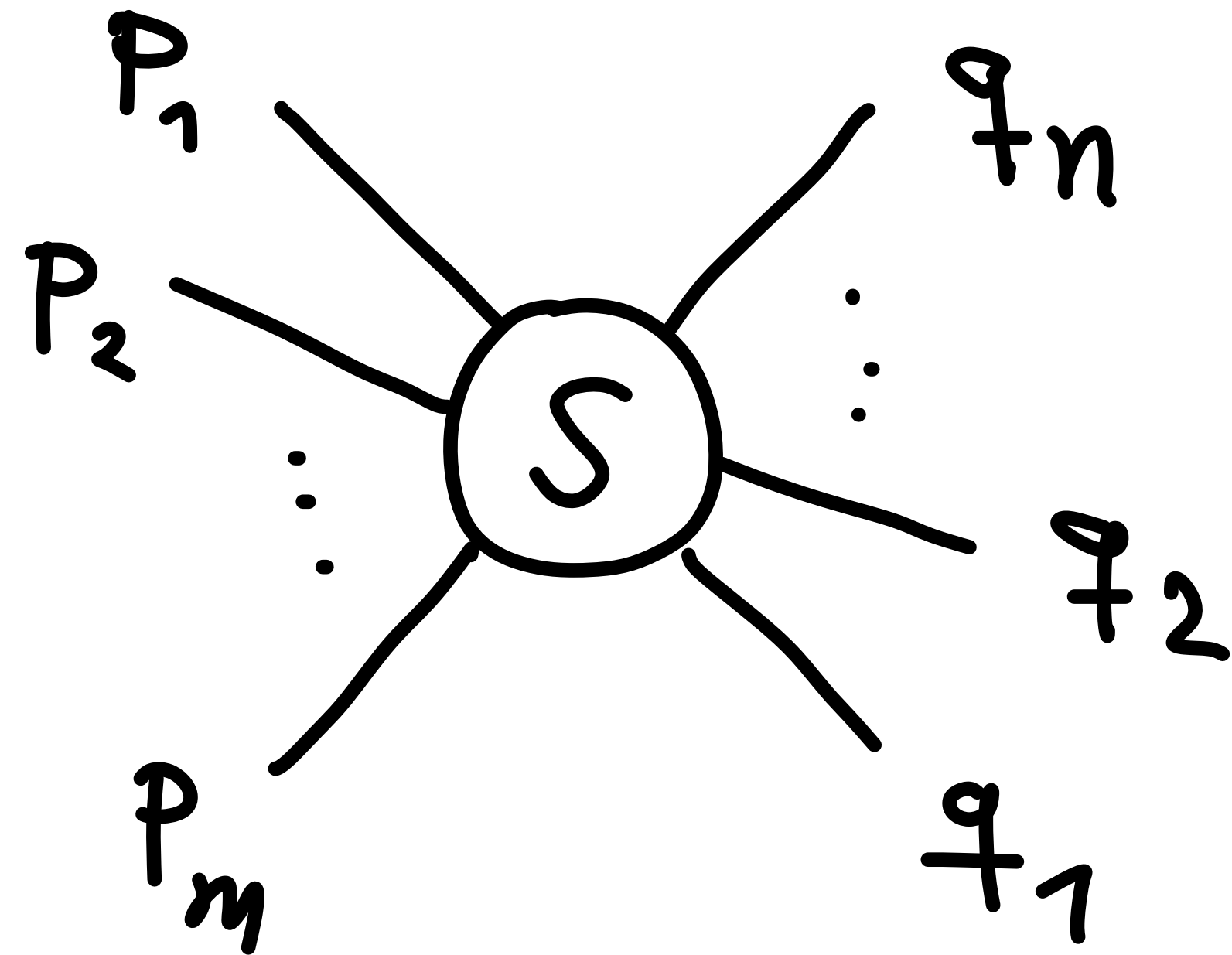


- Expected deviations @  $O(100 \text{ GeV})$
- Avoid mistaking signal for BSM physics



Figures: [Jenny, Maas, Riederer, Phys. Rev. D, 2022, 2204.02756]

# Computing a transition probability



$$\leadsto P = |\text{out} \langle \{p_f\} | \{q_i\} \rangle_{\text{in}}|^2$$

With definite momenta  $p_f$  and  $q_i$

Physical theory encoded in S-Matrix

# Bulava/Hansen method

[Bulava, Hansen, 2019]

$m$  particles out,  $n$  particles in

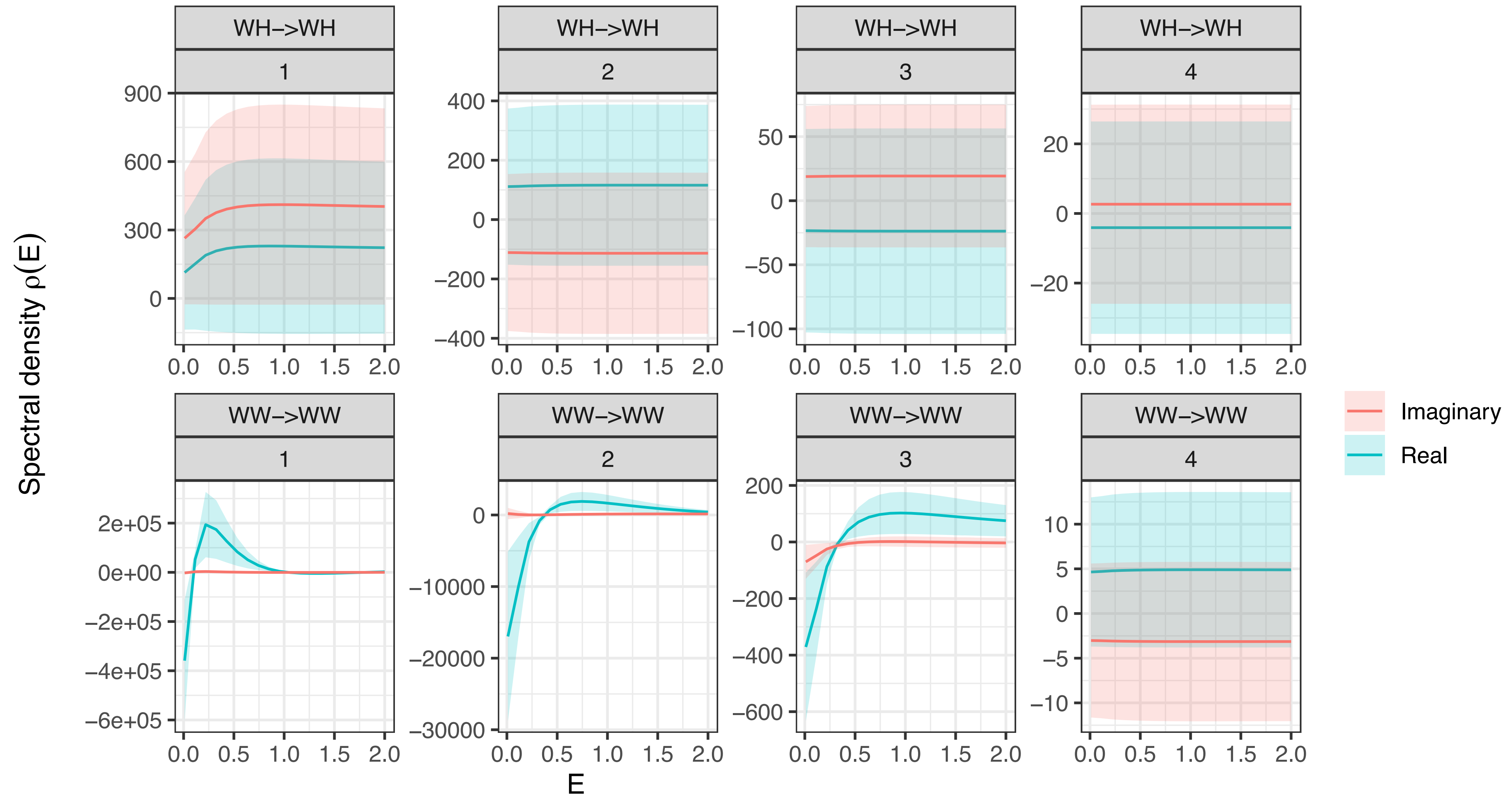
- Disassemble amplitude spectrally

$$i\mathcal{M}_c(\{P_f\}|\{q_i\}) \propto S_{P_m q_1}(\{P_f\}|P_m, \{q_i\}|q_1)$$

- Compute this spectral function by solving inverse problem

$$C_{P_m q_1}(\{P_f\}|P_m, \{q_i\}|q_1) = \int \frac{d^r E}{\pi^r} e^{-\Sigma E \Delta t} S_{P_m q_1}(\{P_f\}|P_m, \{q_i\}|q_1)$$

# Matrix element spectral decomposition



# Summary

- PDFs show internal structure that can be probed at the LHC
- Scalar + Gauge + Fermionic theory shows the expected phases and asymptotic states predicted by the FMS mechanism
- Augmented perturbation theory agrees with lattice spectral densities

# Outlook

- Compare cross sections with analytical results for simulations with dynamical fermions
- Yukawa couplings if we find excited states consistent with generations/flavors
- Mixing angles between excited leptonic states

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**Thank you for your attention!**