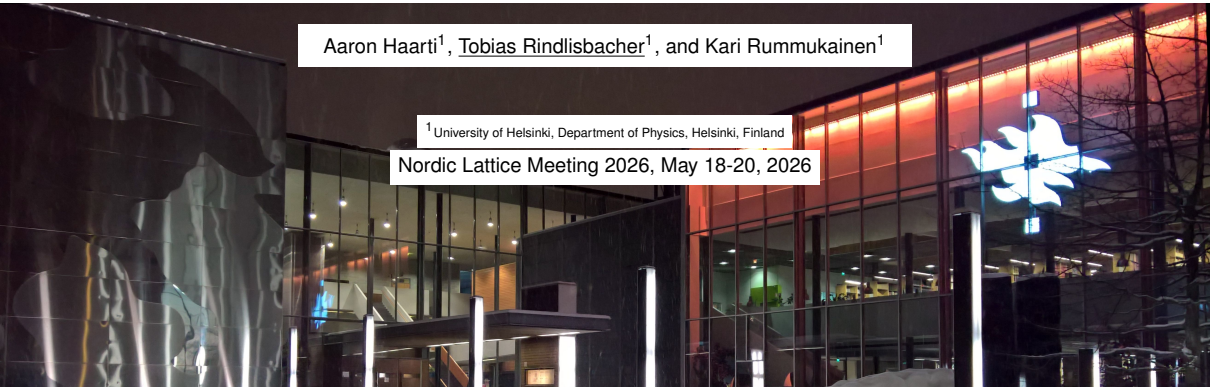


Measuring d-d interface tensions in $SU(N)$ gauge theory with the capillary wave method

Aaron Haarti¹, Tobias Rindlisbacher¹, and Kari Rummukainen¹

¹University of Helsinki, Department of Physics, Helsinki, Finland

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Capillary wave method for confined-deconfined interface tension

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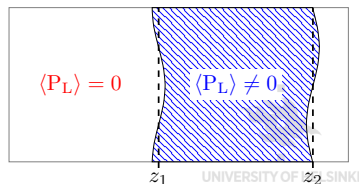
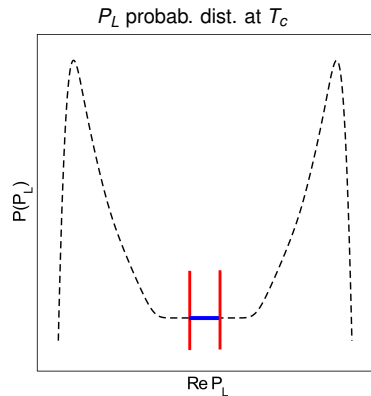
[Moore & Turok (1997)]

→ avoid supercritical slowing down by sampling mixed phase ensemble.

■ z-elongated lattice: $V = \underbrace{N_{x,y}^2}_{V_s} \times N_z \times N_t$ with $N_t \ll N_{x,y} \ll N_z$

■ pure $SU(N)$ gauge theory: $S_G[U] = \frac{\beta}{N} \sum_{x, \mu < \nu} \text{Re tr}(\mathbb{1} - U_{\mu\nu}(x))$

■ Monte Carlo (heatbath + overrelaxation)



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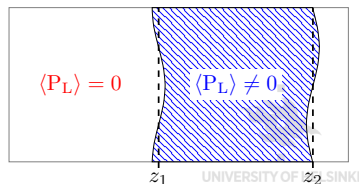
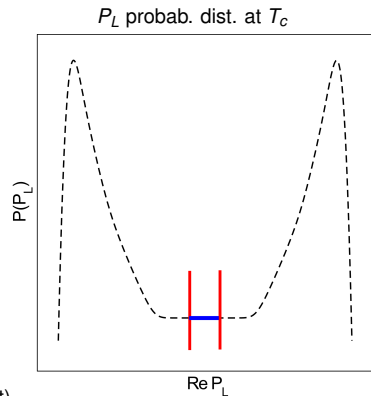
■ Initialize simulation:

1) induce mixed phase state: $\beta = \begin{cases} \tilde{\beta}_c - \Delta\beta & \text{if } z < N_z/2 \\ \tilde{\beta}_c + \Delta\beta & \text{if } z \geq N_z/2 \end{cases}$, $\tilde{\beta}_c \approx \beta_c$ -estimate

2) set $\beta = \tilde{\beta}_c \forall z$ and impose constraint $\langle P_L \rangle \in [(1/2 - c)P_0, (1/2 + c)P_0]$
where $P_0 = |\langle P_L \rangle_{\text{deconf.}}|$, $c \sim 10^{-2}$. (constr. implemented via accept/reject)

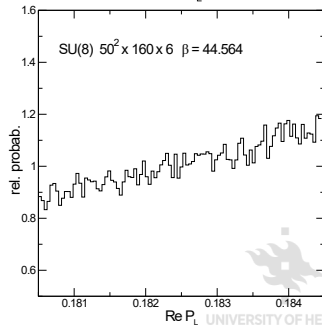
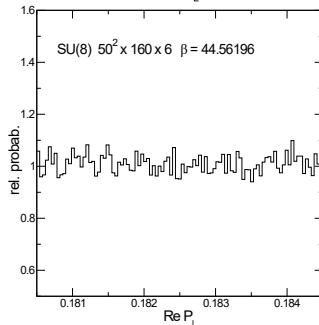
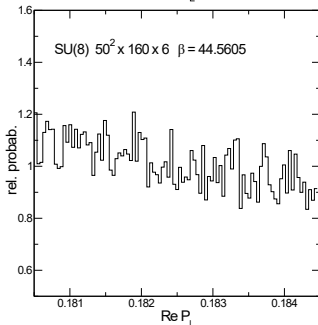
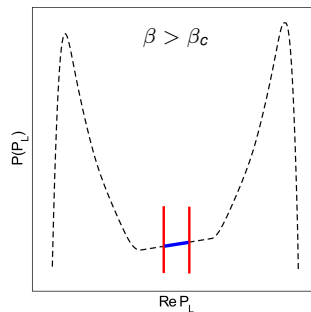
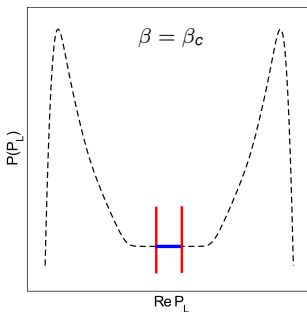
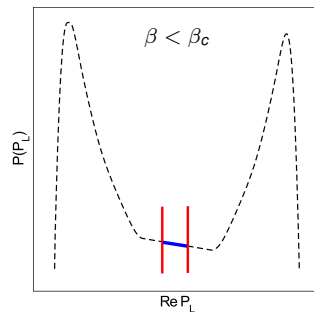
3) thermalize (with constraint)

■ Generate mixed phased configurations for measurements (with constraint)



Capillary wave method for confined-deconfined interface tension

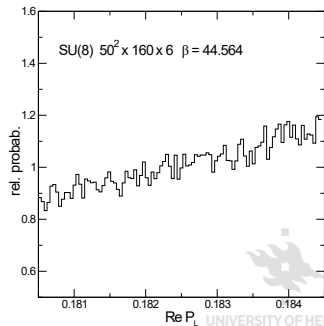
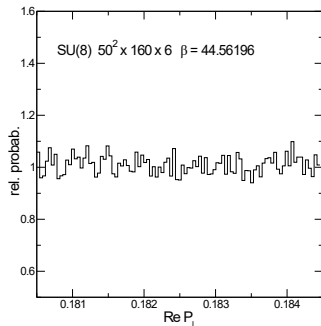
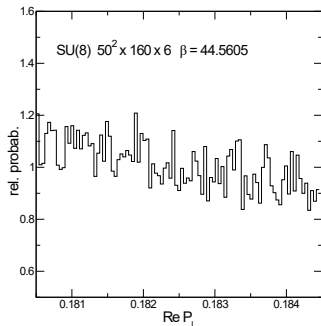
Determining the critical inverse coupling β_c



Capillary wave method for confined-deconfined interface tension

Determining the critical inverse coupling β_c

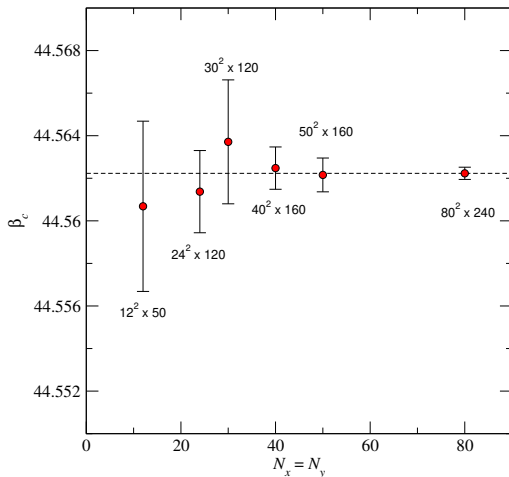
→ can use reweighting to find precise $\beta = \beta_c$ for which Polyakov loop histogram is flat.



Capillary wave method for confined-deconfined interface tension

Determining the critical inverse coupling β_c

- can use reweighting to find precise $\beta = \beta_c$ for which Polyakov loop histogram is flat.
- exponentially suppressed finite volume effects



Surface tension from surface fluctuation spectrum

[Moore & Turok (1997)]

Fluctuating thin surface of size L^2 , parametrized over x-y-plane.

Surface fluctuations $z(x, y)$ at finite temp. $T \rightarrow$ two-dimensional field theory: $Z = \int \prod_{x,y} dz(x, y) e^{-H/T}$

with Hamiltonian:

$$H = \sigma \int_0^L dx dy \sqrt{1 + |\nabla z|^2} = \sigma \int_0^L dx dy \left[1 + \frac{1}{2} |\nabla z|^2 - \frac{1}{8} |\nabla z|^4 + \dots \right].$$



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$\rightarrow Z \propto \prod_{\substack{n_x, n_y \\ n_x^2 + n_y^2 \neq 0}} \int d\hat{z}(n_x, n_y) e^{-|\hat{z}(n_x, n_y)|^2 4\pi^2 n^2 \sigma / T}$ s.t. $\langle |\hat{z}(n_x, n_y)|^2 \rangle = \frac{T}{4\pi^2 n^2 \sigma}$, $n^2 \equiv n_x^2 + n_y^2 \neq 0$



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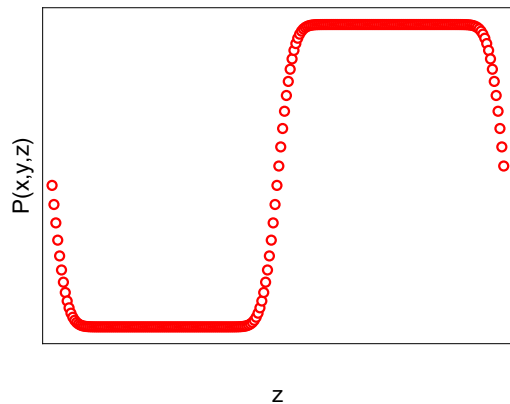
\rightarrow Surface tension from surface fluctuation spectrum: $\sigma = \frac{T}{4\pi^2 n^2 \langle |\hat{z}(n_x, n_y)|^2 \rangle}$, $n^2 \equiv n_x^2 + n_y^2 \neq 0$, $n^2 \rightarrow 0$

\rightarrow Effects from non-renormalizable terms in H expected to vanish as $O(n^2/L^2)$.



Interface detection and smearing

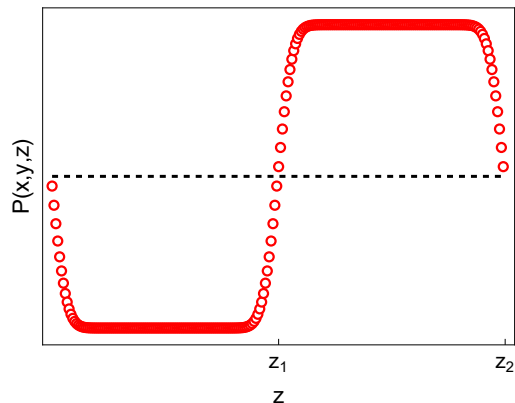
For each (x, y) detect interface from Polyakov loop profile as function of z



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Expect to find $z_i(x, y)$, $i = 1, 2$ so that $P(x, y, z_i(x, y)) = P_0$ (with $P_0 = |\langle P \rangle_{\text{deconf.}}|/2$)



Confined-deconfined interface tension

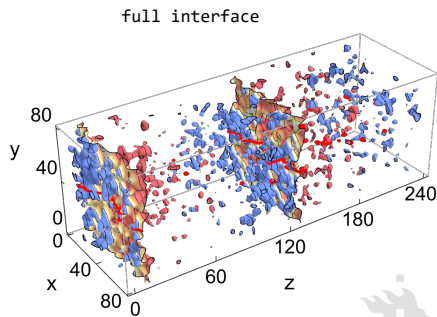
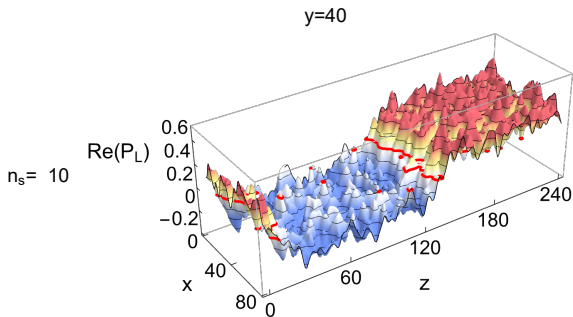
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Polyakov loop field very noisy \Rightarrow perform n_s iter. of recursive smearing $P_L \rightarrow P_L^{(n_s)}$

Smearing iter.: $P_L^{(n+1)}(\bar{x}) = \sum_{\bar{y}} S(\bar{x}, \bar{y}) P_L^{(n)}(\bar{y})$, kernel: $S(\bar{x}, \bar{y}) = \frac{1}{1 + 6\rho} \left(\delta_{\bar{x}, \bar{y}} + \rho \sum_{i=1}^3 (\delta_{\bar{x}+\hat{i}, \bar{y}} + \delta_{\bar{x}-\hat{i}, \bar{y}}) \right)$.



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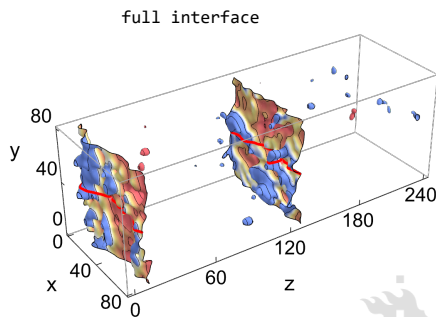
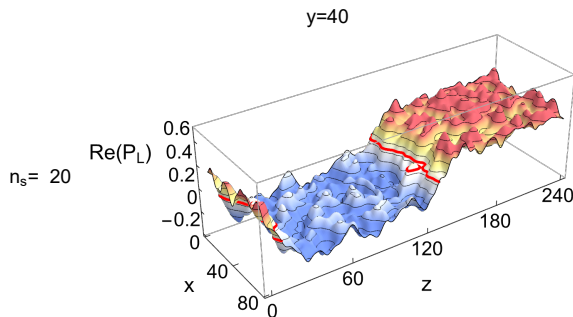
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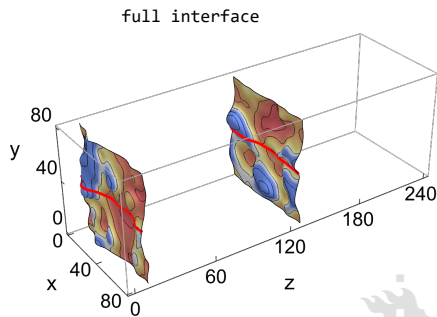
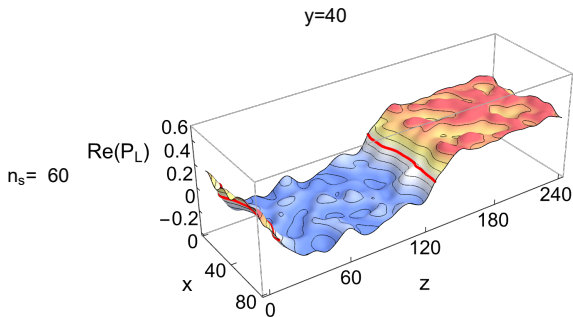
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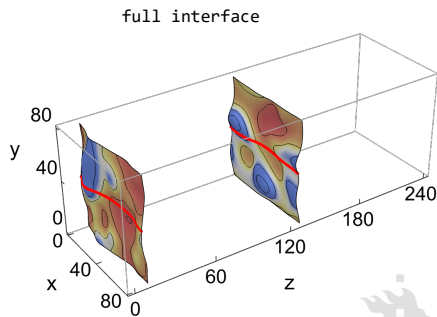
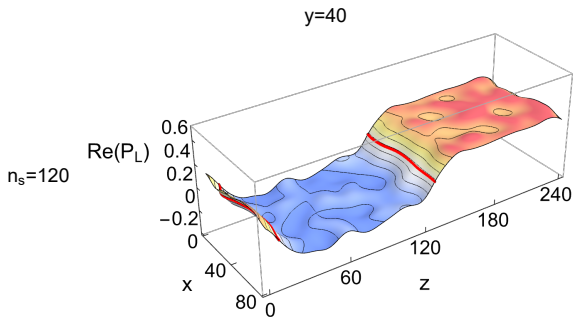
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Effect of smearing on surface fluctuation spectrum?



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→ idealized situation: Polyakov loop field $P(x, y, z)$ given by:

$$A(x, y, z) = \begin{cases} A_0 & \text{if } z_1(x, y) \leq z < z_2(x, y) \\ 0 & \text{else} \end{cases}$$

with interface parametrizations $z_i(x, y)$, $i = 1, 2$, $A_0 = |\langle P_L \rangle_{\text{deconf.}}|$ (neglecting bulk fluctuations)



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Fourier transform in z : $\hat{A}(x, y, k_z) = \frac{i A_0}{2 \pi k_z} \left(e^{i k_z z_1(x, y)} - e^{i k_z z_2(x, y)} \right)$



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→
$$\hat{A}(x, y, k_z) = \frac{A_0}{2 \pi} \left(-\frac{e^{i k_z z_{0,1}} - e^{i k_z z_{0,2}}}{i k_z} - e^{i k_z z_{0,1}} \Delta z_1(x, y) + e^{i k_z z_{0,2}} \Delta z_2(x, y) + \mathcal{O}(k_z \Delta z_i^2(x, y)) \right)$$



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Fourier transform also in (x, y) :

→
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→
$$\hat{z}_i(k_x, k_y)^{(\bar{k}^2 > 0)} \Delta \hat{z}_i(k_x, k_y) \propto \hat{A}(k_x, k_y, k_z) \quad , \quad i = 1, 2$$

→
$$\hat{A}^{(n_s)}(k_x, k_y, k_z) = \hat{S}^{n_s}(k_x, k_y, k_z) \hat{A}(k_x, k_y, k_z) \quad \Rightarrow \quad \hat{z}_i^{(n_s)}(k_x, k_y) = \hat{S}^{n_s}(k_x, k_y, k_z \approx 0) \hat{z}_i(k_x, k_y)$$



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→ Kernel corrected estimator for
$$\langle |\hat{z}(k_x, k_y)|^2 \rangle = \frac{1}{\hat{S}^{2 n_s}(k_x, k_y, k_z \approx 0)} \langle |\hat{z}^{(n_s)}(k_x, k_y)|^2 \rangle$$



Zero-mode extrapolation

Fit ansatz: $f(x) = c_1 \exp(c_2 x + c_3 x^2)$ with $x = k^2/T^2$



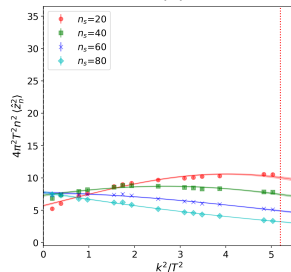
Confined-deconfined interface tension

Zero-mode extrapolation

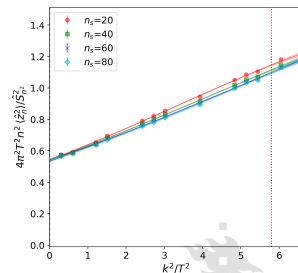
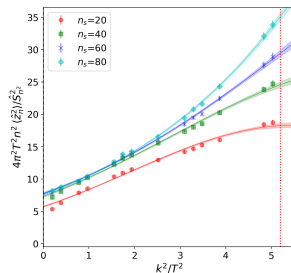
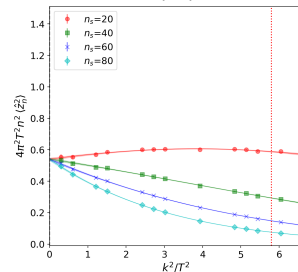
Fit ansatz: $f(x) = c_1 \exp(c_2 x + c_3 x^2)$ with $x = k^2/T^2$

N	k.-corr.	n_s	c_1	c_2	c_3	$\chi^2/\text{#dof}$
4	no	20	5.71(11)	0.308(16)	-0.0383(28)	3.792
		40	7.31(15)	0.132(17)	-0.0247(30)	1.597
		60	7.71(16)	-0.014(17)	-0.0140(30)	1.047
		80	7.79(18)	-0.146(18)	-0.0046(31)	0.737
	yes	20	5.71(11)	0.423(16)	-0.0384(28)	3.787
		40	7.30(15)	0.362(17)	-0.0247(30)	1.573
		60	7.71(16)	0.331(17)	-0.0141(30)	1.007
		80	7.78(17)	0.315(18)	-0.0048(31)	0.695
10	no	20	0.5397(75)	0.0607(84)	-0.0079(11)	1.070
		40	0.5404(76)	-0.0682(85)	-0.0064(11)	1.462
		60	0.5397(77)	-0.1913(86)	-0.0053(11)	1.809
		80	0.5397(77)	-0.3110(86)	-0.0045(11)	2.155
	yes	20	0.5400(75)	0.1751(84)	-0.0078(11)	0.850
		40	0.5409(76)	0.1607(85)	-0.0062(11)	0.961
		60	0.5404(77)	0.1522(86)	-0.0052(11)	0.984
		80	0.5407(77)	0.1469(86)	-0.0043(11)	0.985

SU(4)



SU(10)

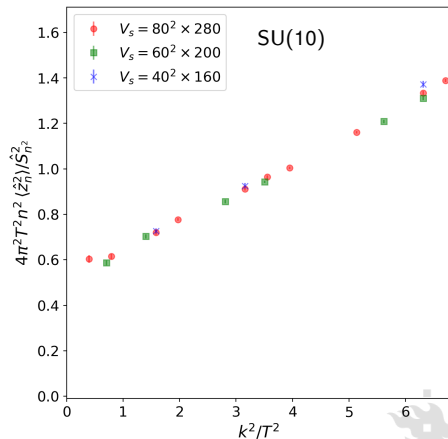
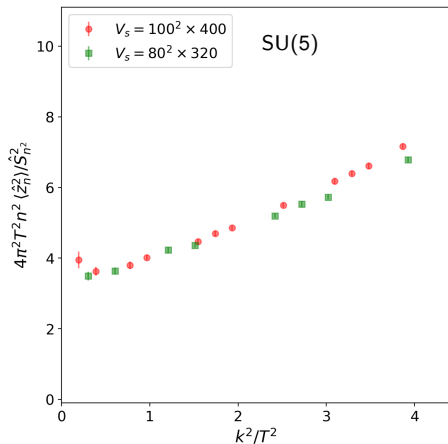


Confined-deconfined interface tension

Zero-mode extrapolation

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finite-volume effects?

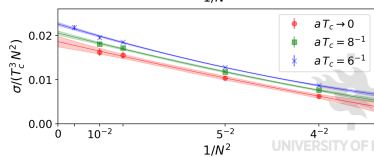
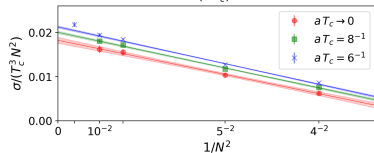
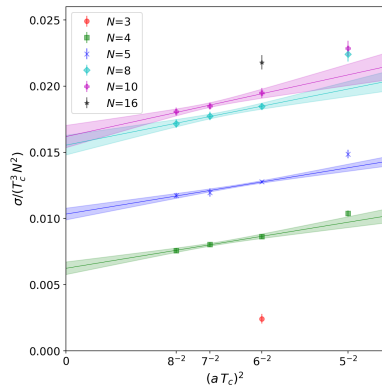


Confined-deconfined interface tension

Continuum limit

Fit ansatz: $f(x) = c_1 + c_2 x$ with $x = (a T_c)^2 = 1/N_t^2$, $N_t > 5$

(error bands: $\delta f(x) = \sqrt{\sum_{i,j} \sigma^{ij} \frac{\partial f(x)}{\partial c_i} \frac{\partial f(x)}{\partial c_j}}$, with covariance matrix σ^{ij})



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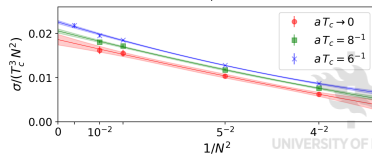
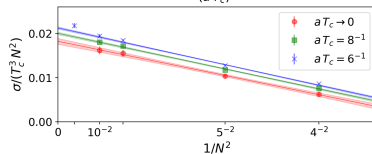
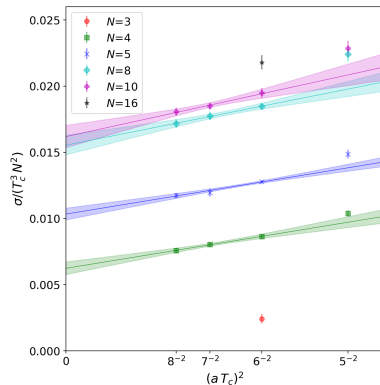
Large-N limit

Fit ansatz (linear): $f(x) = c_1 + c_2 x$ with $x = 1/N^2$,

→ $\frac{\sigma}{T^3} = 0.0182(7)N^2 - 0.194(15)$ (continuum)
($\chi^2/\text{\#dof} = 0.2$)

Fit ansatz (quadratic): $f(x) = c_1 + c_2 x + c_3 x^2$ with $x = 1/N^2$

→ $\frac{\sigma}{T^3} = 0.0186(12)N^2 - 0.220(71) + 0.35(92)/N^2$ (continuum)
($\chi^2/\text{\#dof} = 0.25$)



Deconfined-deconfined interface tension

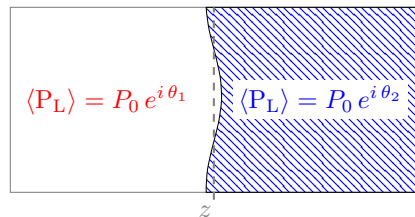
Capillary wave method for deconfined-deconfined interface tension

- $\beta > \beta_c \rightarrow$ Interface between two deconfined Z_N -vacua, corresponding to $\langle P_L \rangle = P_0 e^{i\theta_i}$, $i = 1, 2$

$$\theta_i = 2\pi n_i/N, \quad n_i \in \{0, \dots, N-1\}, \quad n_1 \neq n_2.$$

- periodic boundary conditions with Z_N -twist Δn in z -direction.

\rightarrow interface always across x - y -plane



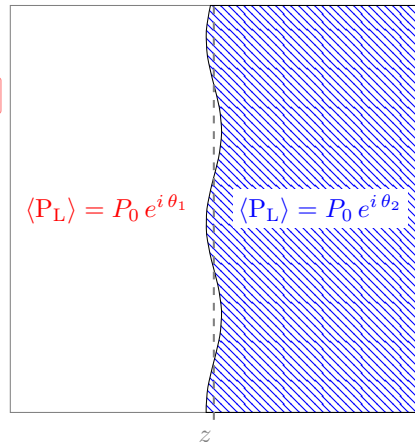
"Capillary" wave method for deconfined-deconfined interface tension

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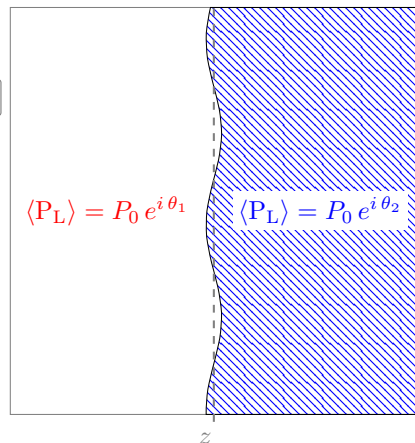
- Monte Carlo (heatbath + overrelaxation)

- Initialize simulation:

$$1) \text{ start with "mixed phase" state: } U_4(x_1, x_2, x_3, 0) = \begin{cases} \mathbb{1} e^{i\theta_1} & \text{if } z < N_z/2 \\ \mathbb{1} e^{i\theta_2} & \text{if } z \geq N_z/2 \end{cases}$$

2) thermalize

- Sample configurations for measurements



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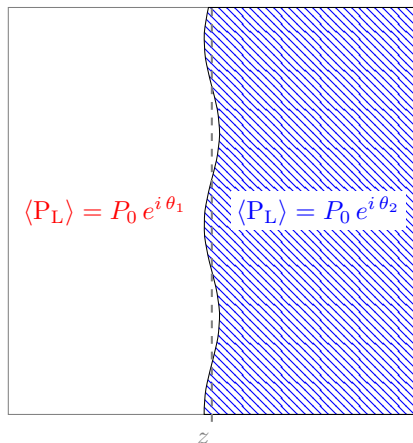
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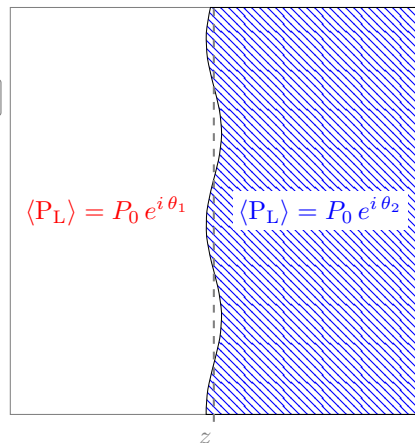
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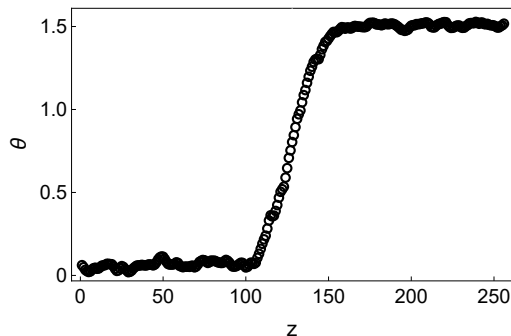
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- Can periodically recenter interface in z -direction to simplify detection.



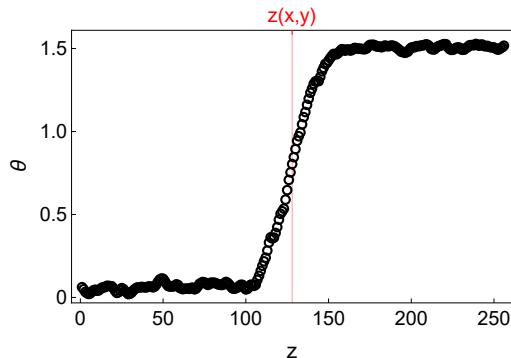
(Naive) interface detection and smearing

- For each (x, y) detect interface from profile of Polyakov loop phase $\theta(z) = \arg(P(x, y, z))$ as function of z



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Polyakov loop field (and its phase) very noisy \Rightarrow perform $n_s = n_{\text{pre}} + n$ iter. of smearing $P_L \rightarrow P_L^{(n_s)}$

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n_{pre} pre-smearing steps:

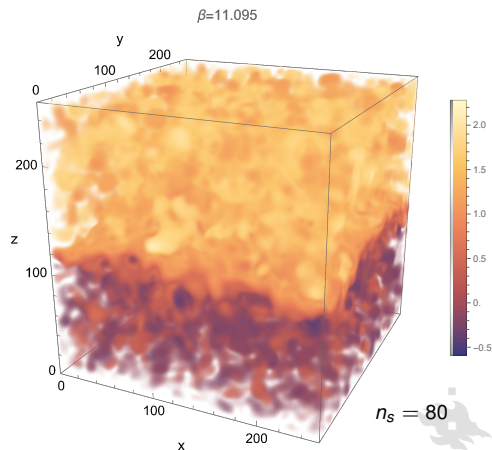
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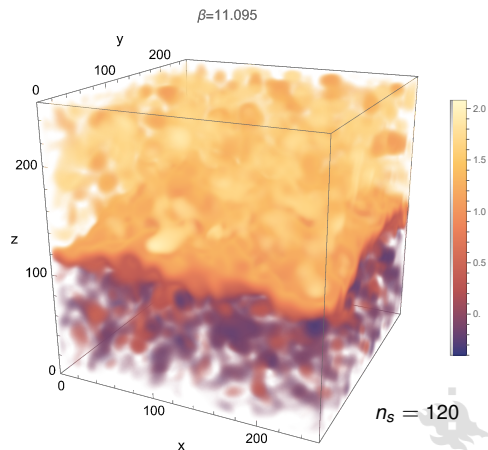
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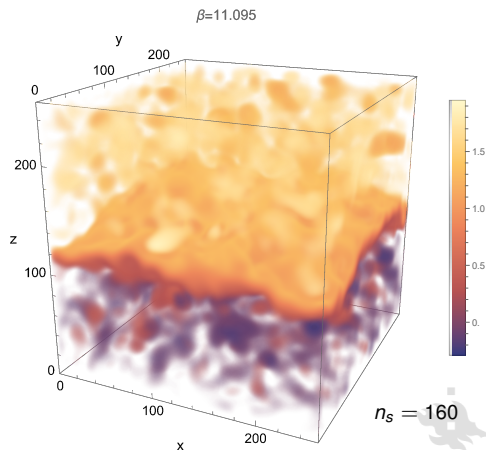
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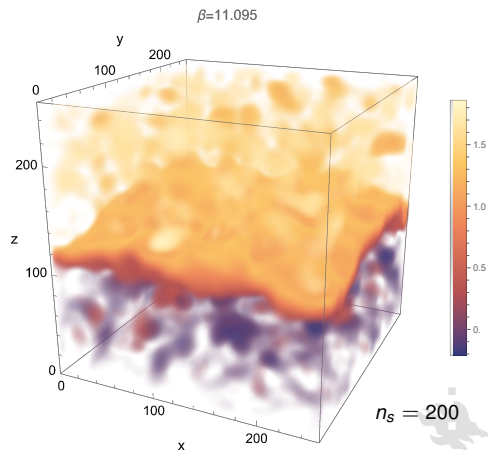
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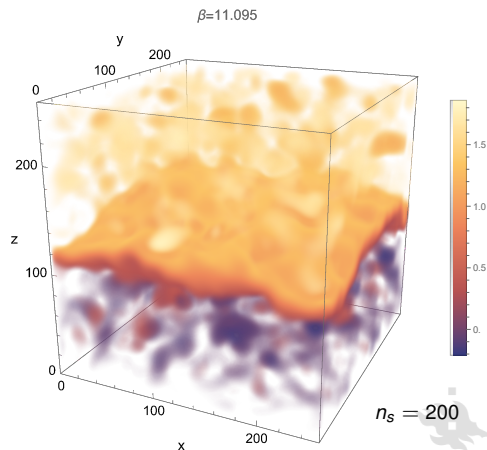
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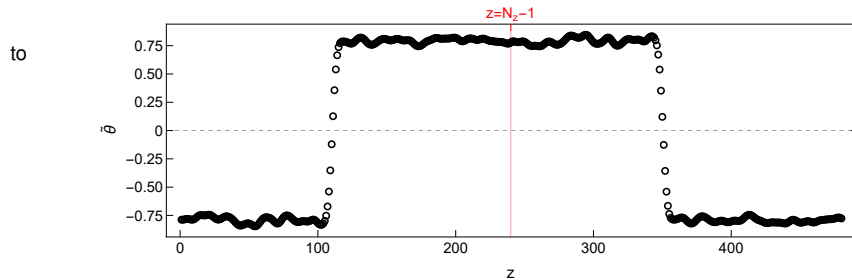
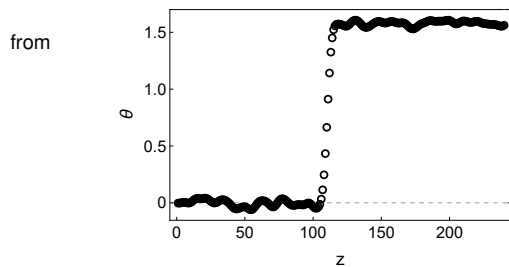
\rightarrow Naive interface detection requires lots of smearing!



Deconfined-deconfined interface tension

Interface detection from Fourier mode phase shift

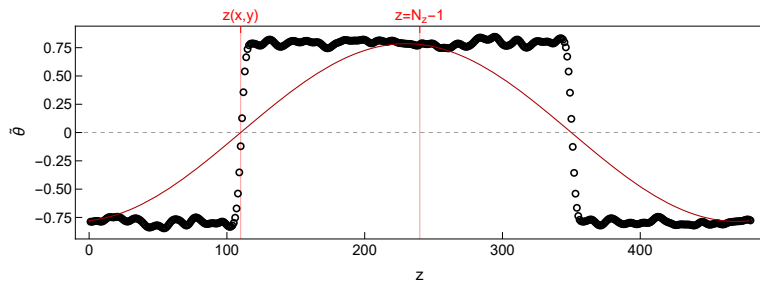
→ For each (x, y) transform profile of Polyakov loop phase $\theta(z) = \arg(P(x, y, z))$



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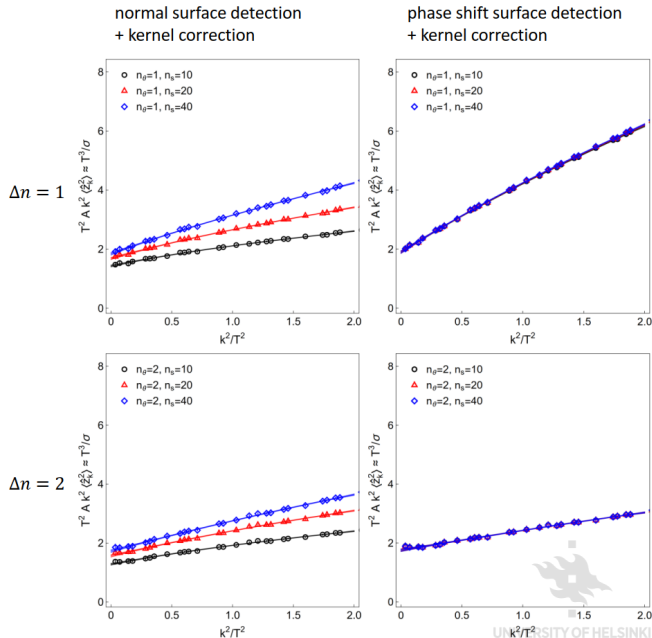
→ For each (x, y) transform profile of Polyakov loop phase $\theta(z) = \arg(P(x, y, z))$

→ extract $z(x, y)$ from phase shift of lowest non-trivial Fourier mode:



Comparison of normal and phase-shift interface detection method

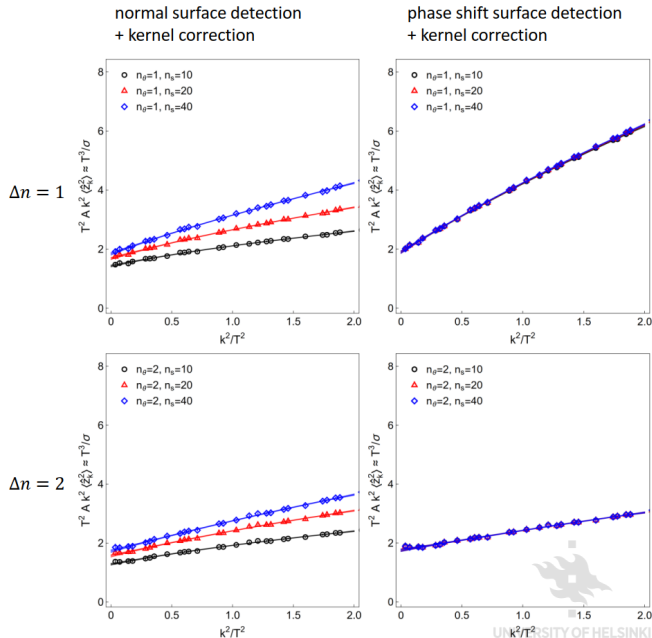
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Deconfined-deconfined interface tension

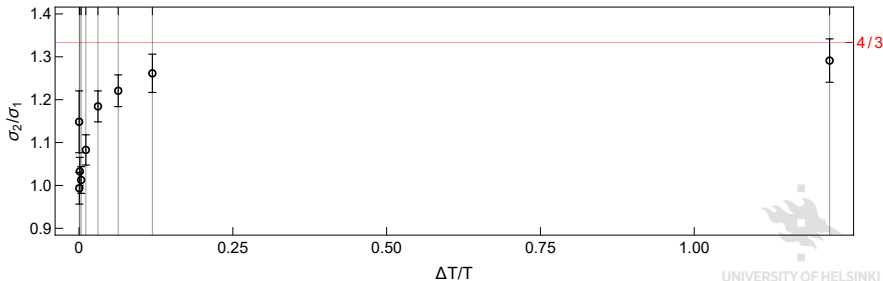
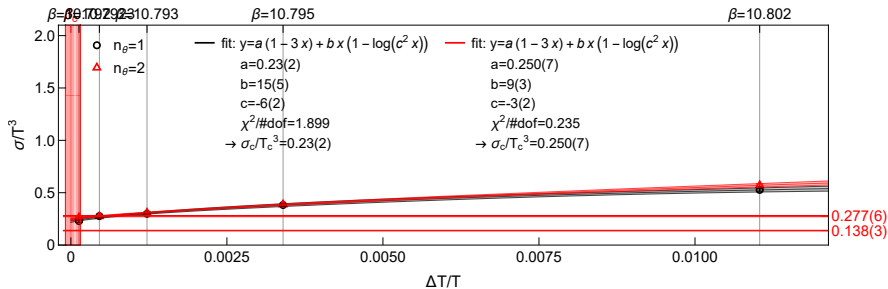
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- interface fluctuation spectra obtained from normally detected interface would require more smearing to collapse after kernel correction.
 - interface fluctuation spectra obtained from interface located by phase shift method collapse perfectly after kernel correction.
- With sufficient smearing (and in sufficiently big boxes) the two methods give the same results.



Deconfined-deconfined interface tension

Results for SU(4)



Deconfined-deconfined interface tension

Results for SU(8)

