LLP Theory overview*

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*actually, my own biased pheno POV...

Long-Lived Particles and the Third Generation, Edinbourgh, 20.11.2019

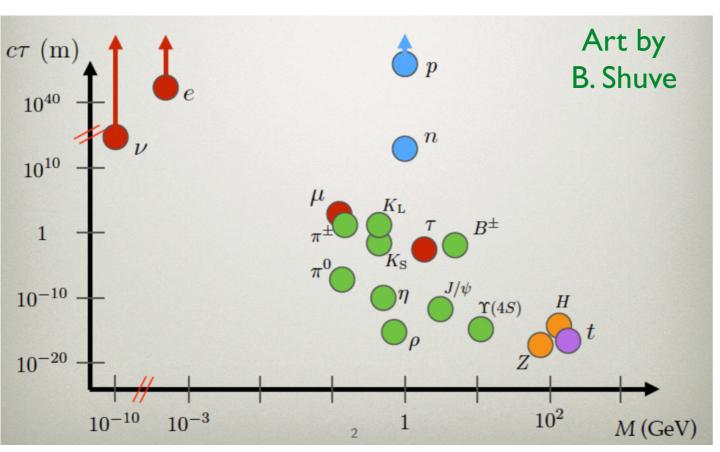
The Quest for New Physics

- Electroweak naturalness problem solved by New Physics (NP) at the TeV scale.
- Other fundamental questions (dark matter, CP asymmetry, neutrino masses, flavor, etc) can also be solved if the NP scale, Λ_{NP} is around the TeV scale.
- No New Physics at the LHC yet (modulo flavour anomalies...)
 - 1) *collider-phobic* (axions, dark photons, sub-GeV dark matter, sterile neutrinos, ...): "we'll need <another kind of experiment>" (e.g. FASER, MATHUSLA, ADMX, DUNE)
 - 2) Λ_{NP} higher than expected: "let's build a new collider!" [BSM-doer, energy] "let's compute more loops!" [QCD-doer, precision]
 - 3) $\Lambda_{\text{NP}} \sim 0.1$ -1 TeV, but it operates in *stealth mode*: heavy mediators, tiny couplings, compressed spectra, sequestered sectors, large backgrounds, ...)
 - LLPs in all of them!!! (note that #1 might need a collider as a source!)

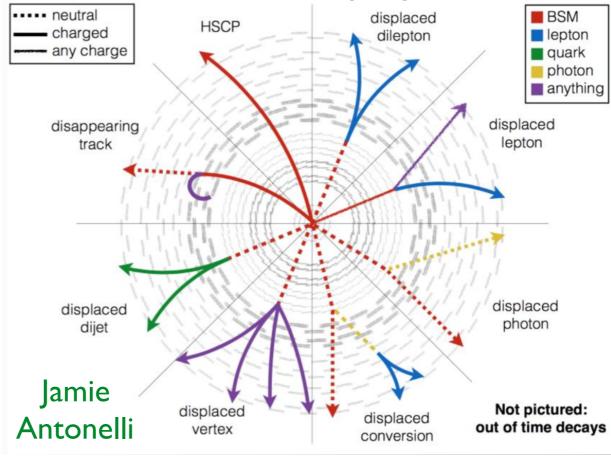
Long-Lived Particles

LLPs: BSM states with macroscopic lifetimes (ns), theoretically well motivated.

Exist in the SM!



A lot of interesting signatures!



large cτ, small Γ

- Large mass hierarchies
- Compressed spectra
- Small couplings

EW Baryogenesis
Dark Matter
Hierarchy Problem
Neutrino Masses

BSM Models: RH neutrinos, dark QCD, stealth SUSY, Neutral Naturalness, Higgs Portal, Z' Portal, Hidden Valleys, ...



MATHUSLA Physics case Curtin et al, 1806.07396 Signatures

LLP@LHC White Paper: Alimena et al, 1903.04497

LLP lessons from the SM

3 ways to get large $c\tau$ / small Γ (correlated with LHC limitations)

•Large mass hierarchies / off-shell mediator —— heavy E scale

$$c\tau(\mu \to e\nu) = \frac{1.2 \text{ fm}}{g_X^4} \left(\frac{m_e}{m_\mu}\right)^4 \left(\frac{1 \text{ TeV}}{m_\mu}\right) \sim 1 \text{ cm} \begin{cases} m_e = 10 \text{GeV}, \\ m_\mu = 10 \text{ GeV}, \\ m_\mu = 10 \text{ GeV}, \\ m_\mu = 1 \text{ TeV}, \\ m_\mu = 1 \text{ TeV}, \\ m_\mu = 10 \text{ GeV}, \\$$

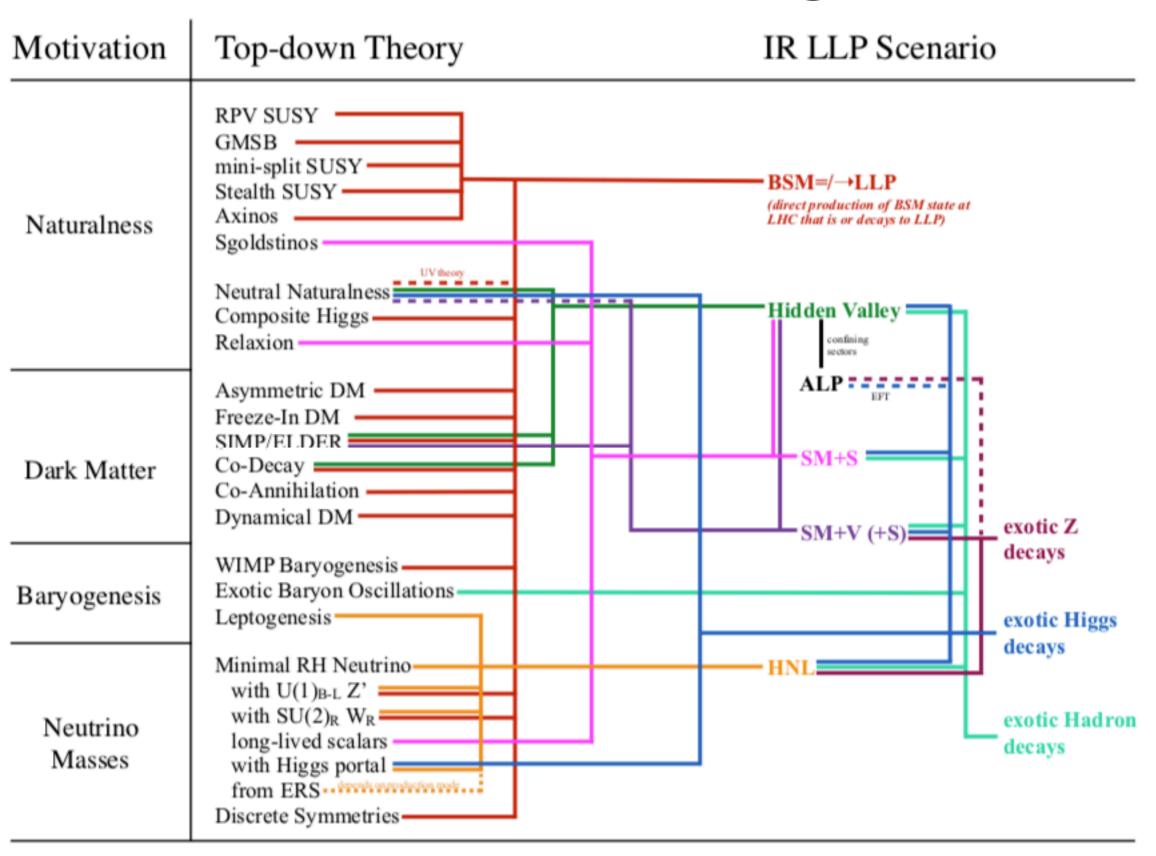
•Compressed spectra — Object reconstruction, thresholds

$$c\tau(n \to pe\nu) \sim \frac{1.2 \text{ fm}}{g_X^4} \left(\frac{m_p}{m_n - m_p}\right)^4 \left(\frac{1 \text{ TeV}}{m_n - m_p}\right) \sim 1 \text{ cm}$$
 $m_n = 101 \text{ GeV}, m_p = 100 \text{ GeV}, g_X^4 = 10^{-2}$

•Small coupling — Low rates

$$c au(Z o
u
u) \sim rac{0.02~\mathrm{fm}}{g_Z^4} \left(rac{1~\mathrm{TeV}}{m_Z}
ight) \sim 1~\mathrm{cm}$$
 $m_Z=1~\mathrm{GeV},~\mathrm{g}_\mathrm{Z}^2=10^{-12}$ Z_D models

From theories to signatures



1 Introduction 7	5 Detector Upgrades 81 5.1 The ATLAS and CMS experiments 83 5.2 LHCb Upgrade 123
 Simplified Models Yielding Long-Lived Particles Goals of the Present Simplified Model Framework 	5.2 LHCb Upgrade 123 5.3 Dedicated Detectors for LLPs 135
2.2 Existing Well-Motivated Theories for LLPs 17 2.3 The Simplified Model Building Blocks 19 2.4 A Simplified Model Proposal 26 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Models & Future Opportunities 37 Experimental Coverage of Long-Lived Particle Signatures 39 3.1 All-Hadronic Decays 41	Reinterpretation and Recommendations for the Presentation of Search Results 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Simplified Models 163 6.4 Recasting Examples for Specific Searches 167 6.5 Handling Long-Lived Particles in Delphes-Based Detector Simulations 187 6.6 Recasting Inside the Experimental Collaborations 192 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198
 3.2 Leptonic Decays 46 3.3 Semi-Leptonic Decays 52 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 	7 New Frontiers: Dark Showers 201 7.1 Introduction: The Anatomy of a Dark Shower 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213
4 Common Sources of Backgrounds for LLP Searches 71 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 4.4 Real Particles Originating from Outside the Detector 75	7.5 Trigger Strategies 220 7.6 Off-Line Analysis 227 7.7 Executive Summary 237 7.8 Appendix: Example Models 238
 4.4 Real Particles Originating from Outside the Detector 75 4.5 Fake-Particle Signatures 77 4.6 Algorithmically Induced Fakes 77 	8 Conclusions 245 A Appendix: Simplified Model Library 249
4.7 Summary 79	A.1 Instructions for the Simplified Model Library 249

1	Introduction 7	5	Detector Upgrades 81 5.1 The ATLAS and CMS experiments 83
2	Simplified Models Yielding Long-Lived Particles 2.1 Goals of the Present Simplified Model Framework 16		5.2 LHCb Upgrade 123 5.3 Dedicated Detectors for LLPs 135
	2.2 Existing Well-Motivated Theories for LLPs 17 Any exp. 2.3 The Simplified Model Building Blocks 19 search using 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Models & Future Opportunities 37	6	Reinterpretation and Recommendations for the Presentation of Search Results 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Simplified Models 163 6.4 Recasting Examples for Specific Searches 167 6.5 Handling Long-Lived Particles in Delphes-Based Detector Simulations 187
3	Experimental Coverage of Long-Lived Particle Signatures 39 3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46		 6.6 Recasting Inside the Experimental Collaborations 192 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198
	 3.3 Semi-Leptonic Decays 52 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 	7	New Frontiers: Dark Showers 201 7.1 Introduction: The Anatomy of a Dark Shower 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213
4	Common Sources of Backgrounds for LLP Searches 71 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72		7.5 Trigger Strategies 220 7.6 Off-Line Analysis 227 7.7 Executive Summary 237 7.8 Appendix: Example Models 238
	4.4 Real Particles Originating from Outside the Detector 754.5 Fake-Particle Signatures 77	8	Conclusions 245
	4.6 Algorithmically Induced Fakes 77 4.7 Summary 79	Α	Appendix: Simplified Model Library 249 A.1 Instructions for the Simplified Model Library 249

1 Introduction 7	5 Detector Upgrades 81 5.1 The ATLAS and CMS experiments 83
2 Eimplified Models Yielding Long-Lived Particles 2.1 Goals of the Present Simplified Model Framework 16	5.2 LHCb Upgrade 123 5.3 Dedicated Detectors for LLPs 135
2.2 Existing Well-Motivated Theories for LLPs 17 Any exp. 2.3 The Simplified Model Building Blocks 19 search using 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Models & Future Opportunities 37	Reinterpretation and Recommendations for the Presentation of Search Results 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Simplified Models 163 6.4 Recasting Examples for Specific Searches 167 6.5 Handling Long-Lived Particles in Delphes-Based Detector Simulations 187
Experimental Coverage of Long-Lived Particle Signatures 39 3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46 Closer look in	 6.6 Recasting Inside the Experimental Collaborations 192 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198
3.3 Semi-Leptonic Decays 52 next slide 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 4 Common Sources of Backgrounds for LLP Searches 71	7 New Frontiers: Dark Showers 201 7.1 Introduction: The Anatomy of a Dark Shower 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213 7.5 Trigger Strategies 220
 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 	7.6 Off-Line Analysis 227 7.7 Executive Summary 237 7.8 Appendix: Example Models 238
4.4 Real Particles Originating from Outside the Detector 754.5 Fake-Particle Signatures 77	8 Conclusions 245
4.6 Algorithmically Induced Fakes 77 4.7 Summary 79	A Appendix: Simplified Model Library 249 A.1 Instructions for the Simplified Model Library 249

2. Existing Well-Motivated Theories for LLPs 17 Any exp. 2.3 The Simplified Model Building Blocks 19 search using 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Model Library 32 3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46 Closer look in 3.3 Semi-Leptonic Decays 56 3.4 Observery Opportunities: Overroiew of Gaps in Coverage 67 4 Common Sources of Backgrounds for LLP Searches 71 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 75 4.4 Real Particles Produced via Interactions with the Detector 75 4.5 Fake-Particle Signatures 77 4.6 Algorithmically Induced Fakes 77 4.7 Summary 79 Summary 79 Any exp. 5.2 LHICL Introduction 133 5. Dedicated Detectors for LLPs 135 Reinterpretation and Recommendations for the Presentation of Search Results 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Experimental Results 160 6.3 Reinterpretation using Experimental Results 160 6.3 Remander Presenting Simplified Models 163 6.4 Recasting Institute Experimental Collovations 187 6.5 Hardling Long-Lived Particles in Detrutes-Based Detector Simulations 187 6.6 Recasting Institute in the Experimental Collovations 192 6.7 Reinterpretation using Experimental Results 160 6.8 Reinterpretation using Experimental Results 160 6.8 Recasting Examples for Specific Searches 167 6.9 Handling Long-Lived Particles in Detrutes-Based Detector Simulations 187 6.6 Recasting Examples for Specific Searches 167 6.8 Cour Proposals for the Presentation of Results 198 6.7 New Frontiers: Dark Showers 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213 7.5 Trigger Strategies 220 7.6 Off-Line Analysis 227 7.7 Executive Summary 237 7.8 Appendix: Example Models 238 7.8 Conclusions 245 7.9 Examples Model Library 249 7.9 Summary 79 7.1 Introduction 168 7.0 Recasting Examples for Searches 167 7.0 Recasting Examples for Searches 167 7.0 Recasting Examples for	1	Introduction 7	5	Detector Upgrades 81 5.1 The ATLAS and CMS experiments 83 Impact of LLP Community
2.2 Existing Well-Motivated Theories for LLPs 17 Any exp. 2.3 The Simplified Model Building Blocks 19 search using them? 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Models & Future Opportunities 37 3 Experimental Coverage of Long-Lived Particle Signatures 39 3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46 3.3 Semi-Leptonic Decays 52 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 4.1 Introduction 71 4.2 Long-Lived Particles or the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 4.4 Real Particles Originating from Outside the Detector 75 4.5 Reheardite Signatures 77 4.6 Algorithmically Induced Fakes 77 Appendix: Simplified Model Library 249	2			5.2 LHCb Upgrade 123
3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46 3.3 Semi-Leptonic Decays 52 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 4.4 Real Particles Signatures 77 4.6 Algorithmically Induced Fakes 77 4.7 Summons 50 4.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198 6.7 Reinterpretation with Proposal for the Presentation of Results 198 6.7 Reinterpretation with Proposal for the Presentation of Results 198 6.7 Reinterpretation with Proposal for the Presentation of Results 198 6.7 Reinterpretation of Results 198 6.7 Reinterp		2.2 Existing Well-Motivated Theories for LLPs 17 2.3 The Simplified Model Building Blocks 19 search using 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32	6	 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Simplified Models 163 6.4 Recasting Examples for Specific Searches 167
3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67 4 Common Sources of Backgrounds for LLP Searches 71 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 4.4 Real Particle Signatures 77 4.5 Fake-Particle Signatures 77 4.6 Algorithmically Induced Fakes 77 A Appendix: Simplified Model Library 249	3	3.1 All-Hadronic Decays 41		6.7 Reinterpretation with Prompt Analyses 197
4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72 4.4 Real Particles Originating from Outside the Detector 75 4.5 Fake-Particle Signatures 77 4.6 Algorithmically Induced Fakes 77 A Appendix: Simplified Model Library 249		3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58	7	7.1 Introduction: The Anatomy of a Dark Shower 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213
4.5 Fake-Particle Signatures 77 4.6 Algorithmically Induced Fakes 77 A Appendix: Simplified Model Library 249	4	4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71		7.6 Off-Line Analysis 227 7.7 Executive Summary 237
4.7 Summany 50			8	Conclusions 245
			A	

1 Introduction 7	5 Detector Upgrades 81 5.1 The ATLAS and CMS experiments 83 Impact of LLP Community
2 Simplified Models Yielding Long-Lived Particles 2.1 Goals of the Present Simplified Model Framework 16	5.2 LHCb Upgrade 123 5. Dedicated Detectors for LLPs 135
2.2 Existing Well-Motivated Theories for LLPs 17 Any exp. 2.3 The Simplified Model Building Blocks 19 search using 2.4 A Simplified Model Proposal 26 them? 2.5 Proposal for a Simplified Model Library 32 2.6 Limitations of Simplified Models & Future Opportunities 37	Reinterpretation and Recommendations for the Presentation of Search Results 157 6.1 Introduction 158 6.2 Options for Presenting Experimental Results 160 6.3 Reinterpretation using Simplified Models 163 6.4 Recasting Examples for Specific Searches 167 6.5 Handling Long-Lived Particles in Delphes-Based Detector Simulations 187
Experimental Coverage of Long-Lived Particle Signatures 3.1 All-Hadronic Decays 41 3.2 Leptonic Decays 46 Closer look in	 6.6 Recasting Inside the Experimental Collaborations 192 6.7 Reinterpretation with Prompt Analyses 197 6.8 Our Proposals for the Presentation of Results 198
3.3 Semi-Leptonic Decays 52 next slide 3.4 Photonic Decays 56 3.5 Other Exotic Long-Lived Signatures 58 3.6 Discovery Opportunities: Overview of Gaps in Coverage 67	7 New Frontiers: Dark Showers 201 7.1 Introduction: The Anatomy of a Dark Shower 201 7.2 Production 205 7.3 Shower 207 7.4 Decay 213 Ask Nishita or Giovanna for details!
4 Common Sources of Backgrounds for LLP Searches 71 4.1 Introduction 71 4.2 Long-Lived Particles in the Standard Model 71 4.3 Real Particles Produced via Interactions with the Detector 72	7.5 Trigger Strategies 220 7.6 Off-Line Analysis 227 7.7 Executive Summary 237 7.8 Appendix: Example Models 238
4.4 Real Particles Originating from Outside the Detector 754.5 Fake-Particle Signatures 77	8 Conclusions 245
4.6 Algorithmically Induced Fakes 774.7 Summary 79	A Appendix: Simplified Model Library 249 A.1 Instructions for the Simplified Model Library 249

Gaps in coverage

1. All-hadronic LLP decays

- Associated-object triggers (especially motivated by Higgs-like VBF and VH production modes) need to be more comprehensively used to improve sensitivity to low- p_T objects
- Improvements are needed in sensitivity at lower masses & lifetimes (e.g., for LLPs produced in Higgs decays)
- Single hadronic DVs need to be looked for in searches that currently use two (such as decays in ATLAS HCAL and MS)
- Possibilities need to be explored for online reconstruction of hadronic displaced objects, as the inclusive H_T triggers of ATLAS and CMS miss these objects unless they have a large p_T . Note that LHCb can actually trigger on a displaced hadronic vertex.
- Low-mass hadronically decaying LLPs can look somewhat like tau leptons, so the question remains as to whether there is any possibility of using, for example, L1 tau triggers to seed displaced jet triggers at HLT and improve trigger efficiency; studies need to be performed by the experimental collaborations
- The prospects for dedicated searches for displaced hadronic taus need to be investigated, since no dedicated searches currently exist. By the same token, it would be interesting to explore the possibility to flavor-tag displaced jets (b-displaced jets, c-displaced jets, etc).

2. Leptonic

- Coverage needs to be provided for the intermediate region between boosted, low-mass LLPs (lepton jets) and high-mass, resolved LLPs (resolved ATLAS/CMS searches)
- Improvements need to be made to extend coverage to lower masses and to lower p_T thresholds, though currently it's unknown how this can be accomplished, and dedicated studies need to be done
- Searches need to be done for different combinations of charge and flavor of displaced leptons (e.g., same-sign vs. opposite-sign, oppositeflavor vs. same-flavor)
- Searches need to be done for tau leptons in LLP decays, in particular
 if they come from the ID; an unanswered question remains as to
 whether displaced-jet triggers can be used

3. Semi-Leptonic

- Searches do not exist and need to be done for masses below or about 30 GeV, theoretically well motivated by Majorana neutrinos.
- Searches need to be performed for all flavor combinations (for example, one CMS search only covers $e^{\pm}\mu^{\mp}$), as well as same-sign vs. opposite-sign leptons
- Currently unknown improvements need to be made to relax or modify isolation criteria wherever possible to recover sensitivity to boosted semi-leptonic decays
- Searches need to be done that better exploit triggering on associated objects for improved sensitivity to low-mass objects, or to employ high-multiplicity lepton triggers if there are multiple LLPs

4. Photonic

- There is currently no coverage for LLPs decaying into $l\gamma$, $j\gamma$, or without missing energy, and searches urgently need to be performed for this decay topology
- There is currently poor coverage (i.e., there exists no dedicated search) for single-γ topologies; the only searches with sensitivity require two jets to be present.
- There is currently no coverage for softer non-pointing or delayed photons, and searches need to be performed here
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5. Other exotic long-lived signatures

- Disappearing tracks with $c\tau \sim \text{mm}$ are very hard to probe, and new ideas and detector components are needed to extend sensitivity to this potential discovery regime. It's unclear if the ATLAS insertable B-layer will be present in HL-LHC run and how sensitivity to the disappearing track topology will improve with the replacement of the current inner detector with the new ITk (Inner Tracker), or whether new tracking layers very close to the beam line can be added. It's an open question as to what is the lowest distance at which new layers (or double layers) can be inserted. Another open question that needs to be answered is whether there are any prospects for disappearing tracks at LHCb with an upgraded detector.
- No dedicated searches for quirks exist at the LHC, a huge, open discovery possibility for ATLAS, CMS, and LHCb. Some LHC constraints exist by reinterpreting heavy stable charged particle searches, but dedicated searches need to be performed. There are significant challenges in modeling the propagation and interaction of quirks with the detector, as well as in fitting tracks to their trajectories, but there new ideas have been proposed that need to be explored by the experimental collaborations that might allow improved sensitivity to quirks with less ambitious analysis methods.

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= lighter and softer objects

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5. Other exotic long-lived signatures

- Disappearing tracks with $c\tau \sim \text{mm}$ are very hard to probe, and new ideas and detector components are needed to extend sensitivity to this potential discovery regime. It's unclear if the ATLAS insertable B-layer will be present in HL-LHC run and how sensitivity to the disappearing track topology will improve with the replacement of the current inner detector with the new ITk (Inner Tracker), or whether new tracking layers very close to the beam line can be added. It's an open question as to what is the lowest distance at which new layers (or double layers) can be inserted. Another open question that needs to be answered is whether there are any prospects for disappearing tracks at LHCb with an upgraded detector.
- No dedicated searches for quirks exist at the LHC, a huge, open discovery possibility for ATLAS, CMS, and LHCb. Some LHC constraints exist by reinterpreting heavy stable charged particle searches, but dedicated searches need to be performed. There are significant challenges in modeling the propagation and interaction of quirks with the detector, as well as in fitting tracks to their trajectories, but there new ideas have been proposed that need to be explored by the experimental collaborations that might allow improved sensitivity to quirks with less ambitious analysis methods.

LLPs in concrete BSM scenarios

LLPs and Electroweak baryogenesis

- Assume 2 neutral WIMPs that freeze out. One of them works as dark matter
 via the standard WIMP miracle, and the other one is metastable and decays
 after freeze-out, triggering baryogenesis [Cui, Sundrum, 1212.2973]
- Baryogenesis better happens before BBN, then

$$1 \text{ cm} \left(\frac{100 \text{ GeV}}{\text{m}_{\chi}}\right)^2 \lesssim \text{c}\tau \lesssim 10^8 \text{ m}$$

 Interestingly, this span various orders of magnitude! If a neutral particle lives too long (more than a few meters) it will escape the detector unscathed and will only appear as missing energy, but then it can be studied with a dedicated detector, such as MATHUSLA, FASER, etc...

LLPs and Dark matter (I): EW multiples (minimal DM)

- Assume DM candidate is the neutral component of one EW multiplet.
- Then EWSB (W,γ-loops) split charged components from the DM candidate.
 Charged lifetime pretty much independent of DM mass!

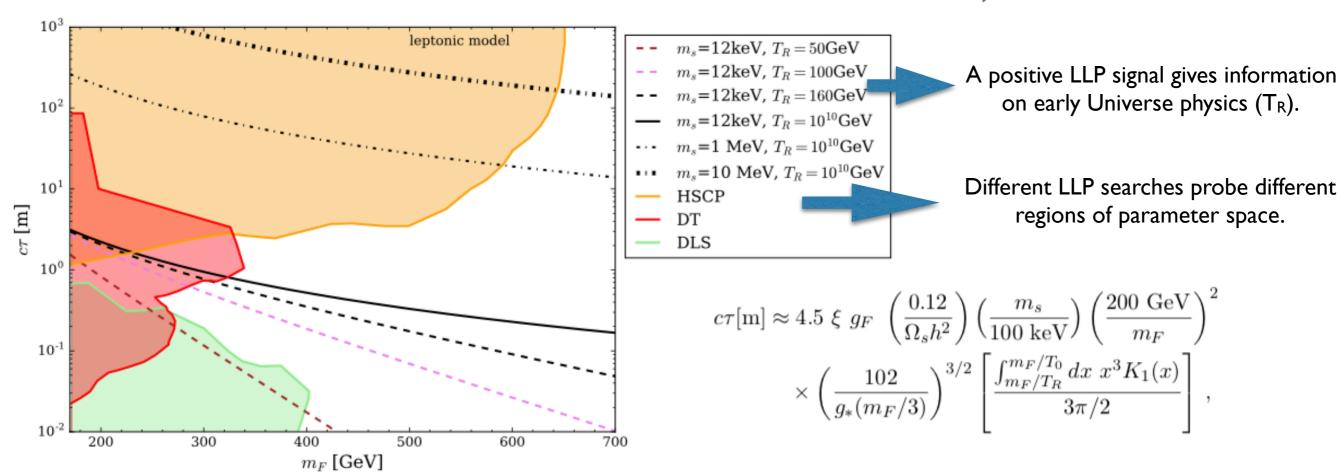
$$c\tau \approx 0.7 \,\mathrm{cm} \times \left(\frac{\Delta m}{340 \,\mathrm{MeV}}\right)^3 + O(m_Z^2 \,/\, m^2)$$

- Lifetimes of 6.6 mm (n-even, Y=1/2), 6.0 cm (n-odd, Y=0).
- Above cases correspond to pure Higgsino and pure Wino if DM is fermion.
- Main search: disappearing tracks (more to discuss in Nishita's talk Thursday)
- Note:mixing with additional states (e.g. singlets) can shorten the lifetime.

LLPs and Dark matter (II): freeze-in (FIMP DM)

tiny Y: DM never in thermal equilibrium. Slowly produced, reaches measured relic.

New fields: scalar DM (s)
$$\mathcal{L} = \mathcal{L}_{SM} + \partial_{\mu} s \ \partial^{\mu} s - \frac{\mu_s^2}{2} s^2 + \frac{\lambda_s}{4} s^4 + \lambda_{sh} s^2 \left(H^{\dagger} H\right)$$
 + VL-fermion (F).
$$+ \bar{F} \left(i \not \!\!D\right) F - m_F \bar{F} F - \sum y_s^f \left(s \bar{F} \left(\frac{1+\gamma^5}{2}\right) f + \text{h.c.}\right)$$
 Details in Nishita's talk on Thursday (right before tea...)

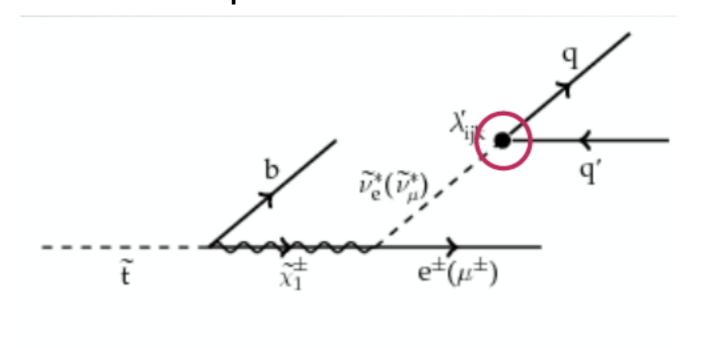


Belanger, Desai, Goudelis, Harz, Lessa, No, Pukhov, Sekmen, Sengupta, Zaldivar, JZ, 1811.05478, 1910.00117.

See Co, D'Eramo, Hall, Pappadopulo, 1506.07532; Evans, Shelton: 1601.01326; Calibbi, Lopez-Honorez, Lowette, Mariotti, 1805.04423, also: Junius, Lopez-Honorez, Mariotti 1904.07513, No, Tunney, Zaldivar 1908.11387

LLPs and Hierarchy Problem (I): RPV SUSY

- Spontaneous breaking of a symmetry naturally generates a small coupling: symmetry gets restored if the coupling vanishes (technical naturalness)
- RPV operators induce proton decay, hence the coupling must be bounded. Indeed, $|\lambda'| \lesssim 10^{-8}$ (rough estimate)
- RPV SUSY also explain null results in MET + X searches!

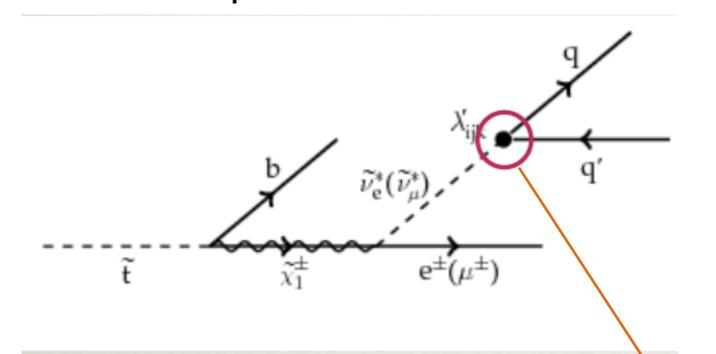


For details:

Z. Liu, B. Tweedie 1503.05923; C. Csaki, E. Kuflik, S. Lombardo, O. Slone, T. Volansky, 1505.00784

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Flavor structure of λ_{ijk} is arbitrary!

LLPs and Hierarchy Problem (I): Neutral Naturalness

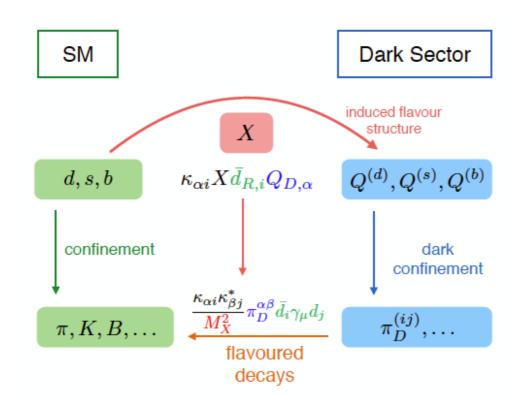
- Expected "supersymmetric" LHC outcome: produce coloured particles (stops, gluinos), then electroweak ones (charginos, neutralinos) and finally the Higgs (because it goes via a dimension 5 operator).
- However, the particles canceling the top loops could simply be uncoloured (Twin Higgs, Orbifold Higgs, Quirky Little Higgs, Folded SUSY) or even fully SM singlets: Hyperbolic Higgs [Cohen, Craig, Giudice, McCullough, 1803.03647], Top Siblings [Cheng, Li, Salvioni, Verhaeren, 1803.03651].
- In all cases, glueballs are present and tend to be the leading signal

$$\mathbf{T} \sim 4 \cdot 10^{-13} \text{ s} \left(\frac{5 \text{ GeV}}{\Lambda_{\text{QCD}_{B,C}}} \right)^{9} \left(\frac{m_{\tilde{s}_{\Delta}^c}}{300 \text{ GeV}} \right)^{4}$$
Cheng, Li, Salvioni, Verhaeren, 1803.03651

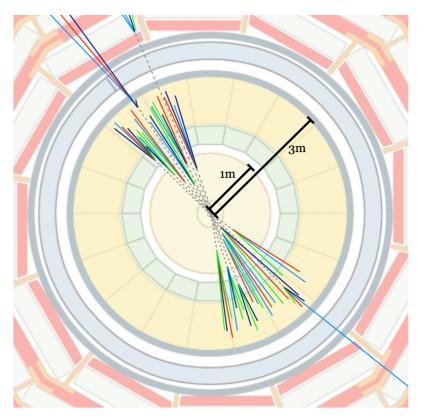
Very sensitive to precise values of $\Lambda_{QCD(B,C)}$ and m_{Δ} !

LLPs and ... String Theory: Hidden Valley models

- String theory models live in 10+1 dimensions with larger gauge groups: SO(32), E8xE8,... so there is plenty of room to have in those constructions some SU(N) confining subgroup (helps to think in terms of QCD).
- The lightest HV meson will drive the pheno (pions if QCD):



Schwaller, Renner 1803.08080

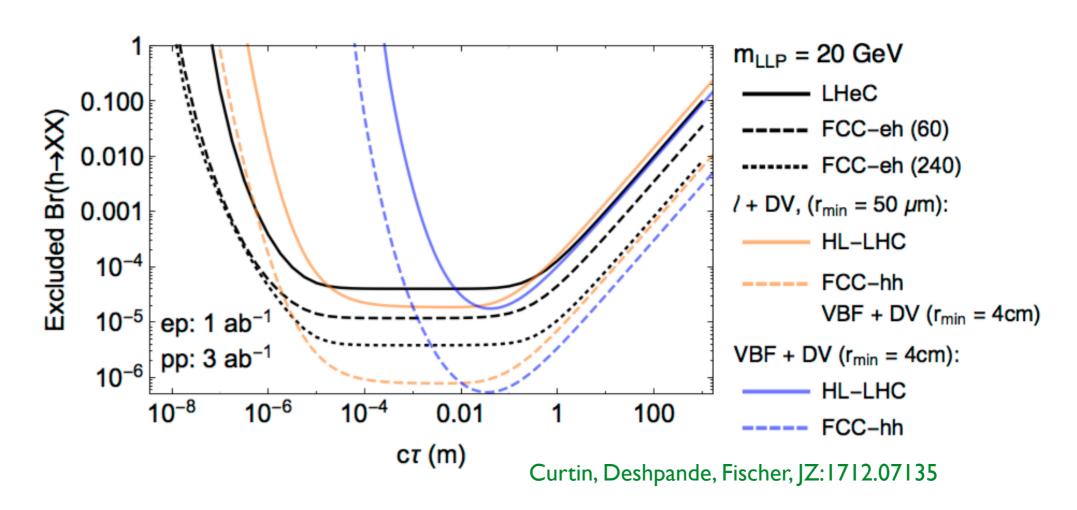


Schwaller, Stolarski, Weiler 1502.05409

• Dark sector can have extra U(I) bosons: current benchmark for displaced jets.

LLPs and ...electron-proton colliders

- Electron-proton colliders provide the clean environment of an e⁺-e⁻ machine with a CME larger by (Ep/Ee)^{1/2} but lower by same amount than p-p machine.
- Hence their window of opportunity is narrow: they can only compete for physics cases where the high pile-up of a p-p collider affects the search, and the energy is on the high side e⁺-e⁻ machine.
- Incredible useful to test $O(\mu m)$ lifetimes; competes with p-p machines!



A few open issues

(with a slight preference for 3rd gen solutions)

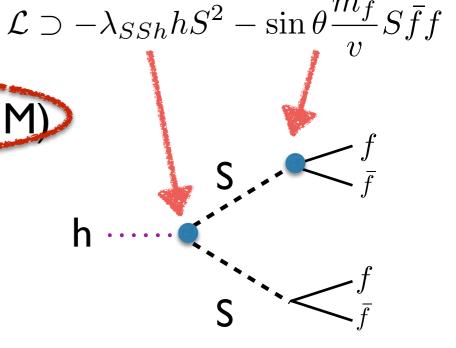
Disclaimer

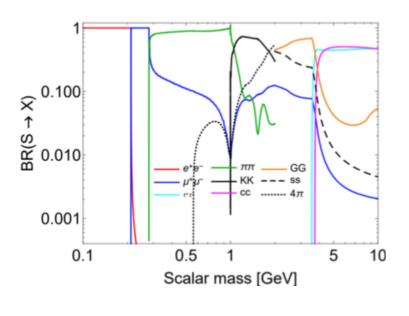
- This part is largely based on my own opinions and ideas, plus interesting discussions with many colleagues.
- I do not intend to point fingers and show a clear and concrete direction, but rather share my own musings... hoping they can trigger some spark!
- Always willing to take any discussion offline!

Why is the 3rd gen special?

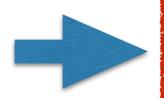
It could be the <u>discovery mode</u>:
 In "Higgs-like" theories (Higgs portal, HAHM) with minimal flavour violation, new scalars couple strongly to tt, bb, T⁺T⁻.

 SM Higgs decays 66% into 3rd gen!





I. Boiarska, K. Bondarenko, A. Boyarsky, V. Gorkavenko, M. Ovchynnikov, and A. Sokolenko, *Phenomenology of GeV-scale scalar portal*, arXiv:1904.10447.



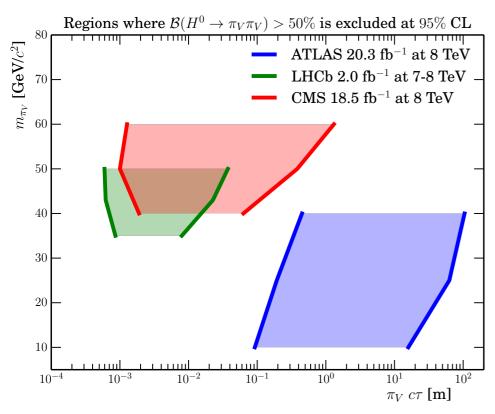
S with $2m_T < m_S < 2m_b$ decays 1:1 into cc: T^+T^- !

• It is a must for <u>signal characterization</u>, even if we find NP before in Ist/2nd gen, e.g. $DV(\mu,\mu)$, DV+MET, etc...

- Hadronic DVs searches do not identify jet flavour (leptonic DVs can distinguish e from μ).
- Naively: b-tagging based on B-meson lifetimes ("SM LLPs"). Hence we expect displaced b-jet ID to be a difficult task (particularly if the LLP lifetime lies in the near a SM LLP).
- However, flavour ID can be important in many cases!
- Higgs-like particles couple strongly to b-quarks, tops, tau leptons (if kinematics allow!)

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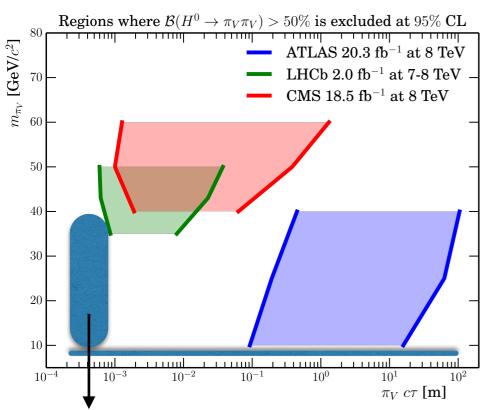
Hadronic DV searches in Hidden Valley Model: Comparison of ATLAS, CMS and LHC.



- 1) B-tagging on top of b-mesons can be complicated ($c\tau \sim 0.15 0.45$ mm).
- 2) For larger cτ, all non-b background: properly identifying displaced b-jets can drastically reduce the background while keeping 90% of the Higgs-like signal.
- 3) Analogous reasoning for H-> $\tau\tau$, 2 m $_{\tau}$ < m $_{H}$ < 2 m $_{b}$

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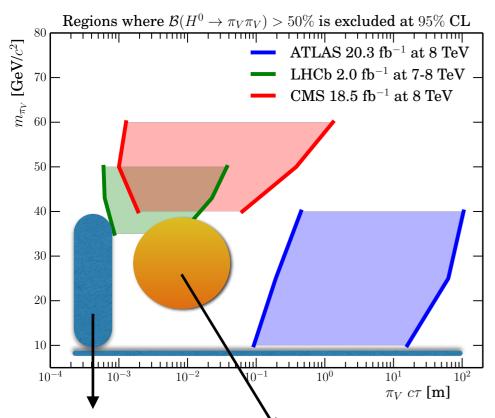
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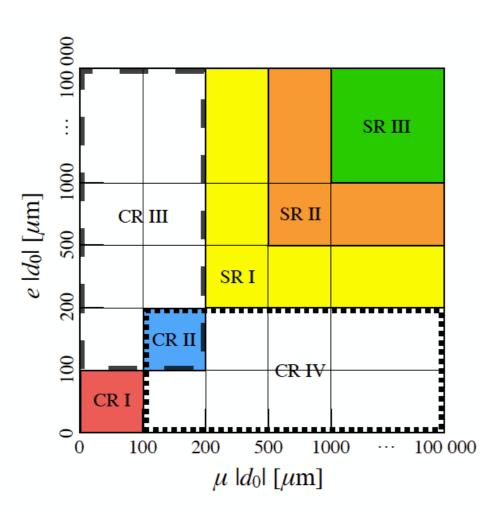
BR(H⁰->bb) ~ 90%. Bgd: non-b mesons. How well can we reject light displaced jets?

2-Soft displaced leptons?

Based on CMS Displaced e-µ search, EXO-16-022-pas

The CMS Displaced Lepton search (DLS) uses exactly 1e and 1µ.

What about Ts?



The plot is valid for pT(e, μ) > 40 GeV.

To have LLPs with Δ~40 GeV we need tiny couplings (phase space too large!). Indeed the search is benchmarked with RPV stops, but it also works for e.g. freeze-in dark matter Evans, Shelton, 1601.01326, Belanger, Desai, Goudelis, Harz, Lessa, No, Pukhov, Sekmen, Sengupta, Zaldivar, JZ, 1803.10379 + in preparation

If we want to consider a compressed phase space, as the one arising from dark matter models, or theories with new electroweak multiplets:

- 1) How would this plot look like to pT = 5, 10, 20, 30 GeV?
- 2) How strict is the d0 < 200 µm region?

LLPs + X:

LLPs and Neutrino Masses LLPs and Machine Learning LLPs and Extra U(1)s LLPs and Dedicated Experiments LLPs and...

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LLPs+Google images -

Conclusions

- Long lived particles (LLPs) are theoretically well motivated: ubiquitous in models trying to solve fundamental problems of the SM.
- There is no "no-lose" theorem for LLPs (also true for BSM), but concrete incarnations do point to LHC testable mass-coupling ranges.
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21

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What about taus???