

Long-lived particles and co-annihilating dark matter with tau-lepton

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[arxiv:1702.00750] [JHEP 06(2017)041]

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Dark Matter

- All evidence for the existence of Dark Matter is purely gravitational.
- The particle physics nature of DM is unknown.
- — is our main ‘new physics’ challenge.
- The energy density of the DM is precisely measured:

$$\Omega_{\text{DM}} h^2 \simeq 0.12$$

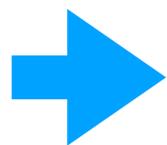
- The thermal freeze-out is attractive scenario for particle dark matter.

WIMP miracle

$$\Omega h^2 \propto \frac{1}{\langle \sigma v \rangle}$$

$$\langle \sigma v \rangle \sim \frac{g^4}{m^2}$$

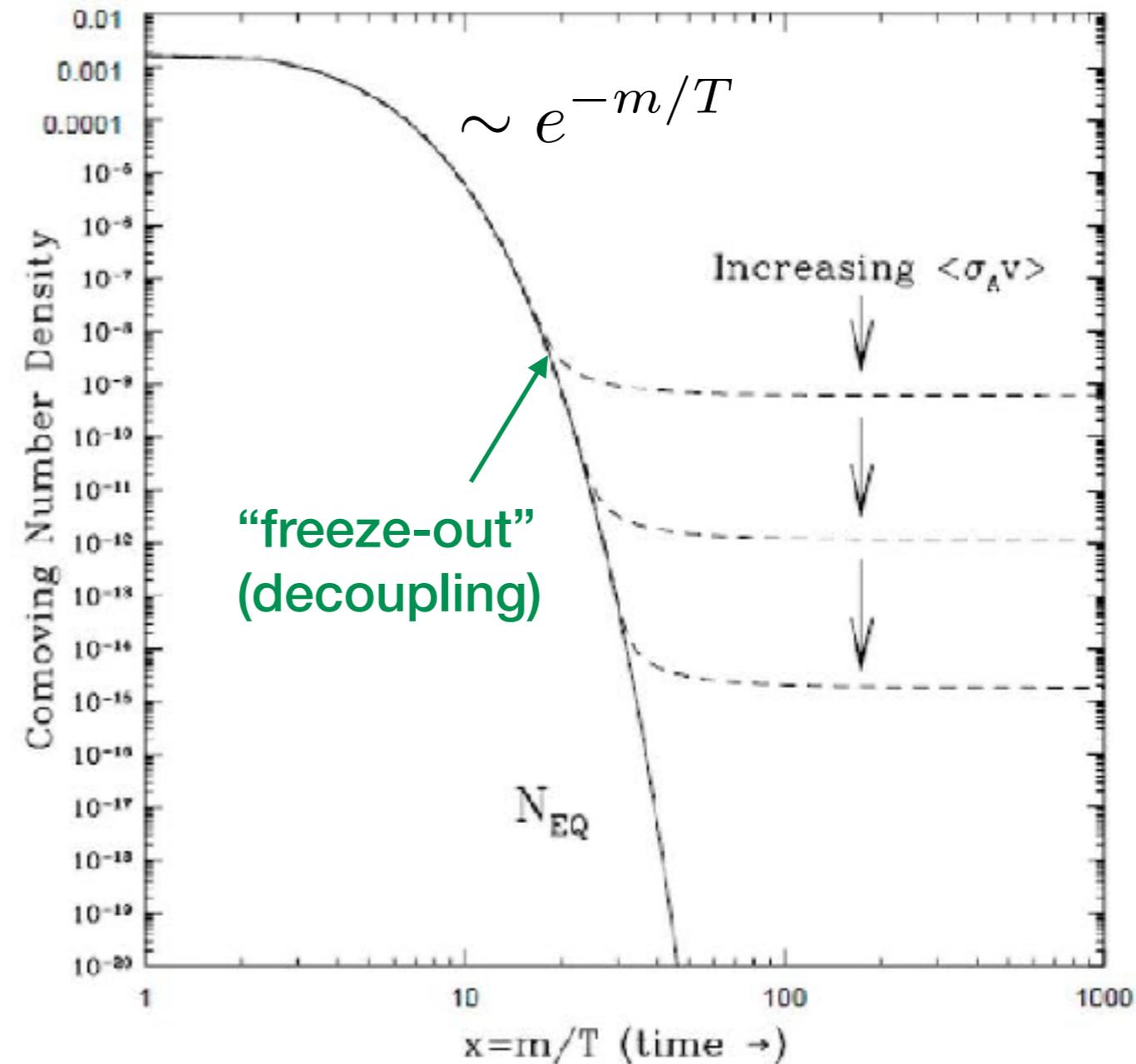
$$g \sim 0.1, \quad \Omega h^2 = 0.12$$



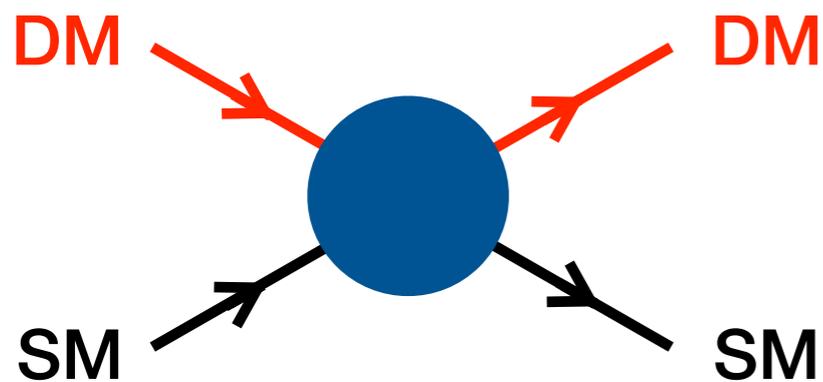
$$m \sim 100 \text{ GeV}$$

within the reach of LHC!

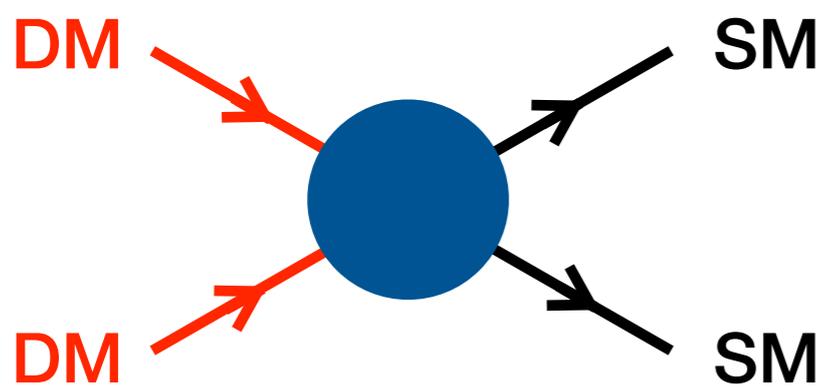
something to do with hierarchy problem?



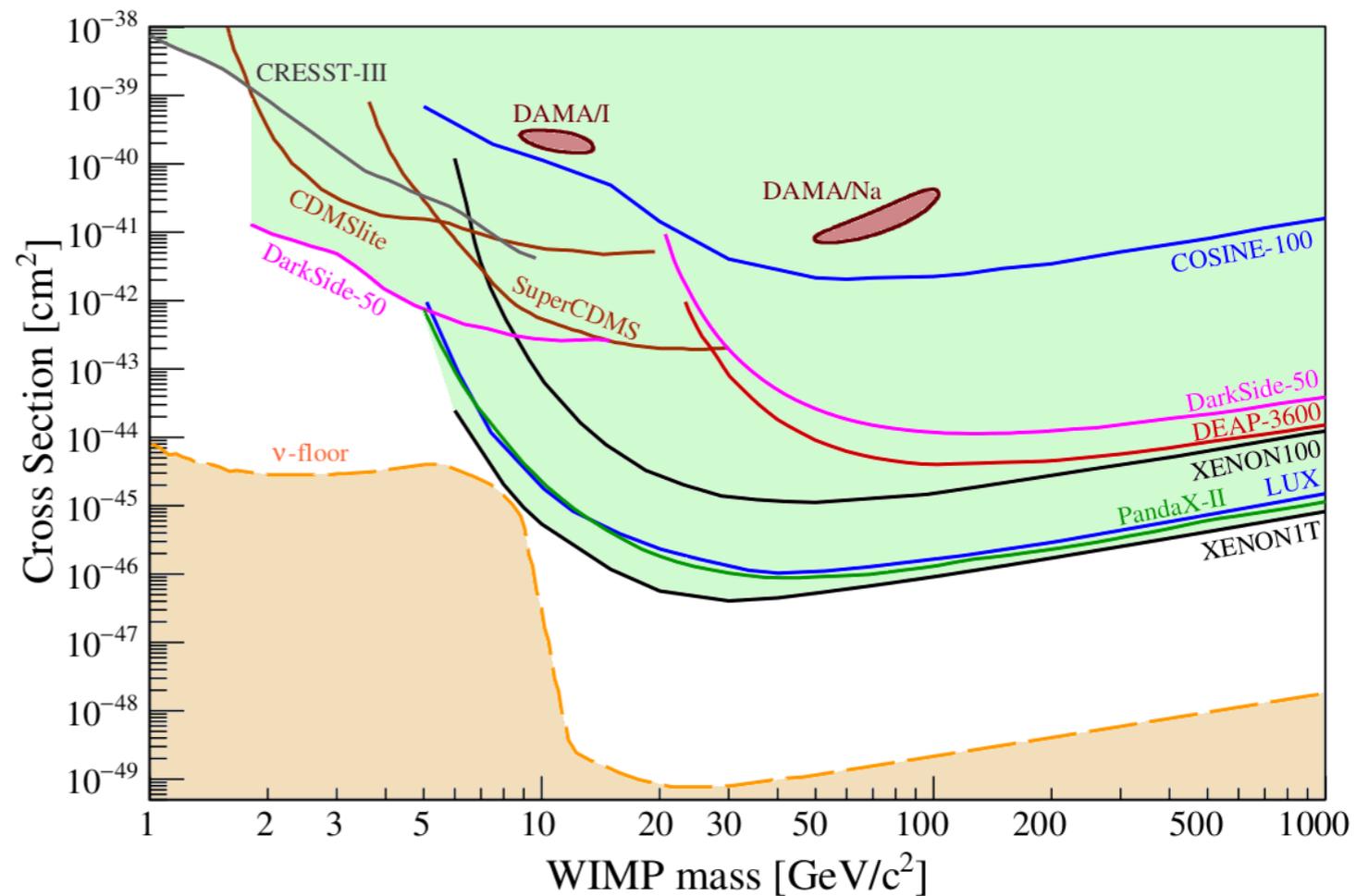
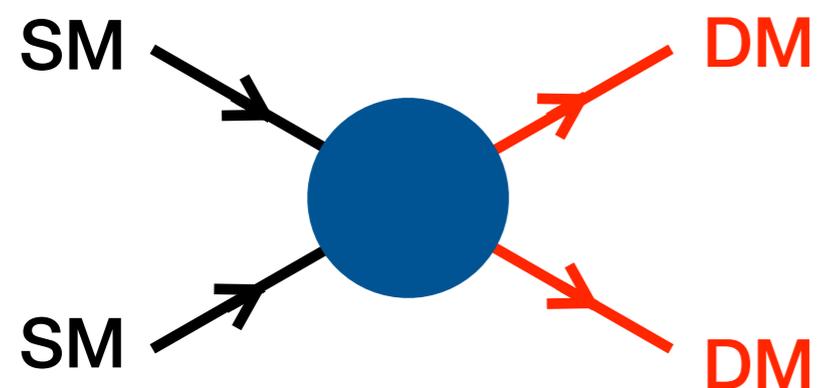
Direct Detection:



Indirect Detection:



Collider (LHC):



No DM signature is found anywhere.

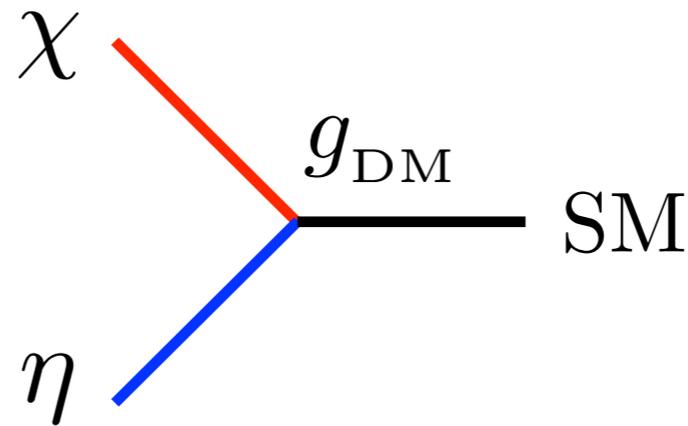
Why?

Co-annihilation

χ : **Dark Matter (DM):** \rightarrow **singlet**

η : **Co-Annihilation Partner (CAP):** \rightarrow **coloured or weakly charged**

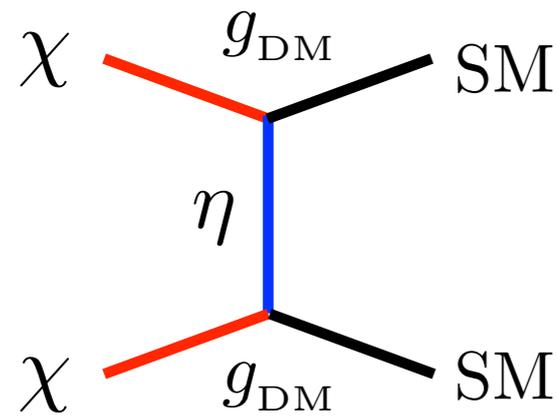
$$\mathcal{L} \supset g_{\text{DM}} \cdot \chi \eta (\text{SM})$$



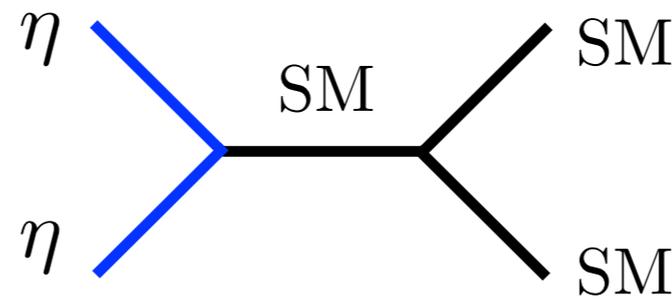
$$\mathbb{Z}_2 \left\{ \begin{array}{l} \eta \rightarrow -\eta \\ \chi \rightarrow -\chi \end{array} \right.$$

For small g_{DM} , $\Delta m = m_\eta - m_\chi$

These two processes are enough to keep χ in the thermal bath

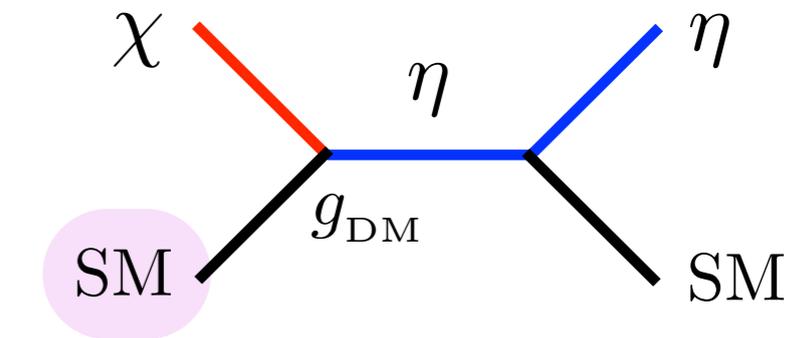


\ll



gauge interaction
(unsuppressed)

\ll



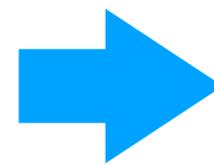
not Boltzmann
suppressed

$$\sigma_{\text{eff}} v = \frac{1}{(g_\chi + \bar{g}_\eta)^2} \left[g_\chi^2 \cdot \sigma(\chi\chi \rightarrow \tau^+\tau^-) + g_\chi \bar{g}_\eta \cdot \sigma(\chi\eta \rightarrow SM \text{ particles}) + \bar{g}_\eta^2 \cdot \sigma(\eta\eta \rightarrow SM \text{ particles}) \right] v$$

$$T_* \sim \frac{m_\chi}{25}$$

$$\bar{g}_\eta = g_\eta \left(\frac{m_\eta}{m_\chi} \right)^{3/2} \exp\left(-\frac{\Delta m}{T}\right)$$

Boltzmann factor of η :
We want this to be order 1



$$\frac{\Delta m}{m_\chi} \ll \frac{1}{25}$$

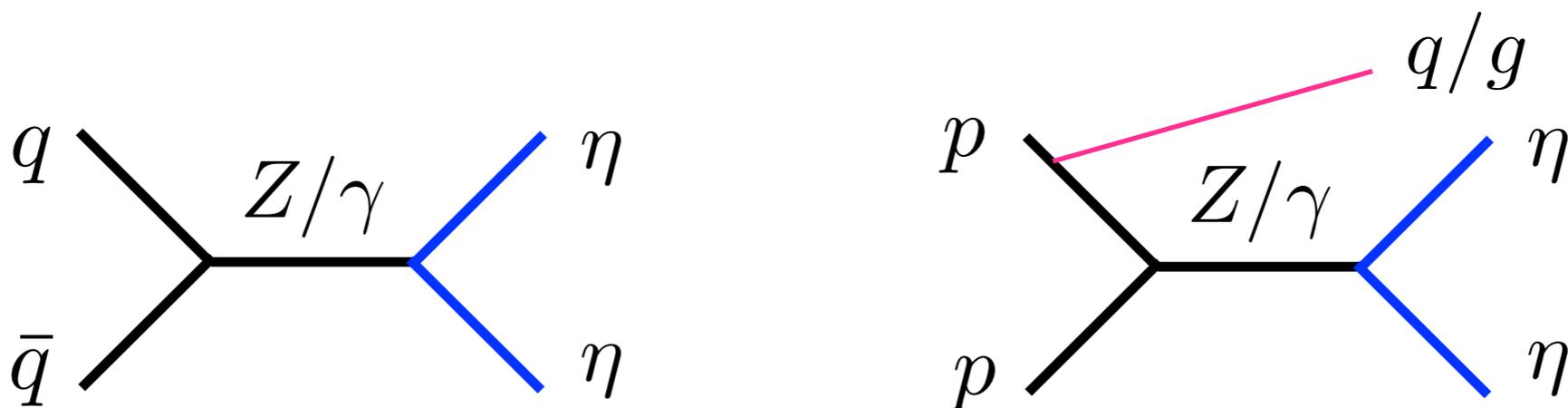
$$g_\chi \sim g_\eta \sim \mathcal{O}(1)$$

mass degeneracy is required

Experimental Signatures

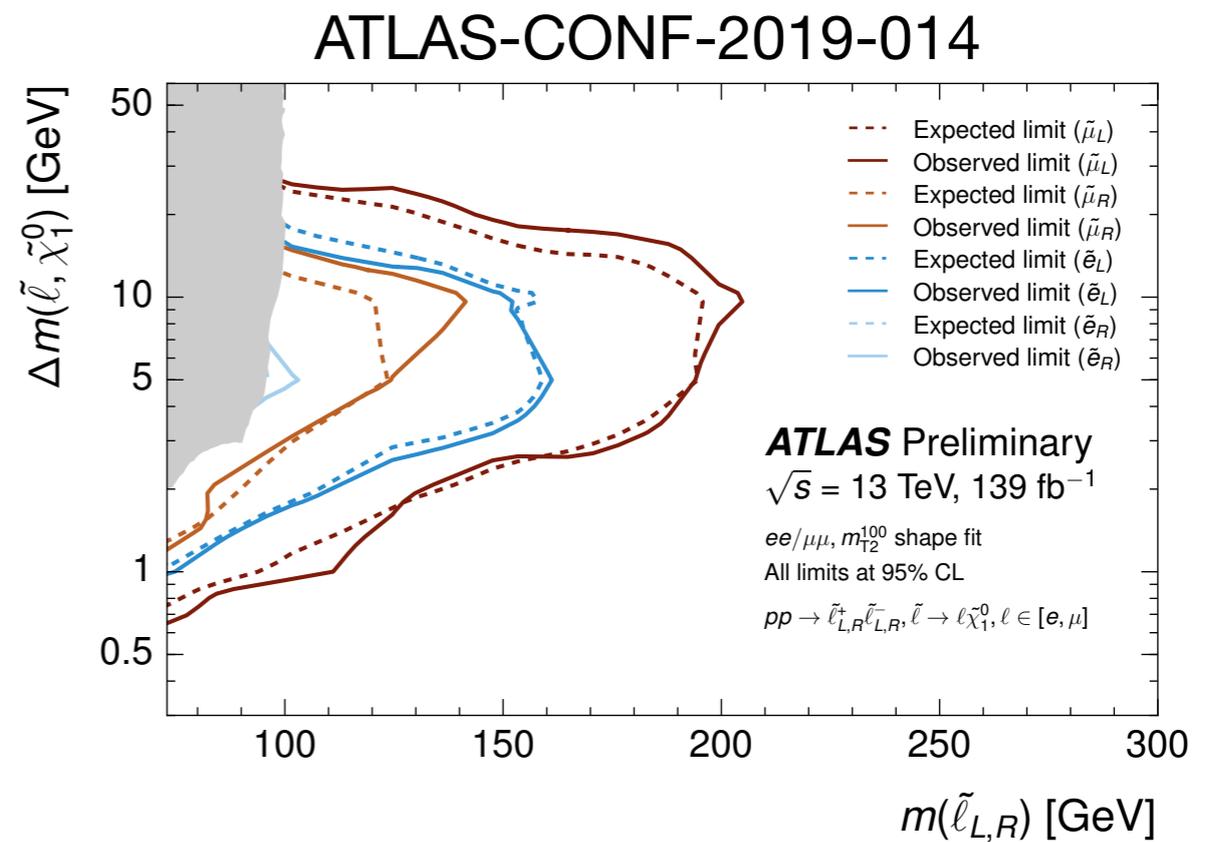
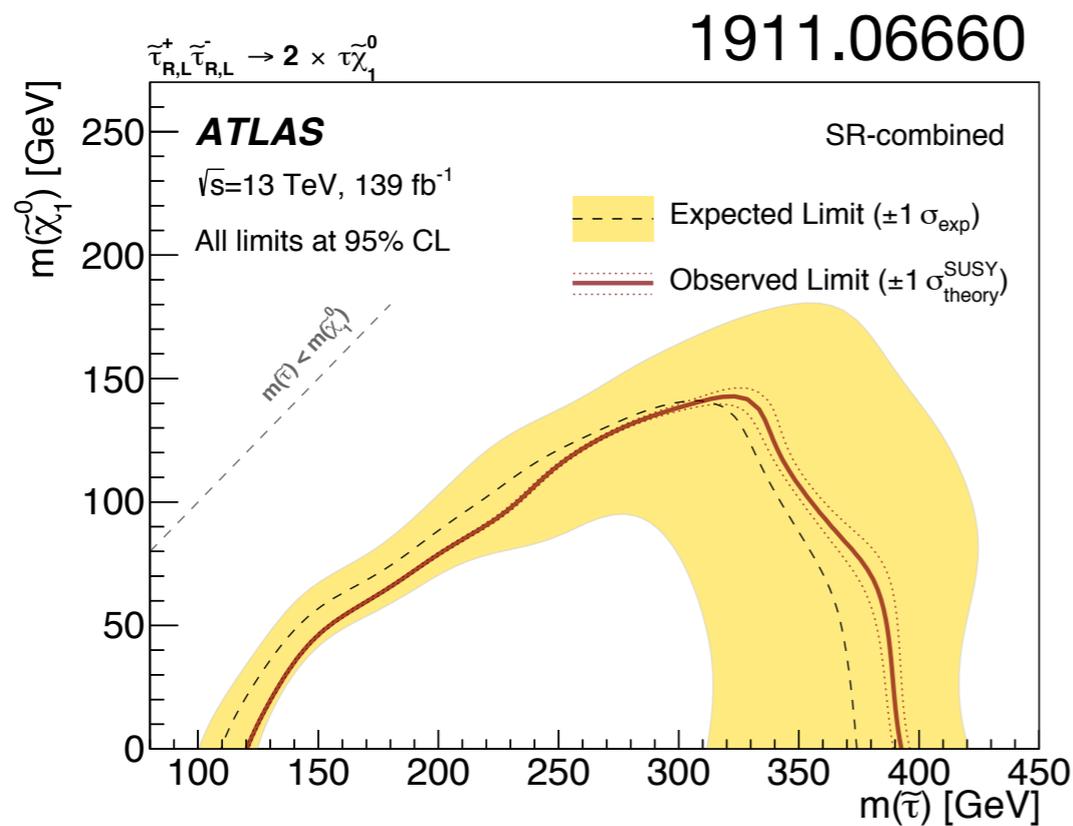
$$\mathcal{L} \supset g_{\text{DM}} \cdot \chi \eta (\text{SM})$$

- For small g_{DM} , the interaction between DM and SM becomes very weak.
- The sensitivities for direct and indirect detections are very low.
- The production rate for direct DM production at the LHC is also very small.
- Since CAP is charged under the SM gauge group, the production rate for CAP is unsuppressed at the LHC.



Collider Signature

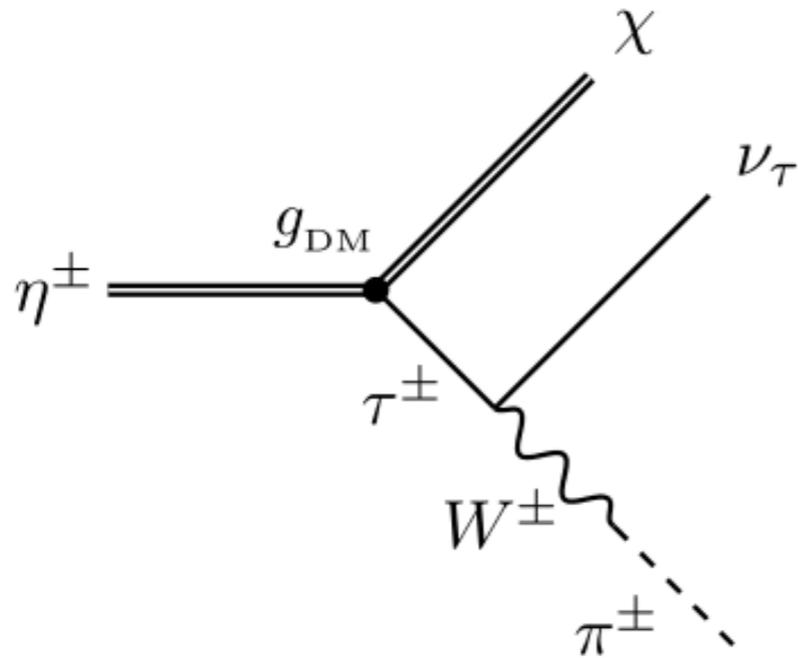
- Once CAP is produced, it will decay into DM + SM with $\mathcal{L} \supset g_{\text{DM}} \cdot \chi \eta (\text{SM})$
- Since co-annihilation requires small mass difference ($\Delta m/m_\chi \ll 4\%$), the SM particles from the decay is very soft.
- In this case, the LHC search relies on the mono-jet channel and sensitivity is very weak in general.



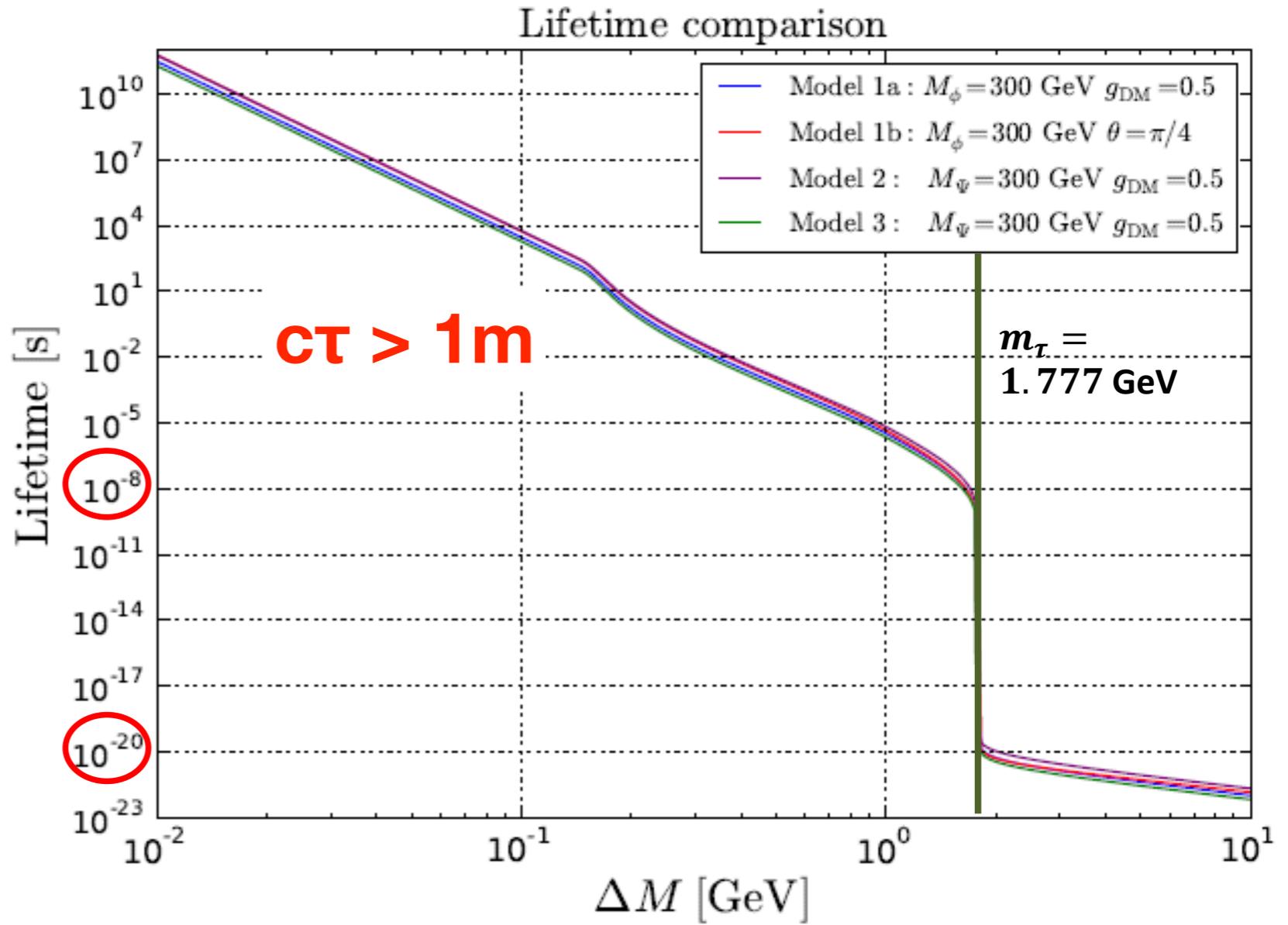
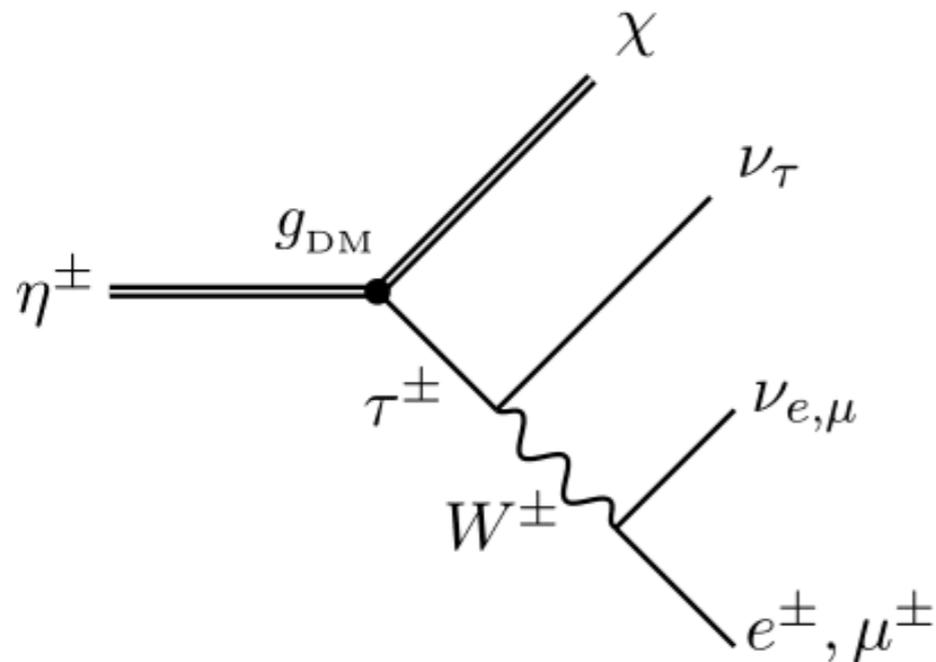
- If DM is exclusively coupled to the **tau-lepton** and $\Delta m < m_\tau$, the decay is further suppressed and CAP becomes **long-lived**.

Lifetime in tau-philic models

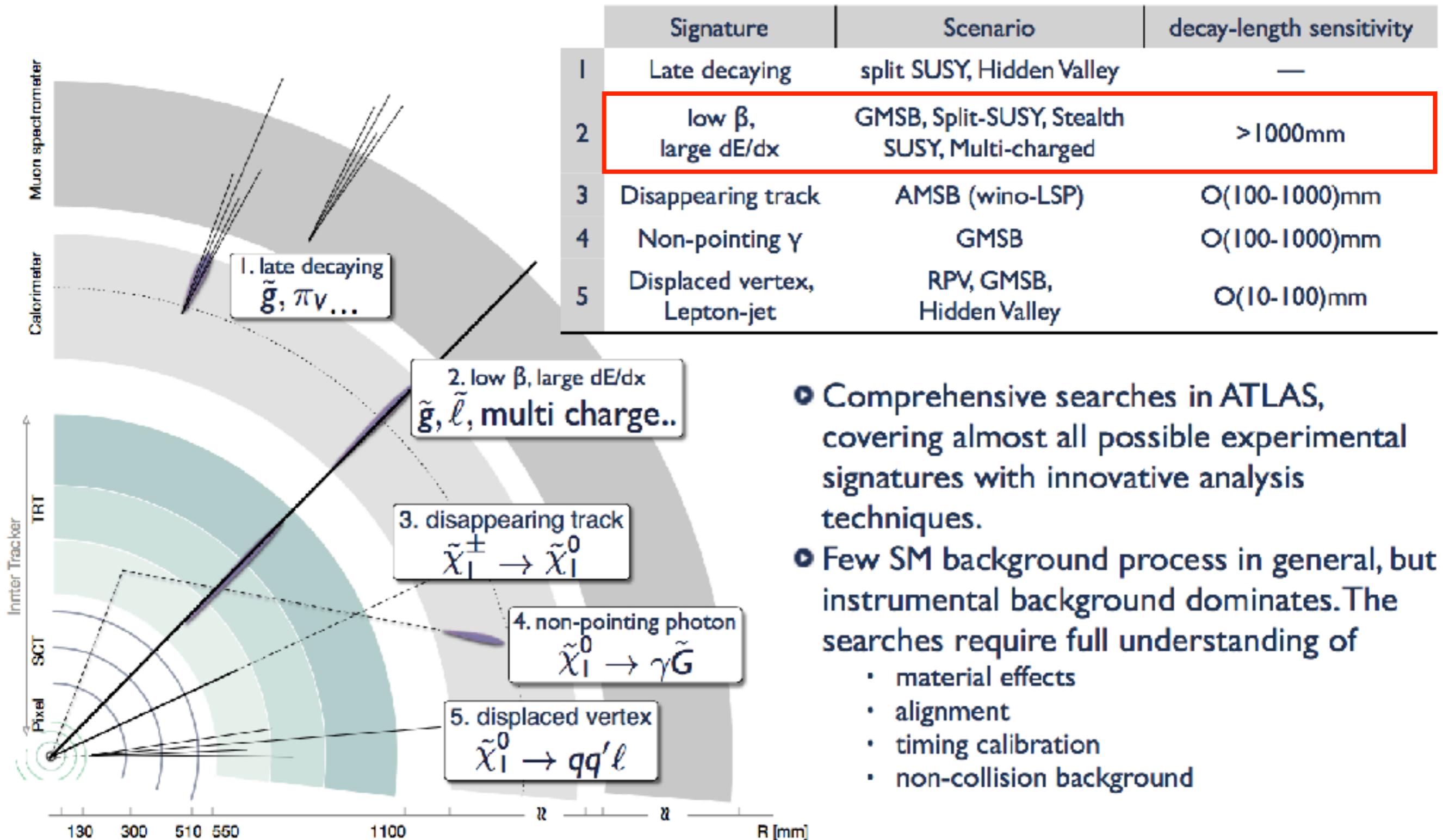
$$\mathcal{L} \supset g_{\text{DM}} \cdot \chi \eta \tau$$



Also ρ and a_1 mesons



Long-Lived Signatures



- Comprehensive searches in ATLAS, covering almost all possible experimental signatures with innovative analysis techniques.
- Few SM background process in general, but instrumental background dominates. The searches require full understanding of
 - material effects
 - alignment
 - timing calibration
 - non-collision background

Simplified Models (SMS)

- A standard signature to search for dark matter at colliders is the mono-X (or multi-jets) plus missing energy.
- These searches are being exploited and interpreted in terms of **simplified dark matter models with mediators**.

Dark Matter + mediator + Standard Model particles

- [A growing number of the analyses are also dedicated to the direct search of the mediators which can decay back to the SM.]
- **We consider instead an alternative DM scenario** characterised by **simplified models without mediators**.

Dark Matter + co-annihilation partner + Standard Model particles

- Our dark sector includes a co-annihilation partner (CAP) particle instead of a mediator (in addition to the cosmologically stable DM).

SMS for long-lived coannihilating DM with tau

Model-1a			
Component	Field	Charge	Interaction
DM	Majorana fermion (χ)	$Y = 0$	$\phi^*(\chi\tau_R) + \text{h.c.}$
CAP	Complex scalar (ϕ)	$Y = -1$	

fermionic DM, scalar CAP
e.g. neutralino-stau (SUSY)

Model-2			
Component	Field	Charge	Interaction
DM	Real scalar (S)	$Y = 0$	$S(\bar{\Psi}P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

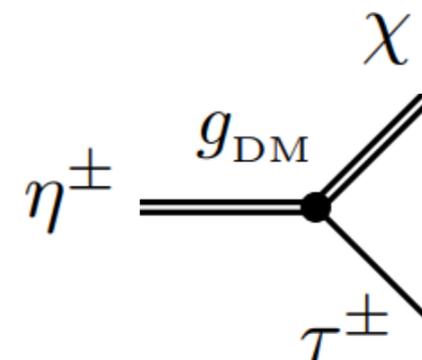
scalar DM, fermionic CAP
e.g. dilaton, KK-tau (extra-dim)

Model-3			
Component	Field	Charge	Interaction
DM	Vector (V_μ)	$Y = 0$	$V_\mu(\bar{\Psi}\gamma^\mu P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

vector DM, fermionic CAP
e.g. KK-photon, KK-tau (UED)

3 free parameters:

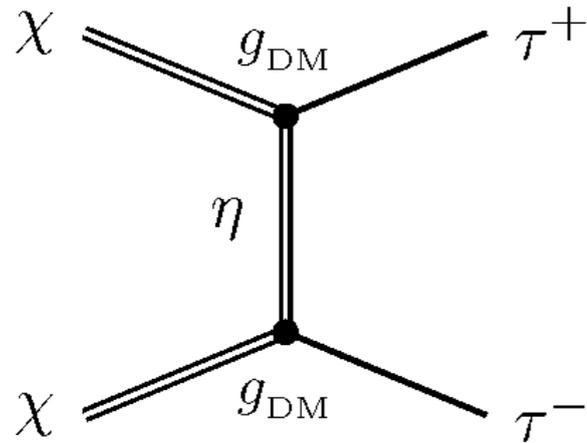
$$m_\chi, \Delta m, g_{\text{DM}}$$



$$\chi = \{\chi, S, V_\mu\}$$

$$\eta = \{\phi, \Psi\}$$

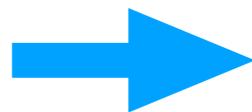
DM annihilation



In model-1 and -2, the DM is a Majorana fermion and a real scalar, respectively.

For those dark matters, the initial state is s-wave (spin-0) (Pauli blocking for the fermion case), therefore the s-wave part of this annihilation cross-section is chiral suppressed by the tau-lepton mass.

$$(\sigma v)_{\text{ann}}^{\text{s wave}} \sim \frac{g_{\text{DM}}^4 m_{\tau}^2}{32\pi m_{\chi}^4}$$



Coannihilation is very important

Other than s-wave:

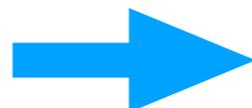
$$\sigma v \propto v^2$$

p-wave suppressed for Majorana DM

$$\sigma v \propto v^4$$

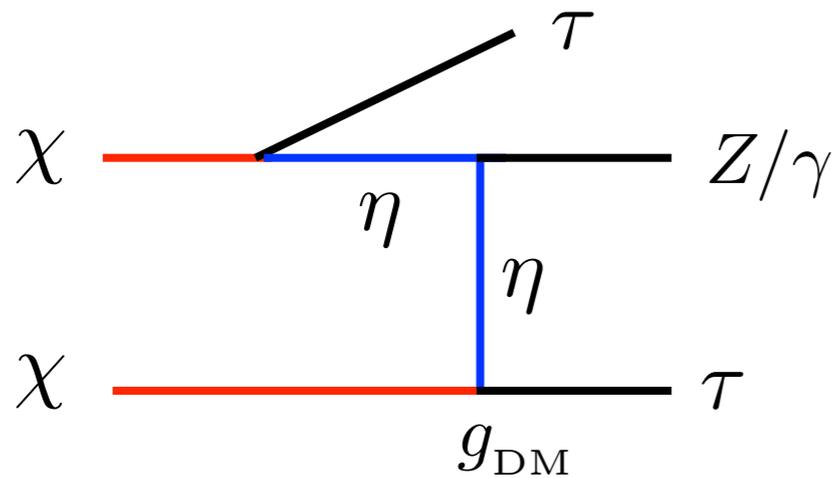
d-wave suppressed for scalar DM

[$v/c \ll 1$ at the present Universe]



indirect detection not promising(?)

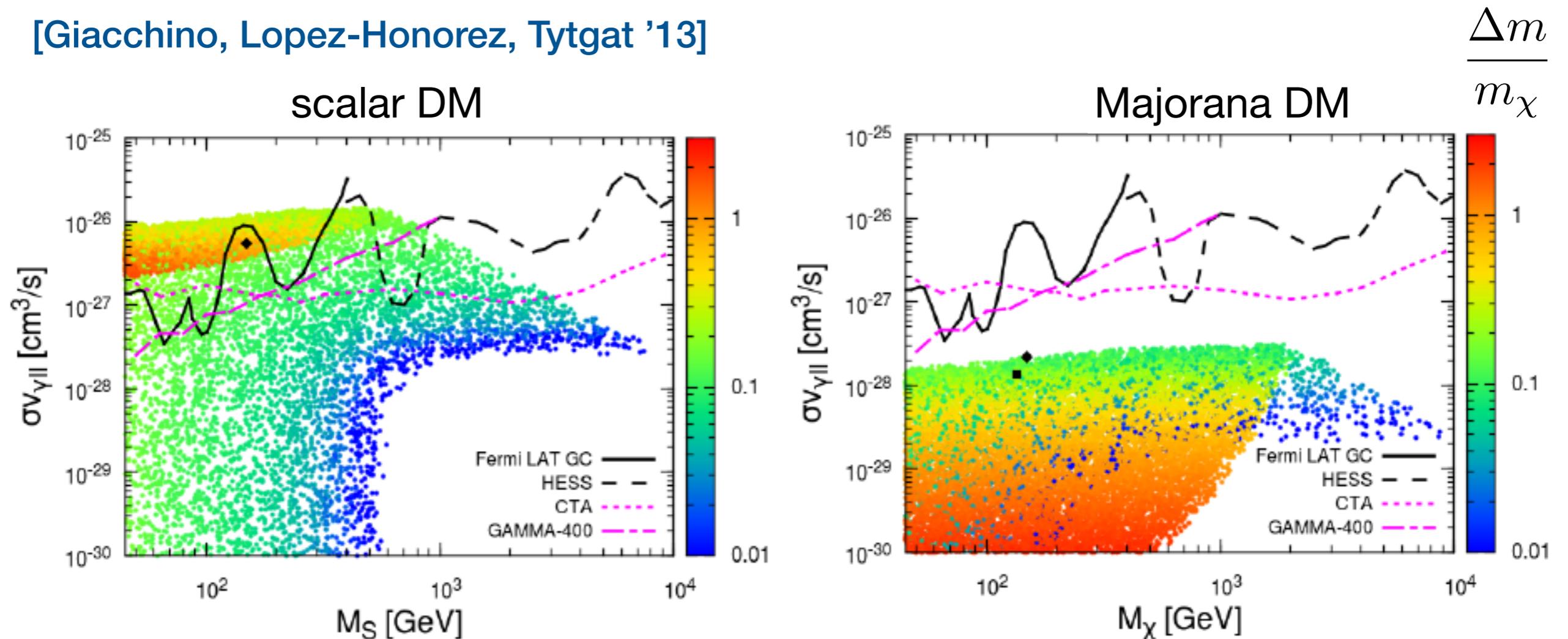
Indirect Detection



The DM can be converted into a slightly off-shell CAP by emitting a soft tau, the CAP then co-annihilates with another DM into a pair of SM particles.

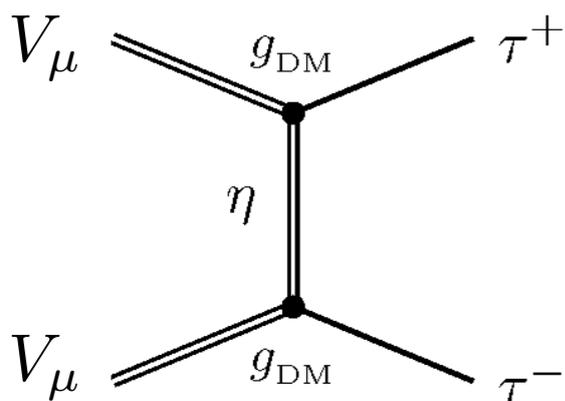
This channel is not chiral suppressed and turns out to be the dominant annihilation channel.

[Giacchino, Lopez-Honorez, Tytgat '13]



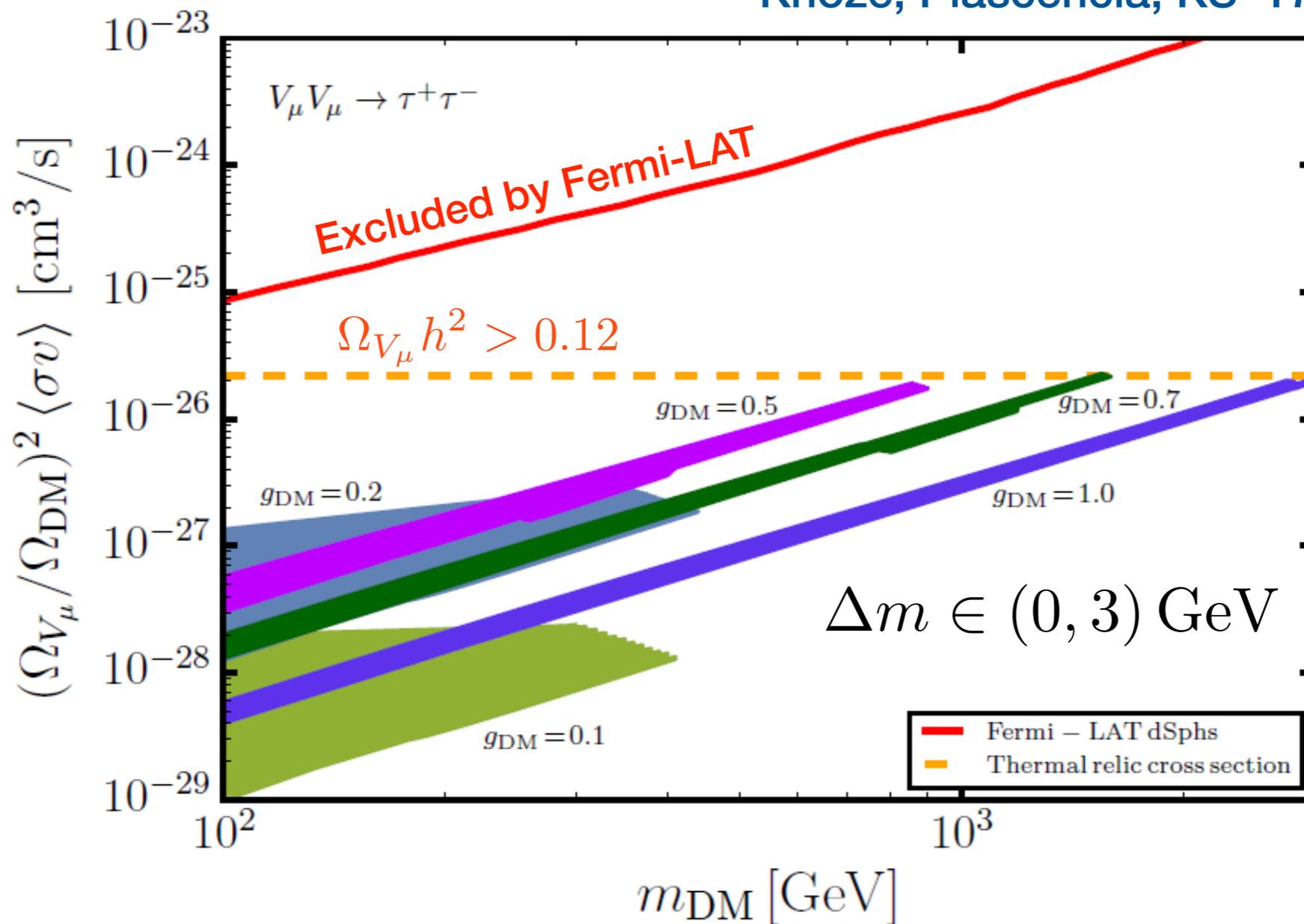
Blue points correspond to the parameter region of interest; $\Delta m/m_\chi \ll 4\%$.

Indirect Detection



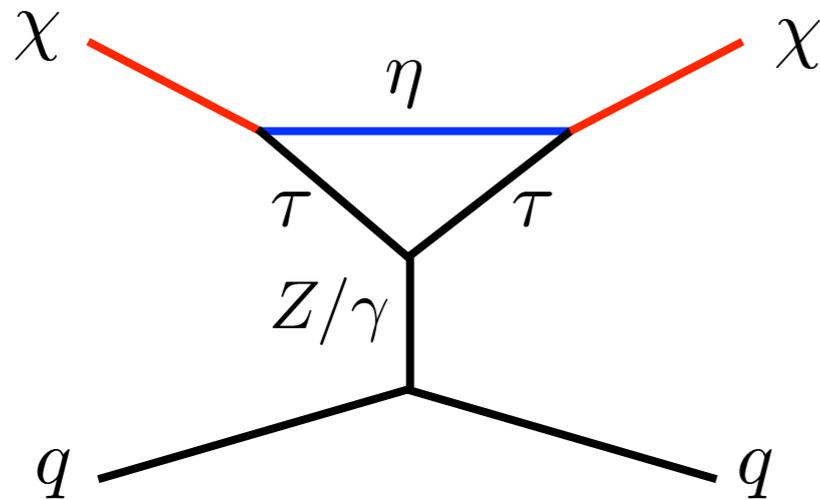
Vector DM annihilation is not chiral suppressed.

Khoze, Plascencia, KS '17



Direct Detection

- The anapole operator for direct detection is generated at 1-loop.



$$\mathcal{A} \bar{\chi} \gamma_{\mu} \gamma_5 \chi \partial^{\nu} F_{\mu\nu}$$

$$m_{\text{DM}} \simeq 500 \text{ GeV and } \Delta M/m_{\tau} < 1,$$

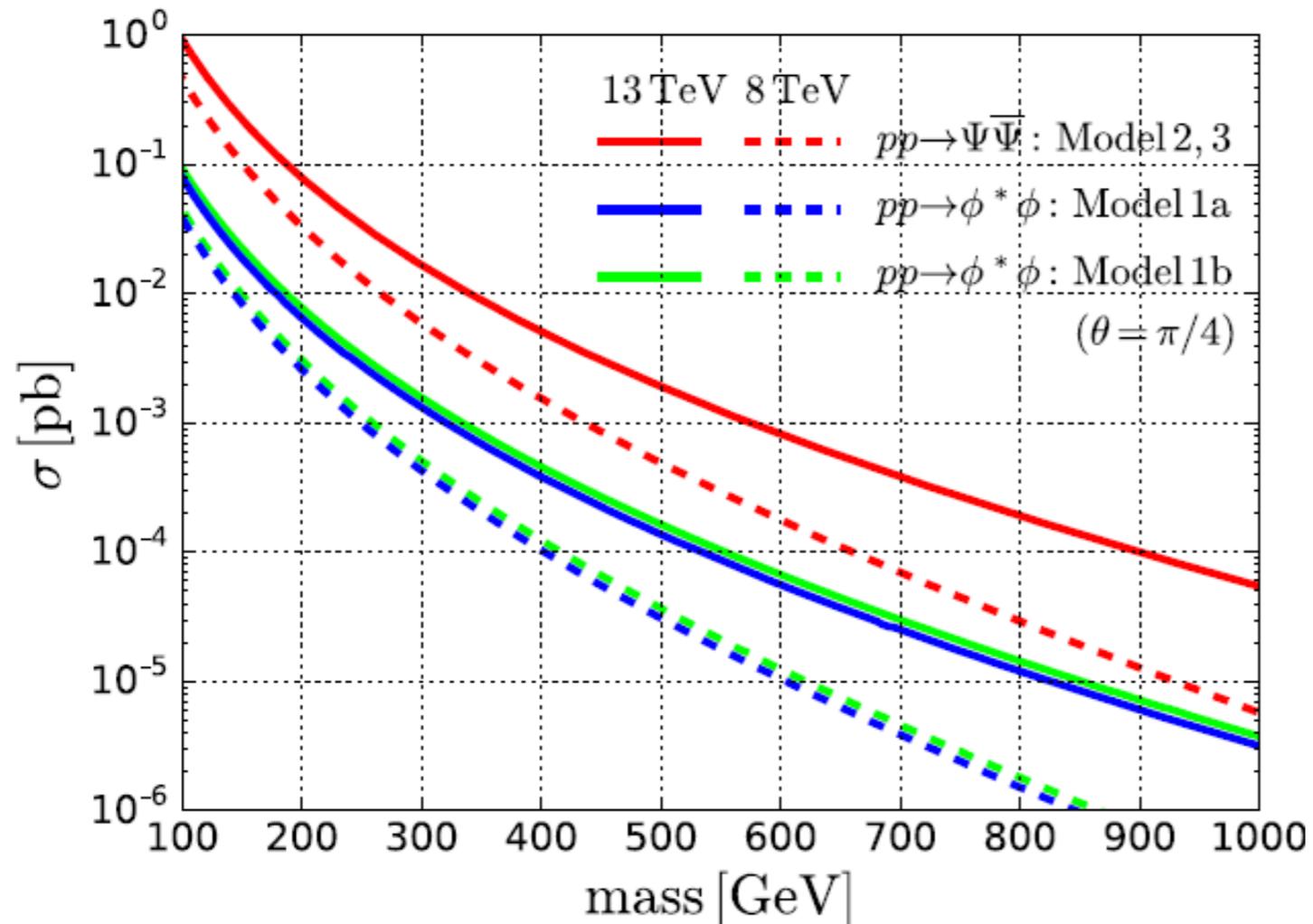
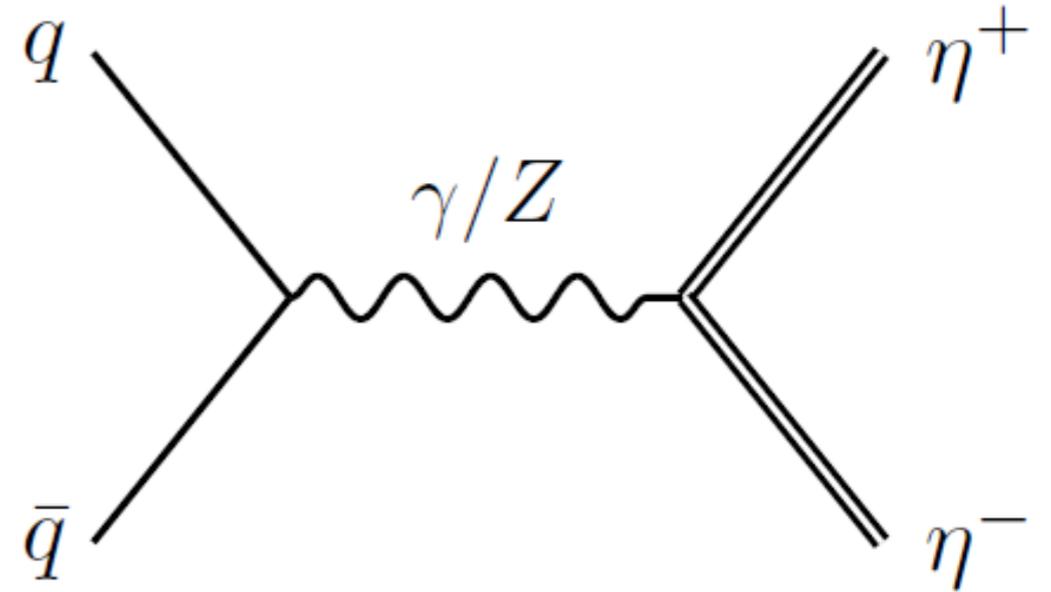
$$\mathcal{A}/g_{\text{DM}}^2 \sim 8 \cdot 10^{-7} [\mu_N \cdot \text{fm}]$$

$$\text{LUX } \mathcal{A} > 2 \times 10^{-5} [\mu_N \text{ fm}]$$

- The current limit is more than one order of magnitude smaller.

Direct Production at LHC

- Drell-Yann pair production of co-annihilation partner



- We study Dirac fermion and complex scalar as co-annihilation partners

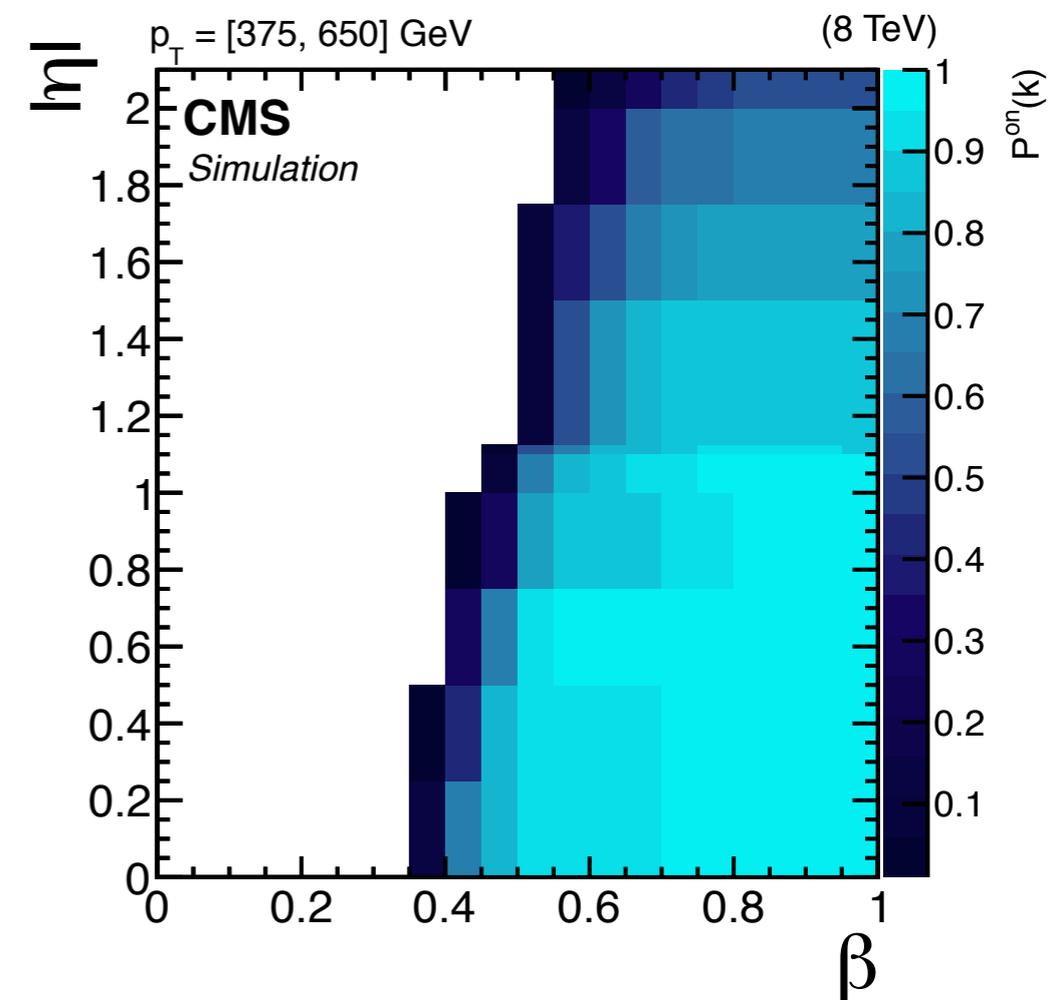
Recasting HSCP analysis

- In order to constrain the long-lived coannihilating DM simplified models, we recast the heavy stable charged particle (HSCP) analysis by CMS (8TeV, 18.8 fb⁻¹) [1305.0491].
- We used the recipe provided by CMS [1502.02522] for recasting the HSCP analysis and used the efficiency maps provided in the paper.

[1305.0491]

	$ Q < 1e$	tracker+TOF	tracker-only	$ Q > 1e$
$ \eta $	< 2.1			
p_T (GeV/c)	> 45			
d_z and d_{xy} (cm)	< 0.5			
σ_{p_T} / p_T	< 0.25			
Track χ^2 / n_d	< 5			
# Pixel hits	> 1			
# Tracker hits	> 7			
Frac. Valid hits	> 0.8			
$\Sigma p_T^{\text{trk}}(\Delta R < 0.3)$ (GeV/c)	< 50			
# dE/dx measurements	> 5			
dE/dx strip shape test	yes		no	
$E_{\text{cal}}(\Delta R < 0.3) / p$	< 0.3		—	
I_h (MeV/cm)	< 2.8	> 3.0		
ΔR to another track	$< \pi - 0.3$	—		

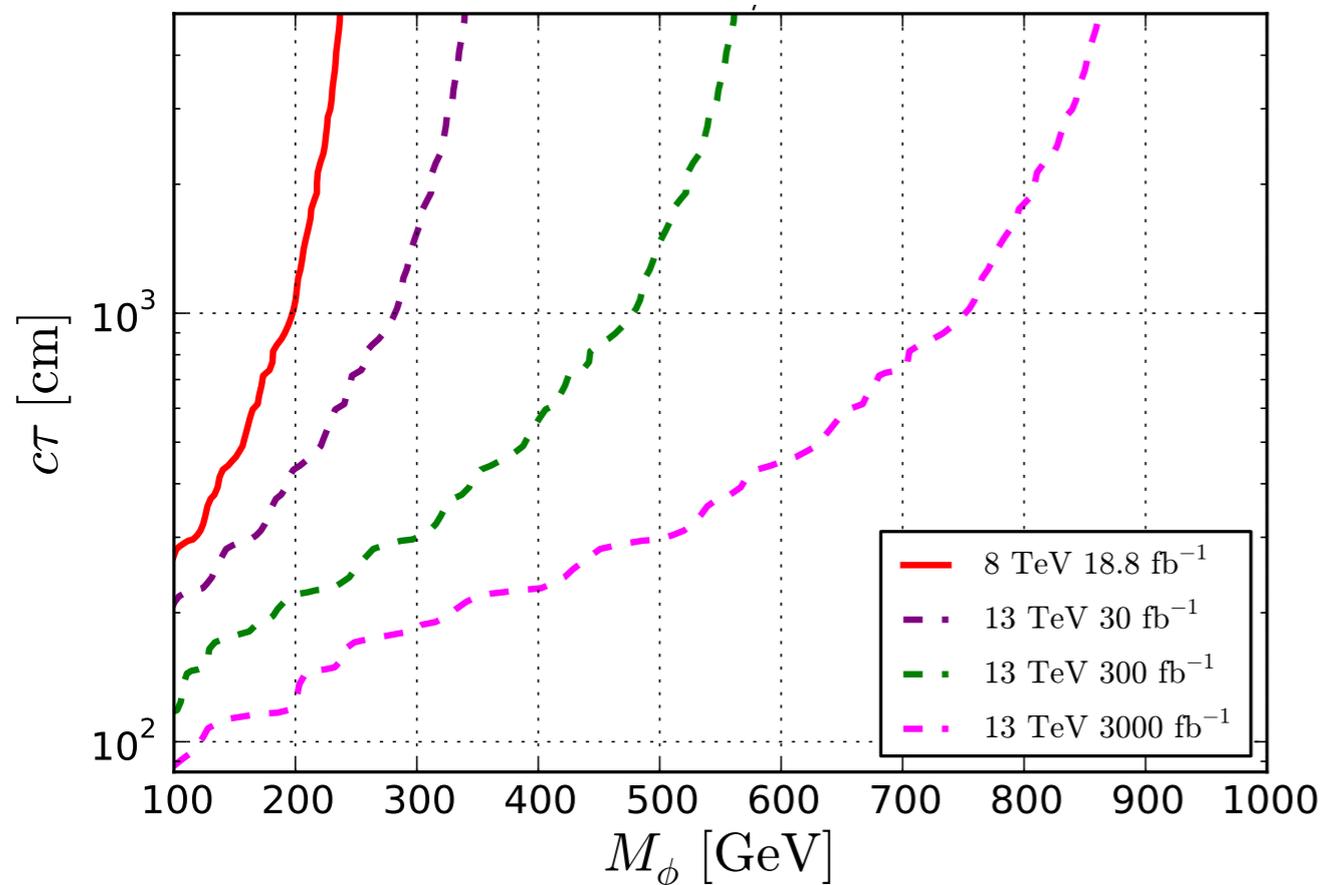
[1502.02522]



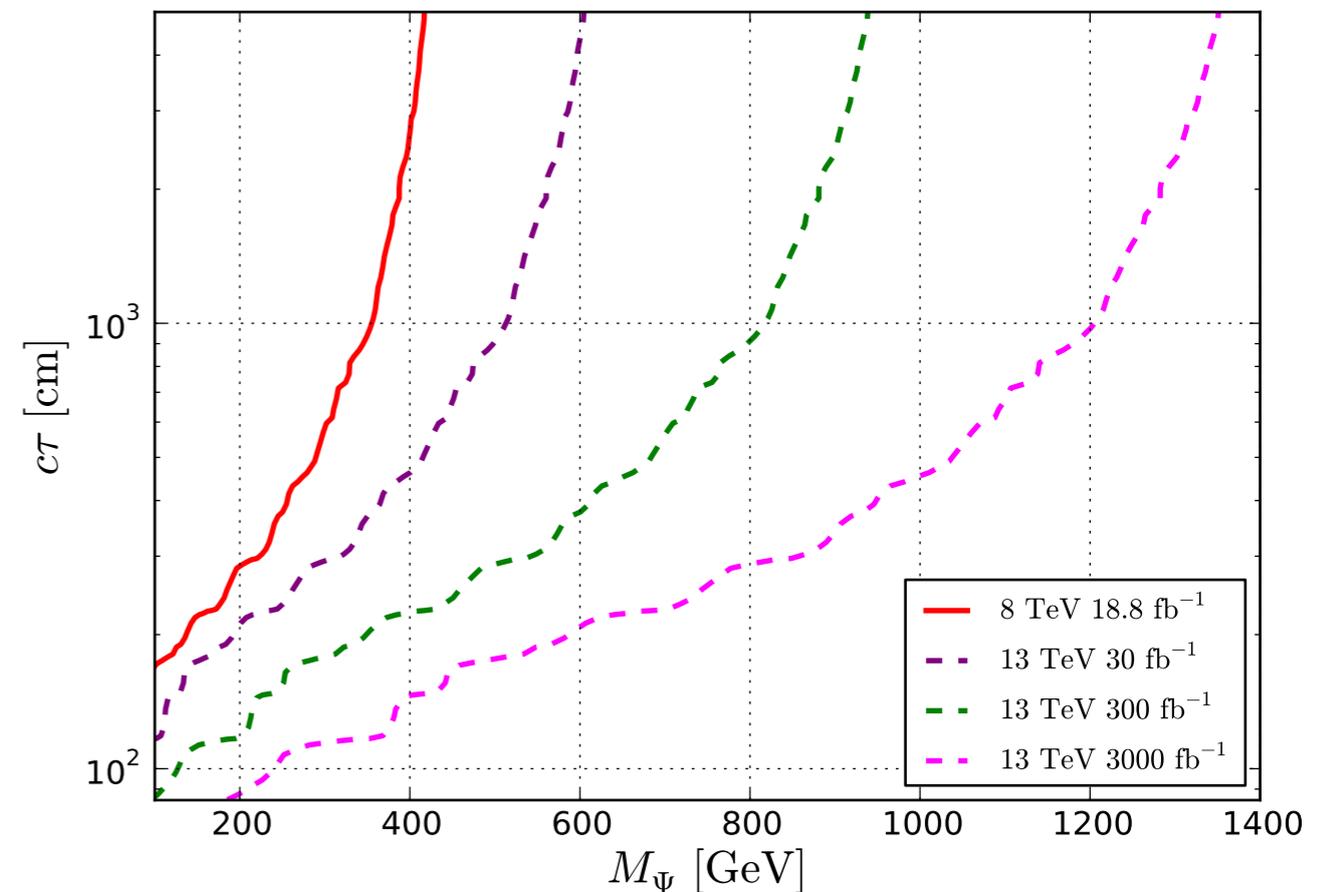
Recasting HSCP analysis

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Scalar CAP (Model-1)



Fermionic CAP (Model-2, 3)



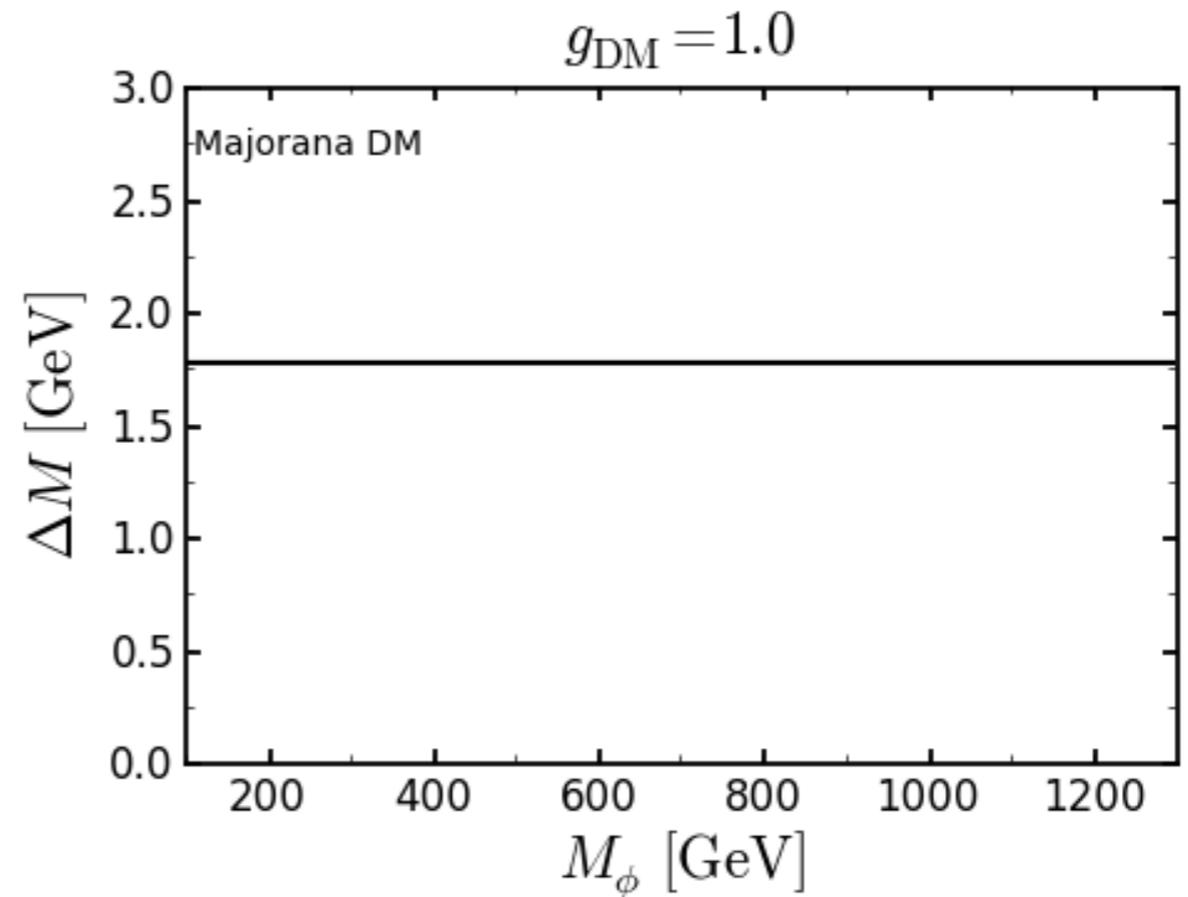
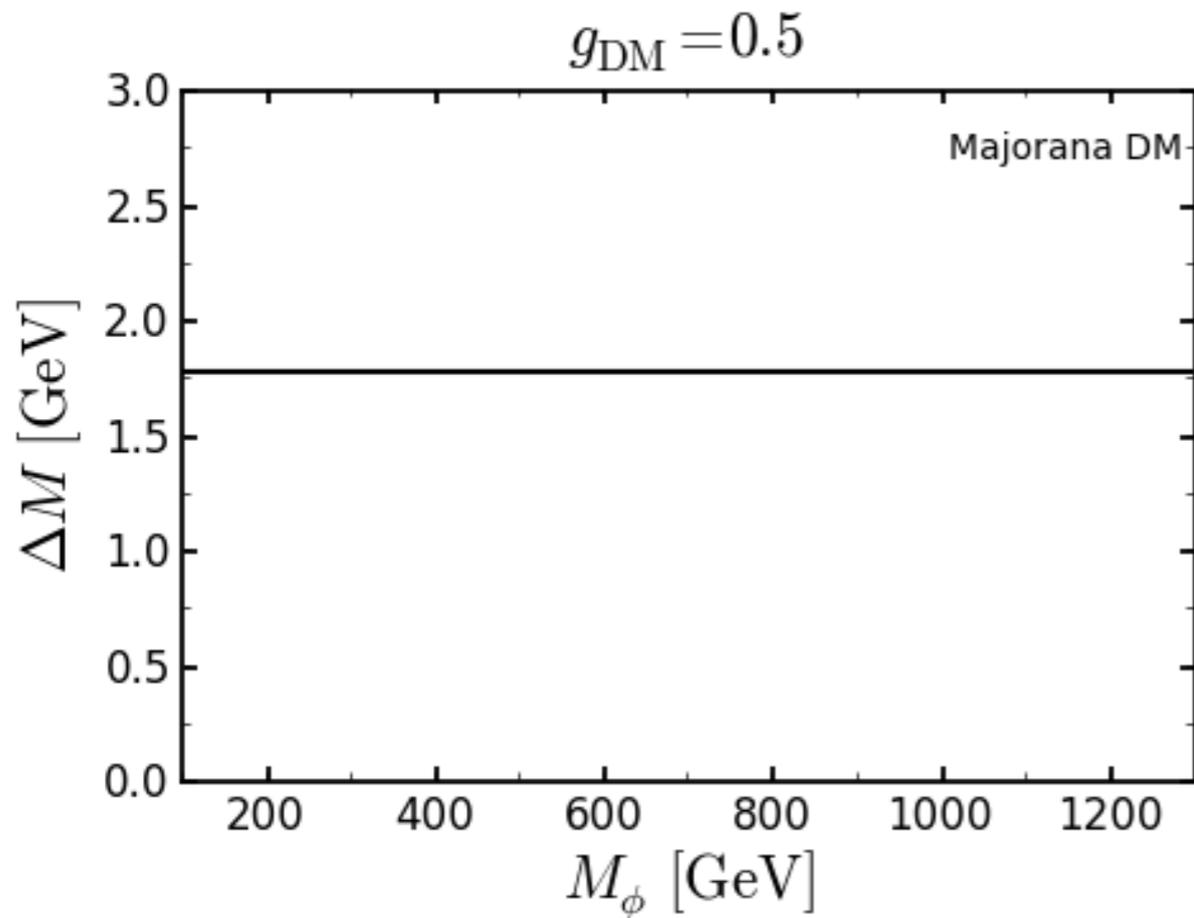
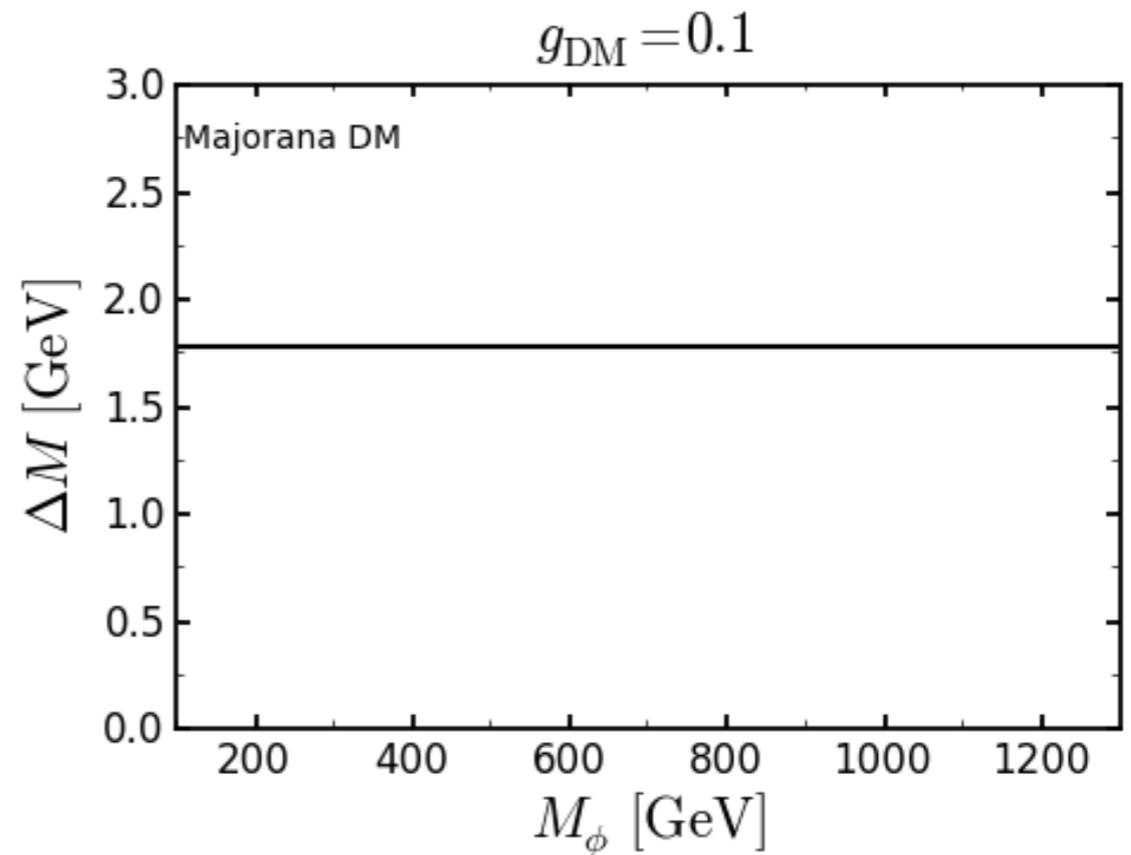
Majorana Dark Matter

DM **CAP** ($Y = 1$ $L_\tau = 1$)

χ ϕ

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

**Gauge-invariant and renormalizable,
no problems of unitarity**



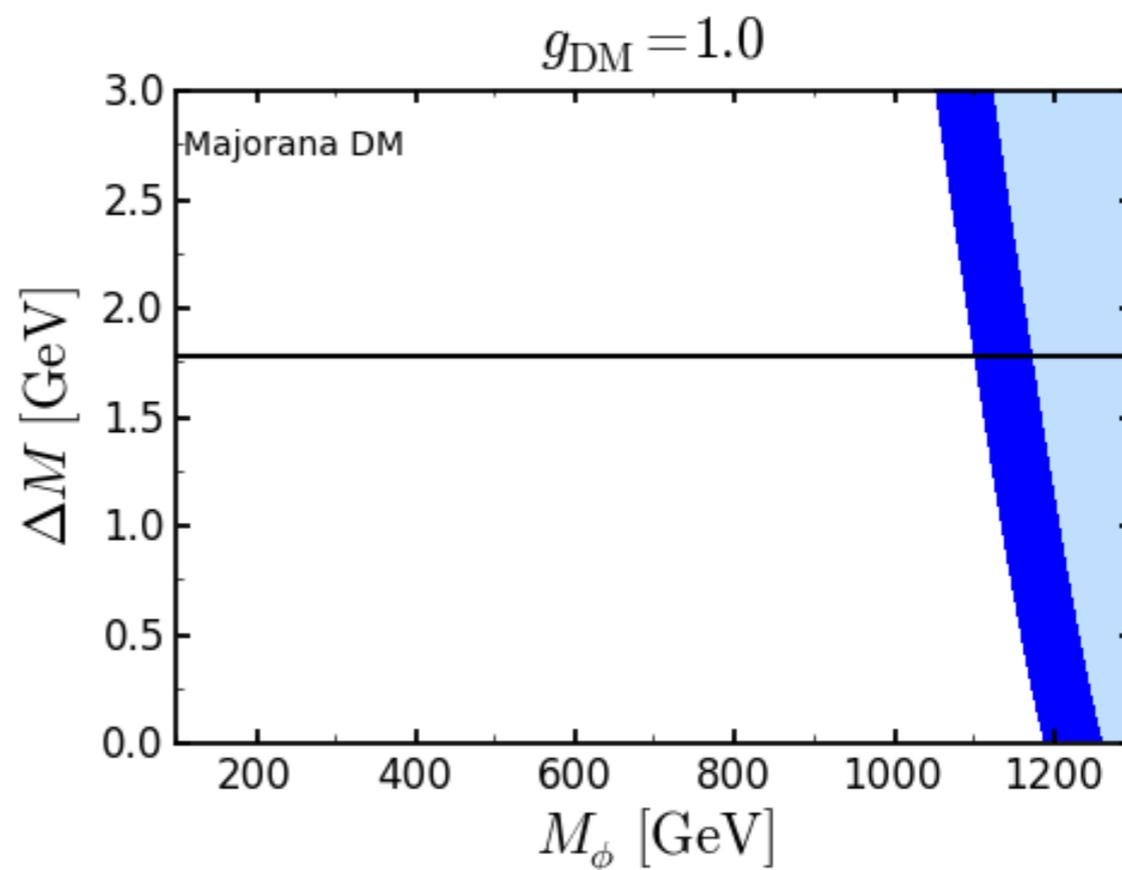
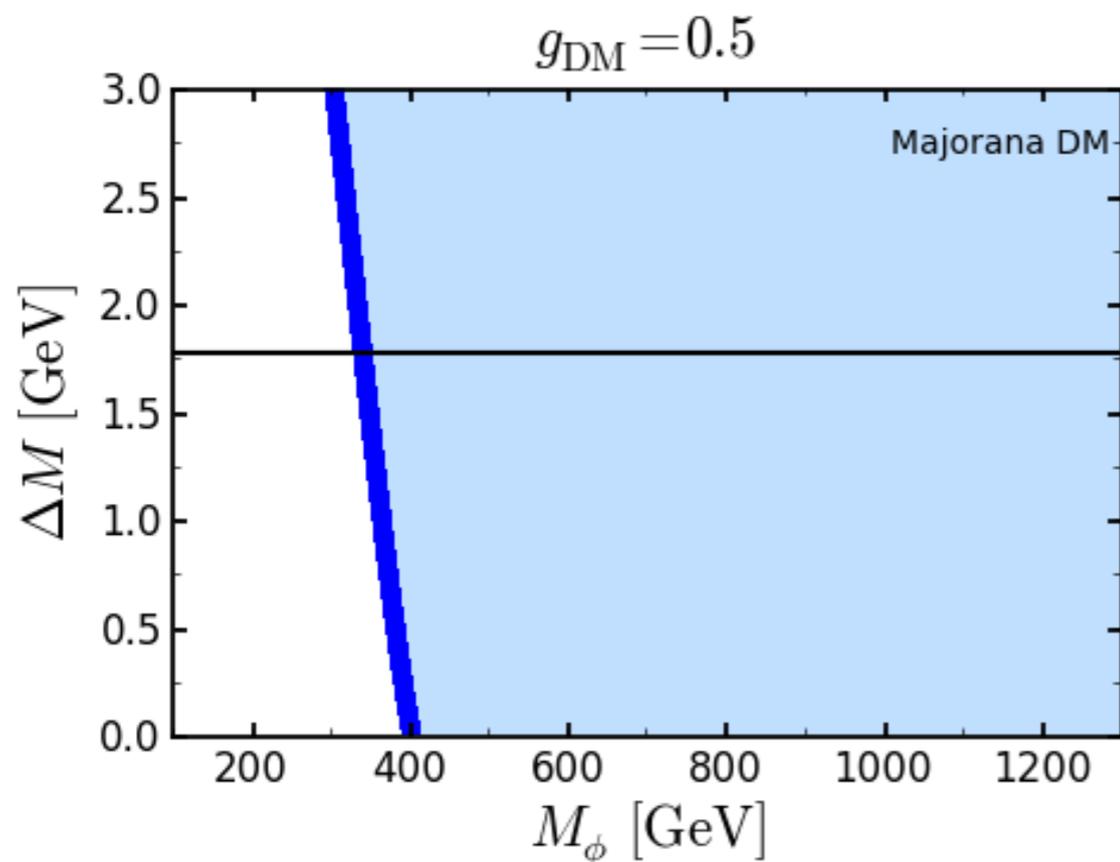
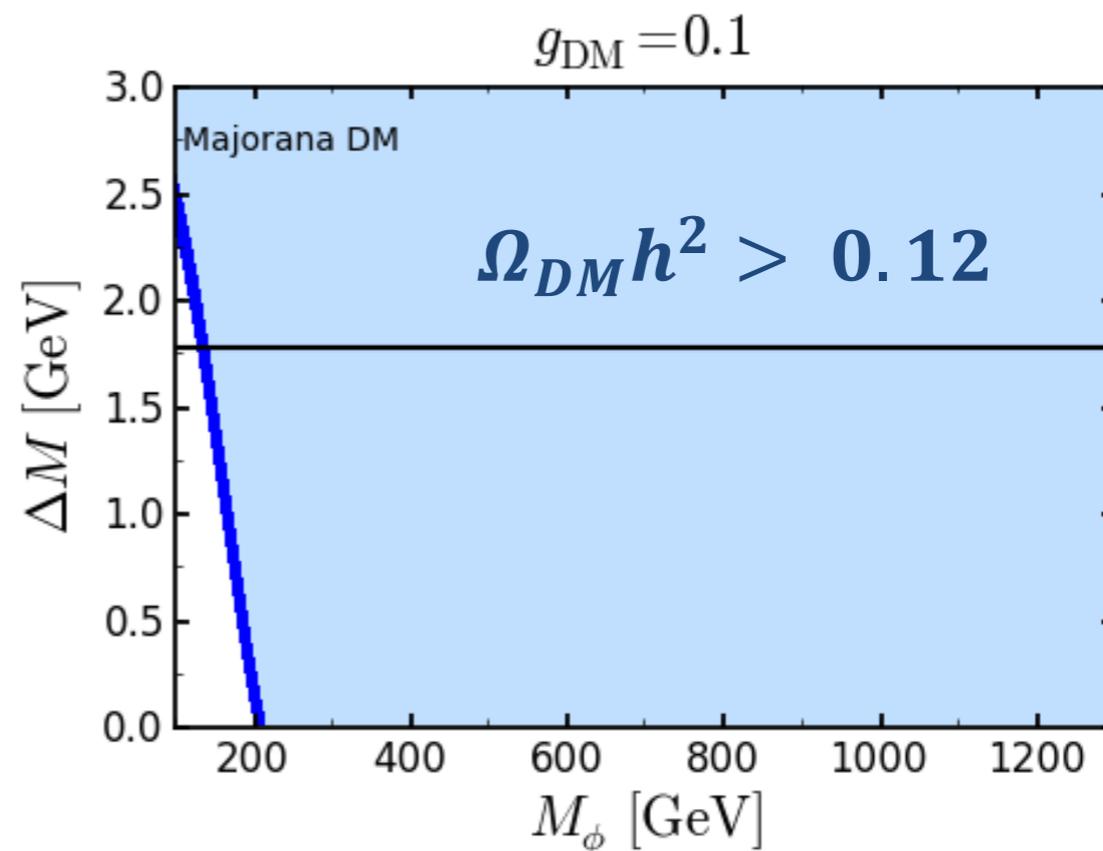
Majorana Dark Matter

DM CAP ($Y=-1$)

χ ϕ

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

Gauge-invariant and renormalizable,
no problems of unitarity



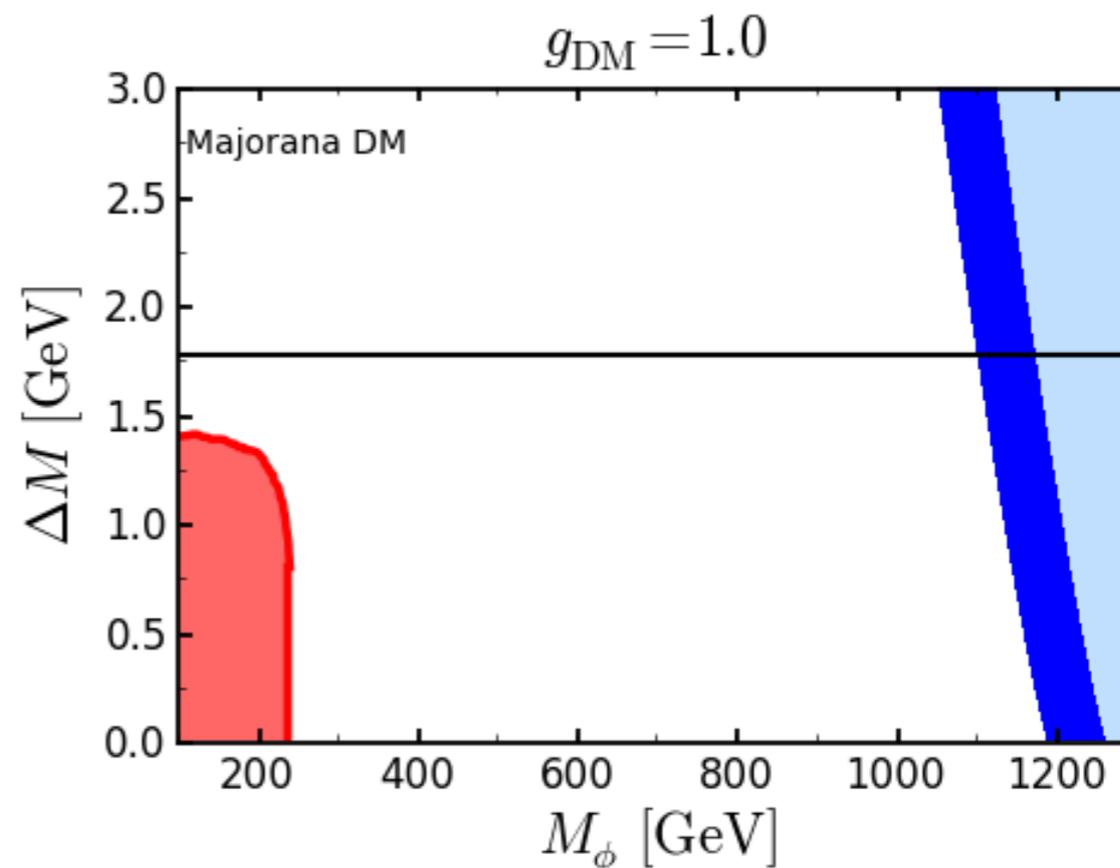
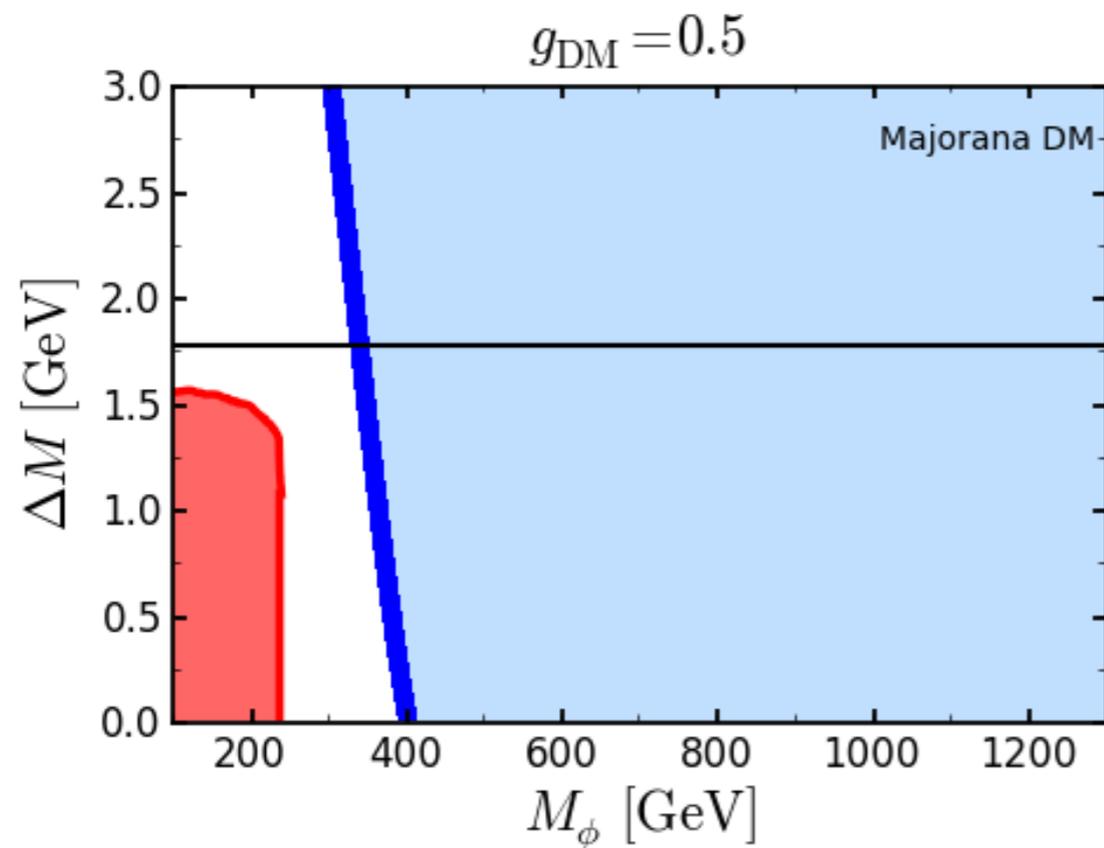
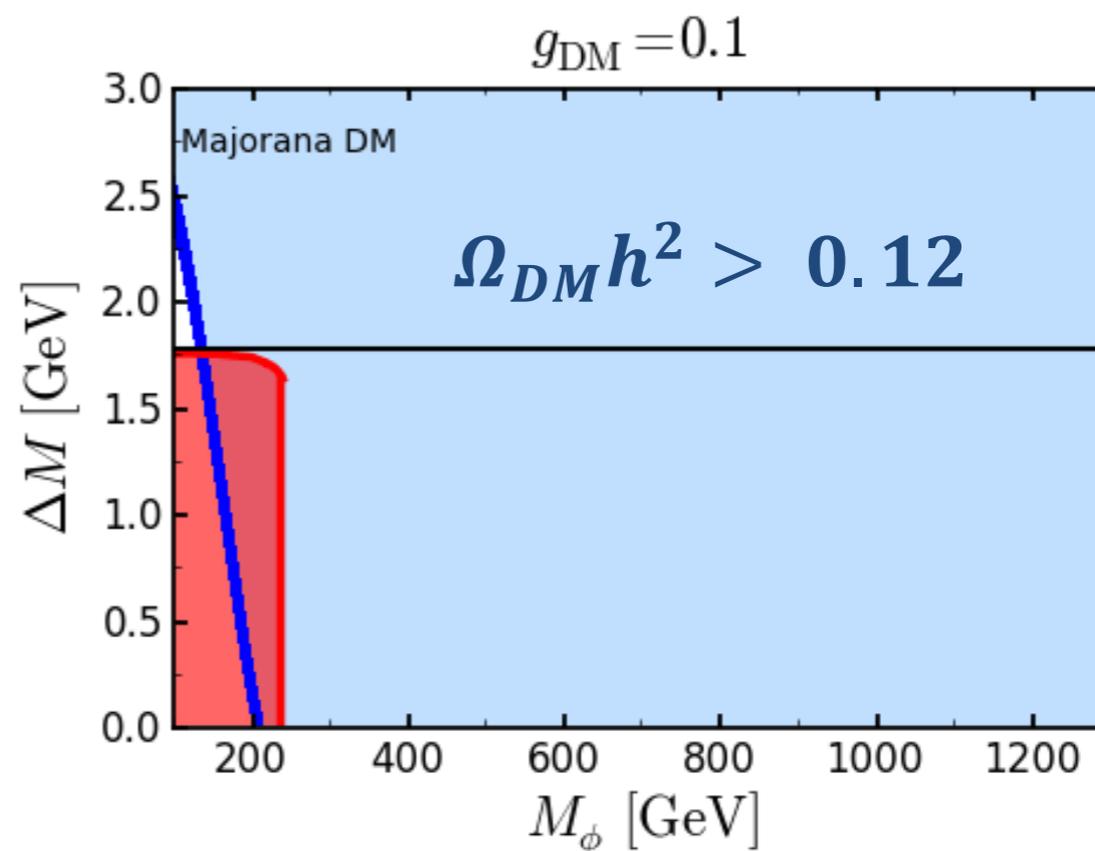
Majorana Dark Matter

DM CAP ($Y = 1 \quad L_\tau = 1$)

χ ϕ

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

CMS 8 TeV 18.8 fb⁻¹



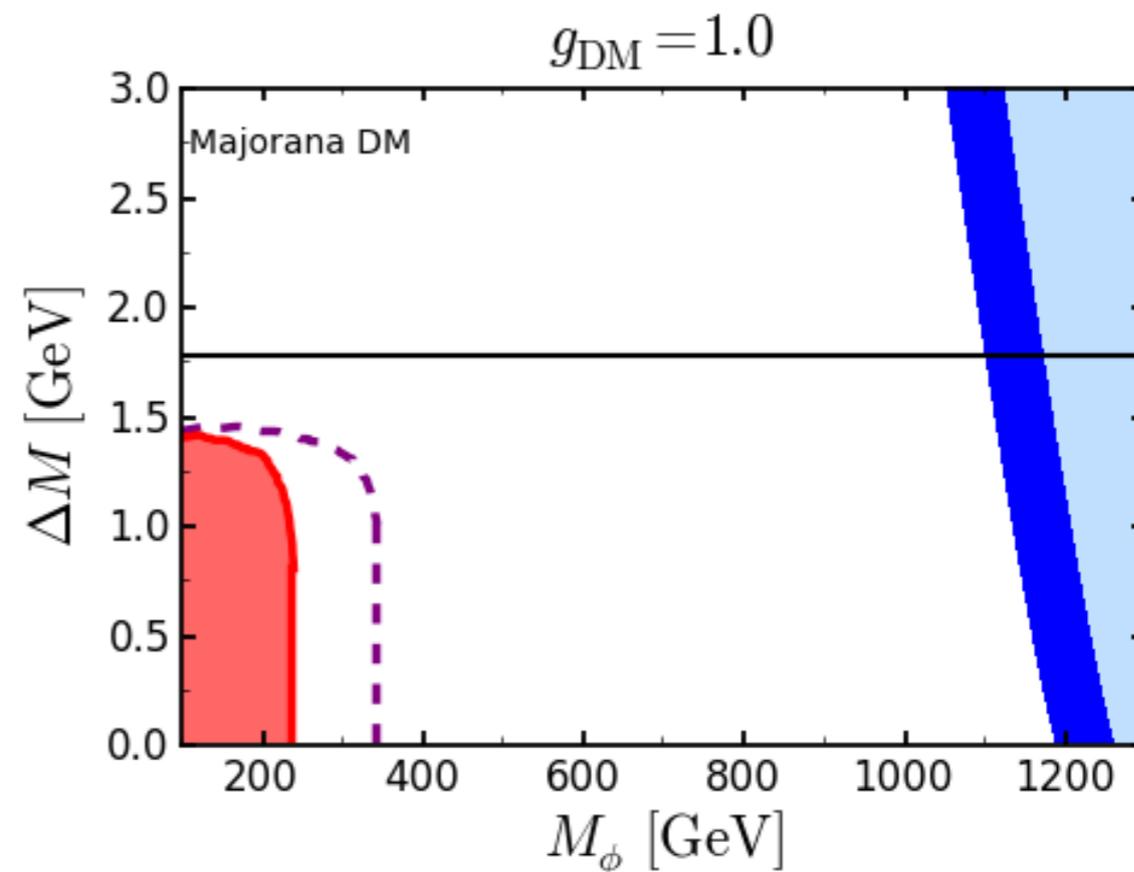
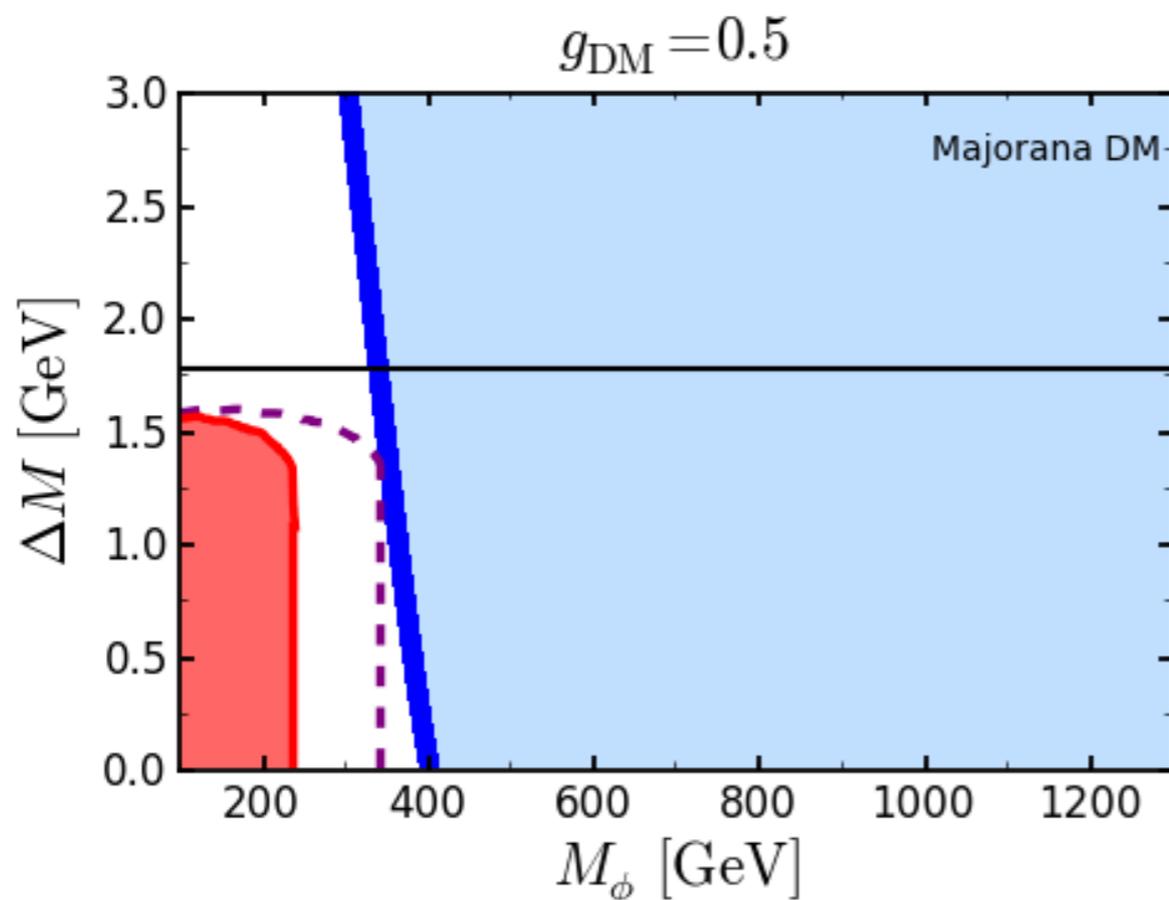
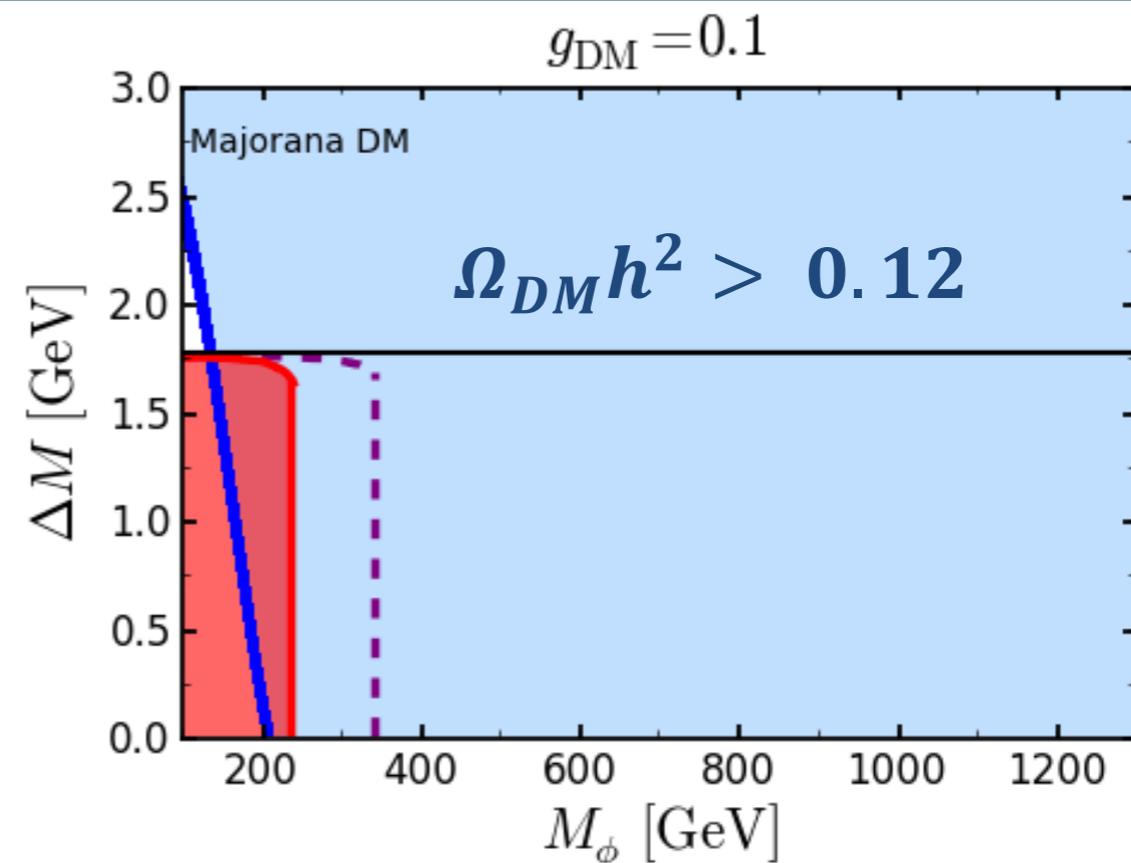
Majorana Dark Matter

DM CAP ($Y = 1 \quad L_\tau = 1$)

$\chi \quad \phi$

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

13 TeV 30 fb^{-1}



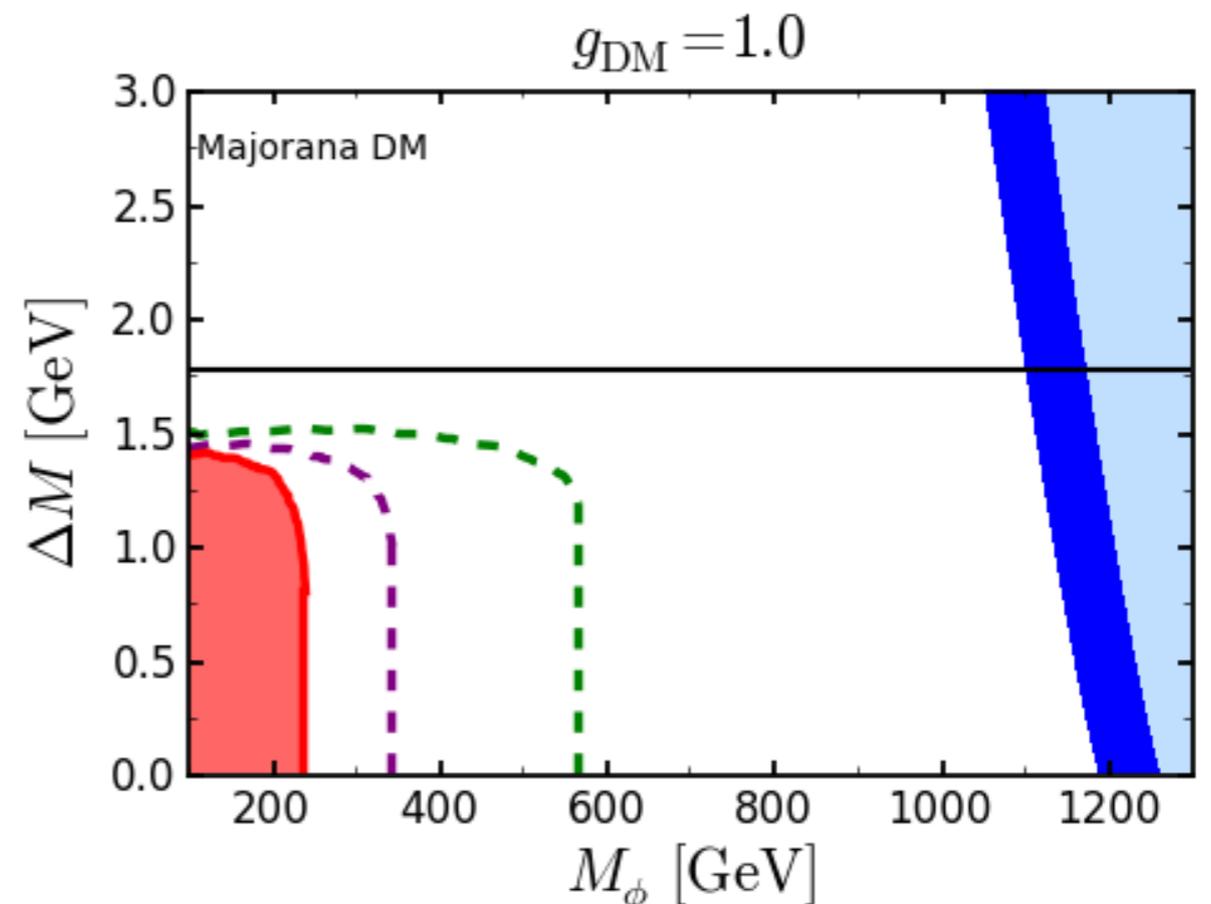
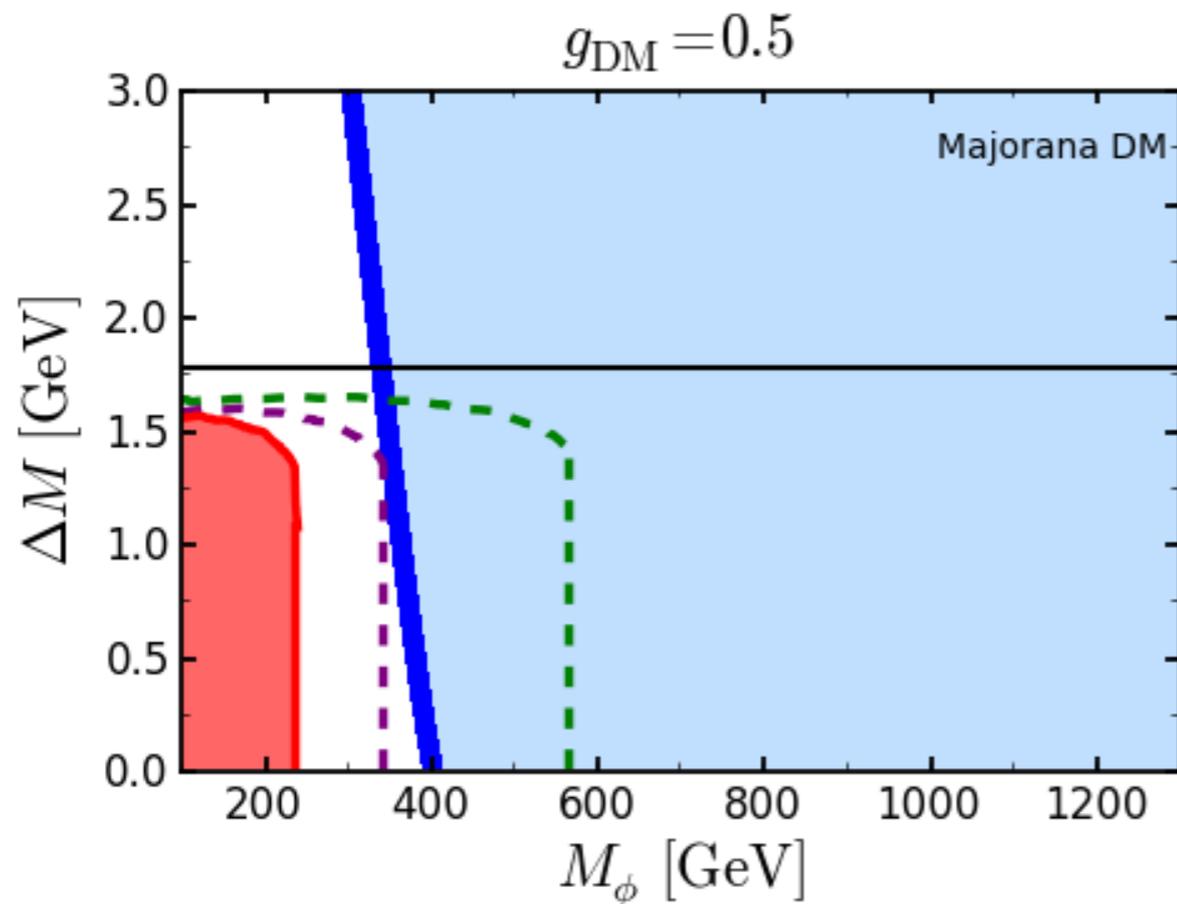
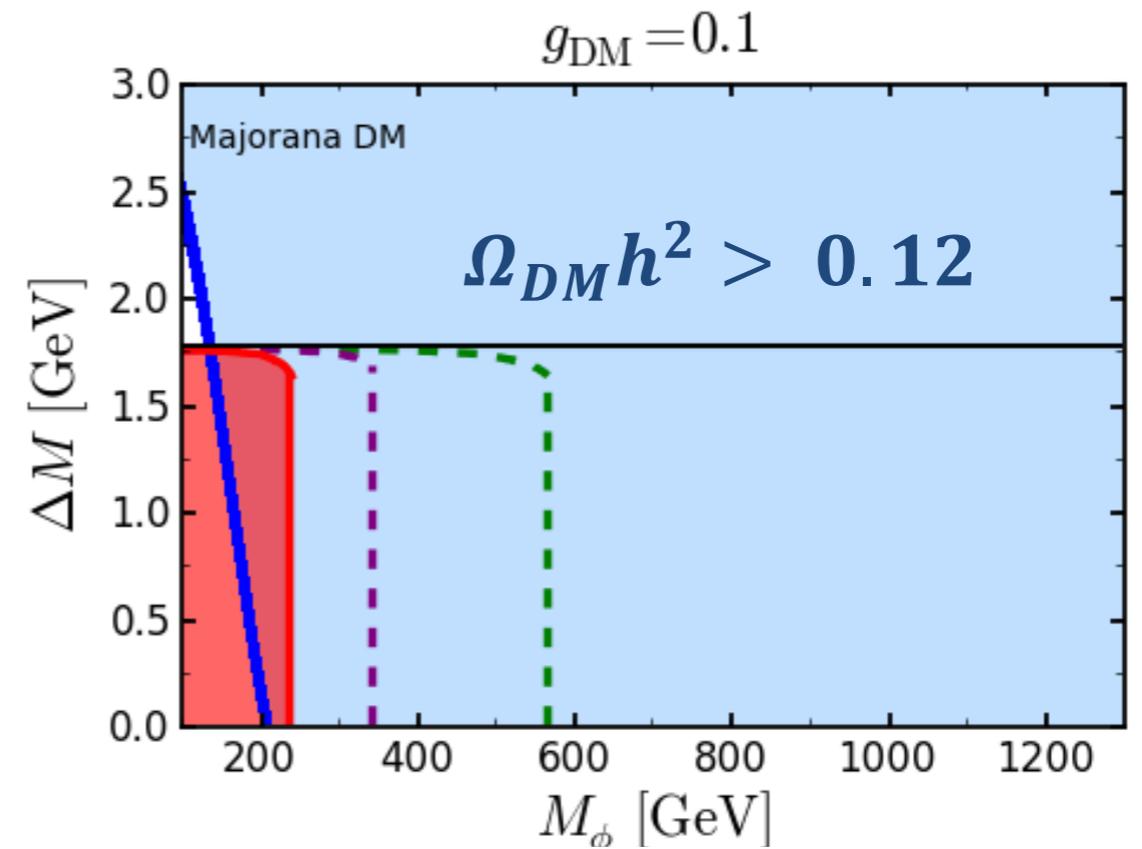
Majorana Dark Matter

DM CAP ($Y = 1 \quad L_\tau = 1$)

$\chi \quad \phi$

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

13 TeV 300 fb^{-1}



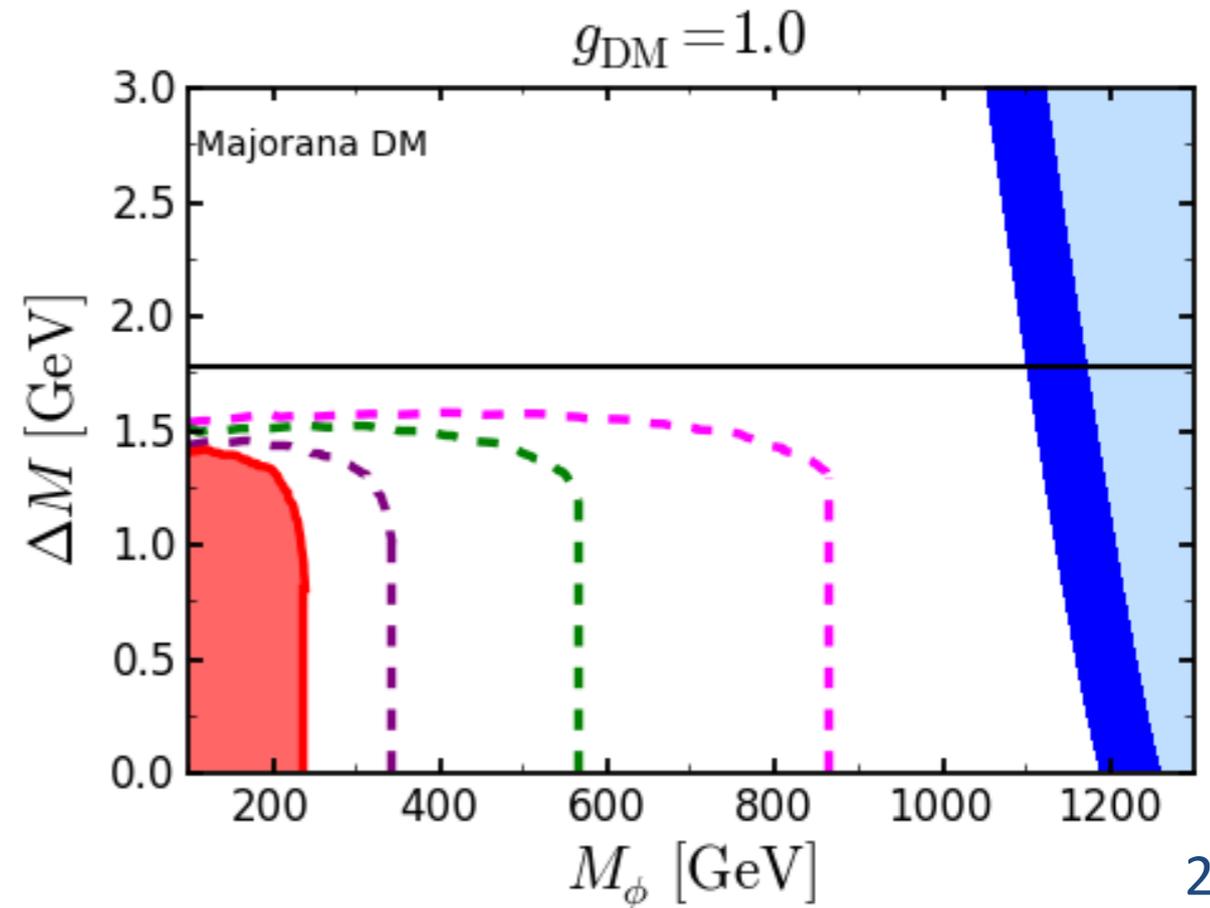
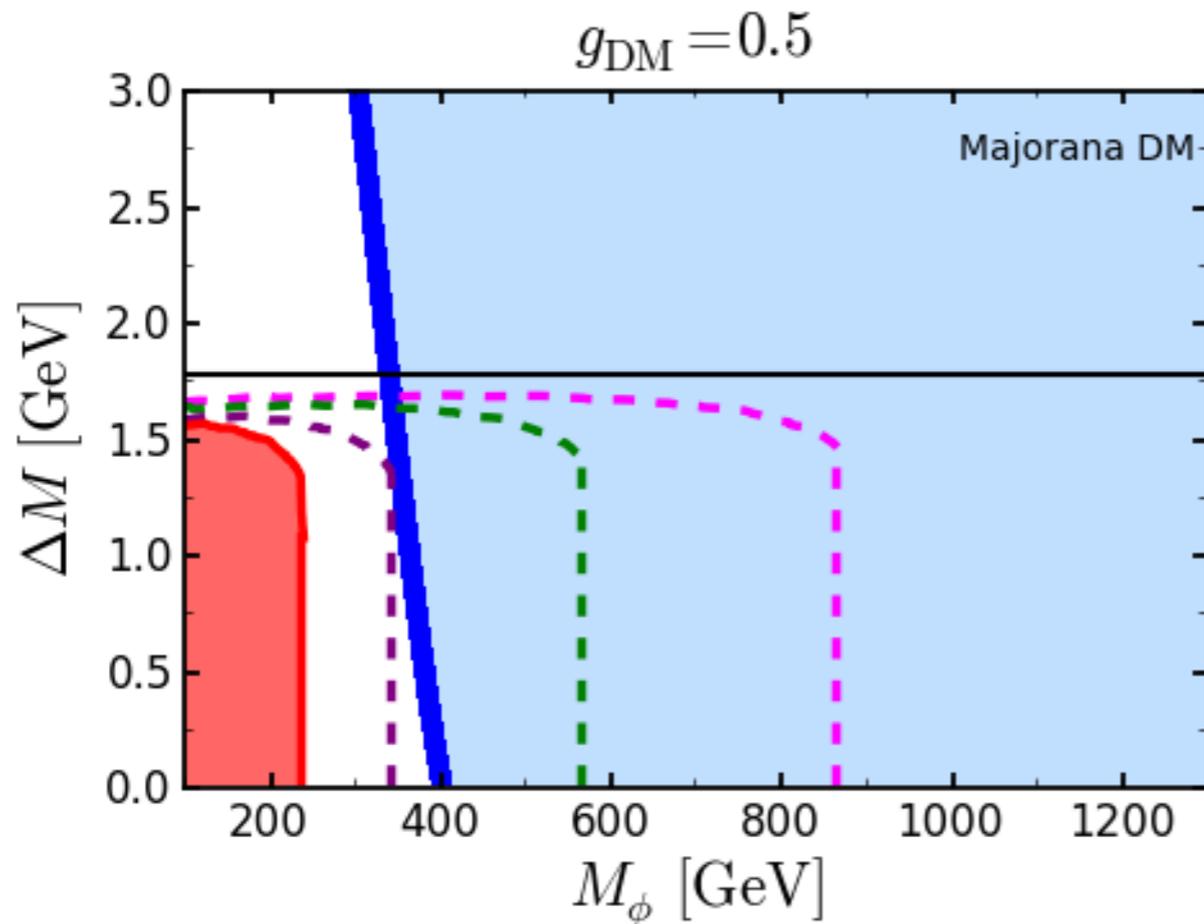
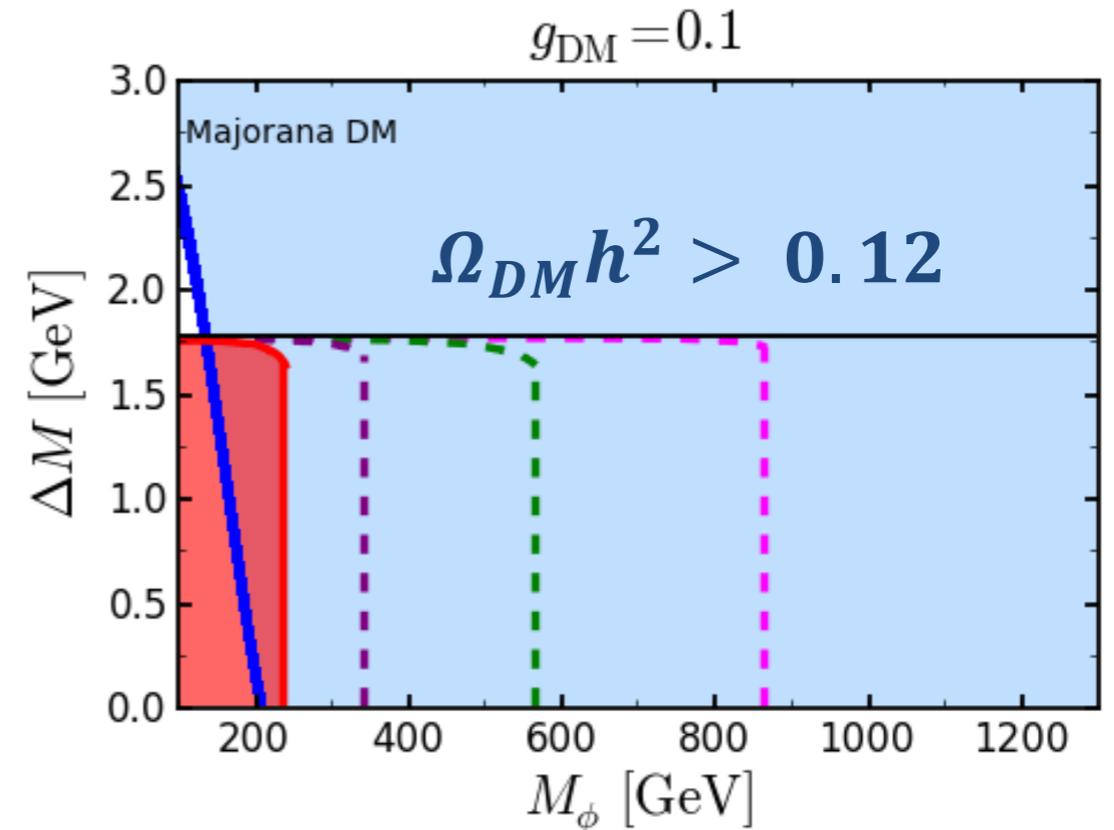
Majorana Dark Matter

DM CAP ($Y = 1 \quad L_\tau = 1$)

$\chi \quad \phi$

$\phi^* (\chi \tau_R) \subset \mathcal{L}$

High Lumi 13 TeV 3000 fb⁻¹



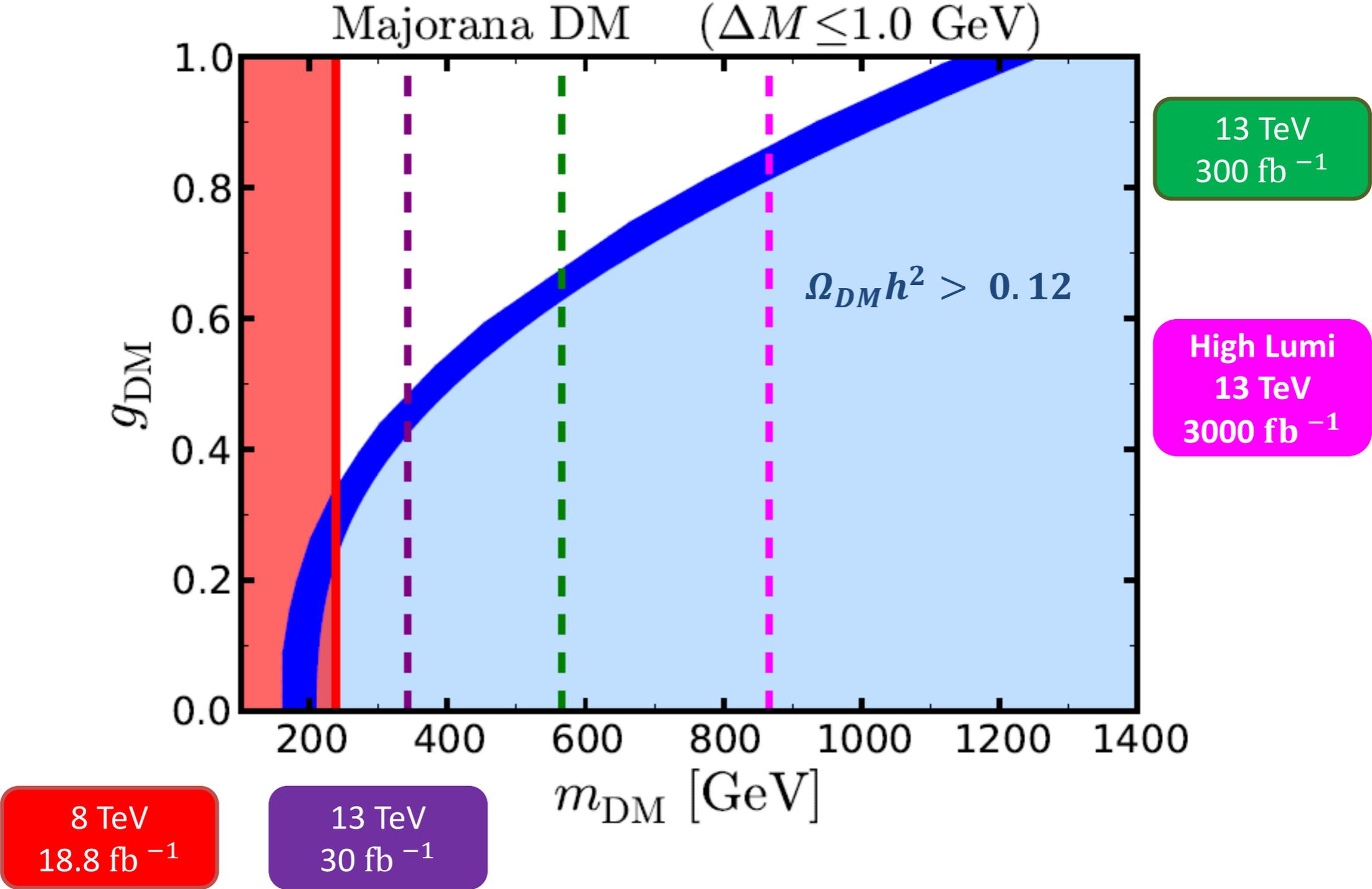
Model 1a

Majorana DM

Scalar CAP

e.g. neutralino-stau in SUSY

Now on the (g_{DM} , m_{DM}) plane:



Model 2

Scalar DM, Fermion CAP

e.g. dilaton, KK-tau

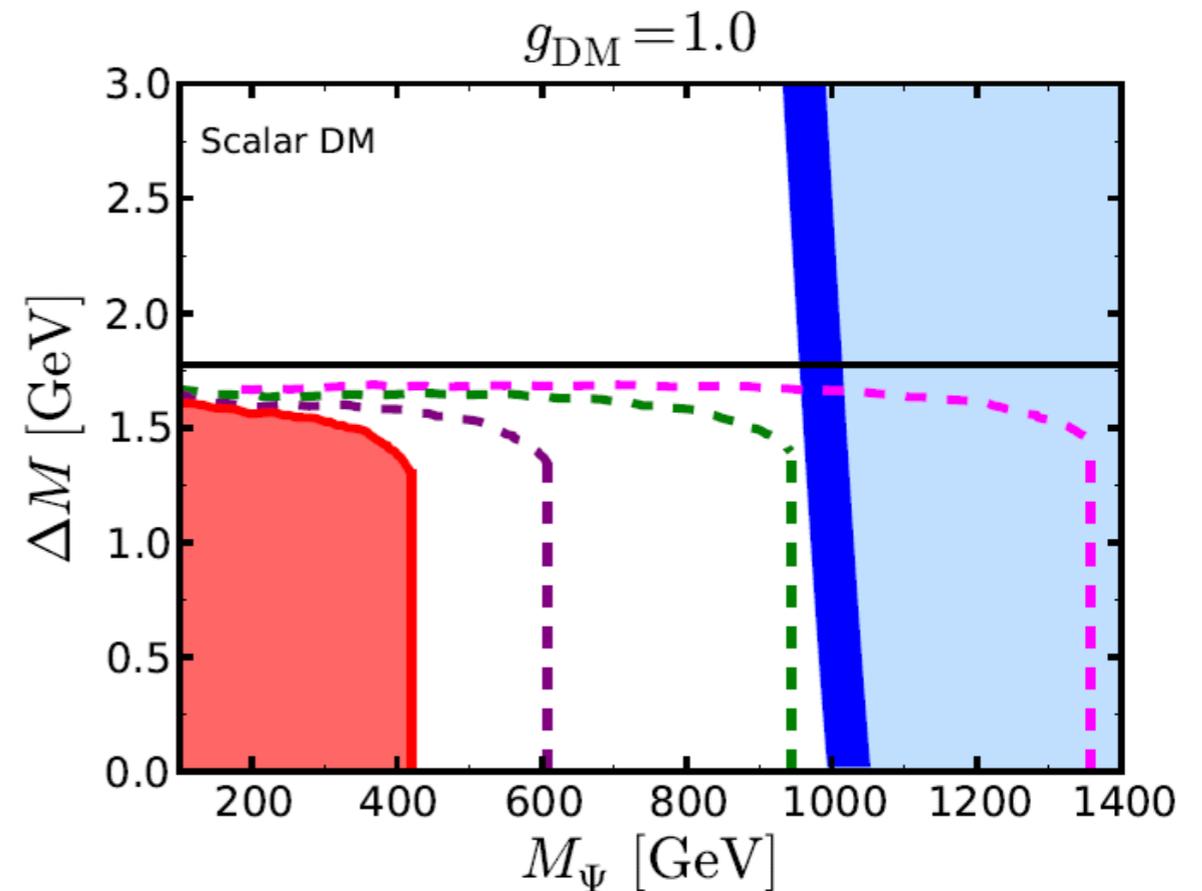
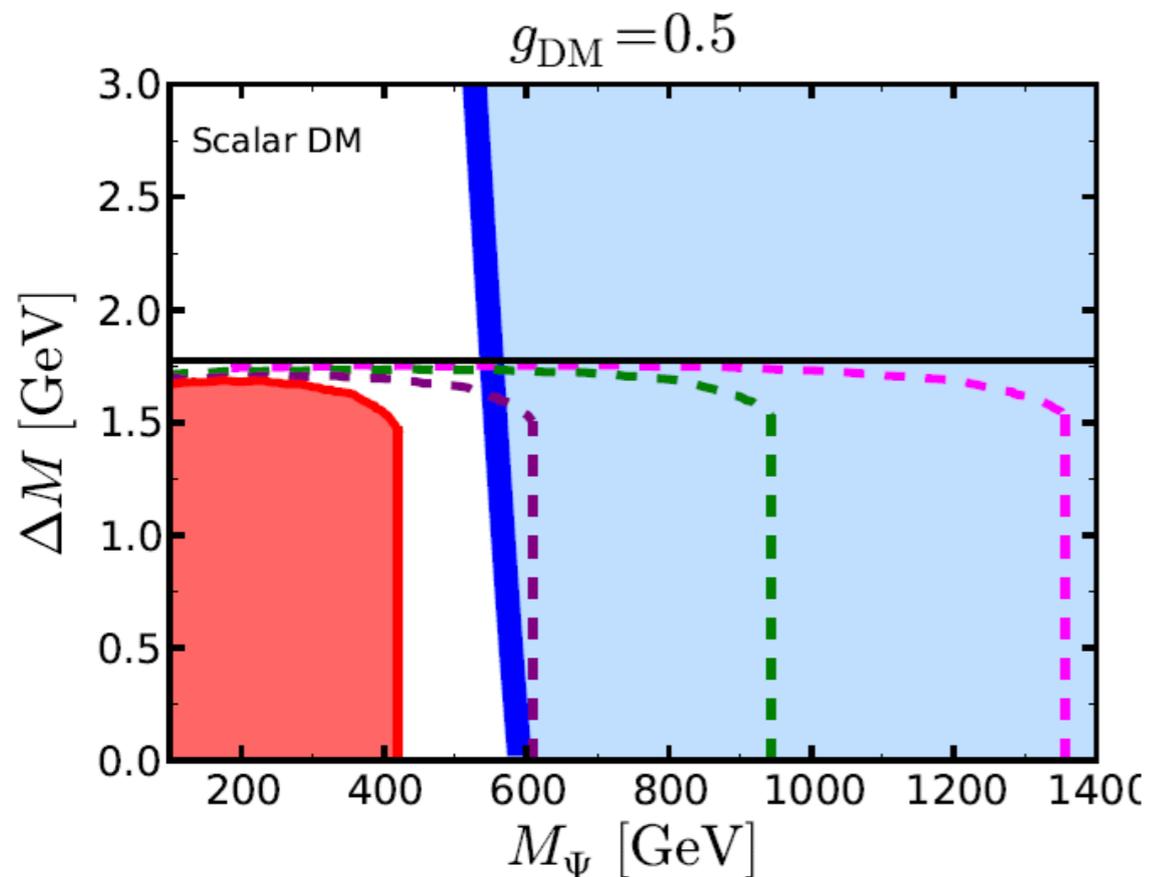
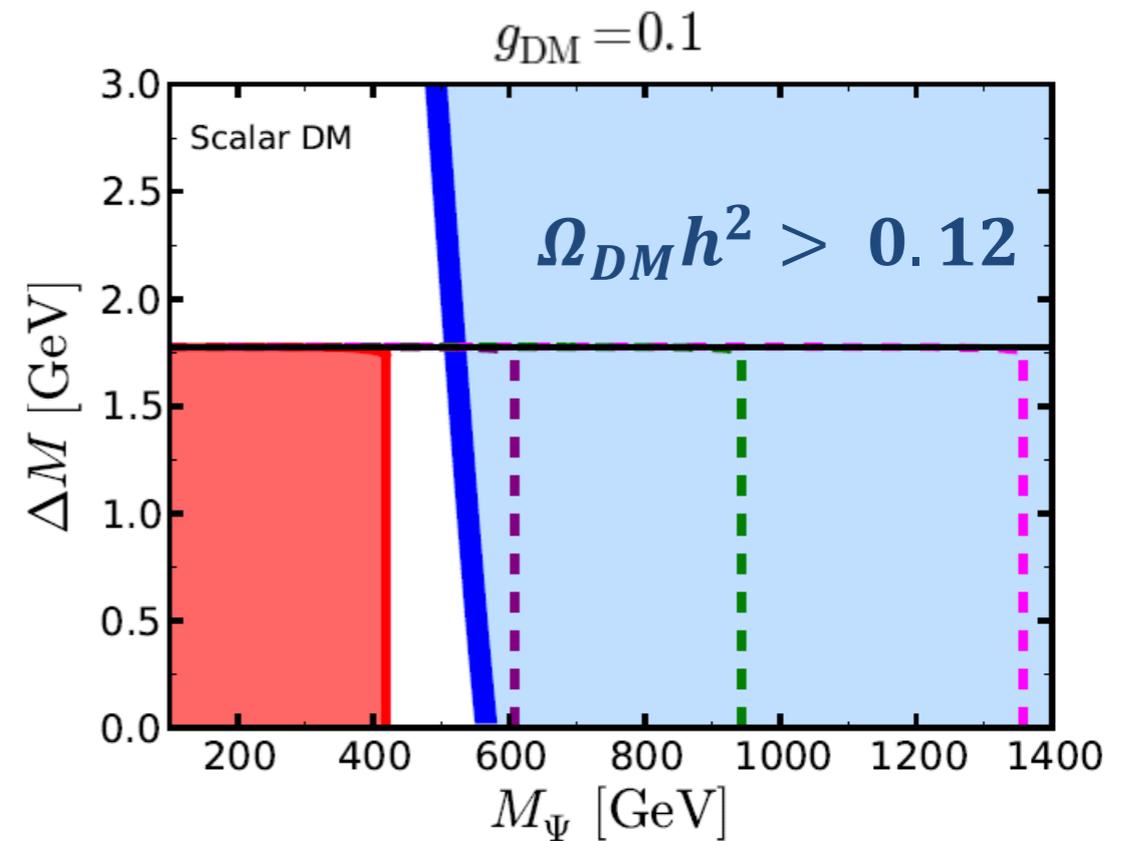
R Scalar Dirac ferm

DM **CAP** ($Y = 1$ $L_\tau = 1$)

S **Ψ**

$\mathcal{S}(\bar{\Psi} \tau_R) \subset \mathcal{L}$

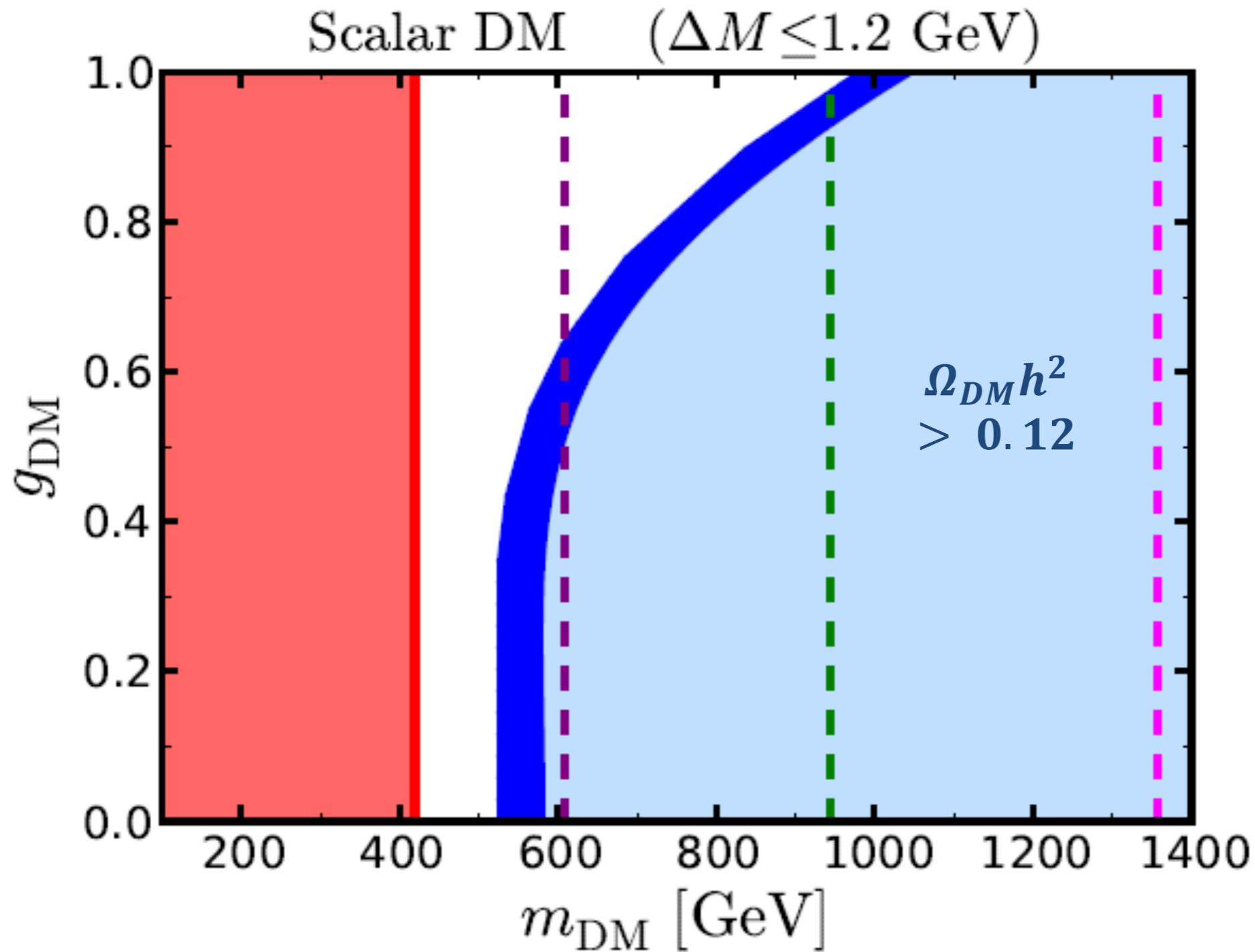
**Gauge-invariant and renormalizable,
no problems of unitarity**



Model 2

Scalar DM, Fermion CAP

e.g. dilaton, KK-tau



8 TeV
 18.8 fb^{-1}

13 TeV
 30 fb^{-1}

13 TeV
 300 fb^{-1}

High Lumi
13 TeV
 3000 fb^{-1}

Model 3

Vector DM, Fermion CAP

e.g. KK-photon, KK-tau

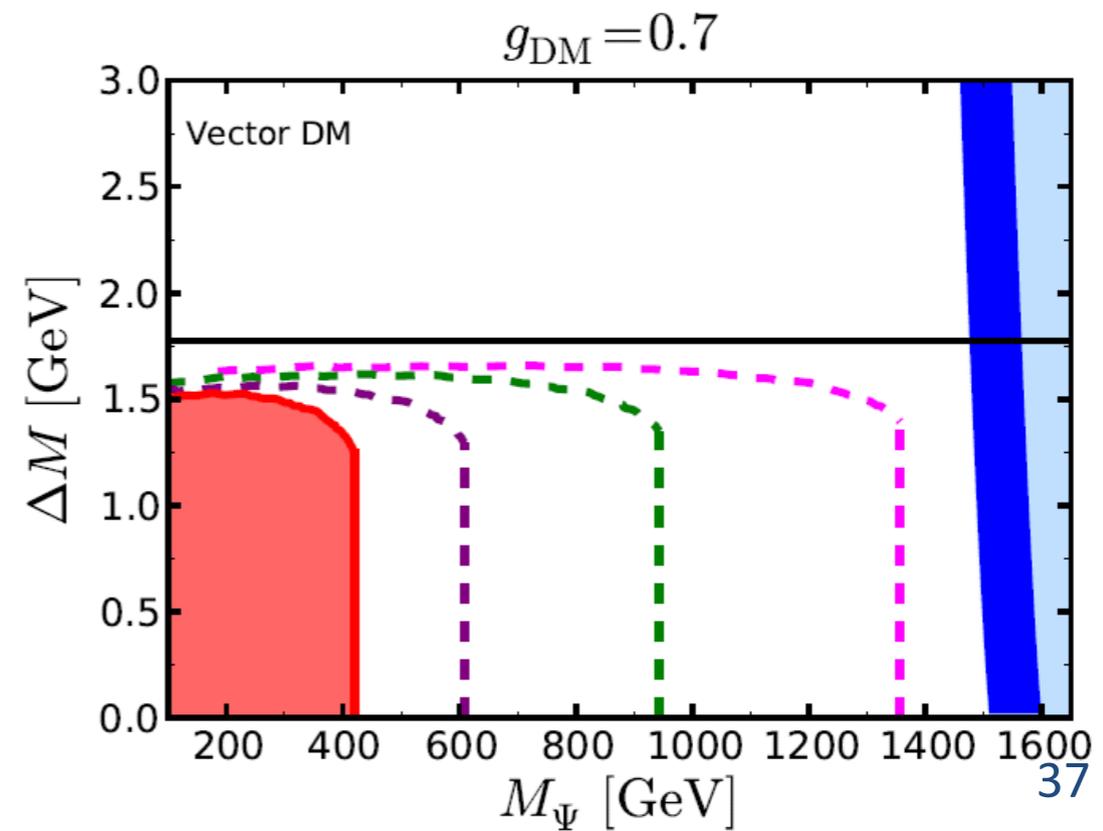
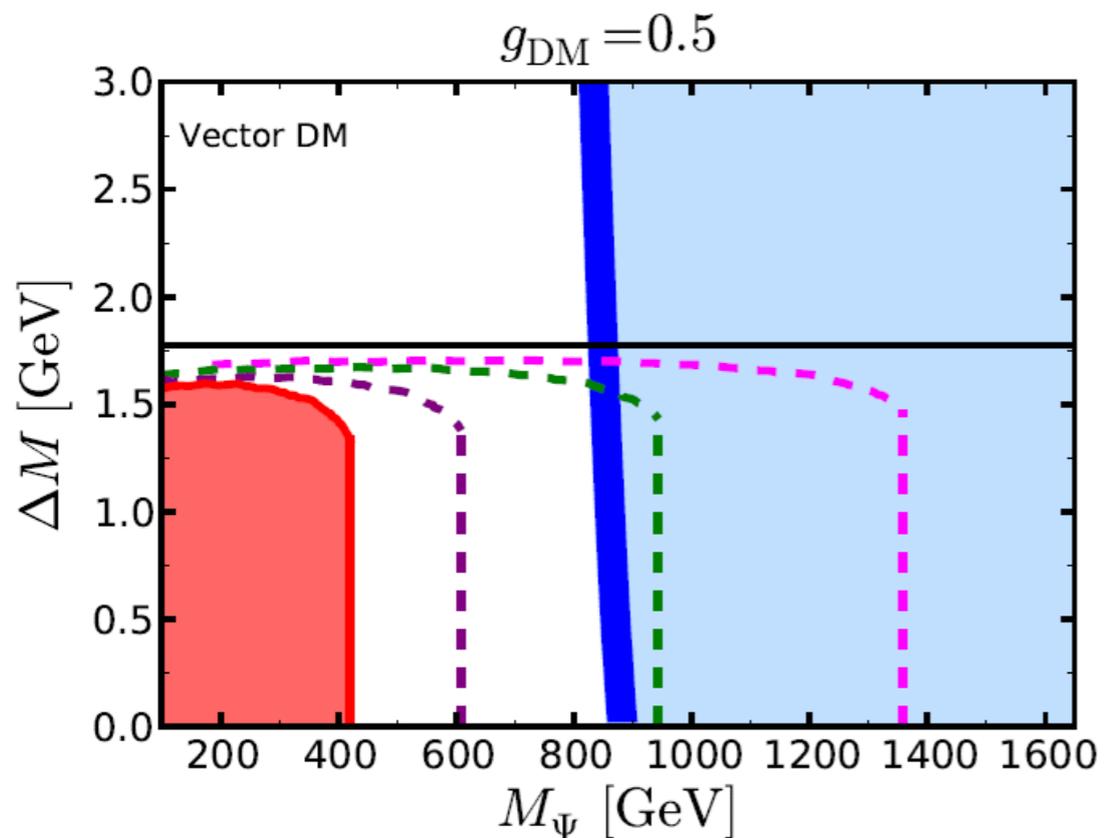
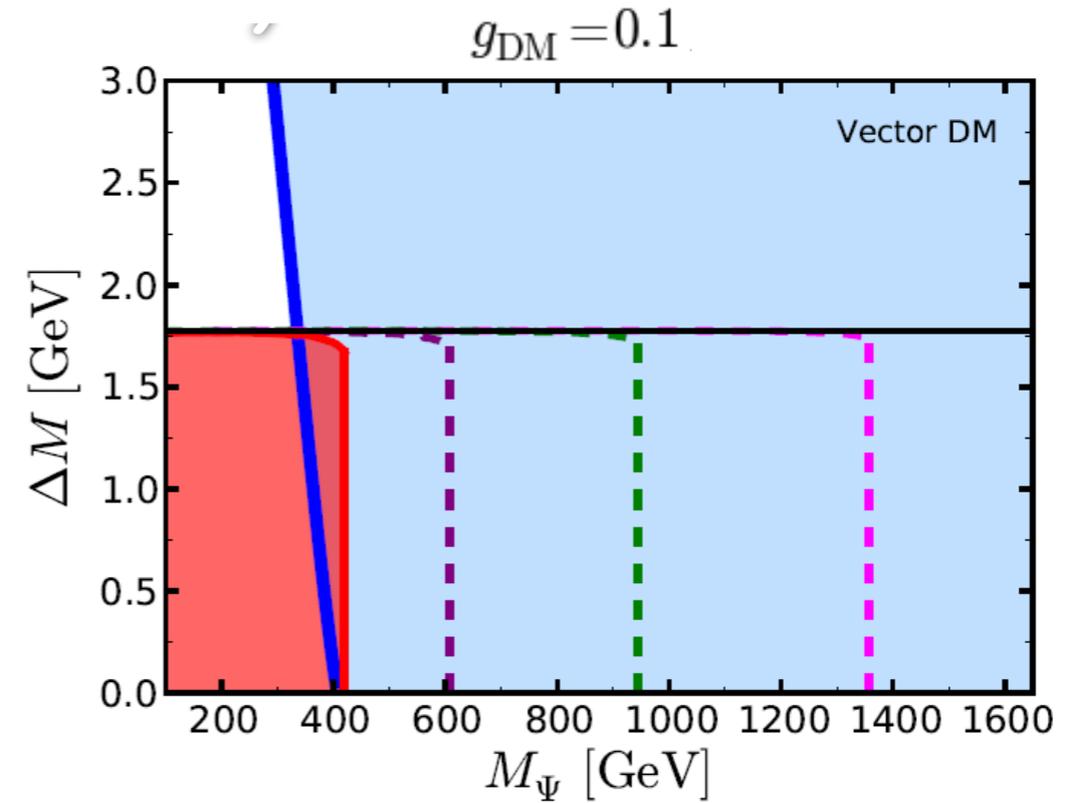
Vector Dirac ferm

DM CAP ($Y = 1$ $L_\tau = 1$)

$A_\mu \quad \Psi$

$$A_\mu (\bar{\Psi} \tau_R) \subset \mathcal{L}$$

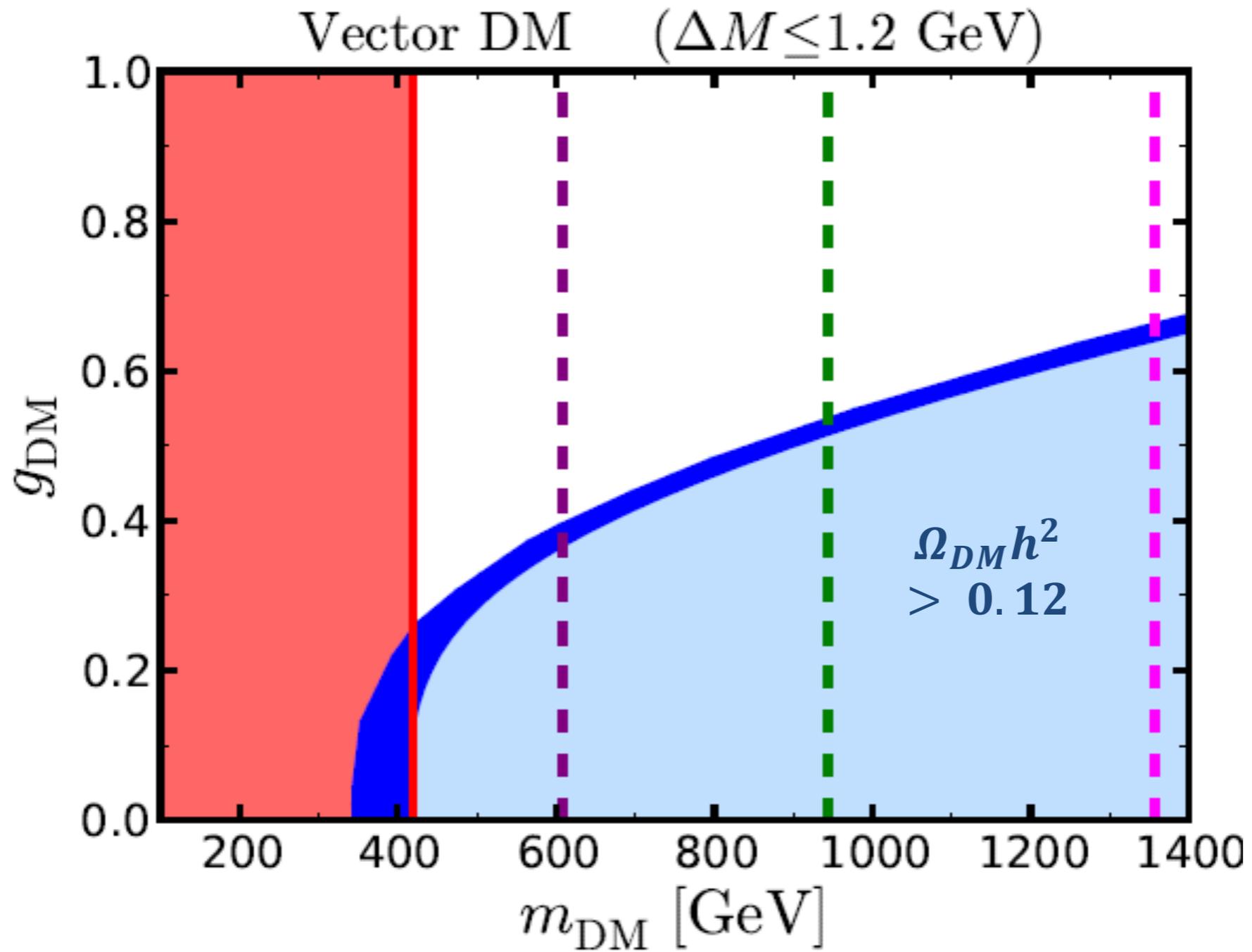
NOT gauge-invariant, requires UV-completion, e.g. Extra-Dimensions



Model 3

Vector DM, Fermion CAP

e.g. KK-photon, KK-tau



13 TeV
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High Lumi
13 TeV
3000 fb⁻¹

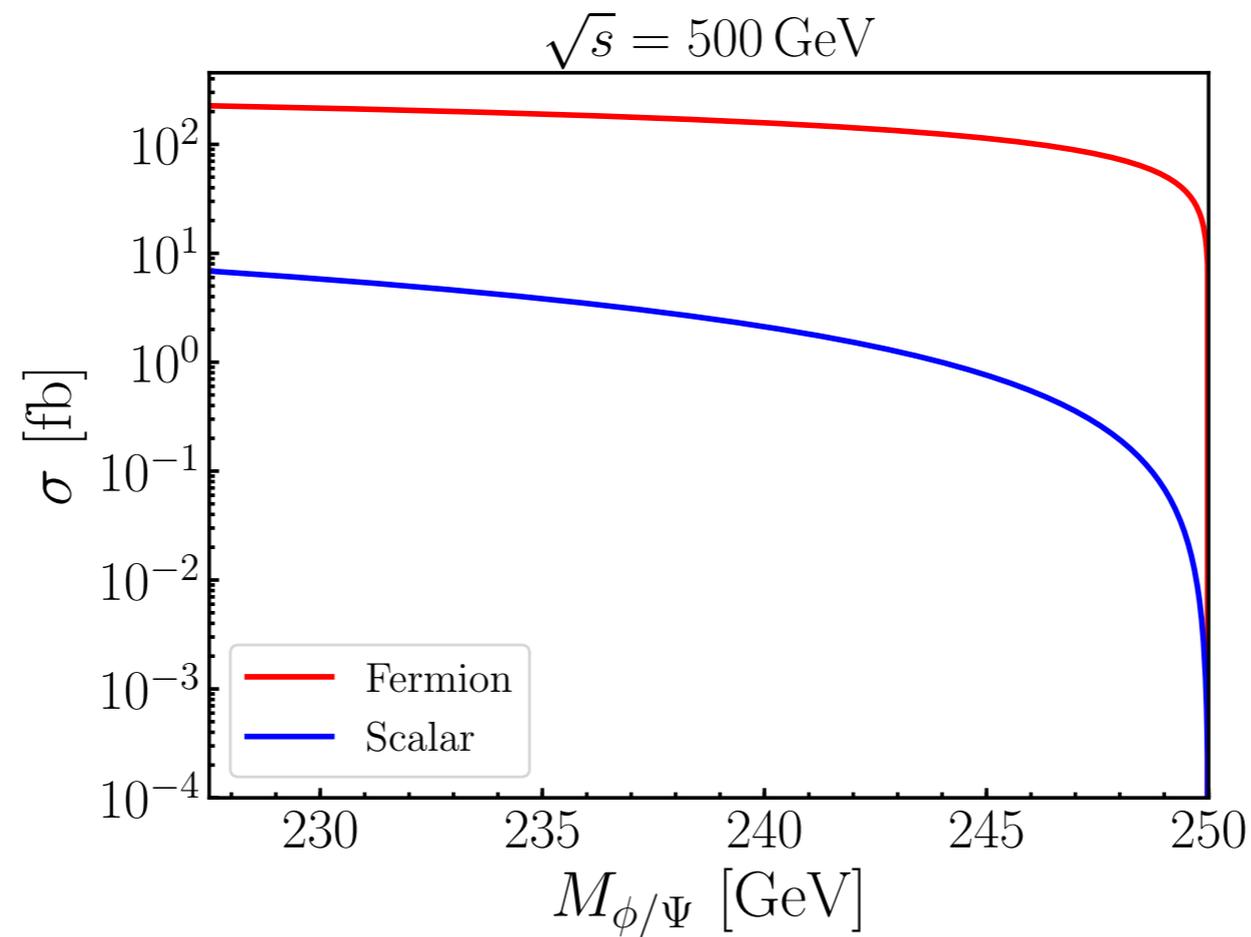
8 TeV
18.8 fb⁻¹

13 TeV
30 fb⁻¹

future e^+e^- collider

Cross-Section

$$e^+e^- \rightarrow \eta^+\eta^- \rightarrow \tau^+\tau^-\chi\chi$$



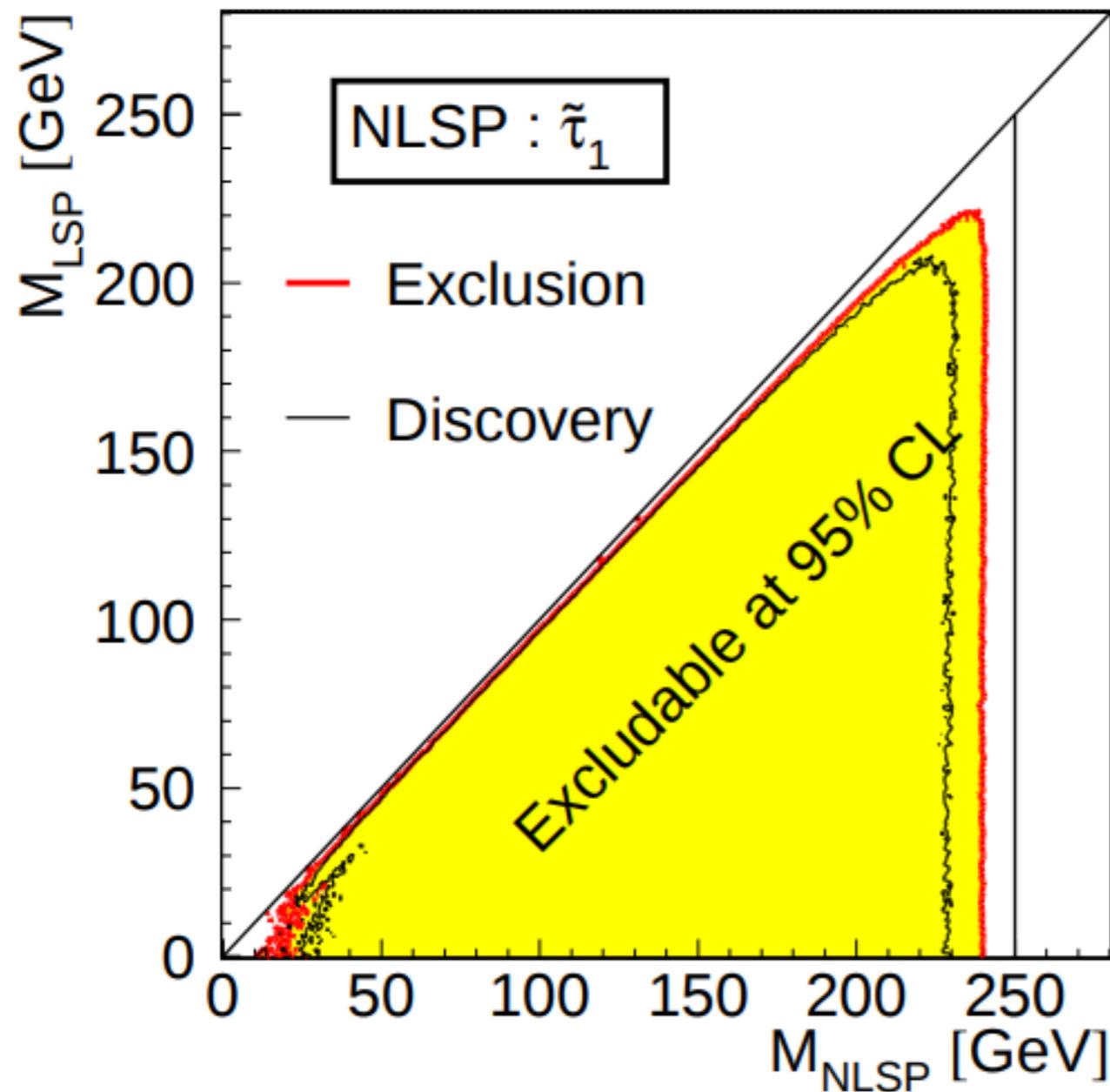
$$\sigma(e^+e^- \rightarrow \phi^+\phi^-) = \alpha^2 \pi s \cdot \mathcal{A} \cdot \frac{1}{6} \beta^3,$$

$$\sigma(e^+e^- \rightarrow \Psi^+\Psi^-) = \alpha^2 \pi s \cdot \mathcal{A} \cdot \beta \left(1 - \frac{1}{3} \beta^2 \right)$$

$$\mathcal{A} = \frac{2}{s^2} + \frac{2}{s} \frac{(g_L + g_R)g_R}{(s - m_Z^2)} + \frac{(g_L^2 + g_R^2)g_R^2}{(s - m_Z^2)^2}, \quad g_L = \frac{-\frac{1}{2} + s_W^2}{s_W c_W}, \quad g_R = \frac{s_W^2}{s_W c_W}, \quad \beta = \sqrt{1 - \frac{4M_{\phi/\Psi}^2}{s}}$$

Prospects for CLIC

- Example: Linear colliders can study compressed spectrum in SUSY



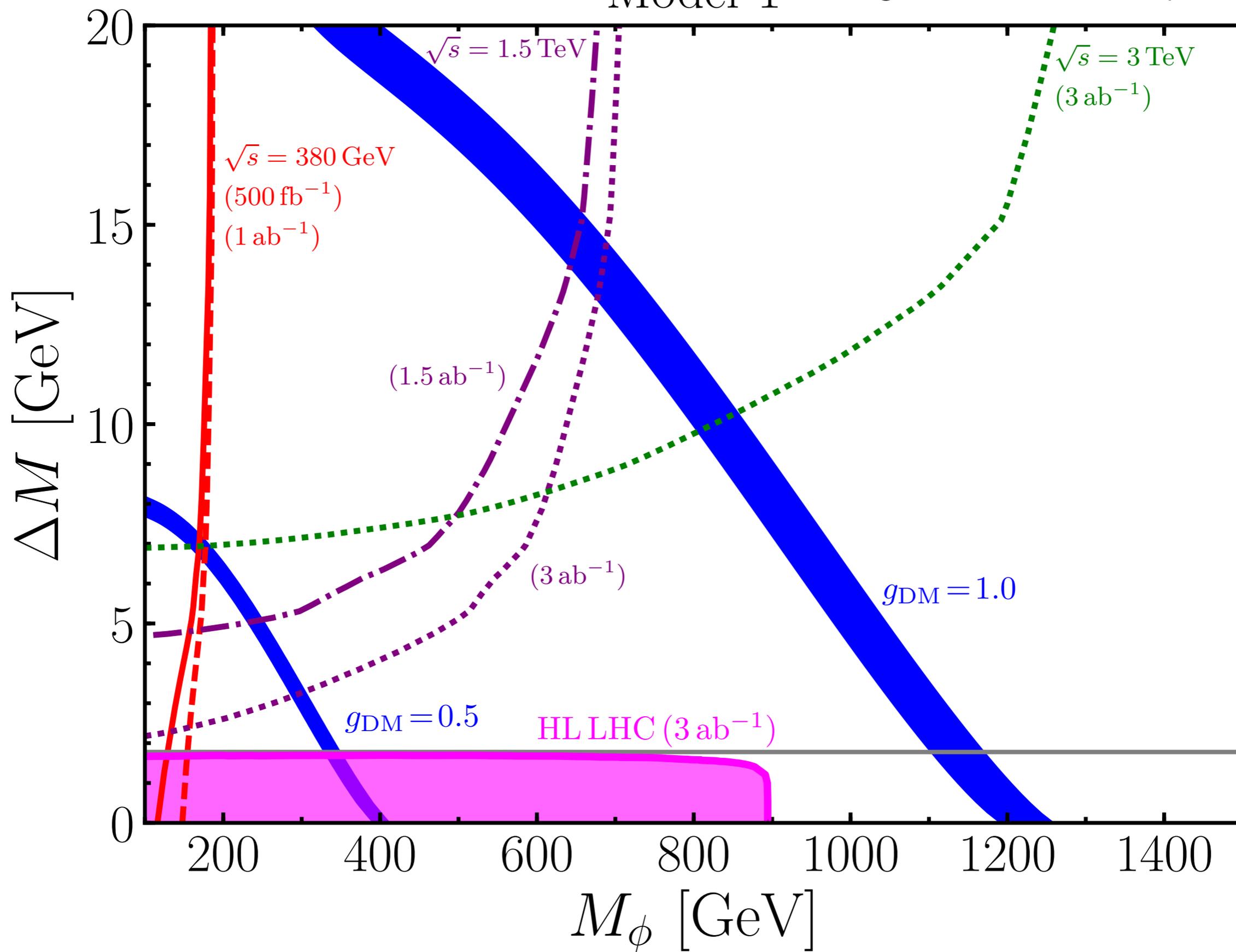
- Stau and neutralino, lightest supersymmetric states
- $\Delta M < 10$ GeV can be tested, virtual $\gamma\gamma$ becomes relevant background

$\sqrt{s} = 500$ GeV

[Berggren 2013]

fermionic DM, scalar CAP
e.g. neutralino-stau (SUSY)

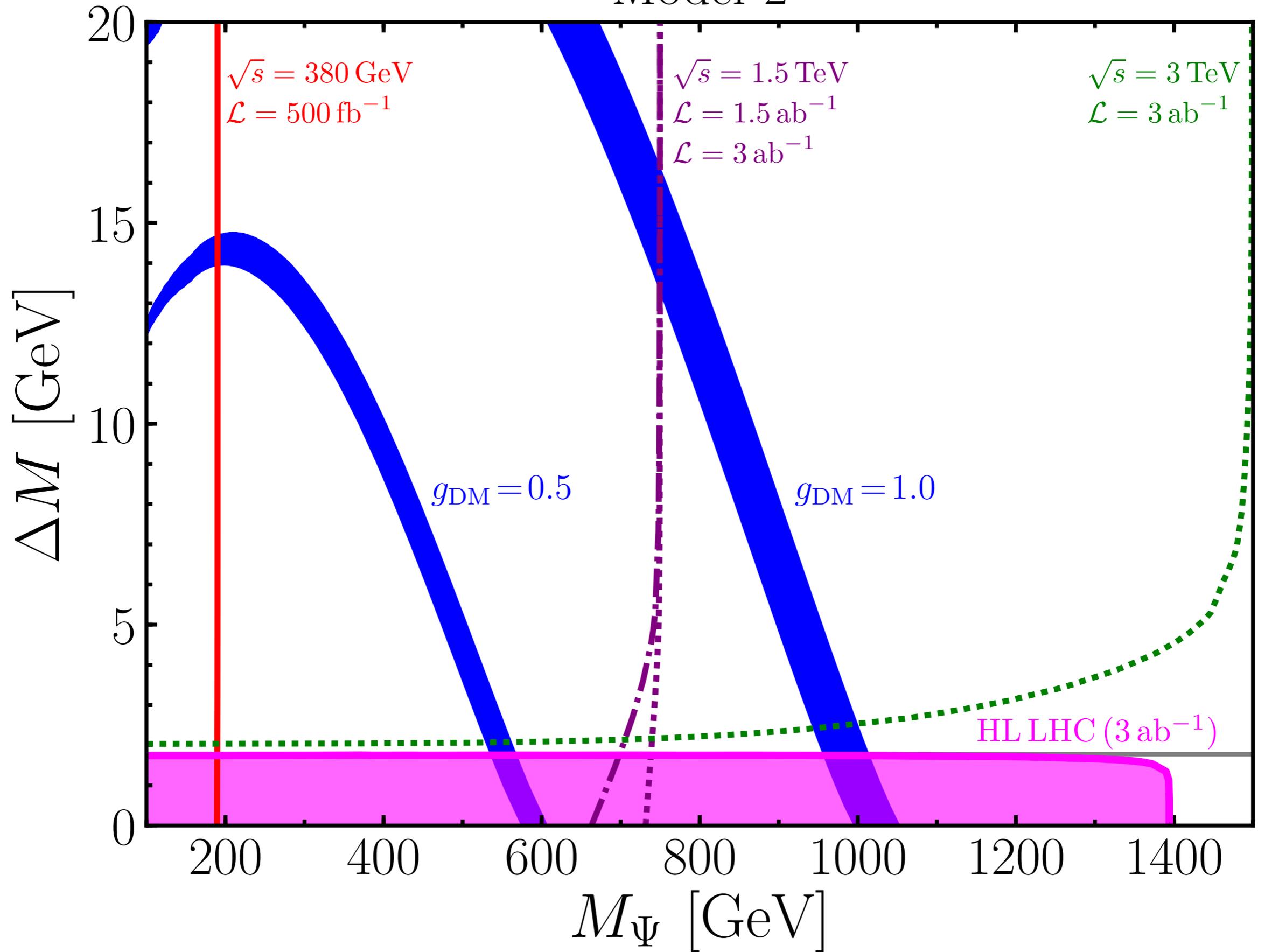
Model 1



scalar DM, fermionic CAP

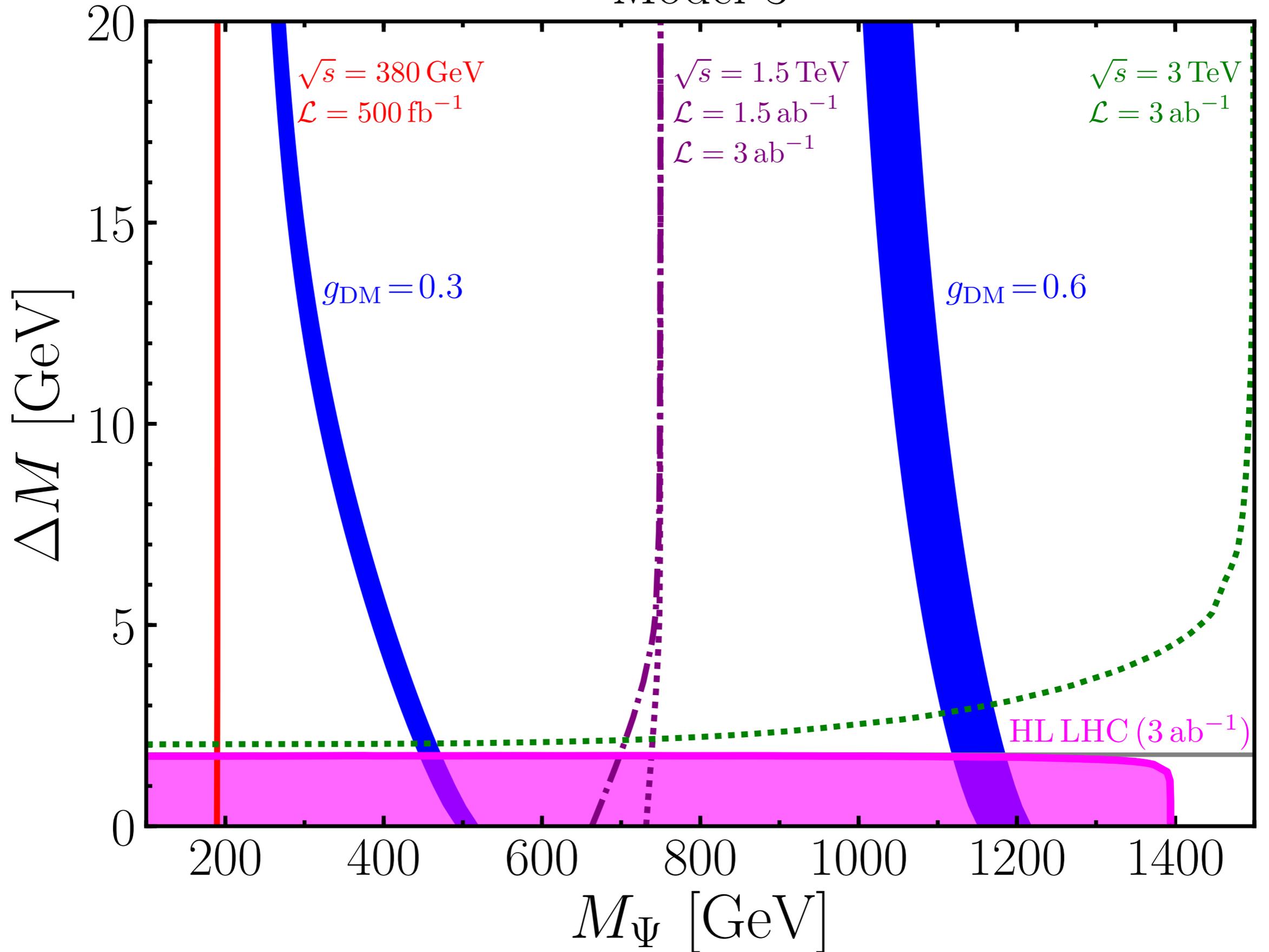
e.g. dilaton, KK-tau (extra-dim)

Model 2



vector DM, fermionic CAP
e.g. KK-photon, KK-tau (UED)

Model 3



Conclusion

- We presented co-annihilating dark matter simplified models with long-lived signatures. There are 3-parameters: m_χ , Δm , g_{DM}

Model-1a			
Component	Field	Charge	Interaction
DM	Majorana fermion (χ)	$Y = 0$	$\phi^*(\chi\tau_R) + \text{h.c.}$
CAP	Complex scalar (ϕ)	$Y = -1$	

fermionic DM, scalar CAP

e.g. neutralino-stau (SUSY)

Model-2			
Component	Field	Charge	Interaction
DM	Real scalar (S)	$Y = 0$	$S(\bar{\Psi}P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

scalar DM, fermionic CAP

e.g. dilaton, KK-tau (extra-dim)

Model-3			
Component	Field	Charge	Interaction
DM	Vector (V_μ)	$Y = 0$	$V_\mu(\bar{\Psi}\gamma^\mu P_R\tau) + \text{h.c.}$
CAP	Dirac fermion (Ψ)	$Y = -1$	

vector DM, fermionic CAP

e.g. KK-photon, KK-tau (UED)

- In this class of models, direct detection is 1-loop suppressed, indirect detection is velocity suppressed (for Model-1 and 2). The collider has the best chance of detecting it, but it needs $\Delta m < m_\tau$ to have long-lived signature.
- These models can be used by ATLAS and CMS to present their results.