

Lambda Polarizing Fragmentation Function from Belle e^+e^- data

Marco Zaccheddu - Università degli Studi di Cagliari & INFN

In collaboration with: Umberto D'Alesio, Francesco Murgia

Based on [Phys. Rev. D 102, 054001 (2020)]

Resummation, Evolution, Factorization 2020



Contents

- Introduction
- TMD FFs with Helicity Formalism
- $e^+e^- \rightarrow h_1^\uparrow h_2 X$
- $e^+e^- \rightarrow h_1(\text{jet})X$
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Similar analysis:

[D. Callos, Z.B. Kang, and J. Terry, Phys. Rev. D 102, 096007 (2020)]

TMD Fragmentation Functions for quarks : Spin - 1/2 hadrons

Helicity density matrix

$$\rho_{\lambda_i, \lambda'_i}^{i, s_i} = \frac{1}{2} \begin{pmatrix} 1 + P_z^i & P_x^i - iP_y^i \\ P_x^i + iP_y^i & 1 - P_z^i \end{pmatrix}$$

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{k}_{\perp h}) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^{q, s_q} \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{k}_{\perp h})$$

8 independent TMD Fragmentation Functions

		Hadron		
		U	L	T
	Pol. States			
Q u a r k	U	$\hat{D}_{h/q}$		$\Delta \hat{D}_{S_Y/q}^h$
	L		$\Delta \hat{D}_{S_Z/s_L}^{h/q}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
	T	$\Delta^N D_{h/q}^\uparrow$	$\Delta \hat{D}_{S_Z/s_T}^{h/q}$	$\Delta \hat{D}_{S_X/s_T}^{h/q} / \Delta - \hat{D}_{S_Y/s_T}^{h/q}$

TMD Fragmentation Functions for quarks : Spin - 1/2 hadrons

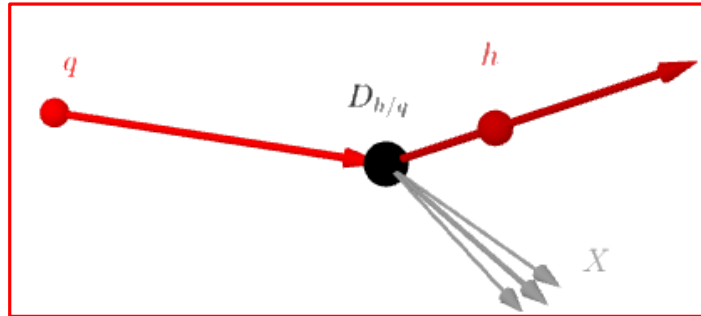
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8 independent TMD Fragmentation Functions

Unpolarized FF



		Hadron		
		U	L	T
Pol. States				
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	L		$\Delta \hat{D}_{S_Z/s_L}^{h/q}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
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TMD Fragmentation Functions for quarks : Spin - 1/2 hadrons

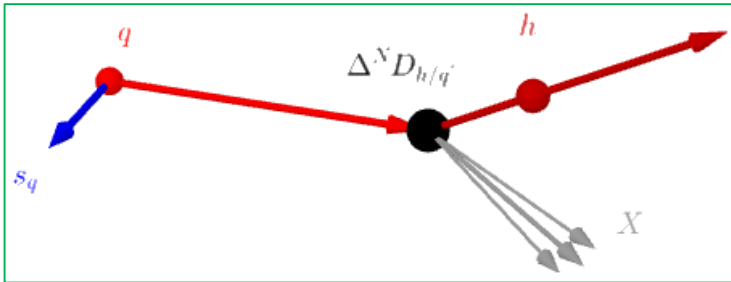
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8 independent TMD Fragmentation Functions

Collins FF



		Hadron		
		U	L	T
Pol. States		U	L	T
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	L		$\Delta \hat{D}_{S_Z/s_L}^{h/q}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
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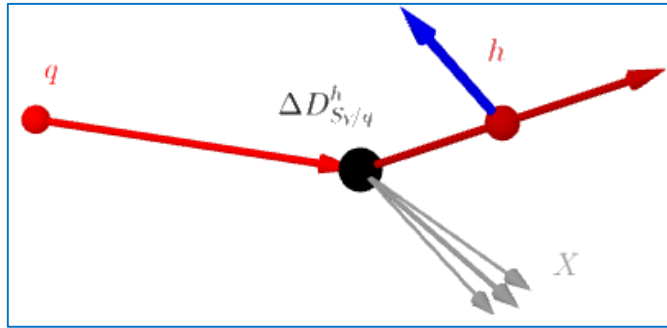
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Polarizing FF



8 independent TMD Fragmentation Functions

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Q u a r k	Pol. States			
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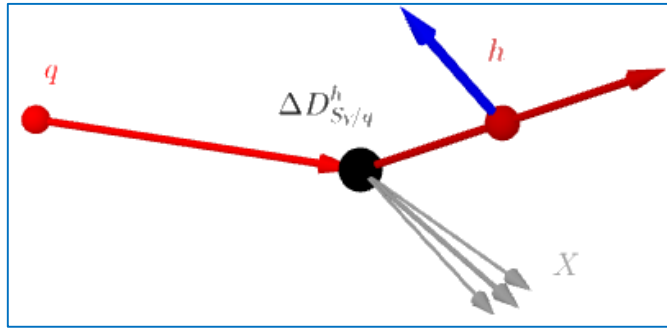
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Polarizing FF



Amsterdam notation:

Collins $\Delta^N D_{h/q^\uparrow} \longleftrightarrow H_1^\perp$

Polarizing $\Delta D_{S_Y/q}^h \longleftrightarrow D_{1T}^\perp$

8 independent TMD Fragmentation Functions

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$e^+e^- \rightarrow h_1h_2X$: Helicity Formalism

$$= \sum_{q_c} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \hat{M}_{\lambda_c \lambda_d, \lambda_a \lambda_b} \hat{M}_{\lambda'_c \lambda'_d, \lambda_a \lambda_b}^* \hat{D}_{\lambda_c, \lambda'_c}^{\lambda_{h_1}, \lambda'_{h_1}}(z_1, \mathbf{p}_{\perp 1}) \hat{D}_{\lambda_d, \lambda'_d}^{\lambda_{h_2}, \lambda'_{h_2}}(z_2, \mathbf{p}_{\perp 2})$$

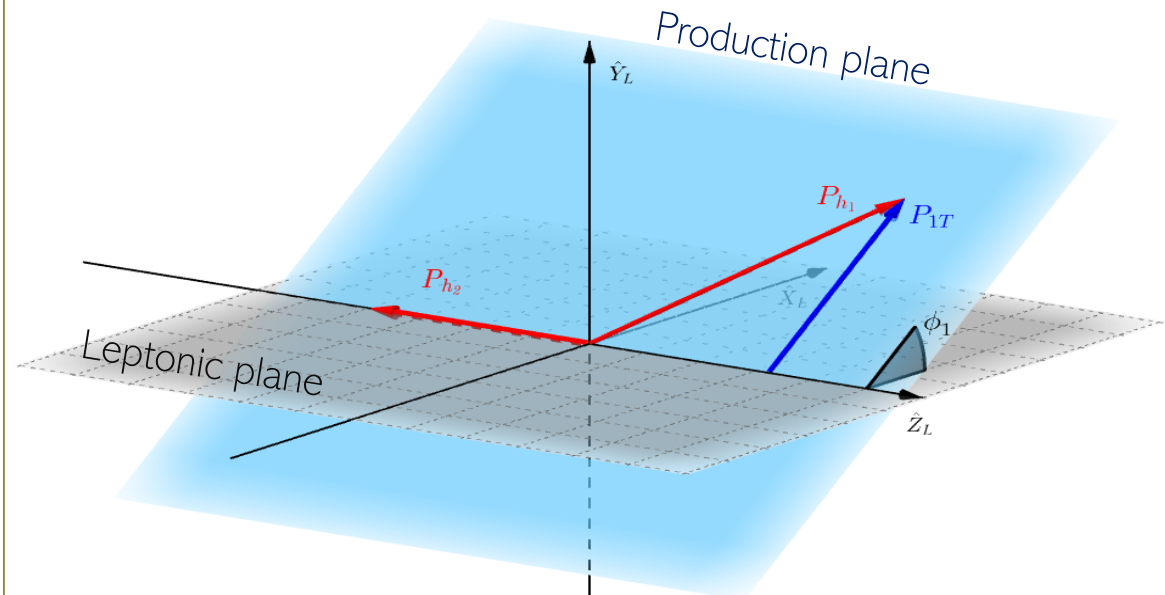
Scaling variables

- Light cone z
- Momentum fraction $z_p = 2|\mathbf{P}_h|/\sqrt{s}$
- Energy fraction $z_h = 2E_h/\sqrt{s}$

$$z_{h,p} \simeq z \left[1 \pm m_h^2 / (z^2 s) \right]$$

Results consistent with [D. Boer, R. Jakob, and P.J. Mulders. Nucl. Phys. B 504 (1997)]

Hadron Frame



$e^+e^- \rightarrow h_1 h_2 X$: Helicity Formalism

$$\begin{aligned}
 & \overbrace{\rho_{\lambda_{h_1}, \lambda'_{h_1}}^{h_1} \rho_{\lambda_{h_2}, \lambda'_{h_2}}^{h_2}}^{\text{Helicity matrices}} \frac{d\sigma^{e^+e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 \mathbf{p}_{\perp 1} dz_2 d^2 \mathbf{p}_{\perp 2}} \\
 = & \sum_{q_c} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \underbrace{\hat{M}_{\lambda_c \lambda_d, \lambda_a \lambda_b} \hat{M}_{\lambda'_c \lambda'_d, \lambda_a \lambda_b}^*}_{\text{Scattering Amplitudes}} \underbrace{\hat{D}_{\lambda_c, \lambda'_c}^{\lambda_{h_1}, \lambda'_{h_1}}(z_1, \mathbf{p}_{\perp 1}) \hat{D}_{\lambda_d, \lambda'_d}^{\lambda_{h_2}, \lambda'_{h_2}}(z_2, \mathbf{p}_{\perp 2})}_{\text{TMD Fragmentation Functions}}
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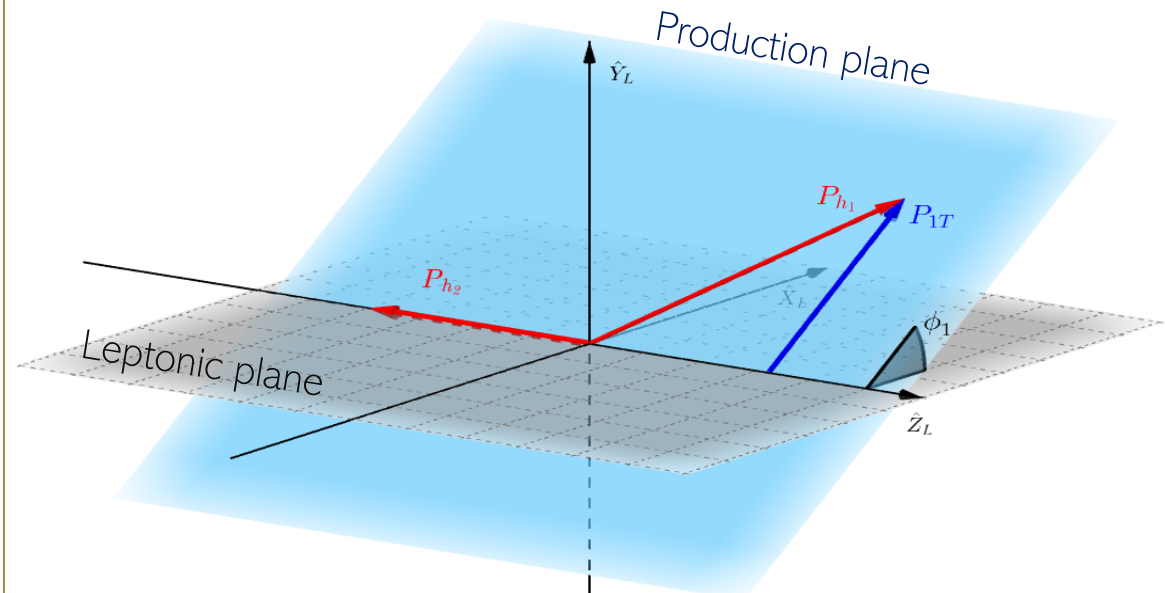
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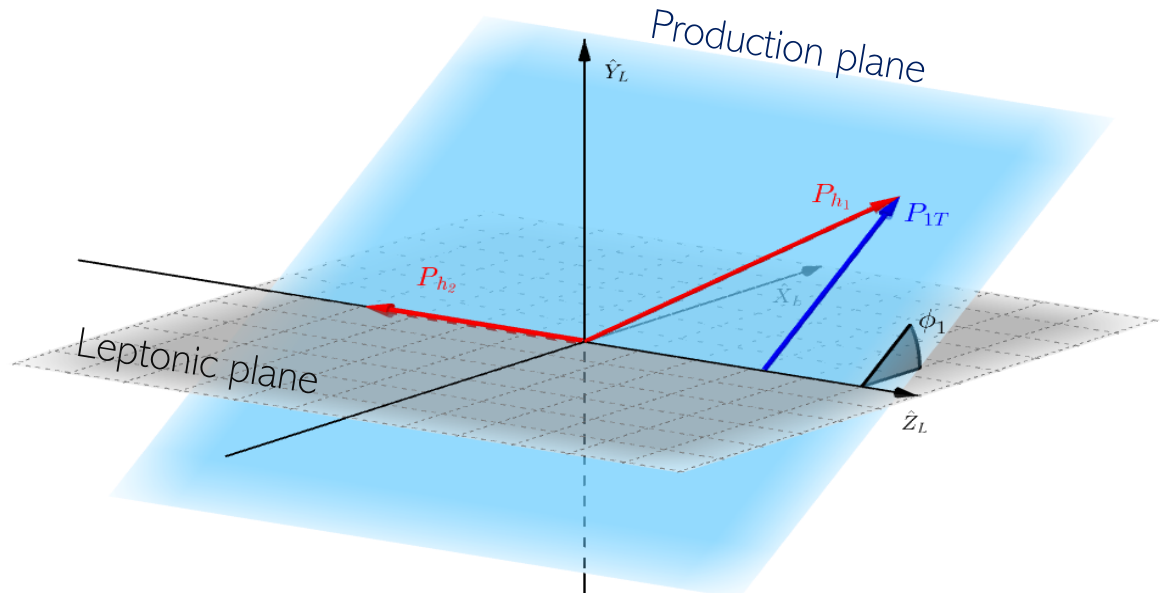
$$z_{h,p} \simeq z \left[1 \pm m_h^2 / (z^2 s) \right]$$

Polarization vector

$$\mathcal{P}^{h_1} = P_x^{h_1} \hat{X}_1 + P_y^{h_1} \hat{Y}_1 + P_z^{h_1} \hat{Z}_1$$

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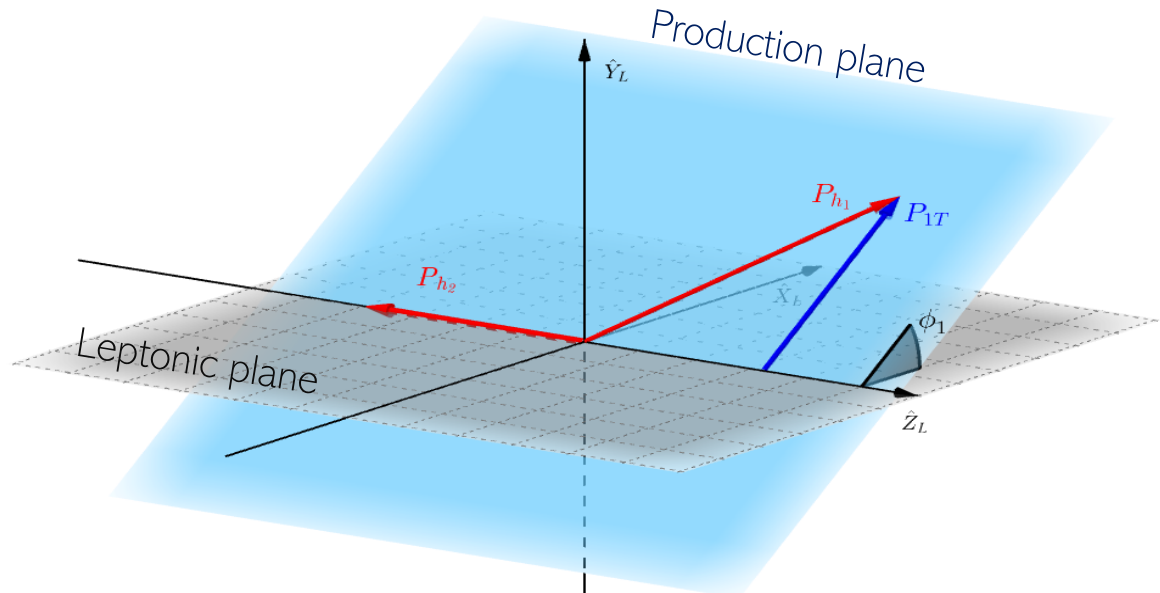
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The polarization is measured along:

$$\hat{n} = -\hat{P}_{h_2} \times \hat{P}_{h_1}$$

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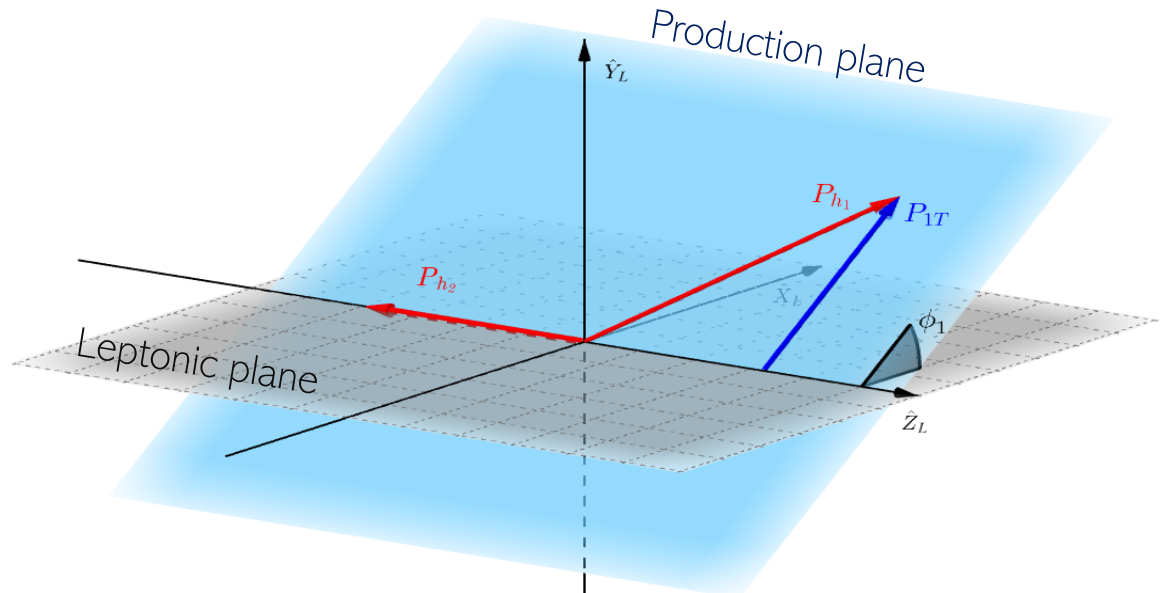
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The polarization projection along \hat{n} :

$$\mathcal{P}^{h_1} \cdot \hat{n} = P_x^{h_1} \cos \tilde{\phi} + P_y^{h_1} \sin \tilde{\phi}$$

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Hadron Frame



$e^+e^- \rightarrow h_1^+h_2^-X$: Hadron Frame

Introduce p_\perp - parameterization for FFs:

$$\Delta D_{S_Y/q}^h(z, p_\perp) = \Delta D_{S_Y/q}^h(z) \sqrt{2} e \frac{p_\perp}{M_{pol}} \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle_{pol}}}{\pi \langle p_\perp^2 \rangle_h}$$

$$D_{h/q}(z, p_\perp) = D_{h/q}(z) \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle_h}}{\pi \langle p_\perp^2 \rangle_h}$$

- Fixed energy scale $\sqrt{s} = 10.58$ GeV
- NO Evolution
- Data depend only on energy fraction $z_\Lambda - z_{\pi,K}$
- $\langle p_\perp^2 \rangle_h = 0.2 \text{ GeV}^2$ width of unp. FF

$e^+e^- \rightarrow h_1^+ h_2^- X$: Hadron Frame

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$e^+e^- \rightarrow h^{\uparrow}_1 h_2 X$: Hadron Frame

Introduce p_{\perp} - parameterization for FFs: Linear term

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- Data depend only on energy fraction $z_\Lambda - z_{\pi,K}$
- $\langle p_\perp^2 \rangle_h = 0.2 \text{ GeV}^2$ width of unp. FF

$$d^2 \mathbf{p}_{\perp 1} \rightarrow d^2 \mathbf{P}_{T1}$$

$$\int d^2 \mathbf{P}_{1T} d^2 \mathbf{p}_{\perp 2}$$

$e^+e^- \rightarrow h_1^\dagger h_2 X$: Hadron Frame

Introduce p_\perp - parameterization for FFs:

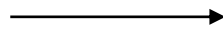
$$\Delta D_{S_Y/q}^h(z, p_\perp) = \underbrace{\Delta D_{S_Y/q}^h(z)}_{z \text{ dependence}} \sqrt{2e} \frac{\underbrace{p_\perp}_{\text{Linear term}}}{M_{pol}} \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle_{pol}}}{\underbrace{\pi \langle p_\perp^2 \rangle_h}_{\text{Gaussian dependence on } p_\perp}}$$

$$D_{h/q}(z, p_\perp) = \underbrace{D_{h/q}(z)}_{z \text{ dependence}} \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle_h}}{\pi \langle p_\perp^2 \rangle_h}$$

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$$d^2\mathbf{p}_{\perp 1} \rightarrow d^2\mathbf{P}_{T1}$$

$$\int d^2\mathbf{P}_{1T} d^2\mathbf{p}_{\perp 2}$$



$$\mathcal{P}_n(z_1, z_2) = \sqrt{\frac{e\pi}{2}} \frac{1}{M_{pol}} \frac{\langle p_\perp^2 \rangle_{pol}^2}{\langle p_\perp^2 \rangle} \frac{z_2}{\{ [z_1(1 - m_{h_1}^2 / (z_1^2 s))]^2 \langle p_\perp^2 \rangle + z_2^2 \langle p_\perp^2 \rangle_{pol} \}^{1/2}}$$

$$\times \frac{\sum_q e_q^2 \Delta D_{h_1^\dagger/q}(z_1) D_{h_2/\bar{q}}(z_2)}{\sum_q e_q^2 D_{h_1/q}(z_1) D_{h_2/\bar{q}}(z_2)}$$

$e^+e^- \rightarrow h_1^\dagger h_2 X$: Hadron Frame

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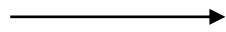
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$$\times \frac{\sum_q e_q^2 \Delta D_{h_1^\dagger/q}(z_1) D_{h_2/\bar{q}}(z_2)}{\sum_q e_q^2 D_{h_1/q}(z_1) D_{h_2/\bar{q}}(z_2)}$$

Consistent with [D. Callos, Z.B. Kang, and J. Terry, Phys. Rev. D 102, 096007 (2020)]

$e^+e^- \rightarrow h_1 + (\text{jet}) X$: Thrust Frame

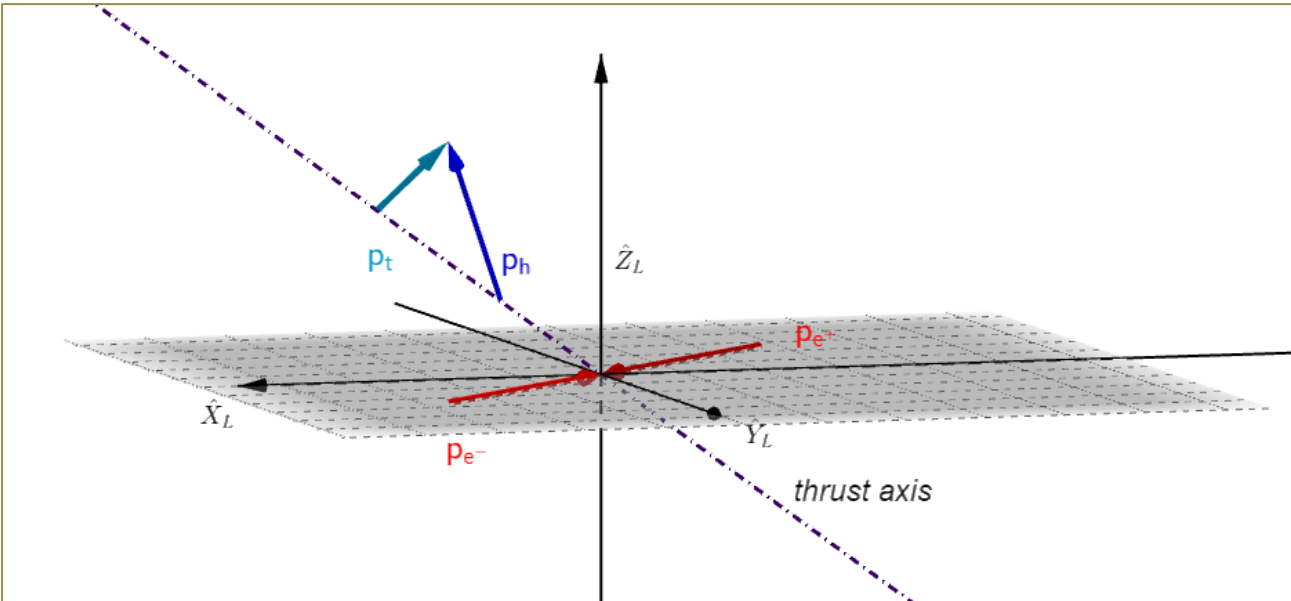
The polarization is measured along:

$$\hat{n} = \hat{T} \times \hat{P}_{h_1}$$

- Data depend only on $z_\Lambda - p_\perp$
- Direct access to p_\perp dependence

Within a phenomenological approach

$$\mathcal{P}_T(z_1, p_{\perp 1}) = \frac{\sum_q e_q^2 \Delta D_{h_1^\uparrow/q}(z_1, p_{\perp 1})}{\sum_q e_q^2 D_{h_1/q}(z_1, p_{\perp 1})}$$



Fit and Results (1)

Belle data: $\sqrt{s} = 10.58$ GeV

- 128 points $\Lambda + h$, in bins of the energy fractions $z_\Lambda - z_{\pi,K}$
- 32 points $\Lambda(\text{jet})$, in bins of $z_\Lambda - p_\perp$

Polarizing FF parametrization:

$$\Delta D_{S_Y/q}^h(z) = \mathcal{N}_q^p(z) \overbrace{D_{h/q}(z)}^{\text{Unpolarized FF}}$$

$$\mathcal{N}_q^p(z) = \mathcal{N}_q^p z^{\alpha_q} (1-z)^{\beta_q} \frac{(\alpha_q + \beta_q)^{(\alpha_q + \beta_q)}}{\alpha_q^{\alpha_q} \beta_q^{\beta_q}}$$

- Normalization factor: \mathcal{N}_q^p , $|\mathcal{N}_q^p| \leq 1$
- Shape for high and low z : α_q β_q

Unpolarized FF set adopted:

- DSS07 for π, K
- AKK08 for $\Lambda + \bar{\Lambda}$

$q\bar{q}$ FF separation for AKK08:

$$D_{\Lambda/\bar{q}}(z_p) = (1 - z_p) D_{\Lambda/q}(z_p)$$

Fit and Results (2)

Data selection:

- $\Lambda + \pi/K$: $z_{\pi.K} = [0.5 - 0.9]$ bin excluded \rightarrow 96 data points
- $\Lambda(jet)$: $z_{\Lambda} = [0.5 - 0.9]$ bin excluded \rightarrow 24 data points

Fit and Results (2)

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- $\Lambda + \pi/K$: $z_{\pi.K} = [0.5 - 0.9]$ bin excluded \rightarrow 96 data points
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Fitted 8 parameters

Flav.	\mathcal{N}_q^p	α_q	β_q	$\langle p_{\perp}^2 \rangle_{pol}$
u	\mathcal{N}_u^p		β_u	
d	\mathcal{N}_d^p			$\langle p_{\perp}^2 \rangle_{pol}$
s	\mathcal{N}_s^p	α_s		
sea	\mathcal{N}_{sea}^p		β_{sea}	

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Fitted 8 parameters

Fitted Parameters Value	
Nu	0.47 ^{+0.32} _{-0.20}
Nd	-0.32 ^{+0.13} _{-0.13}
Ns	-0.57 ^{+0.29} _{-0.43}
Nsea	-0.27 ^{+0.12} _{-0.20}
α_s	2.30 ^{+1.08} _{-0.91}
β_{sea}	2.60 ^{+2.60} _{-1.74}
β_u	3.50 ^{+2.33} _{-1.82}
$\langle p_{\perp}^2 \rangle_{pol}$	0.10 ^{+0.02} _{-0.02}

Full data Fit : $\chi_{dof}^2 = 1.94$

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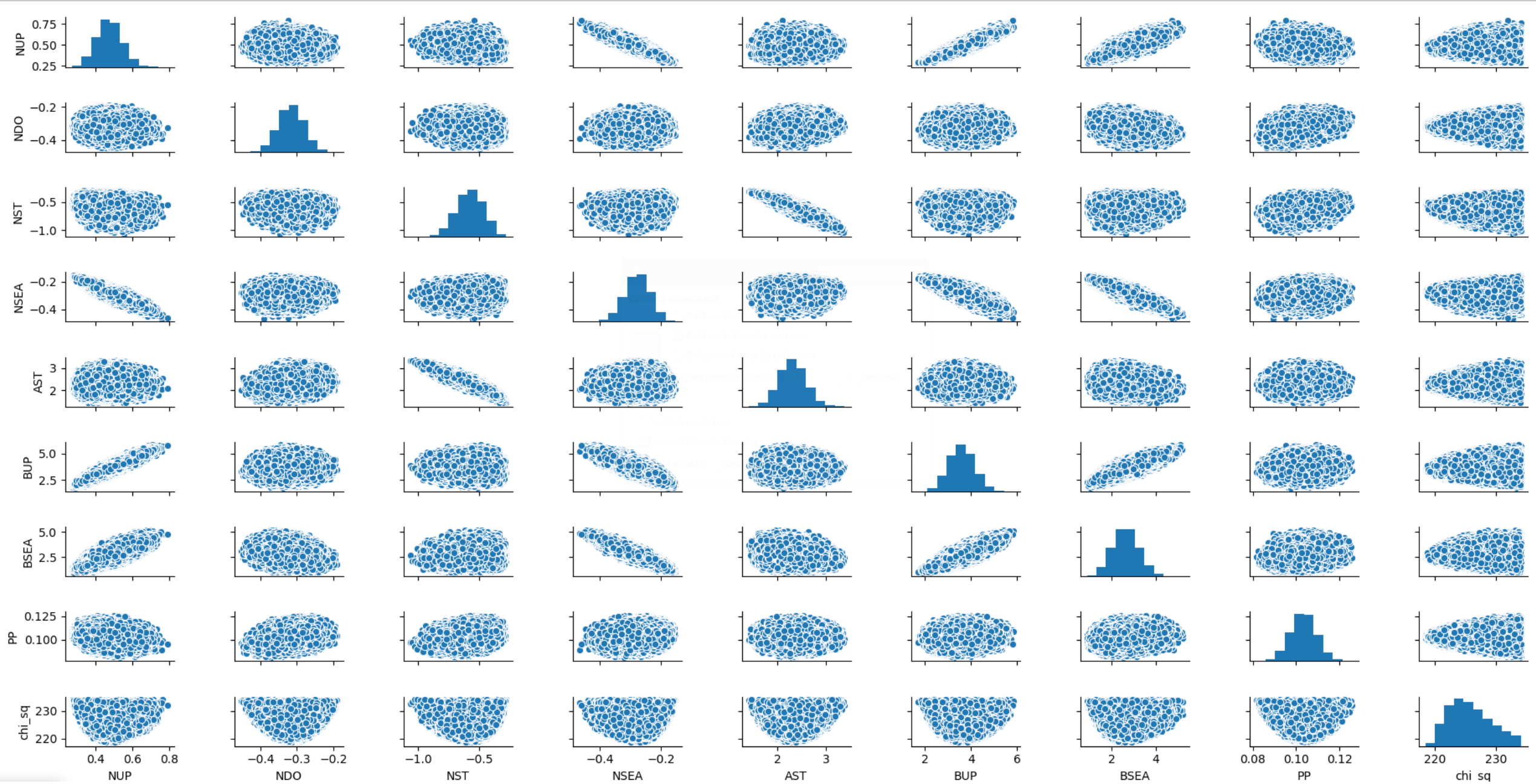
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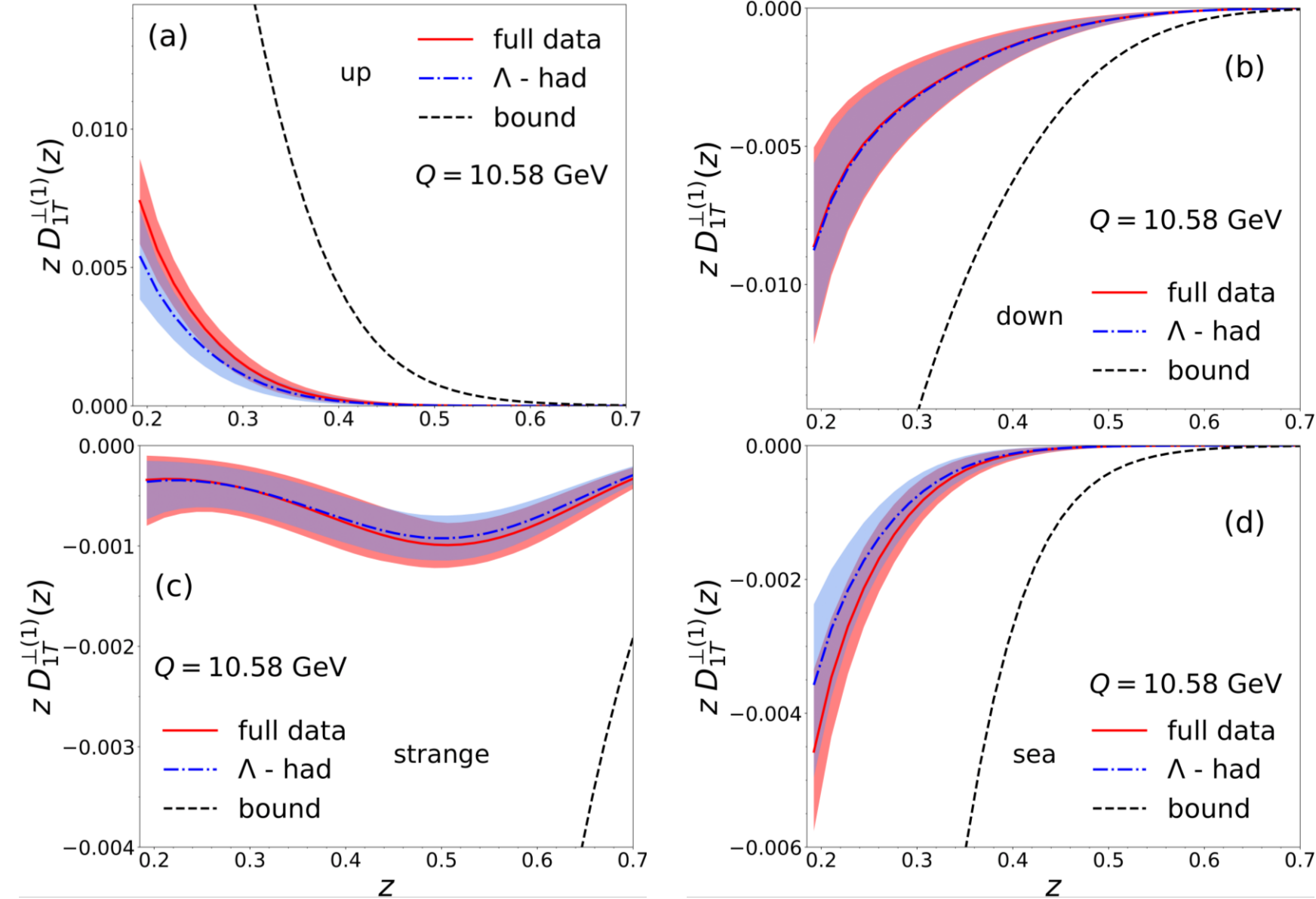
$\Lambda + \pi/K$ data Fit : $\chi_{dof}^2 = 1.26$

The two extractions are consistent

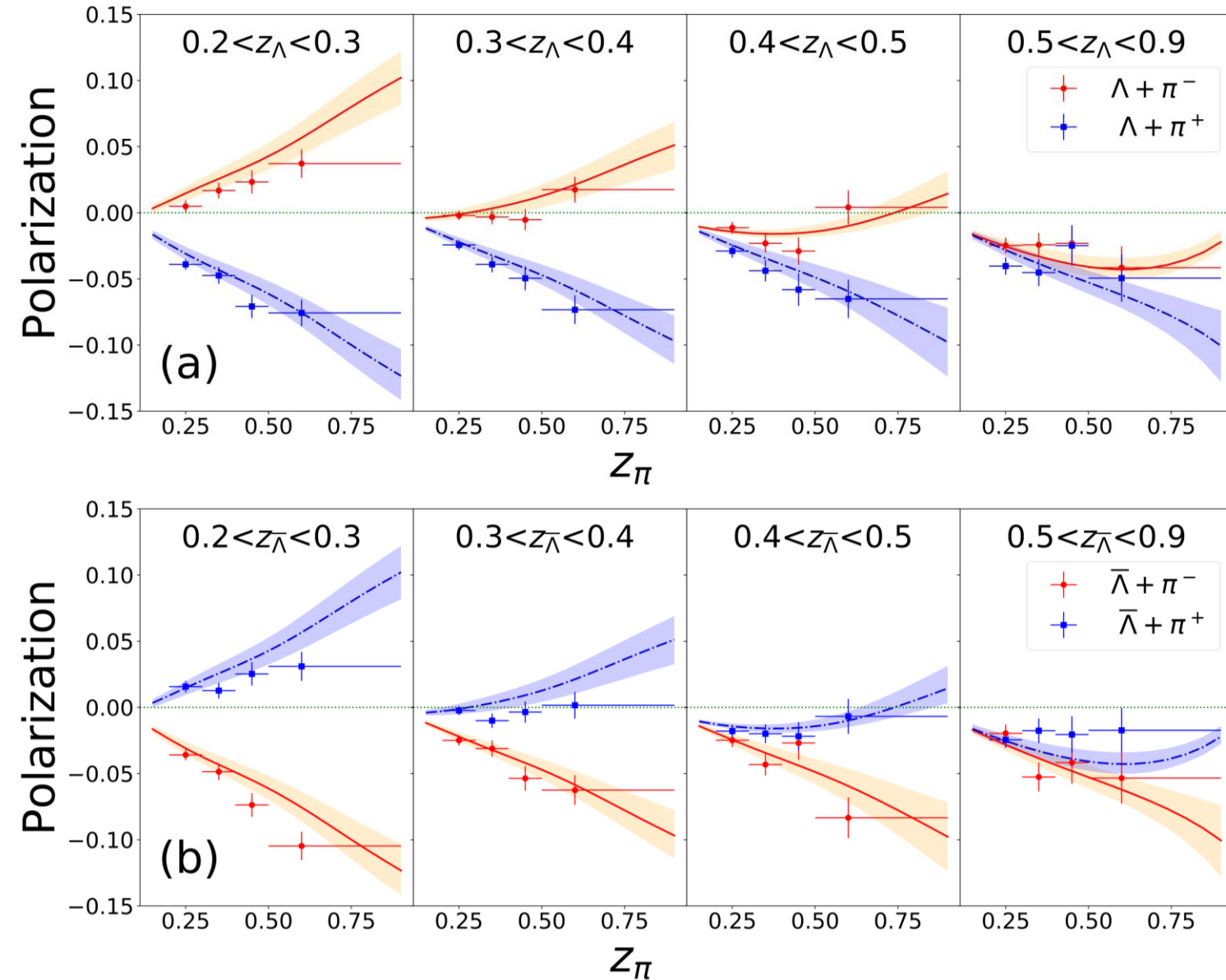


First moments

$$D_{1T}^{\perp(1)}(z) = \int d^2\mathbf{p}_\perp \frac{p_\perp}{2zm_h} \Delta D_{h^\dagger/q}(z, p_\perp)$$



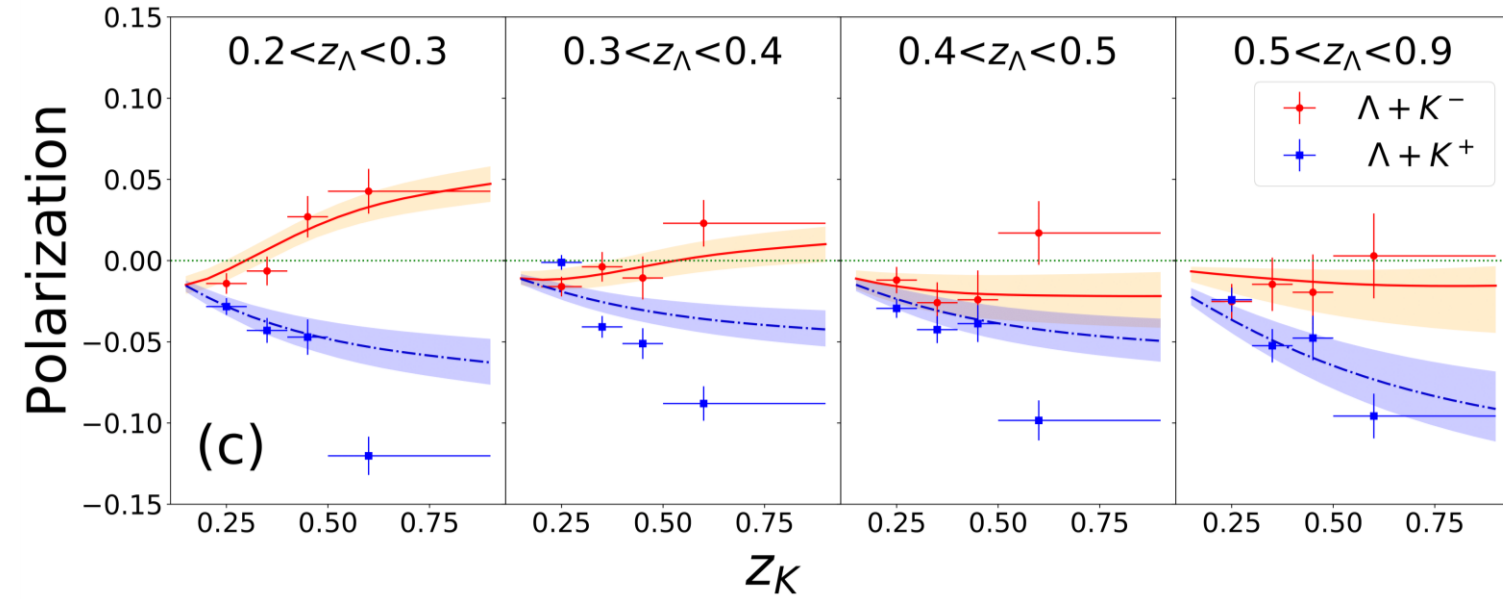
Lambda-pion



Bin excluded
 $z_\pi = [0.5 - 0.9]$

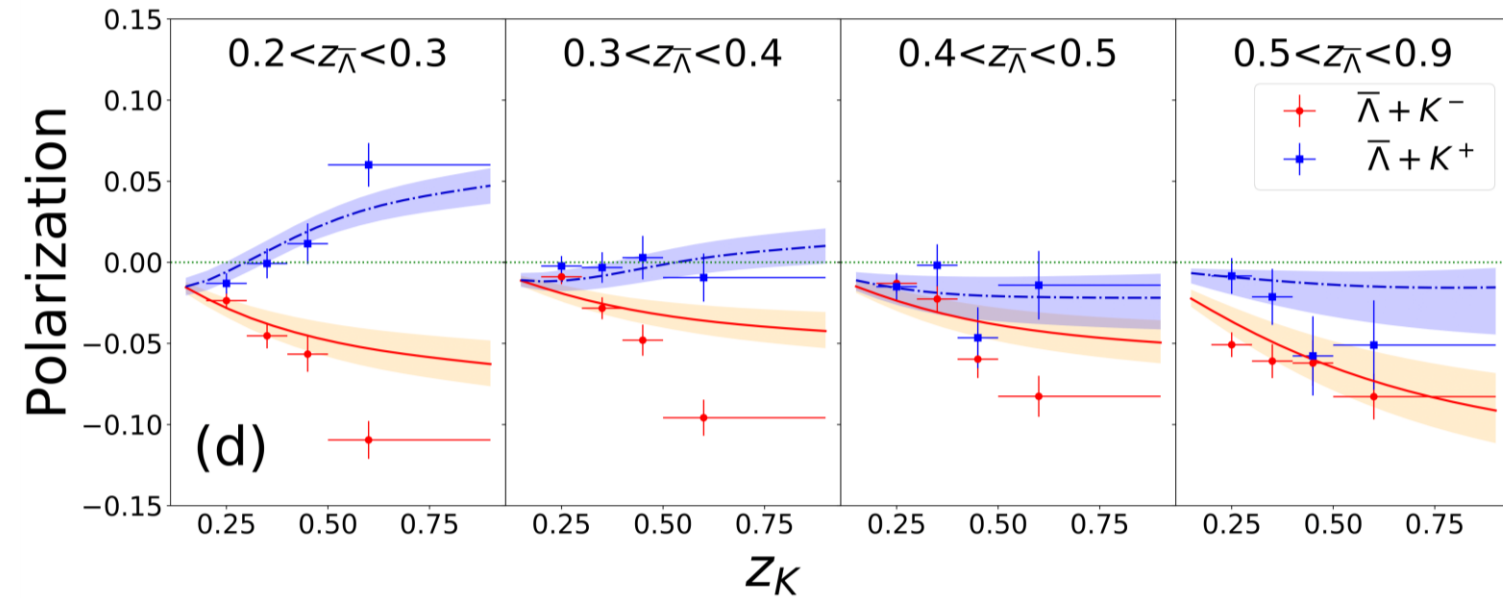
$$\chi_{dof}^2 = 1.94$$

Lambda-kaon

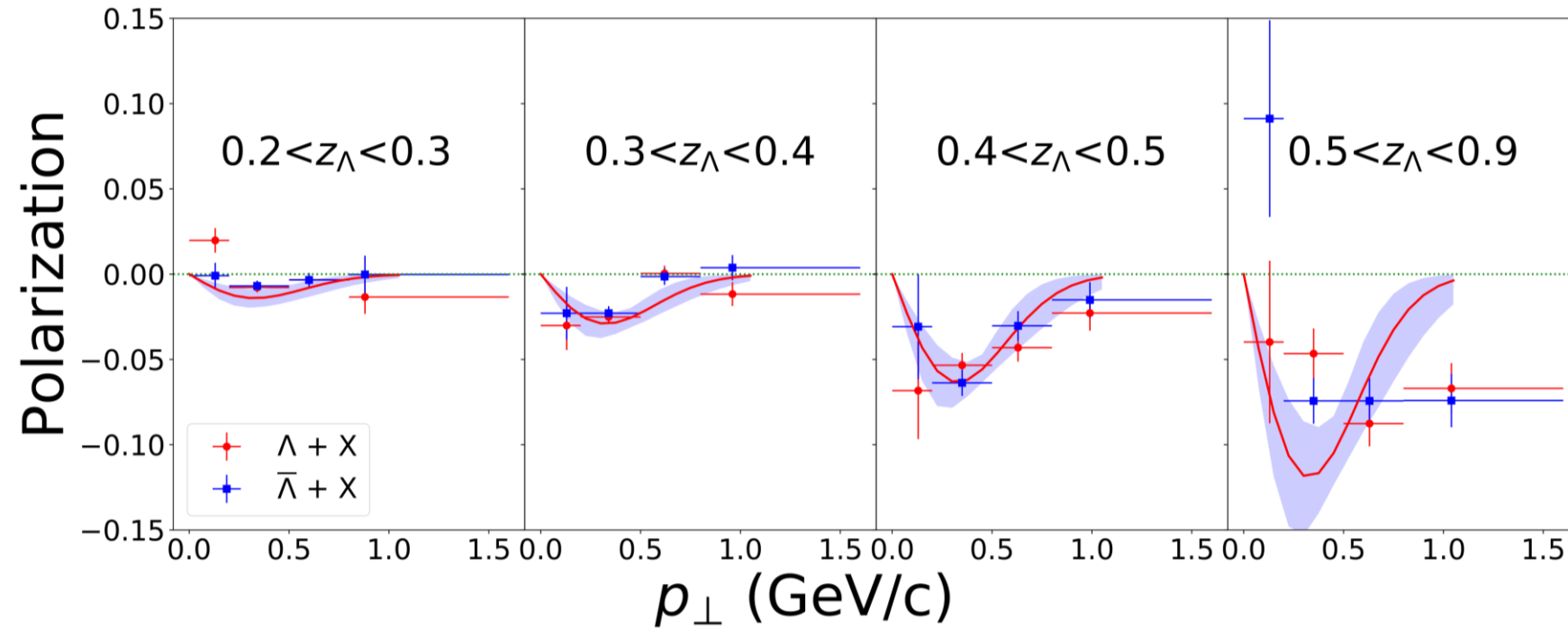


Bin excluded
 $z_K = [0.5 - 0.9]$

$$\chi_{dof}^2 = 1.94$$



Lambda-jet



Bin excluded
 $z_\Lambda = [0.5 - 0.9]$

$$\chi_{dof}^2 = 1.94$$

Expected features:

- $P_T = 0$ when $p_\perp = 0$;
- $P_T(\Lambda) = P_T(\bar{\Lambda})$.

SiDIS – Polarized Lambda Production

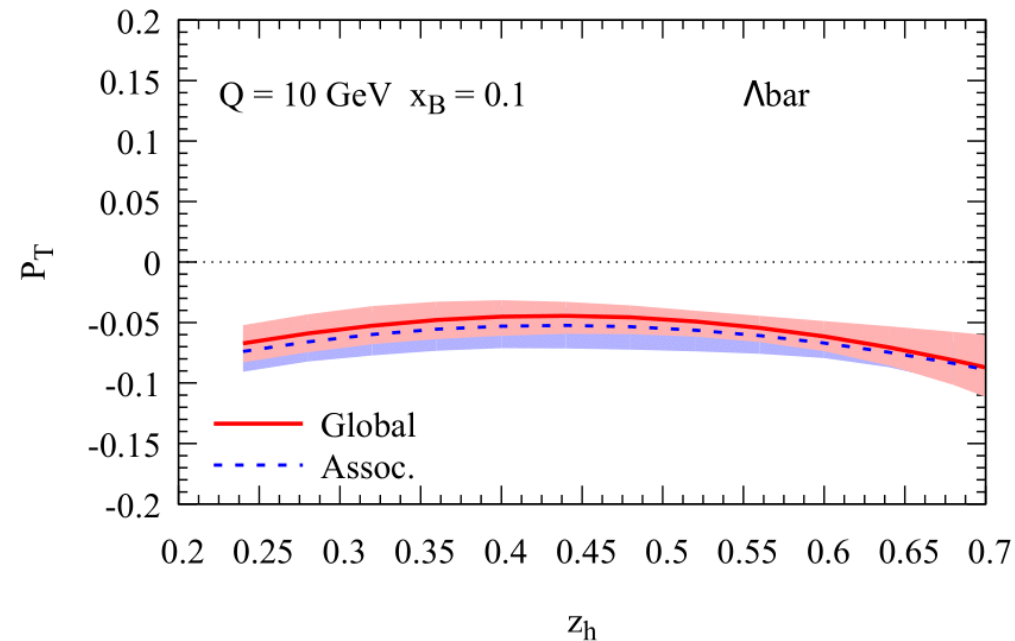
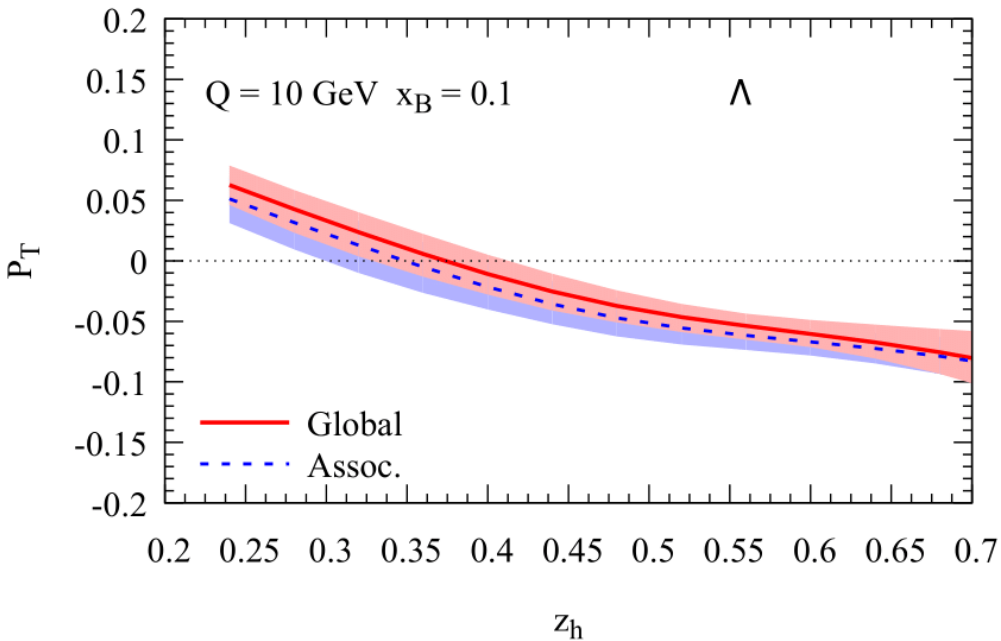
$$P_T(x_B, z_h) = \frac{\sqrt{2e\pi} \langle p_{\perp}^2 \rangle_p^2}{2M_p \langle p_{\perp}^2 \rangle} \frac{1}{\sqrt{\langle p_{\perp}^2 \rangle_p + \xi_p^2 \langle k_{\perp}^2 \rangle}} \times \frac{\sum_q f_{q/P}(x_B) \Delta D_{h\uparrow/q}(z_h)}{\sum_q f_{q/P}(x_B) D_{h/q}(z_h)}$$

Kinematical configuration compatible with EIC

Prediction for the Λ polarization:

- $x_B = 0.1$

$$\xi_p = z_h \left(1 - \frac{m_h^2}{z_h^2 Q^2} \frac{x_B}{1 - x_B} \right)$$



Conclusions

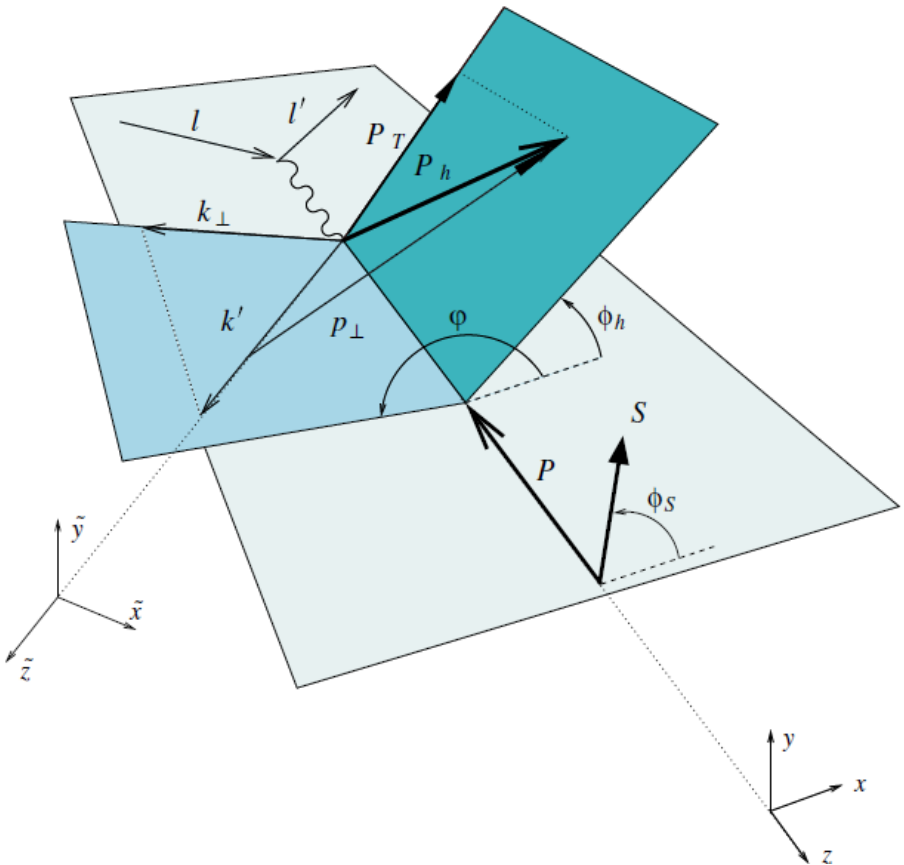
- First extraction of the Λ polarizing fragmentation function from Belle e^+e^- data;
- Clear separation in flavours: three different valence pFF needed;
- Sea-quark pFF also necessary;
- First indication of the p_{\perp} dependence within a Gaussian Ansatz;
- Prediction for transverse Λ polarization in SiDIS at EIC.



Backup Slides

SiDIS – Polarized Lambda Production

$e^-P \rightarrow e^- \Lambda$



$$P_T(x_B, z_h) = \frac{\sqrt{2e\pi} \langle p_{\perp}^2 \rangle_p^2}{2M_p \langle p_{\perp}^2 \rangle} \frac{1}{\sqrt{\langle p_{\perp}^2 \rangle_p + \xi_p^2 \langle k_{\perp}^2 \rangle}} \times \frac{\sum_q f_{q/P}(x_B) \Delta D_{h\uparrow/q}(z_h)}{\sum_q f_{q/P}(x_B) D_{h/q}(z_h)}$$

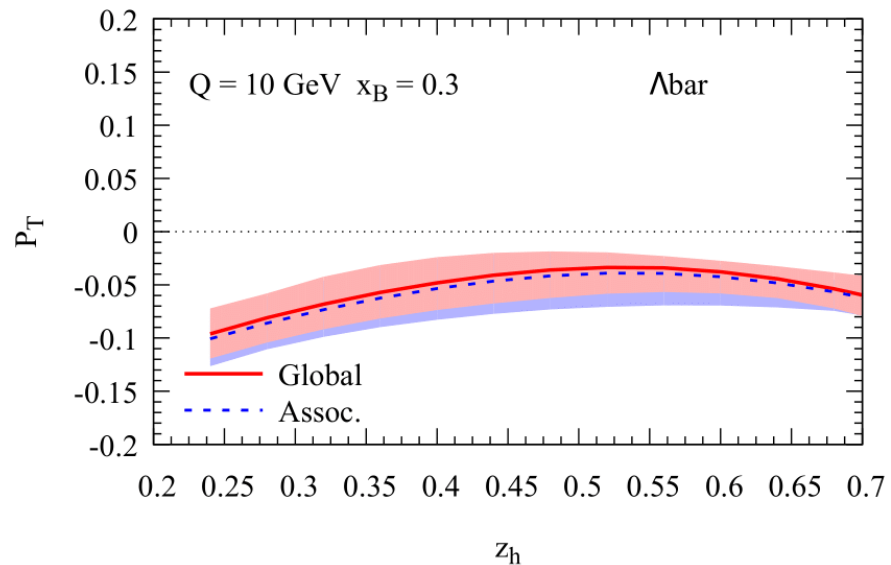
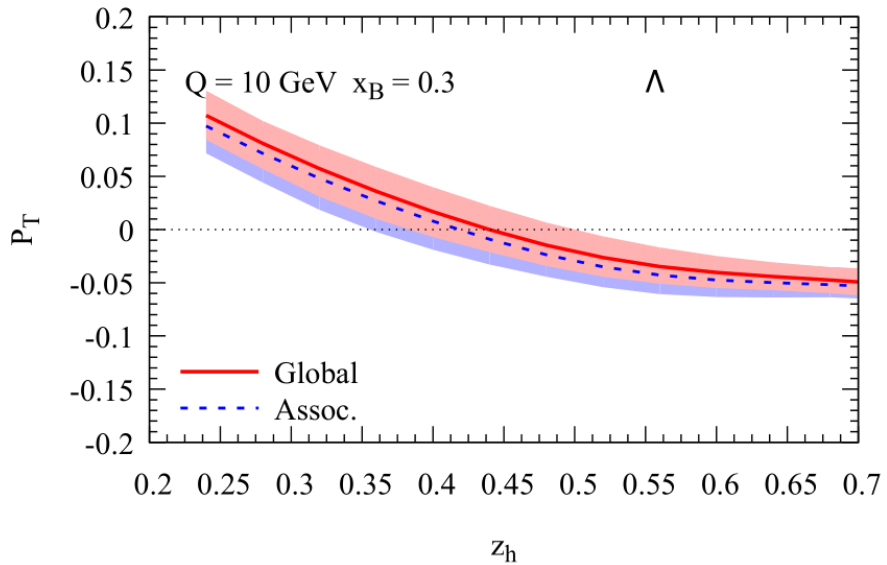
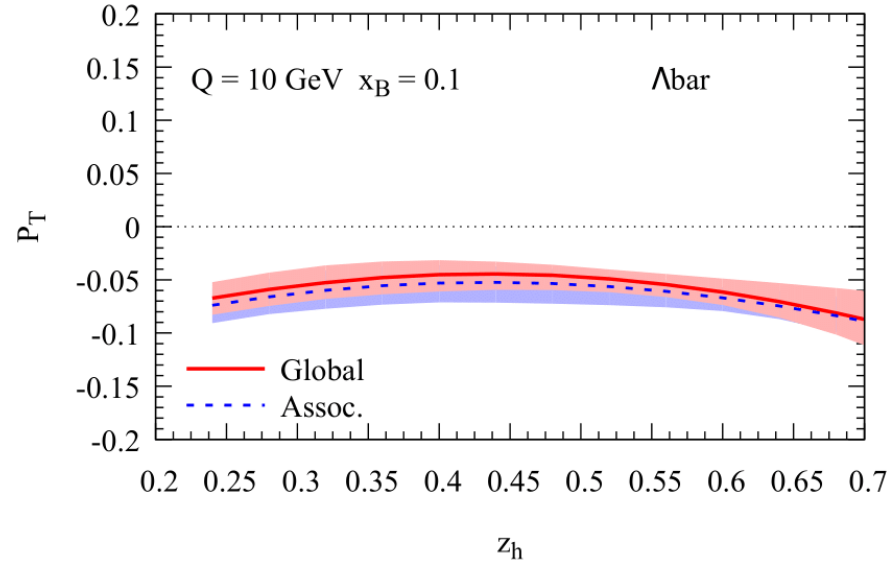
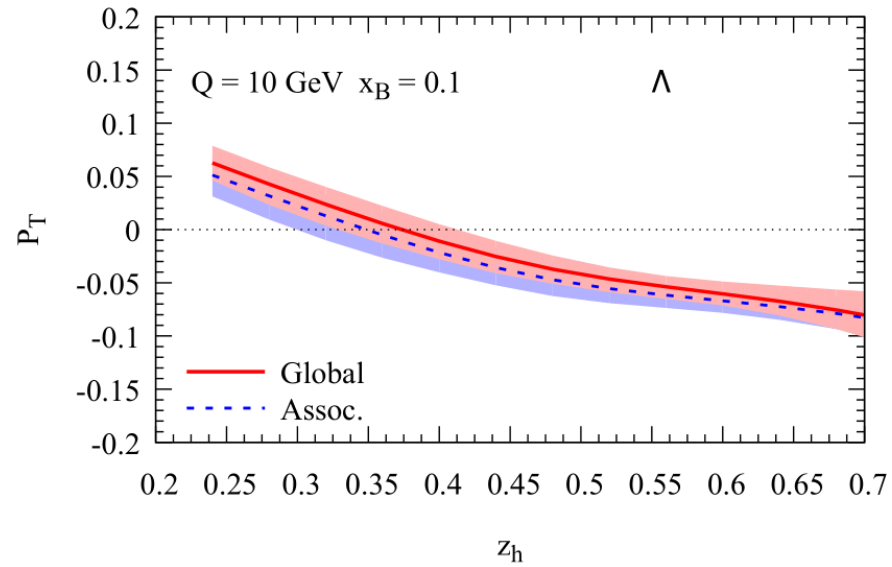
Proton PDF
CTEQ6L1

Lambda Polarizing FF

x_B Bjorken-x
 z_h energy fraction

$$\xi_p = z_h \left(1 - \frac{m_h^2}{z_h^2 Q^2} \frac{x_B}{1 - x_B} \right)$$

SiDIS – Polarized Lambda Production



Prediction for the Λ polarization:

- $x_B = 0.1$
- $x_B = 0.3$

Statistical Uncertainty Band

Multivariate Normal Distribution

MINUIT:

- Best fit parameters
- Covariance matrix
- Minimum Chi-square χ^2

$$\left. \begin{array}{l} \mu: \mathcal{N}_q^p \alpha_q \beta_q \langle p_{\perp}^2 \rangle_p \\ \Sigma \\ \chi^2 \end{array} \right\} \xrightarrow{\mu, \Sigma}$$

$$p(x; \mu, \Sigma) = \frac{1}{(2\pi)^{n/2} |\Sigma|^{1/2}} \exp \left(-\frac{1}{2} (x - \mu)^T \Sigma^{-1} (x - \mu) \right)$$

Generate a random set of parameter

$$x: (\mathcal{N}_q^p \alpha_q \beta_q \langle p_{\perp}^2 \rangle_p)$$

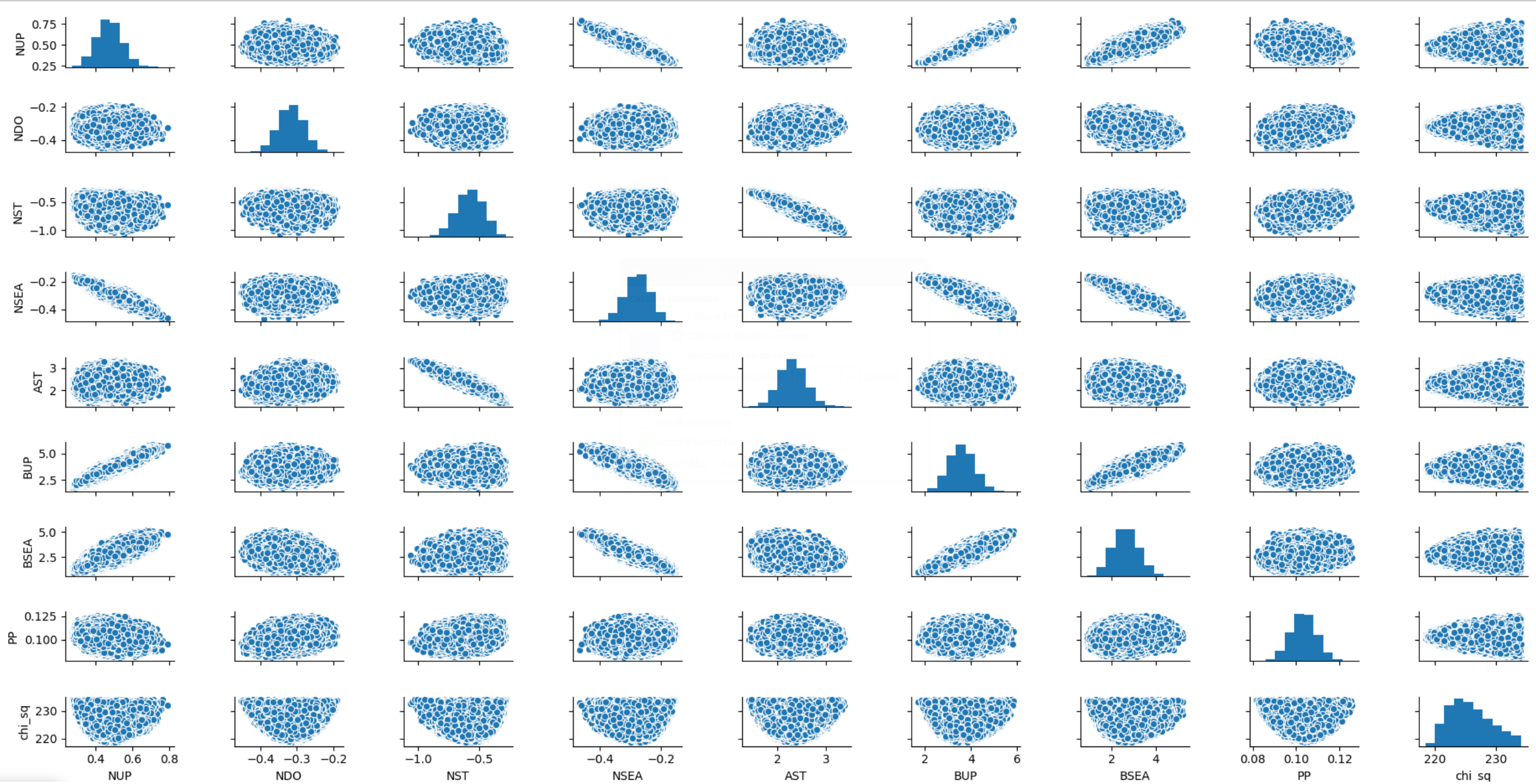
Calculate their own Chi-square

$$\chi'^2$$

Keep set if

$$\chi^2 \leq \chi'^2 \leq \chi^2 + \Delta\chi^2$$

- Minimum Chi-square: $\chi^2 = 232,8$
- Confidence interval $2\sigma \rightarrow 95,5\%$: $\Delta\chi^2 = 15,79$ for 8 parameters (χ^2 -distribution)



$$e^+ e^- \rightarrow h_1^\uparrow h_2 X$$

h_1, h_2 unpolarized

$$\begin{aligned} & \frac{d\sigma^{e^+e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} \\ &= \frac{6e^4 e_q^2}{64\pi \hat{s}} \left\{ \underbrace{D_{h_1/q}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2})}_{\text{Unpolarized FFs}} (1 + \cos^2 \theta) \right. \\ & \quad \left. + \frac{1}{4} \sin^2 \theta \Delta^N D_{h_1/q^\uparrow}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^\uparrow}(z_2, p_{\perp 2}) \cos(2\varphi_2 + \phi_1^{h_1}) \right\} \\ & \qquad \qquad \qquad \text{Collins FFs} \end{aligned}$$

h_1 polarized: Y , h_2 unpolarized

$$\begin{aligned} & P_Y^{h_1} \frac{d\sigma^{e^+e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} \\ &= \frac{6e^4 e_q^2}{64\pi \hat{s}} \left\{ \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2}) (1 + \cos^2 \theta) \right. \\ & \quad \left. + \frac{1}{2} \sin^2 \theta \Delta^- D_{S_Y/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^\uparrow}(z_2, p_{\perp 2}) \cos(2\varphi_2 + \phi_1^{h_1}) \right\} \\ & \qquad \qquad \qquad \text{Collins FF} \\ & \qquad \qquad \qquad \text{FF } h_1 \text{ transv. Pol.} \end{aligned}$$

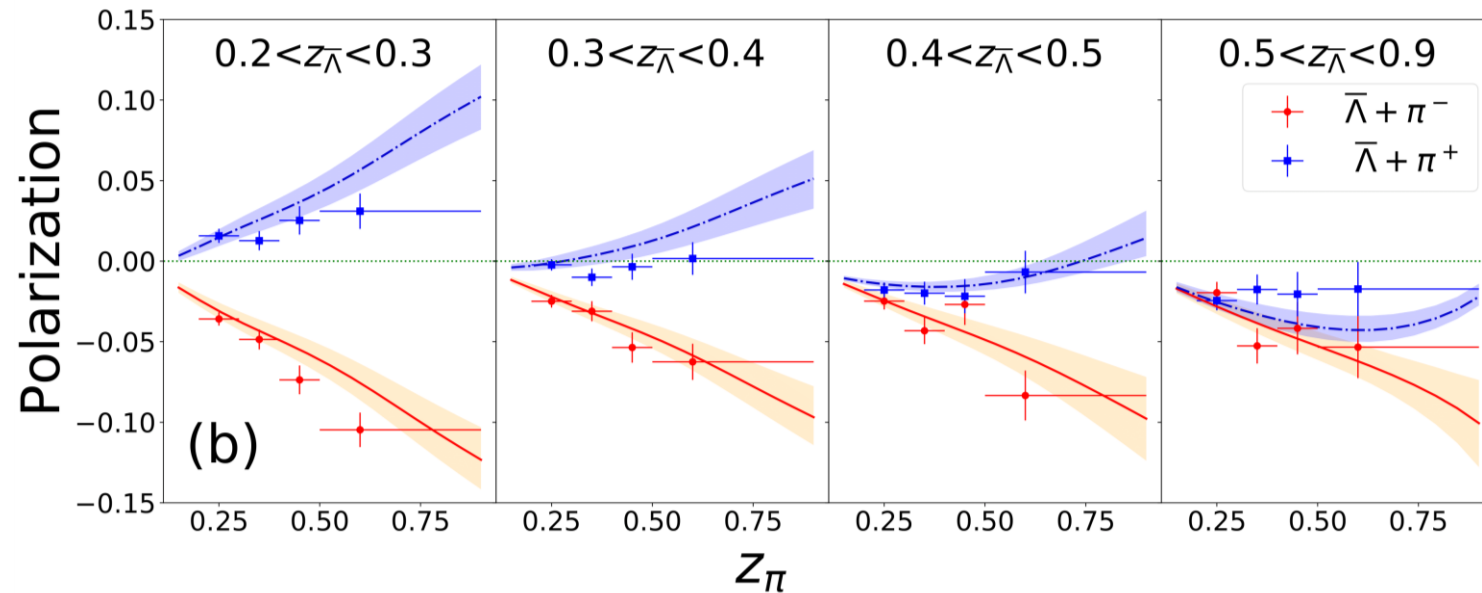
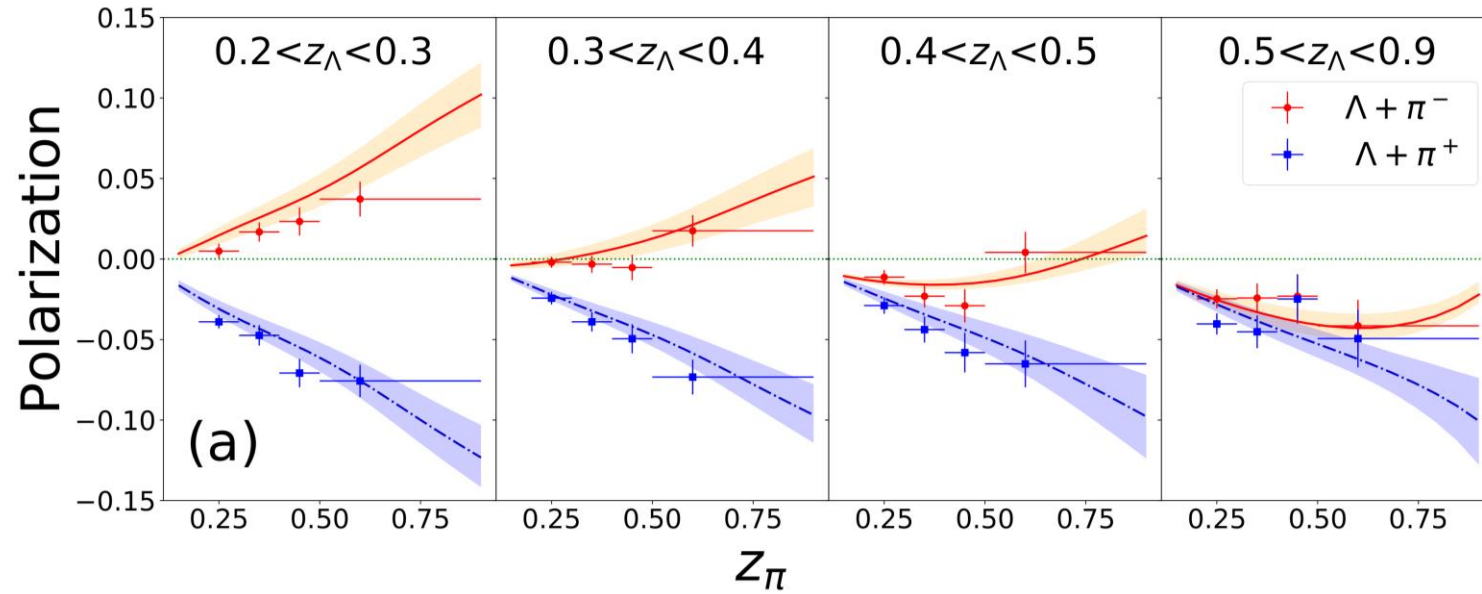
h_1 polarized: X , h_2 unpolarized

$$\begin{aligned} & P_X^{h_1} \frac{d\sigma^{e^+e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} \\ &= \frac{3e^4 e_q^2}{64\pi \hat{s}} \Delta D_{S_X/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^\uparrow}(z_2, p_{\perp 2}) \sin^2 \theta \sin(2\varphi_2 + \phi_1^{h_1}) \\ & \qquad \qquad \qquad \text{FF } h_1 \text{ transv. Pol.} \qquad \qquad \qquad \text{Collins FF} \end{aligned}$$

h_1 polarized: Z , h_2 unpolarized

$$\begin{aligned} & P_Z^{h_1} \frac{d\sigma^{e^+e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 p_{\perp 1} dz_2 d^2 p_{\perp 2}} \\ &= \frac{3e^4 e_q^2}{64\pi \hat{s}} \Delta D_{S_Z/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}^\uparrow}(z_2, p_{\perp 2}) \sin^2 \theta \sin(2\varphi_2 + \phi_1^{h_1}) \\ & \qquad \qquad \qquad \text{FF } h_1 \text{ long. Pol.} \qquad \qquad \qquad \text{Collins FF} \end{aligned}$$

Lambda-pion



Data fitted:

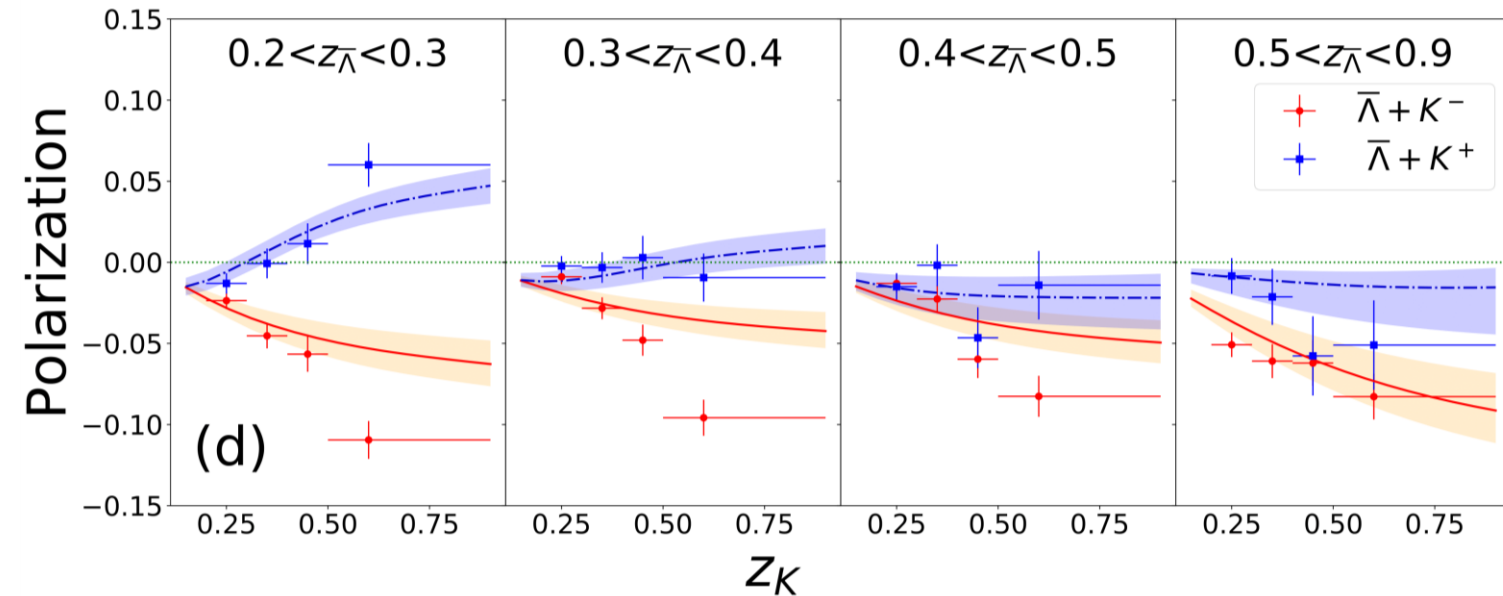
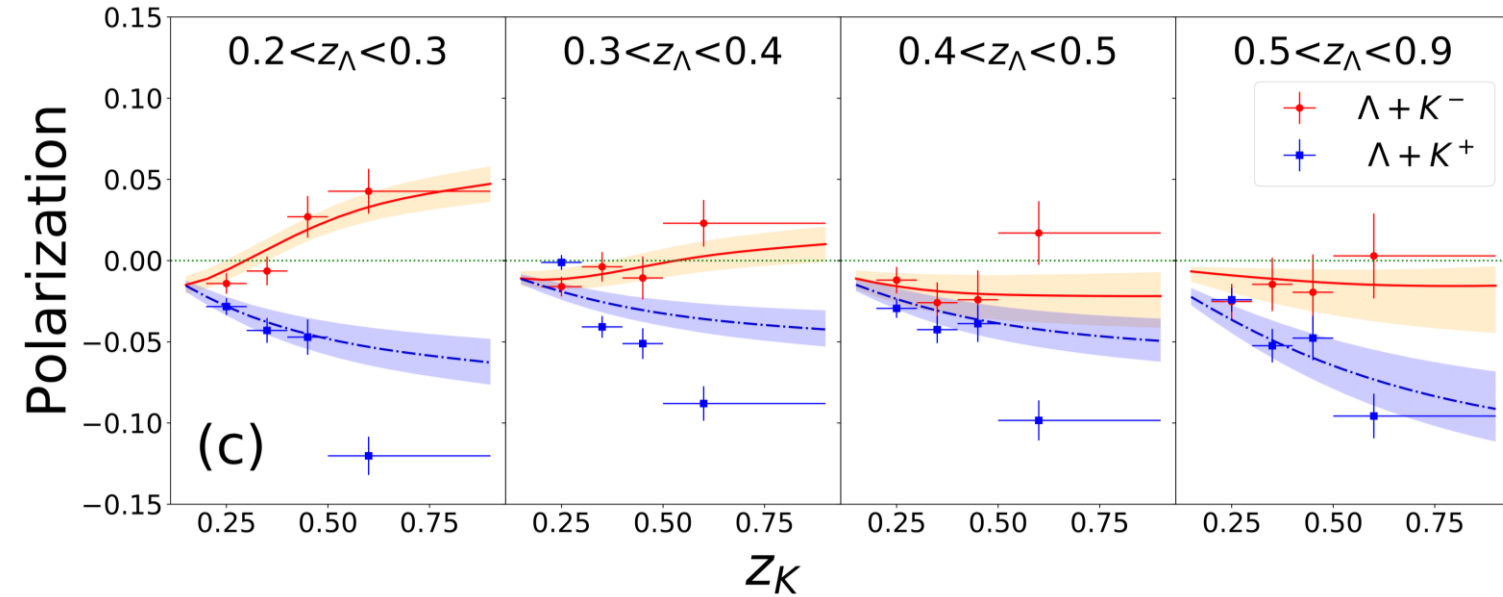
- $z_\pi < 0.50$

$$\chi_{dof}^2 = 1.94$$

Lambda pion

- Charge-conjugation symmetry $P_n(\Lambda\pi^+) = P_n(\bar{\Lambda}\pi^-)$
- Data give information on the pFFs for u and d
- $P_n(\Lambda\pi^+)$ negative, dominated by down pFF
- $P_n(\Lambda\pi^-)$ positive, dominated by up pFF
- $P_n(\Lambda\pi^-)$ strong reduction due to the large suppression of the up pFF
- Small z_a sea pFFs become important and negative,
- $P_n(\Lambda\pi^+)$ up and down cancel each other, sea pFF leads to large, and negative, values of the transverse polarization
- $P_n(\Lambda\pi^-)$ sea pFF partial reduction of the up pFF

Lambda-kaon



Data fitted:

- $z_K < 0.50$

$$\chi_{dof}^2 = 1.94$$

Lambda-kaon

- Charge-conjugation symmetry $P_n(\Lambda K^+) = P_n(\bar{\Lambda} K^-)$
- Similar pattern pion
- Data give information on the pFFs for u and s
- $P_n(\Lambda K^+)$ negative, dominated by strange pFF
- $P_n(\Lambda K^-)$ positive, dominated by up pFF
- $P_n(\Lambda K^-)$ strong reduction due to the large suppression of the up pFF
- Small z_a sea pFFs become important and negative,

$$e^+ e^- \rightarrow h_1(\text{jet})X$$

For Spin-1/2 hadron production, two possible cross sections :

Unpolarised
hadron

$$\frac{d\sigma^{e^+e^- \rightarrow h_1(\text{jet})X}}{d \cos \theta dz_1 d^2 p_{\perp 1}} = \sum_{q_c} \frac{3e^4}{32\pi s} e_q^2 (1 + \cos^2 \theta) \hat{D}_{h/q}(z_1, p_{1\perp h_1})$$

Transversely
polarised hadron

$$P_Y^{h_1} \frac{d\sigma^{e^+e^- \rightarrow h_1(\text{jet})X}}{d \cos \theta dz_1 d^2 p_{\perp 1}} = \sum_{q_c} \frac{3e^4}{32\pi s} e_q^2 (1 + \cos^2 \theta) \Delta D_{S_Y/q}^{h_1}(z_1, p_{1\perp h_1})$$

ratio →

Hadron polarisation

$$\mathcal{P}_T(z_1, p_{\perp 1}) = \frac{\sum_q e_q^2 \Delta D_{h_1^\uparrow/q}(z_1, p_{\perp 1})}{\sum_q e_q^2 D_{h_1/q}(z_1, p_{\perp 1})}$$

$$\hat{X}_{h_1} = \hat{Y}_{h_1} \times \hat{Z}_{h_1}$$

$$\hat{Y}_{h_1} = \frac{\hat{q}_1 \times P_{h_1}}{|\hat{q}_1 \times P_{h_1}|}$$

$$\hat{Z}_{h_1} = \frac{P_{h_1}}{|P_{h_1}|}$$

$$\mathcal{P}^{h_1} = P_x^{h_1} \hat{X}_1 + P_y^{h_1} \hat{Y}_1 + P_z^{h_1} \hat{Z}_1$$

$$\hat{n} = -\hat{P}_{h_2} \times \hat{P}_{h_1}$$

$$\mathcal{P}^{h_1} \cdot \hat{n} = P_x^{h_1} \cos \tilde{\phi} + P_y^{h_1} \sin \tilde{\phi}$$

$$\cos \tilde{\phi} \simeq \frac{z_{h_1} p_{\perp 2}}{z_{h_2} p_{\perp 1}} \sin(\phi_1 - \varphi_2)$$

$$\sin \tilde{\phi} \simeq \frac{P_{1T}}{p_{\perp 1}} - \frac{z_{h_1} p_{\perp 2}}{z_{h_2} p_{\perp 1}} \cos(\phi_1 - \varphi_2)$$

