

# QCD (and QED) in a Euclidean finite volume

*EuroPLEX Summer School 2021*



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August 23rd-27th



THE UNIVERSITY  
of EDINBURGH

# Five-day outline

- Warm-up and definitions
  - Meaning of Euclidean
  - Finite-volume set-up
- $e^{-mL}$  volume effects
  - Mass, matrix element in  $\lambda\phi^4, g\phi^3$
  - Pion mass
  - LO-HVP for  $(g - 2)_\mu$
  - Bethe-Salpeter kernel
- $2 \rightarrow 2$  finite-volume (f.v.) formalism
  - Scattering basics
  - Derivation
  - Generalizations
  - Applications
- $(1+\mathcal{J}) \rightarrow 2$  f.v. formalism
  - Derivation
  - Expansions
  - Applications
- $2 + \mathcal{J} \rightarrow 2$  f.v. formalism
  - Derivation
  - Testing the result
  - Numerical explorations
- Non-local matrix elements
  - Derivation
  - Applications
- $3 \rightarrow 3$  f.v. formalism
  - New complications
  - Derivation ( $E_n(L)$  to  $\mathcal{K}_{df,3}$ )
  - Integral equations ( $\mathcal{K}_{df,3}$  to  $\mathcal{M}_3$ )
  - Testing the result
  - Numerical explorations/calculations
- Finite-volume QED + QCD
  - Different formalisms
  - Basic examples

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## Warm-up and definitions

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## $e^{-mL}$ volume effects

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## $2 \rightarrow 2$ finite-volume (f.v.) formalism

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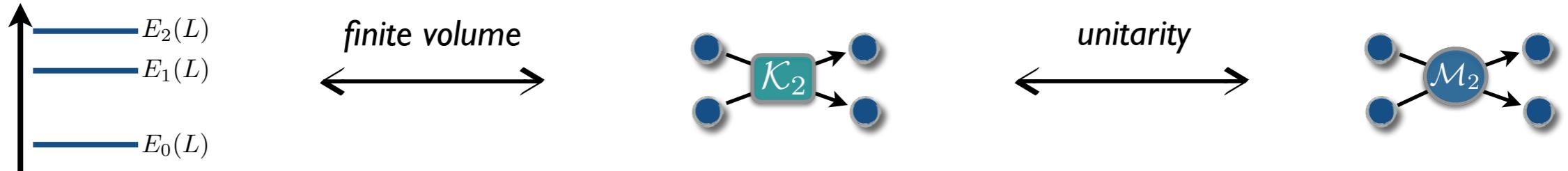
## Finite-volume QED + QCD

- Different formalisms
- Basic examples

# Result

$$\det[\mathcal{K}^{-1}(s) + F(P, L)] = 0$$

$F(P, L) \equiv$  Matrix of known geometric functions



Holds only for two-particle energies  $s < (4m)^2$

Neglects  $e^{-mL}$

Generalized to *non-degenerate masses, multiple channels, spinning particles*

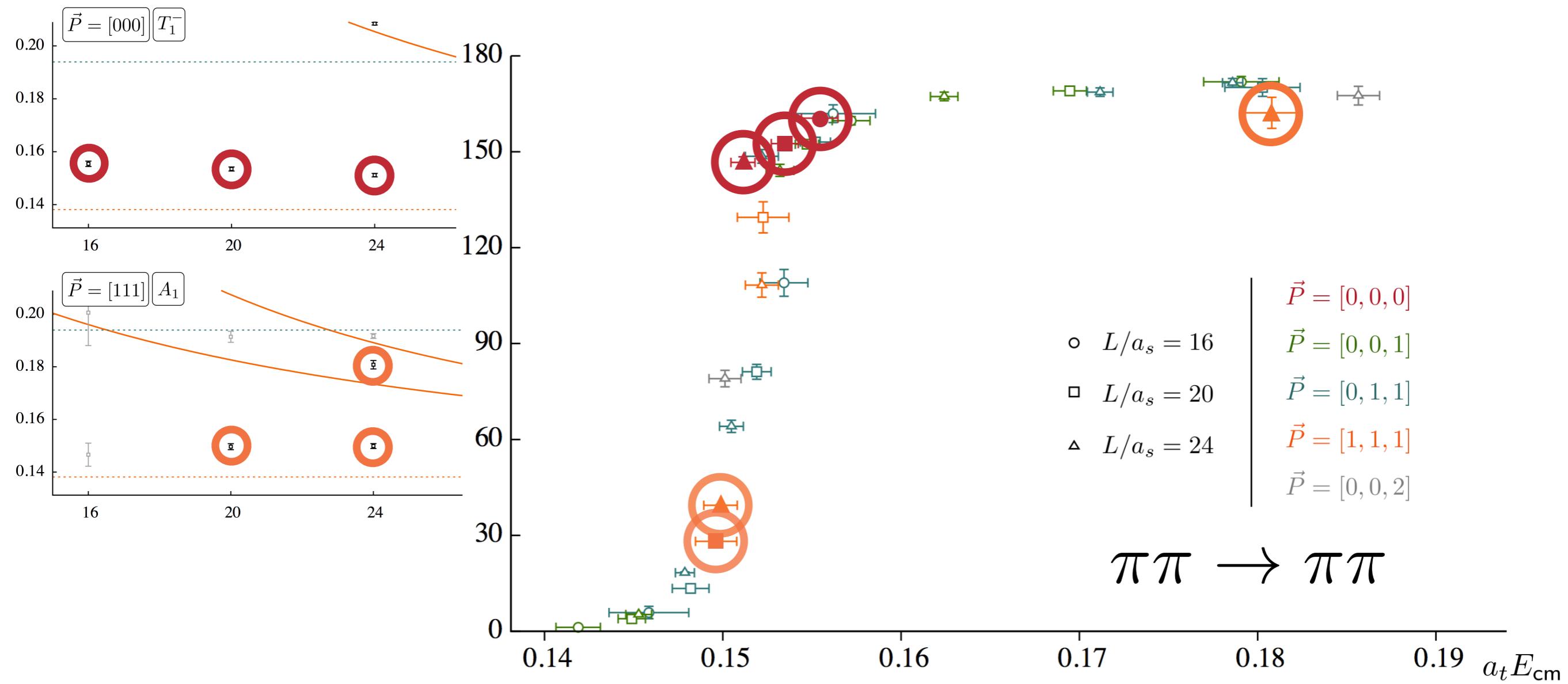
*Encodes angular momentum mixing*

- Huang, Yang (1958) • Lüscher (1986, 1991) • Rummukainen, Gottlieb (1995)  
Kim, Sachrajda, Sharpe (2005) • Christ, Kim, Yamazaki (2005) • He, Feng, Liu (2005)  
Leskovec, Prelovsek (2012) • Bernard *et. al.* (2012) • MTH, Sharpe (2012) • Briceño, Davoudi (2012)  
Li, Liu (2013) • Briceño (2014)

# Using the result

## □ Single-channel case (*pions in a p-wave*)

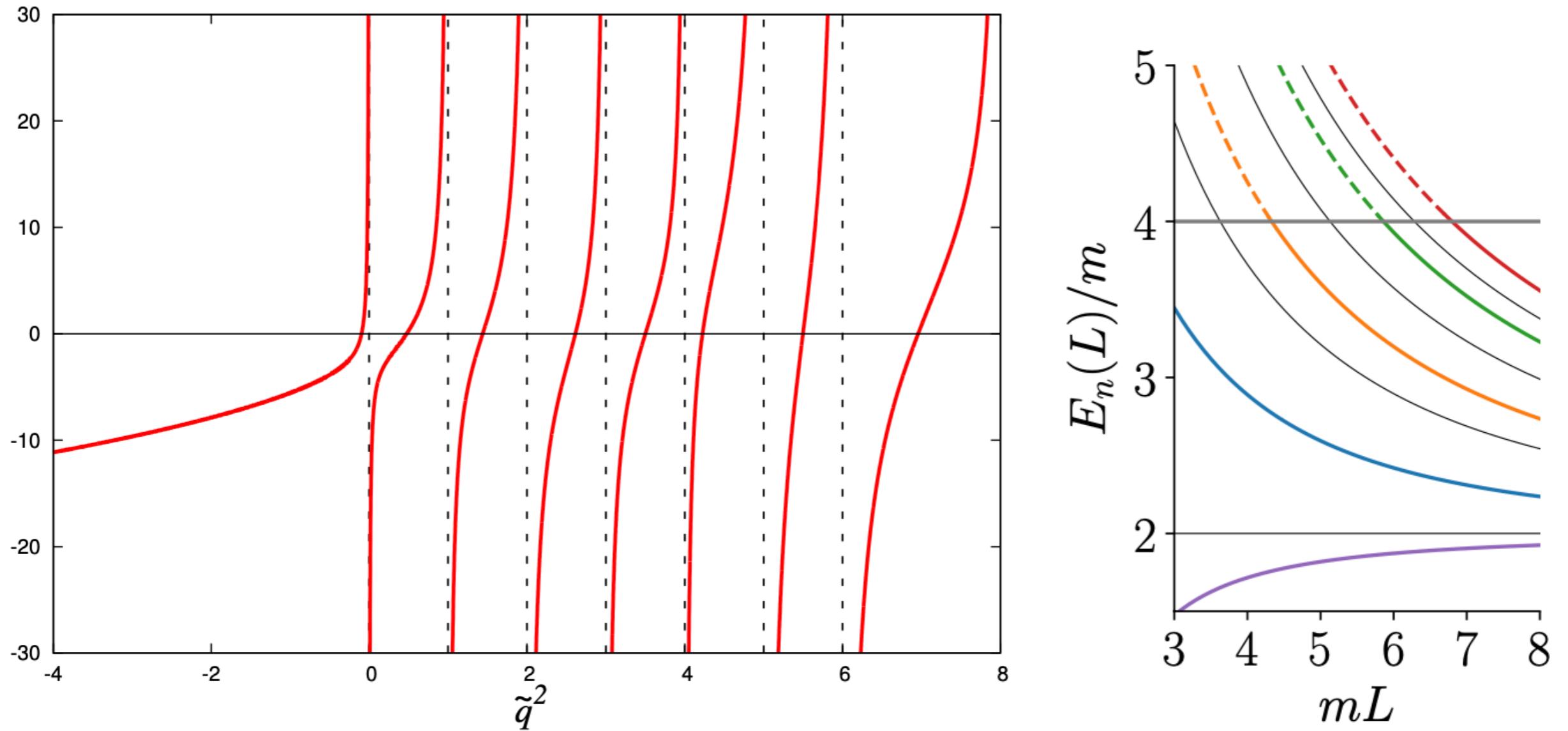
$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$



- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

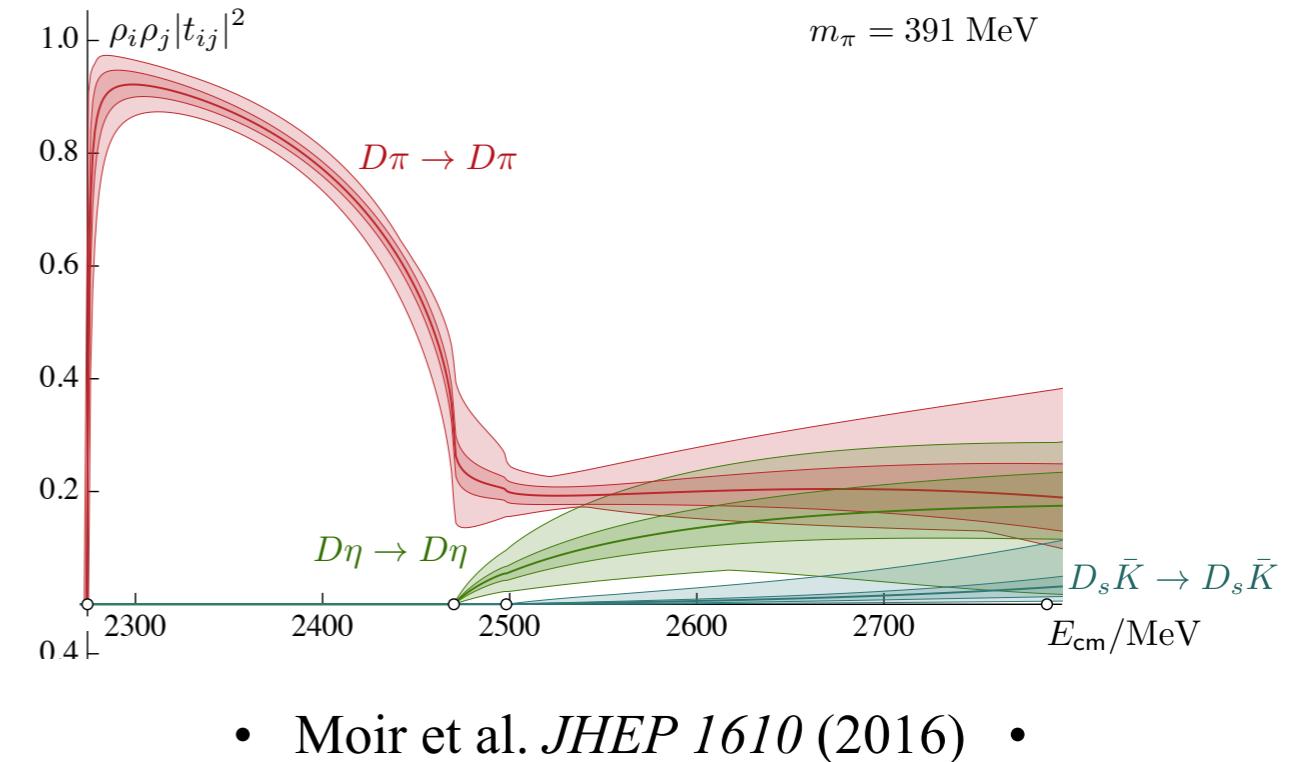
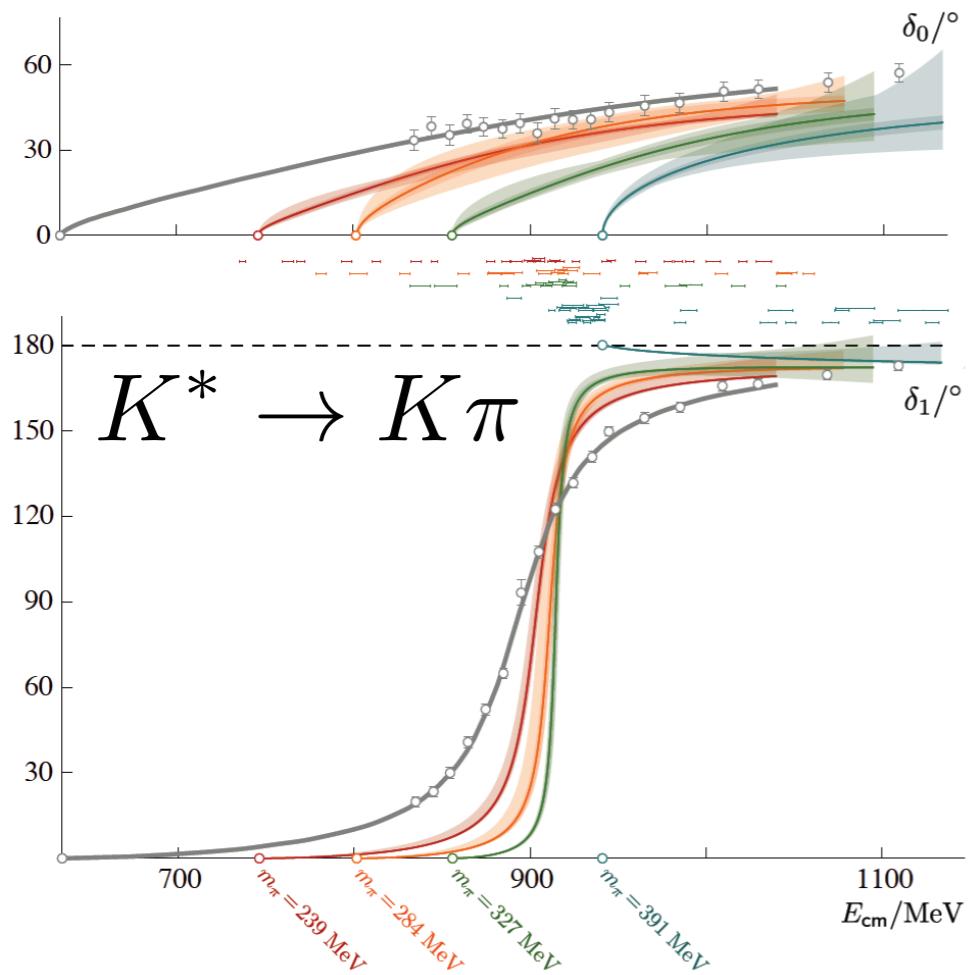
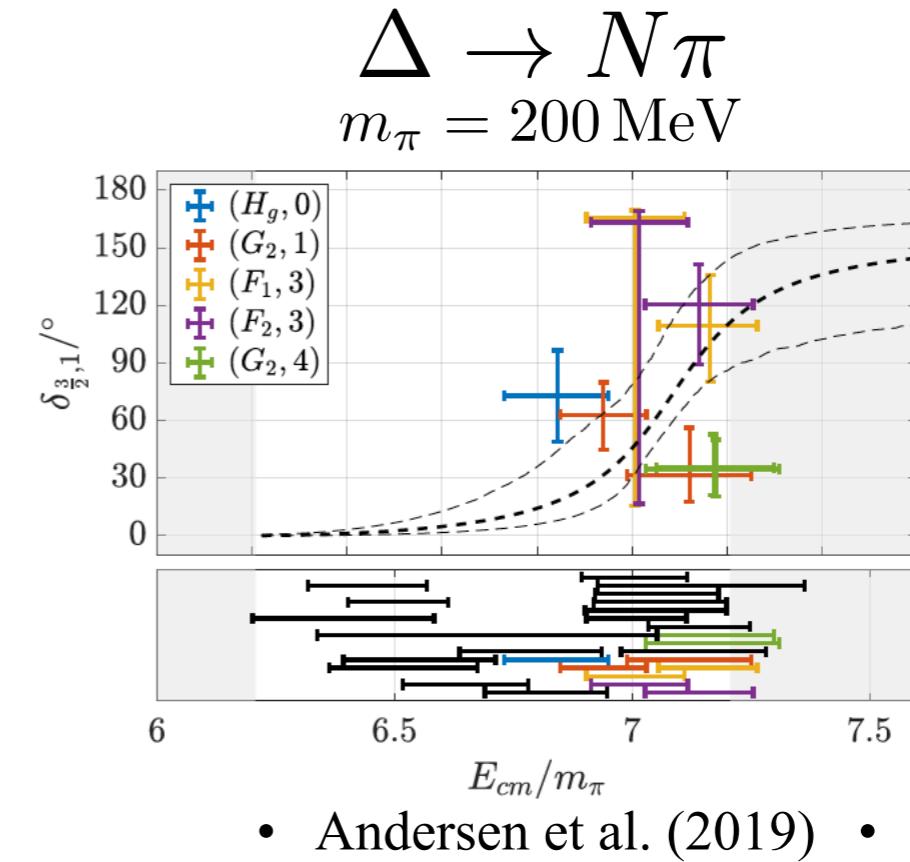
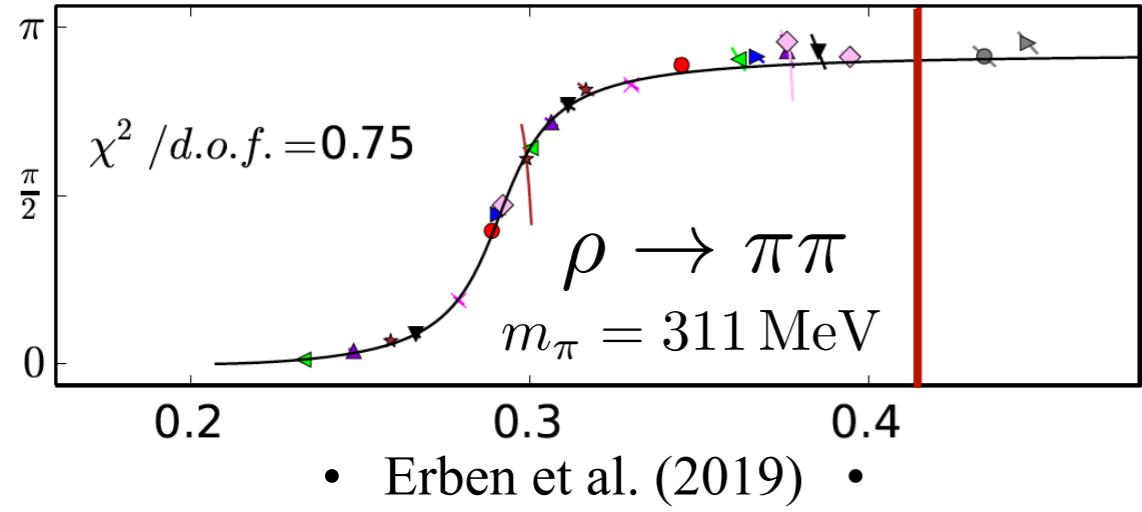
# Visualizing the result

$$p \cot \delta(p) = \frac{1}{\pi L} \left[ \sum_{\mathbf{n}} -\text{p.v.} \int d^3 \mathbf{n} \right] \frac{1}{\mathbf{n}^2 - q^2}$$

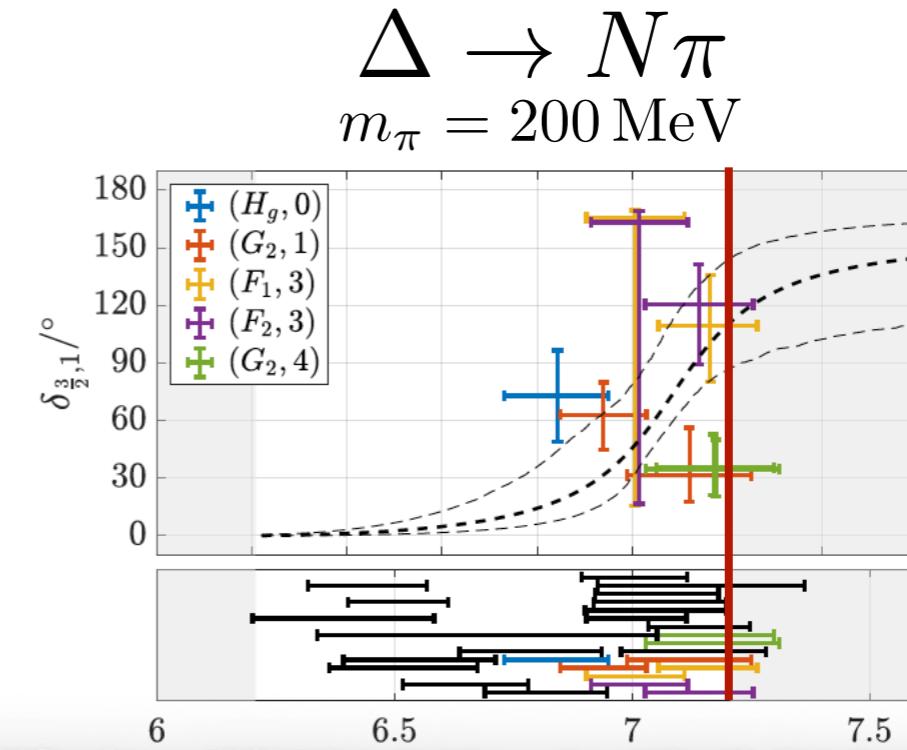
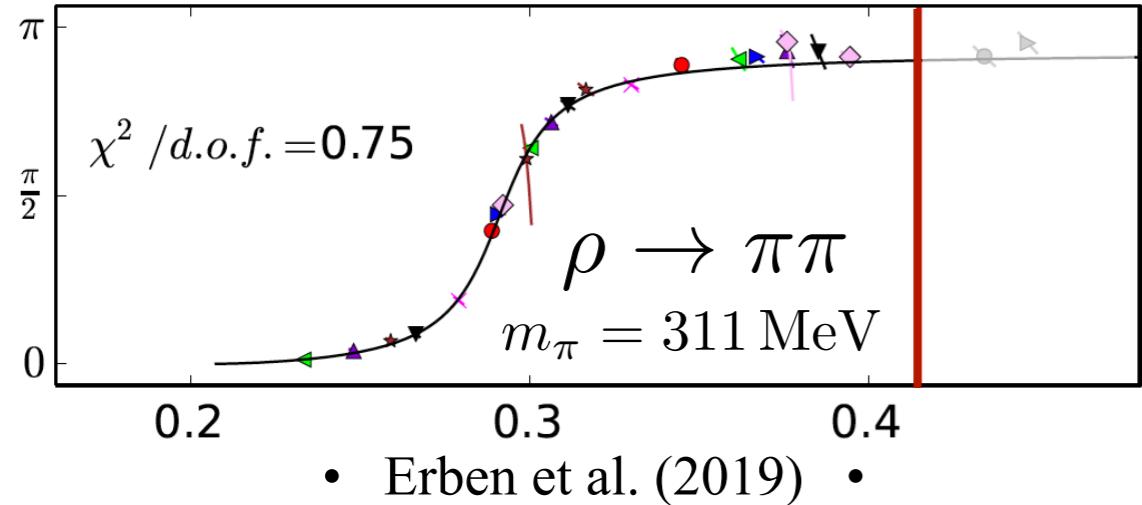


- plot taken from Luu and Savage (2011) •

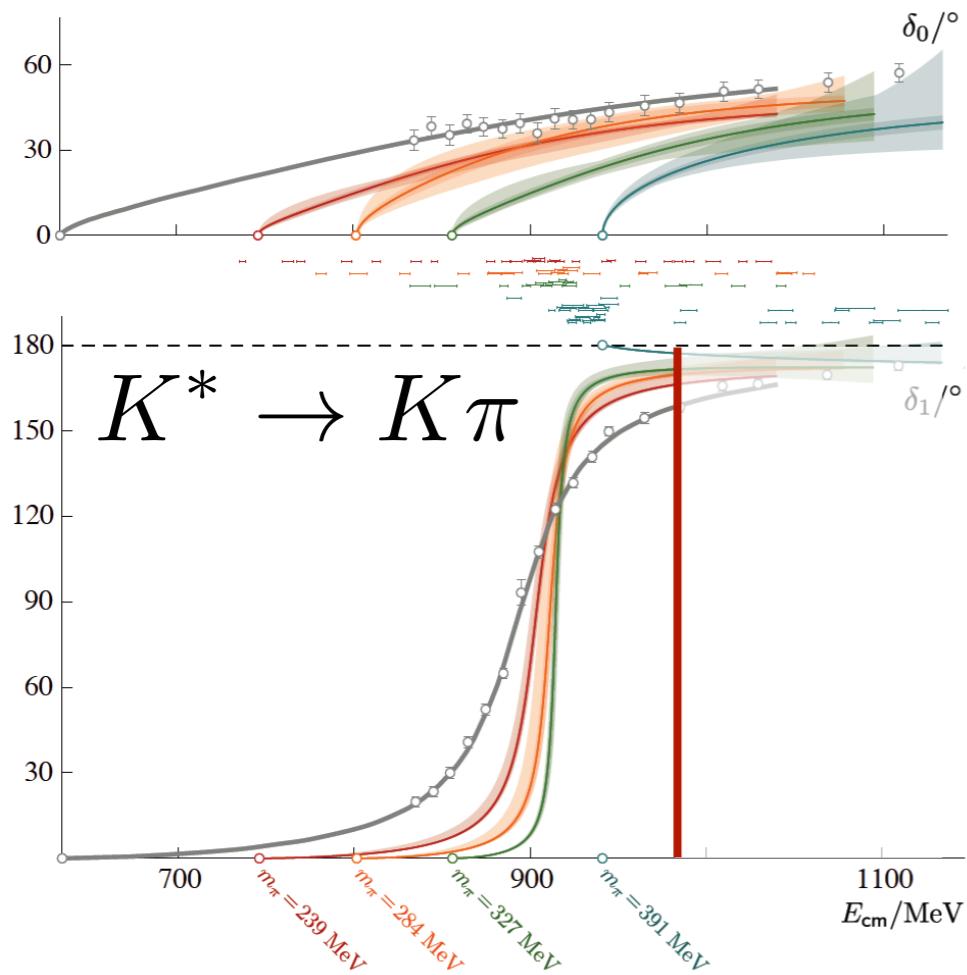
# Many applications



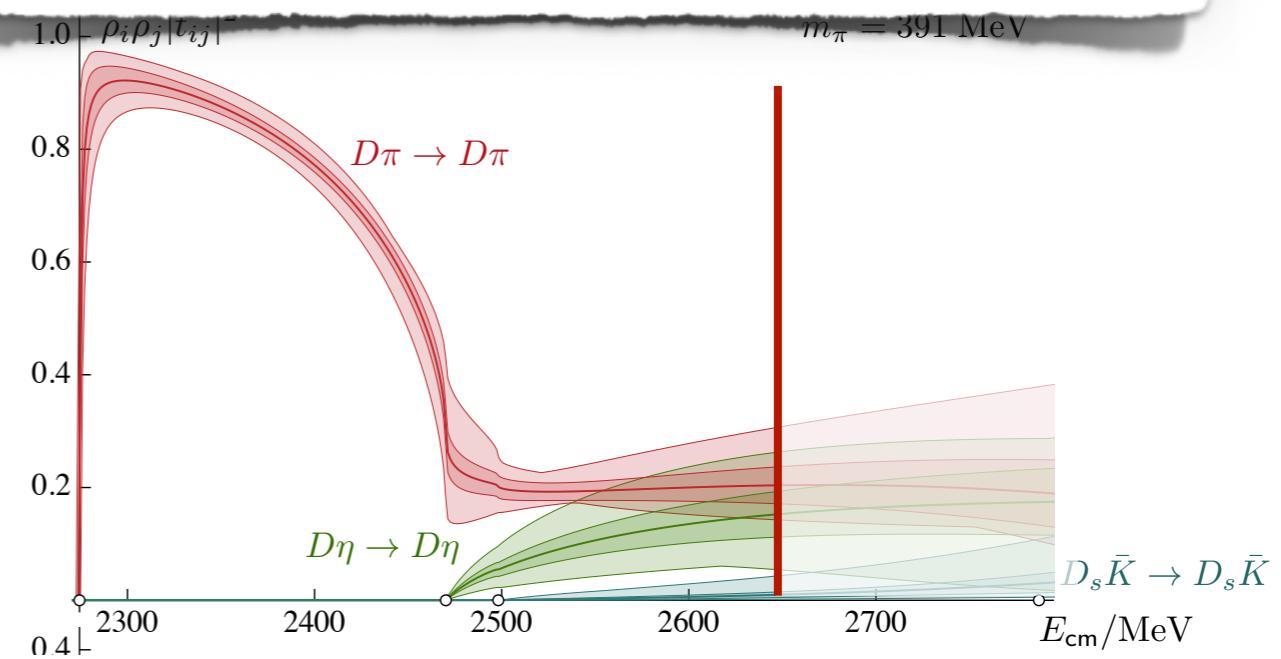
# Many applications



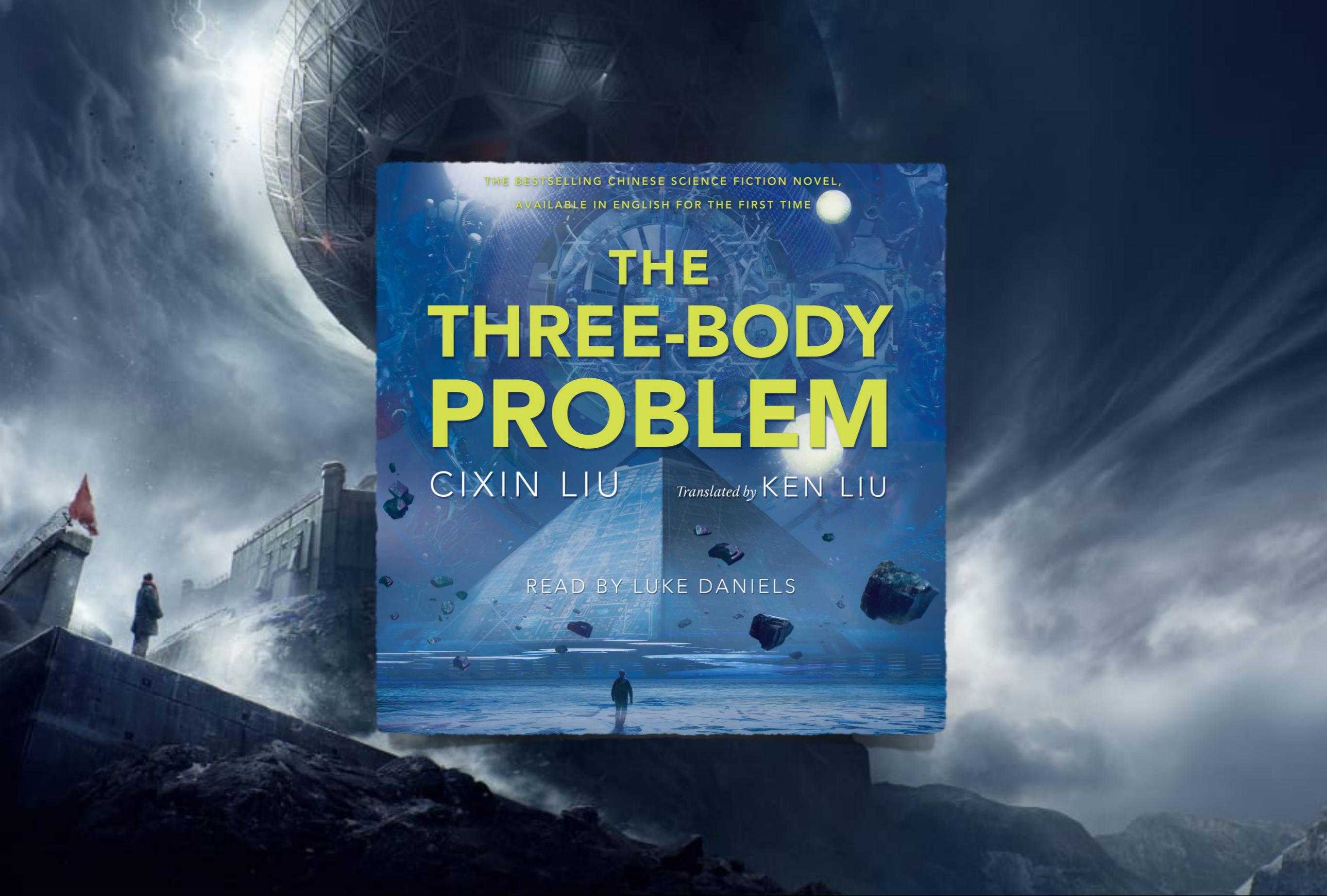
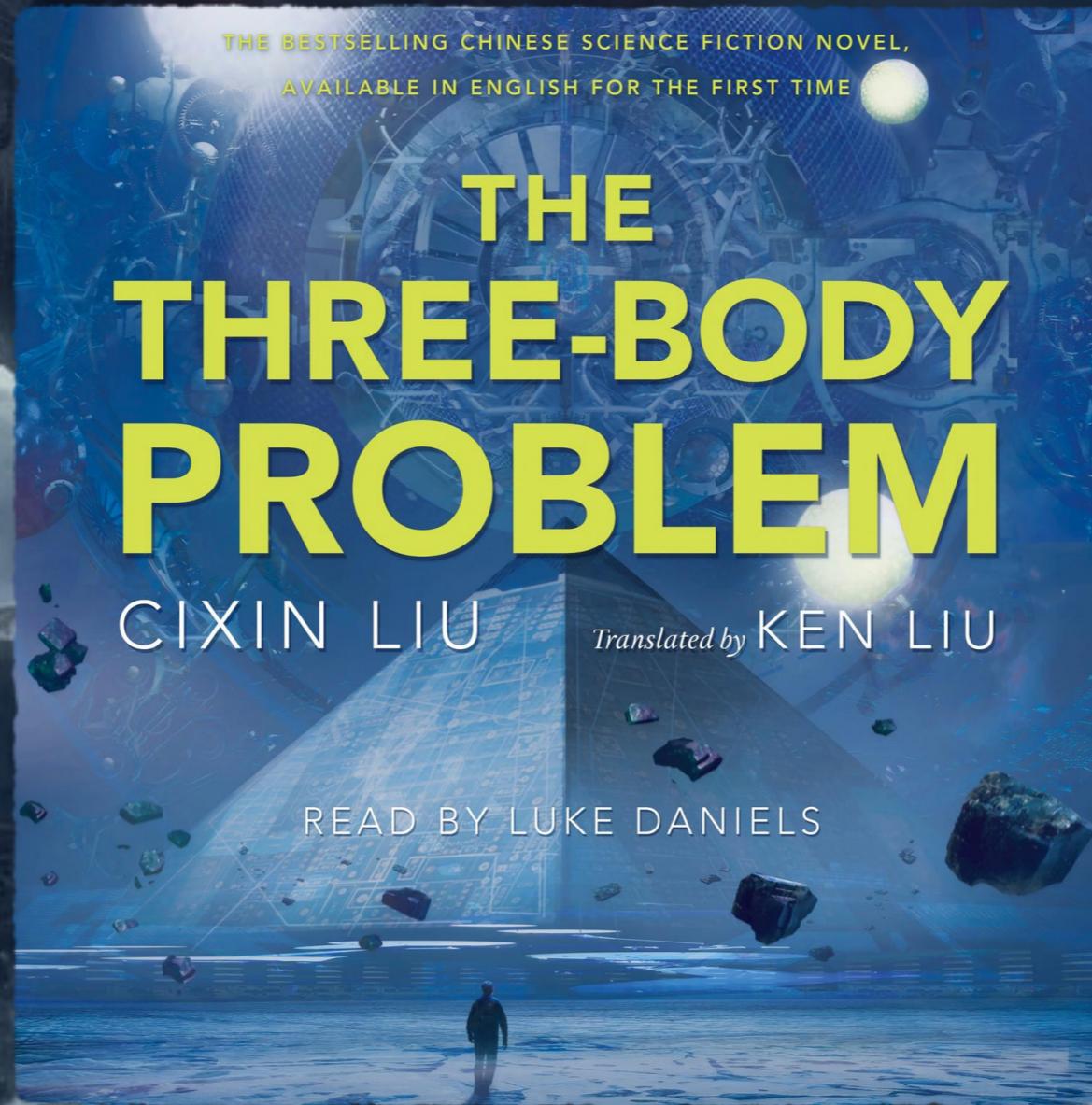
Formalism only strictly holds up to lowest  
3- or 4-particle threshold



• Wilson et al. PRL 123 (2019) •



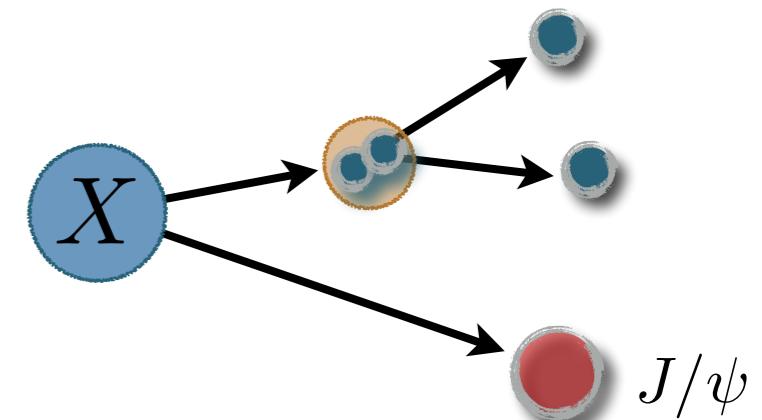
• Moir et al. JHEP 1610 (2016) •



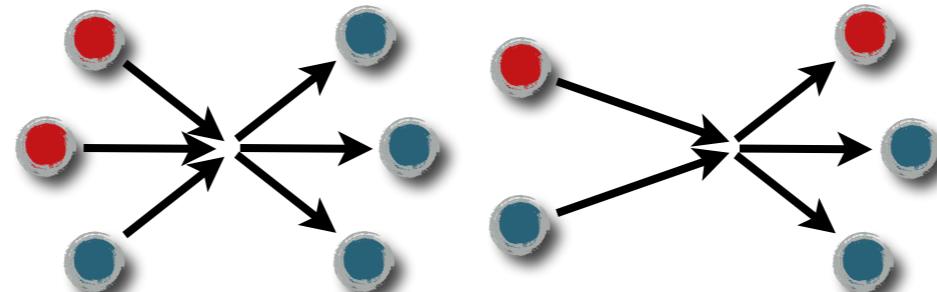
# 3-particle amplitudes

2-to-2 only samples  $J^P \ 0^+ \ 1^- \ 2^+ \dots$

many interesting resonances have significant 3-body decays



**Goal:** finite-volume + unitarity formalism for generic two- and three-particle systems



## Applications...

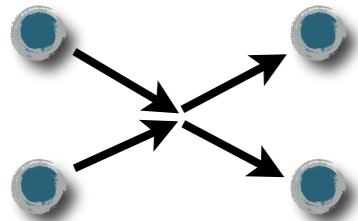
exotic resonance pole positions, couplings, quantum numbers

$\omega(782), a_1(1420) \rightarrow \pi\pi\pi$        $X(3872) \rightarrow J/\psi\pi\pi$        $X(3915)[Y(3940)] \rightarrow J/\psi\pi\pi$

form factors and transitions

and much more!... (3-body forces, weak transitions, gluons content)

# Complication: degrees of freedom

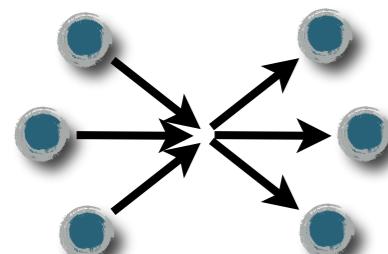


12 momentum components

-10 Poincaré generators

$$\vec{p}_1 + \vec{p}_2 \rightarrow \vec{p}_3 + \vec{p}_4 \longrightarrow \text{Mandelstam } s, t$$

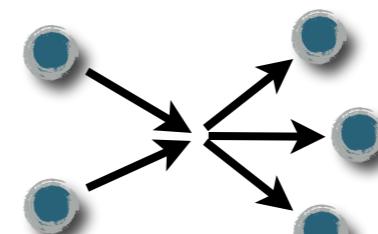
**2 degrees of freedom**



18 momentum components

-10 Poincaré generators

**8 degrees of freedom**



15 momentum components

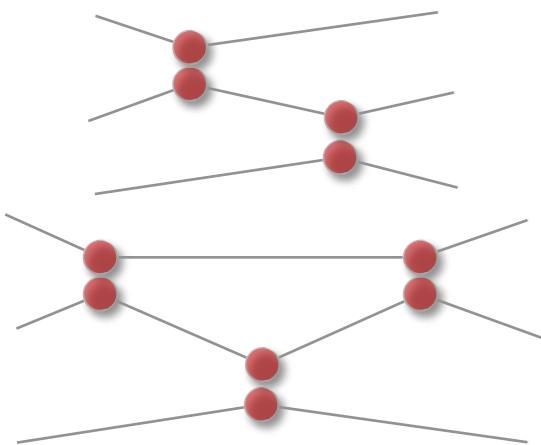
-10 Poincaré generators

**5 degrees of freedom**

# Complication: on-shell states

## □ Classical pairwise scattering

for  $m_1 = m_2 = m_3$  up to 3  
binary collisions are possible



### Dispersion Relations for Three-Particle Scattering Amplitudes. I\*

MORTON RUBIN

Physics Department, University of Wisconsin, Madison, Wisconsin

AND

ROBERT SUGAR

Physics Department, Columbia University, New York, New York

AND

GEORGE TIKTOPOULOS

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

(Received 31 January 1966)

$$b = \frac{(m_1+m_3)(m_2+m_3)}{m_1 m_2}$$

It follows that if

$$b^{n-2}(b-1) > 1, \quad (\text{IV.18})$$

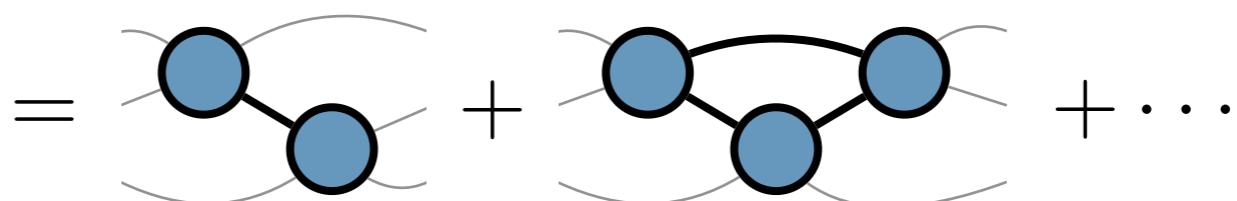
then  $2n+1$  successive binary collisions are kinematically impossible.

$m_1 = m_2 = m_3 - \varepsilon$ :  
4 collisions possible  
 $\pi\pi K$

$b < 2$   
5 collisions possible  
 $\pi K K$

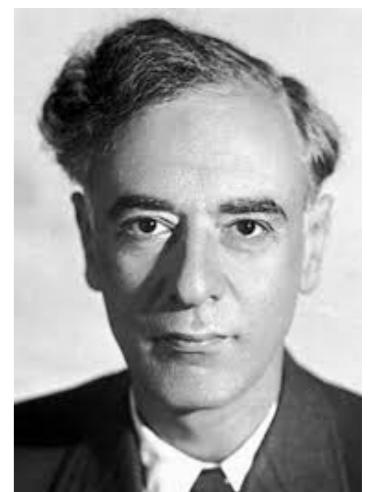
## □ Correspond to Landau singularities

$i\mathcal{M}_{3 \rightarrow 3} \equiv$  fully connected  
correlator



complicate analyticity & unitarity

difficult to disentangle kinematic  
singularities from resonance poles



# Two key observations

- Intermediate  $K_{df,3}$  removes singularities

$$i\mathcal{K}_{df,3} \equiv \begin{array}{l} \text{fully connected diagrams} \\ \text{w/ PV pole prescription} \end{array} - \text{---} + \text{---} + \dots$$

same degrees of freedom as  $M_3$

smooth real function

relation to  $M_3$  = known

- $K_{df,3}$  has a systematic low-energy expansion

$$\mathcal{K}_{df,3}(p_3, p_2, p_1; k_3, k_2, k_1) = \mathcal{K}_{df,3}^{\text{iso},0} + \mathcal{K}_{df,3}^{\text{iso},1} \Delta + \dots \quad \Delta = \frac{s - (3m)^2}{(3m)^2}$$

smooth real function

analogous to effective range expansion

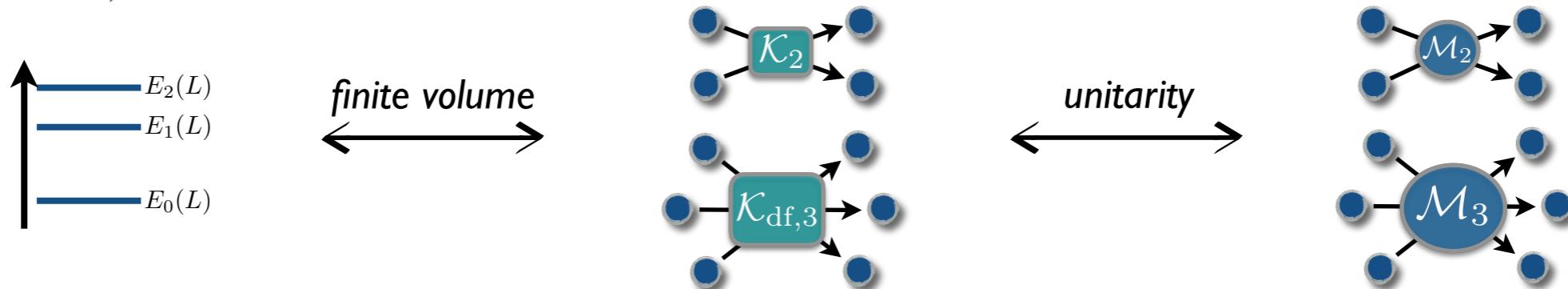
gives handle on many DOFs

$$p \cot \delta = -\frac{1}{a} + \frac{1}{2} r p^2 + \mathcal{O}(p^4)$$

# Status...

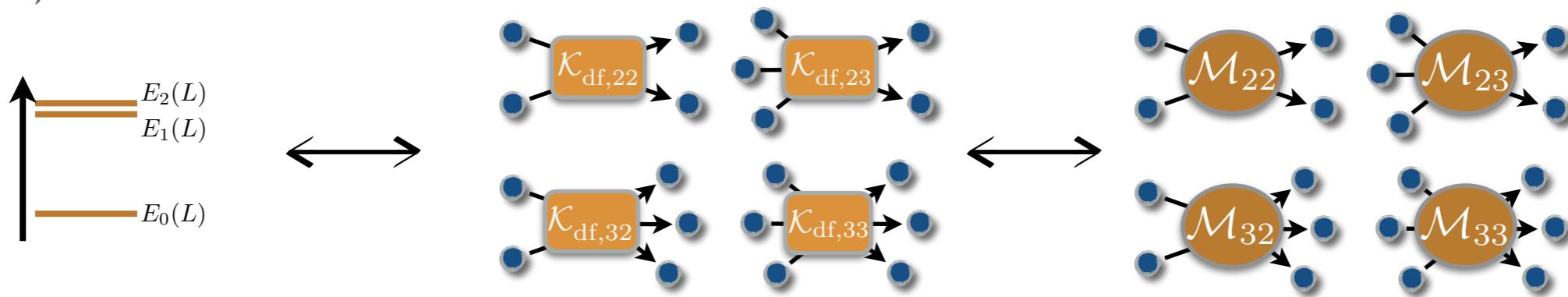
## □ General relation between *energies* and *two-and-three scalar scattering*

No 2-to-3, no sub-channel resonance



• MTH, Sharpe (2014, 2015) •

2-to-3, no sub-channel resonance



• Briceño, MTH, Sharpe (2017) •

Including sub-channel resonances + *different isospins*

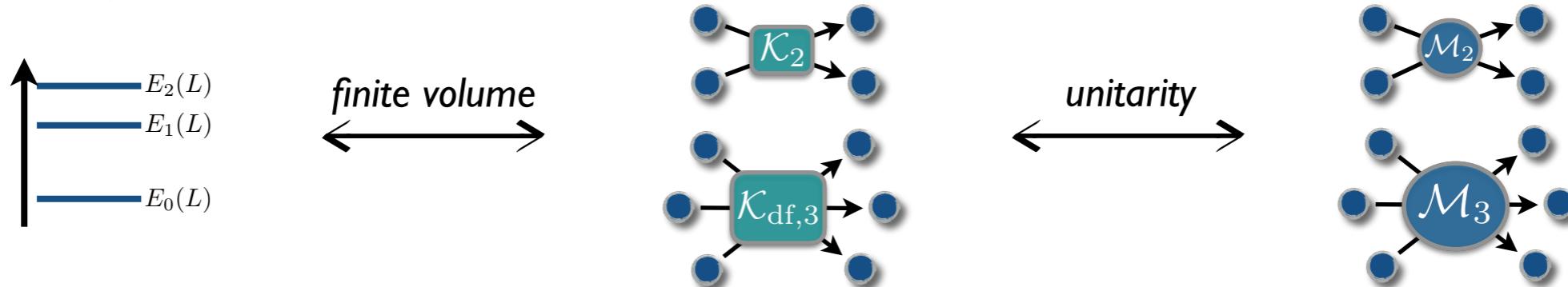
$$\pi\pi\pi \rightarrow \rho\pi \rightarrow \omega \rightarrow \rho\pi \rightarrow \pi\pi\pi$$

• Briceño, MTH, Sharpe (2018) • Blanton, Briceño, MTH, Romero-López, Sharpe (2019 + *in progress*) •

# Status...

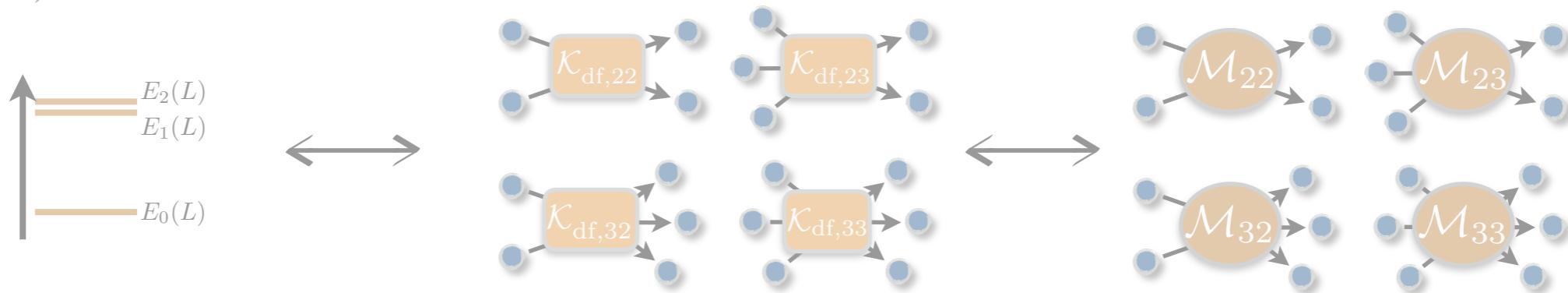
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# Derivation

- Study 3-body correlator in an *all-orders skeleton expansion*

$$\begin{aligned}
 C_L = & \square + \square + \square + \dots \\
 & + \square + \square + \square + \dots \\
 & + \square + \square + \square + \dots \\
 & + \square + \square + \square + \dots \\
 & + \dots \\
 & + \square + \square + \dots \quad \square = \sum_{\mathbf{k}}
 \end{aligned}$$

Here... assume no 2-to-3 allowed

$$\begin{aligned}
 \square &\equiv \star + \star + \star + \dots \\
 \square &\equiv \star + \star + \star + \star + \dots
 \end{aligned}$$

kernels have suppressed  $L$  dependence  
lines = fully dressed hadrons

## Two types of cuts

$$C_L \supseteq \text{Diagram} + \text{Diagram} + \text{Diagram} + \dots$$

$\square = \sum_{\mathbf{k}}$

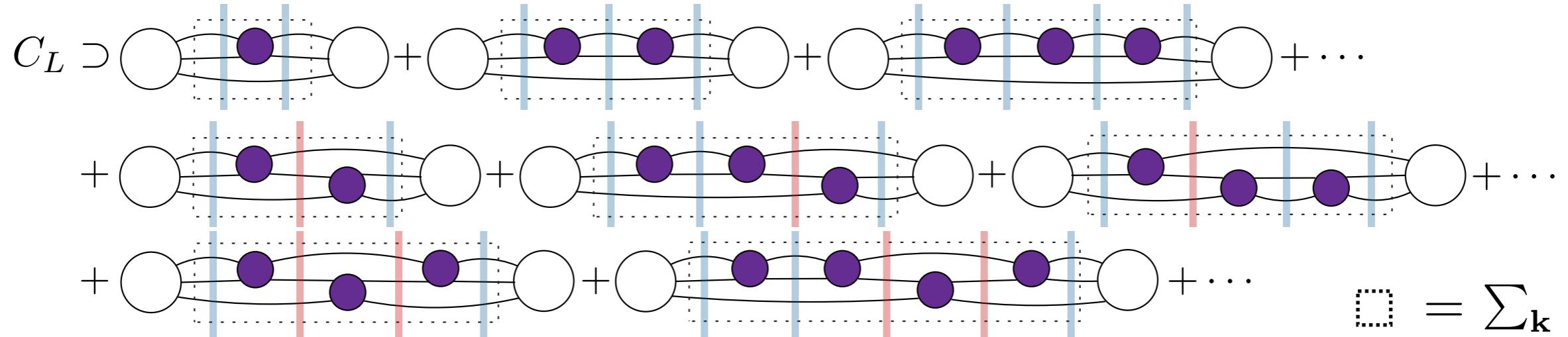
# Two types of cuts

$$\begin{aligned}
 C_L \supset & \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots \\
 + & \text{Diagram 4} + \text{Diagram 5} + \text{Diagram 6} + \dots \\
 + & \text{Diagram 7} + \text{Diagram 8} + \text{Diagram 9} + \dots
 \end{aligned}$$

$\square = \sum_{\mathbf{k}}$

$$\mathbf{A}'_3 \mathbf{F} \mathbf{K}_2 \mathbf{F} \mathbf{A}_3 + \mathbf{A}'_3 \mathbf{F} [\mathbf{K}_2 \mathbf{F}]^2 \mathbf{A}_3 + \mathbf{A}'_3 \mathbf{F} [\mathbf{K}_2 \mathbf{F}]^3 \mathbf{A}_3 + \dots = \mathbf{A}'_3 \mathbf{F} \frac{1}{1 - \mathbf{K}_2 \mathbf{F}} \mathbf{K}_2 \mathbf{F} \mathbf{A}_3$$

# Two types of cuts



$$\mathbf{A}'_3 \mathbf{F} \mathbf{K}_2 \mathbf{F} \mathbf{A}_3 + \mathbf{A}'_3 \mathbf{F} [\mathbf{K}_2 \mathbf{F}]^2 \mathbf{A}_3 + \mathbf{A}'_3 \mathbf{F} [\mathbf{K}_2 \mathbf{F}]^3 \mathbf{A}_3 + \dots = \mathbf{A}'_3 \mathbf{F} \frac{1}{1 - \mathbf{K}_2 \mathbf{F}} \mathbf{K}_2 \mathbf{F} \mathbf{A}_3$$

$$\mathbf{A}'_3 \mathbf{F} \mathbf{K}_2 \mathbf{G} \mathbf{K}_2 \mathbf{F} \mathbf{A}_3 + \dots$$

$$C_L - C_\infty \supset \mathbf{A}'_3 \mathbf{F}_{33} \mathbf{A}_3$$

$$\mathbf{F}_{33} \equiv \frac{1}{3} \mathbf{F} + \mathbf{F} \mathbf{K}_2 \frac{1}{1 - (\mathbf{F} + \mathbf{G}) \mathbf{K}_2} \mathbf{F}$$

single  $\mathbf{F}$  requires special treatment (1/3)

*have not yet considered entire diagram contributions*

missing contributions from off-shellness

missing smooth terms (short-distance parts)

# Short-distance parts & summation

$$C_L = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

+ ...

$$+ \text{Diagram 4} + \text{Diagram 5} + \text{Diagram 6} + \dots$$

+ ...

$$+ \text{Diagram 7} + \text{Diagram 8} + \text{Diagram 9} + \dots$$

+ ...

$$+ \text{Diagram 10} + \text{Diagram 11} + \text{Diagram 12} + \dots$$

+ ...

$$+ \text{Diagram 13} + \text{Diagram 14} + \text{Diagram 15} + \dots$$

$$C_L - C_\infty = \mathbf{A}'_3 \mathbf{F}_{33} \mathbf{A}_3 + \mathbf{A}'_3 \mathbf{F}_{33} \mathbf{K}_{\text{df},3} \mathbf{F}_{33} \mathbf{A}_3 + \cdots$$

$$= \mathbf{A}_3' \frac{1}{\mathbf{F}_{33}^{-1} + \mathbf{K}_{df,3}} \mathbf{A}_3$$

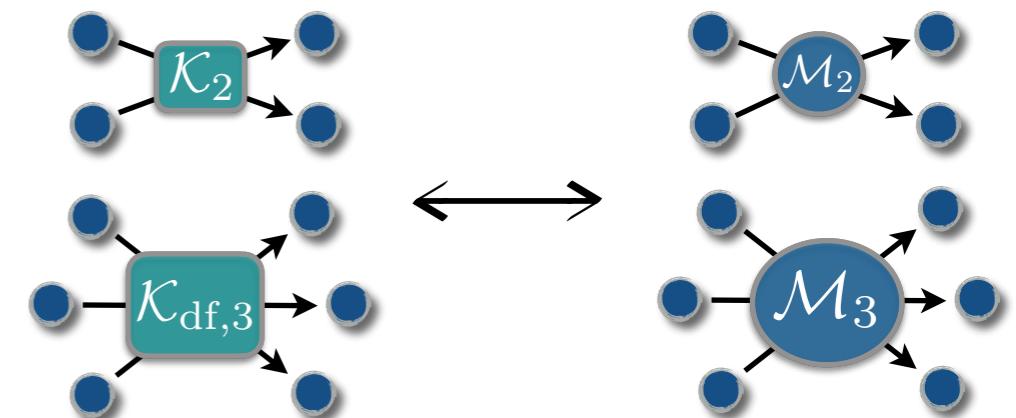
$$\mathbf{F}_{33} \equiv \frac{1}{3}\mathbf{F} + \mathbf{F} \mathbf{K}_2 \frac{1}{1 - (\mathbf{F} + \mathbf{G}) \mathbf{K}_2} \mathbf{F}$$

*no term left behind*

# Relating $\mathcal{K}_{\text{df},3}$ to $\mathcal{M}_3$

## □ Define a new finite-volume correlator

- I. Amputate interpolating fields
2. Drop disconnected diagrams
3. Symmetrize



$$\mathcal{M}_{L,3} \equiv \mathcal{S} \left\{ \begin{array}{c} \text{Diagram with one orange circle} + \text{Diagram with two orange circles connected by a dashed box} + \dots \\ + \text{Diagram with three purple circles connected by a dashed box} + \text{Diagram with four purple circles connected by a dashed box} + \text{Diagram with five purple circles connected by a dashed box} + \dots \\ + \text{Diagram with six purple circles connected by a dashed box} + \text{Diagram with seven purple circles connected by a dashed box} + \dots \\ + \dots \\ + \text{Diagram with one purple circle connected to one orange circle} + \text{Diagram with one orange circle connected to one purple circle} + \dots \end{array} \right\} = \mathcal{F}_L(\mathcal{K}_2, \mathcal{K}_{\text{df},3})$$

*known functional of  $\mathcal{K}_{\text{df},3}$*

$$\mathcal{M}_3 = \lim_{L \rightarrow \infty} \mathcal{F}_L(\mathcal{K}_2, \mathcal{K}_{\text{df},3})$$

- MTH, Sharpe (2015) •

Integral equations, all inputs known

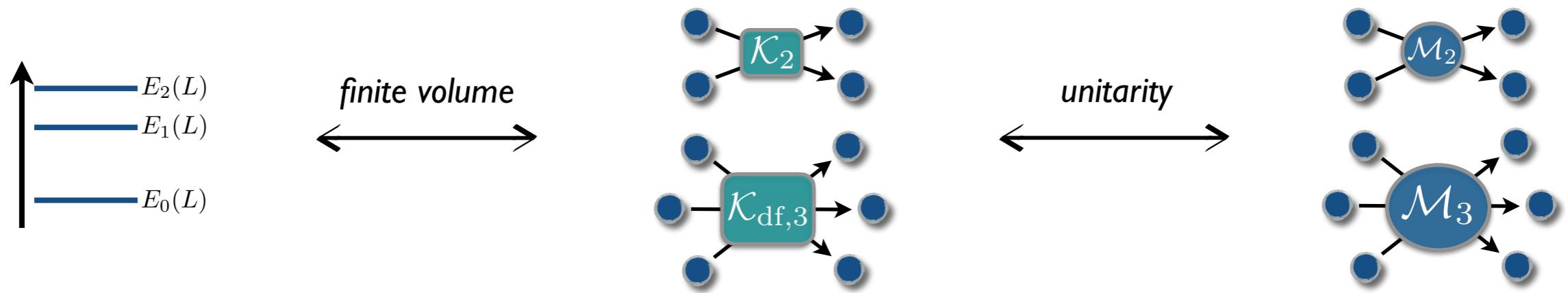
*Exact unitarity solution*

- Briceño, MTH, Sharpe, Szczepaniak (2019) •

# Result

$$\det[\mathcal{K}_{\text{df},3}(s)^{-1} + F_3(\mathcal{M}_2, P, L)] = 0$$

$$\mathcal{M}_3 = \lim_{L \rightarrow \infty} \mathcal{F}_L(\mathcal{K}_2, \mathcal{K}_{\text{df},3})$$

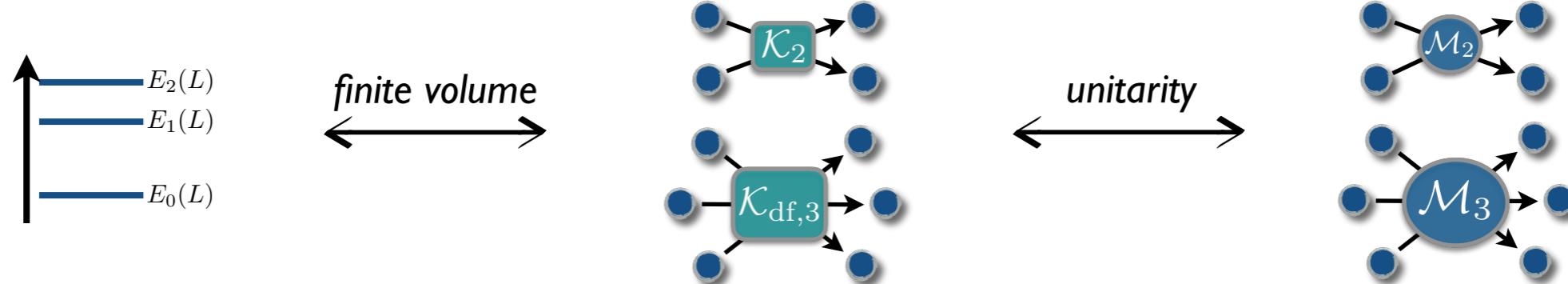


Holds only for three-particle energies  $s < (5m)^2$

Neglects  $e^{-mL}$

- Huang, Yang (1958) • Lüscher (1986, 1991) • Rummukainen, Gottlieb (1995)  
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Leskovec, Prelovsek (2012) • Bernard *et. al.* (2012) • MTH, Sharpe (2012) • Briceño, Davoudi (2012)  
Li, Liu (2013) • Briceño (2014)

# Status...



- Identical spin-zero, no 2-to-3, no K2 poles
  - MTH, Sharpe (2014, 2015) •
- as above... but including 2-to-3
  - Briceño, MTH, Sharpe (2017) •
- as above... but including K2 poles
  - Briceño, MTH, Sharpe (2018) •
- Non-identical, non-degenerate spin-zero
  - MTH, Romero-López, Sharpe (2020) • Blanton, Sharpe (2020, 2021)
- Multiple three-particle channels... Spin!

# Related work

## □ Finite-volume unitarity method

Döring, Mai (2016,2017)

Gives connection to unitarity relations

## □ Non-relativistic EFT method

Hammer, Pang, Rusetsky (2017)

Simplified derivation + integral equations

Do not yet include non-degenerate, 2-to-3

## □ All methods

Rely on intermediate, scheme-dependent quantity

Hold up to  $e^{-mL}$  and for  $E_3^\star < 5m_\pi$

Equivalent where comparable

## □ Review articles

MTH and Sharpe, 1901.00483 • Rusetsky, 1911.01253 • Mai, Döring, Rusetsky, 2103.00577

# Not covered here

## □ Activity extracting and fitting three-hadron energies

- Hörz, Hanlon (2019) • Blanton, Romero-López, Sharpe (2019) •
  - Alexandru, Brett, Culver, Döring, Guo, Lee, Mai (2019) •
  - Fischer, Kostrzewa, Liu, Romero-López, Ueding, Urbach (2020) •
  - Mai, Alexandru, Brett, Culver, Döring, Lee, Sadasivan (2021) •
- Blanton, Hanlon, Hörz, Morningstar, Romero-López, Sharpe (2021) •

## □ Activity connecting and extending formalisms

Relating infinite-volume equations

- Jackura *et al.* (2019) •

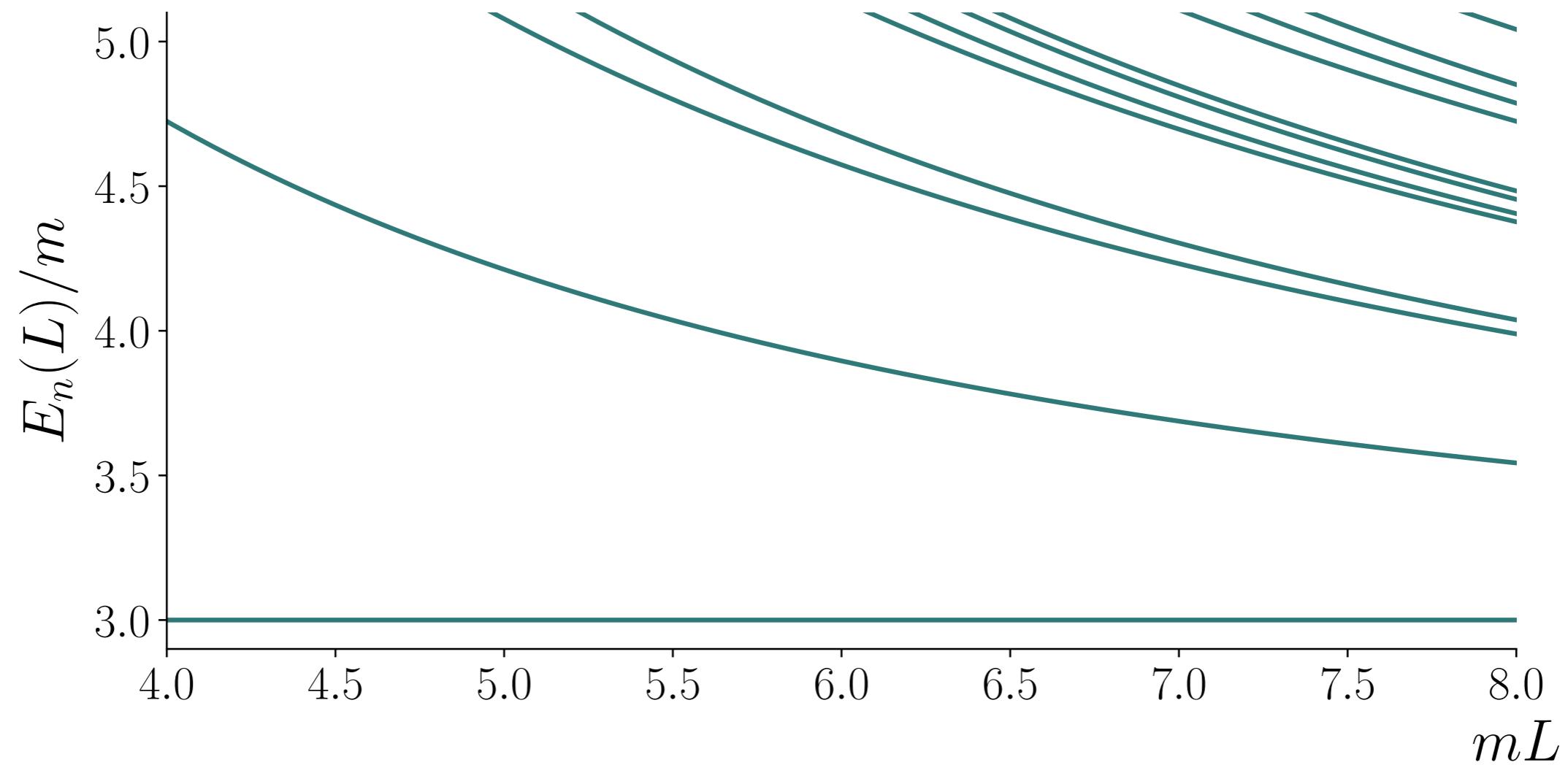
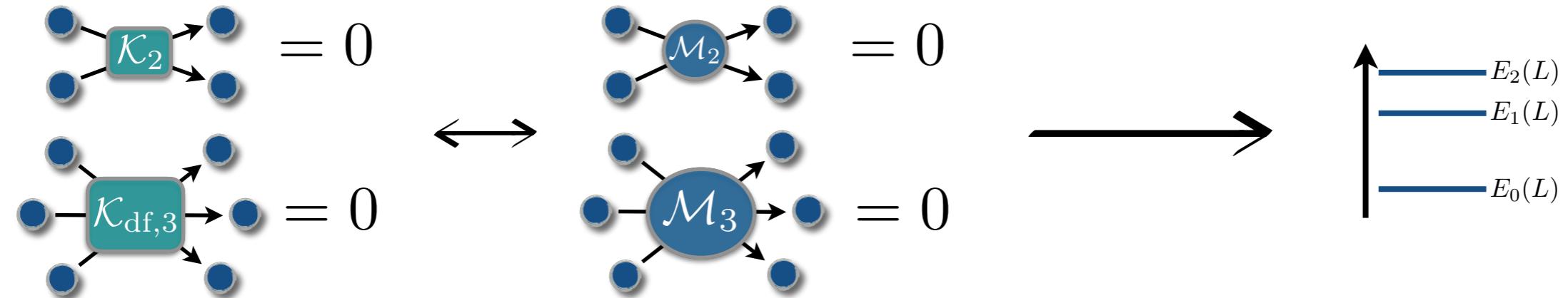
Alternative derivations

- Blanton, Sharpe (2020) •

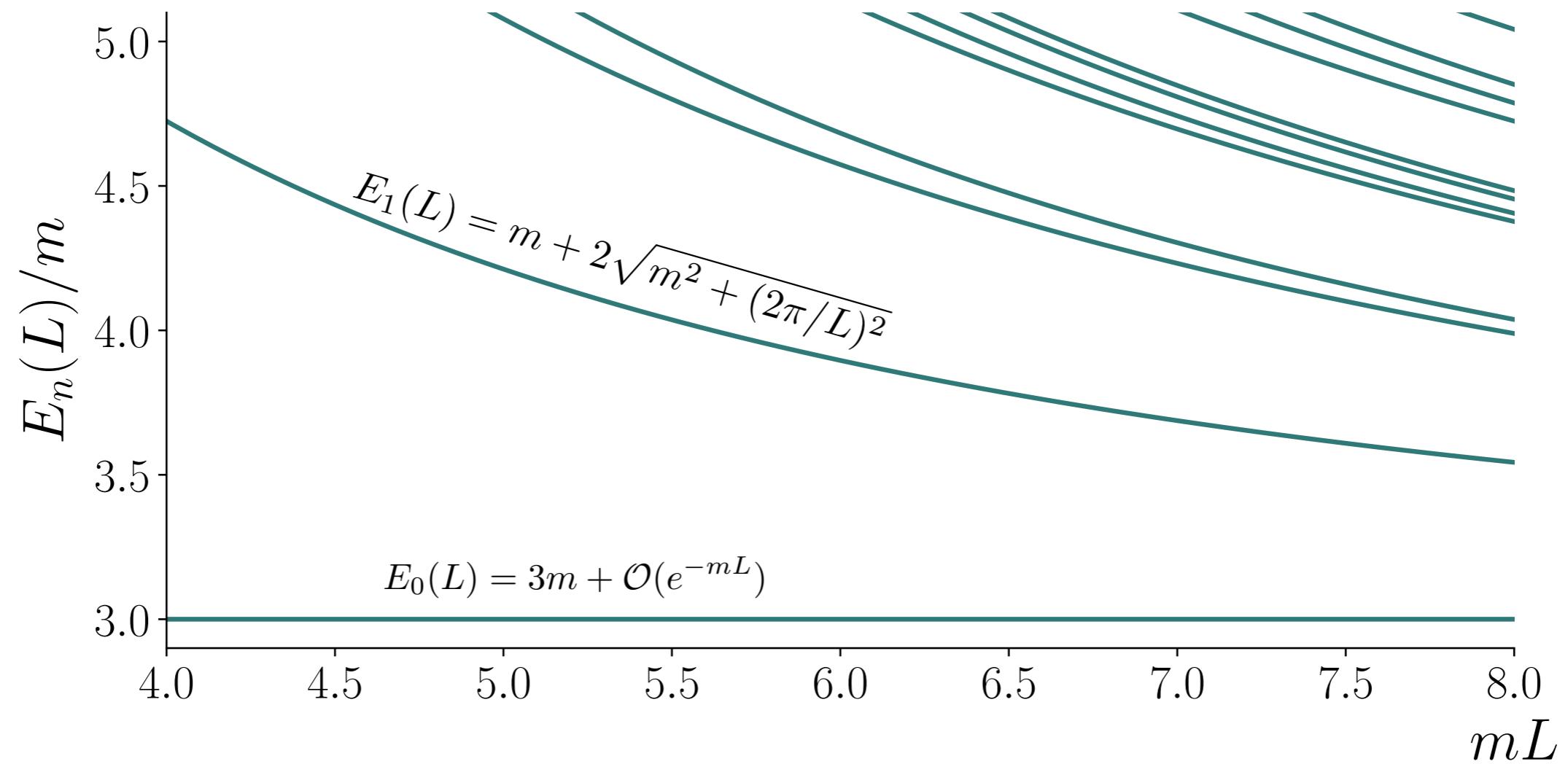
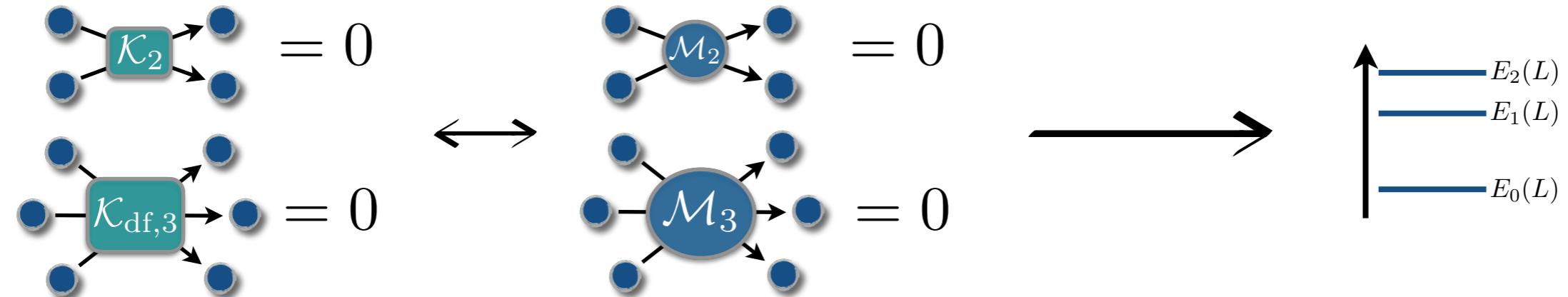
Equivalence of formalisms (where comparable)

- Blanton, Sharpe (2020) •

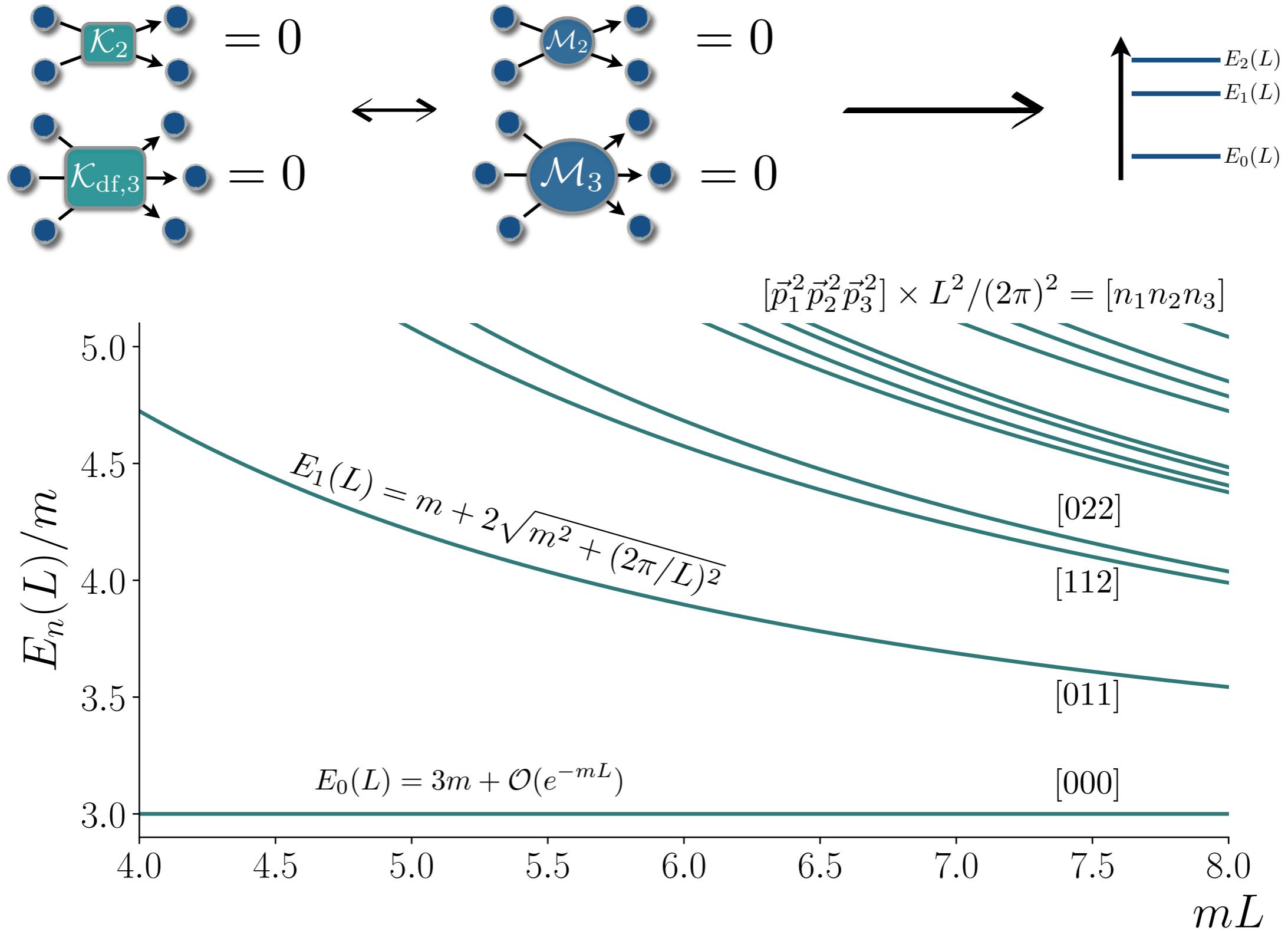
# Non-interacting energies



# Non-interacting energies



# Non-interacting energies



# Two-particle interactions

$\mathcal{K}_2$

$$= -16\pi\sqrt{s} a$$

$\mathcal{K}_{\text{df},3}$

$$= 0$$

$\mathcal{M}_2$

$$= \frac{16\pi\sqrt{s}}{-1/a - ip}$$

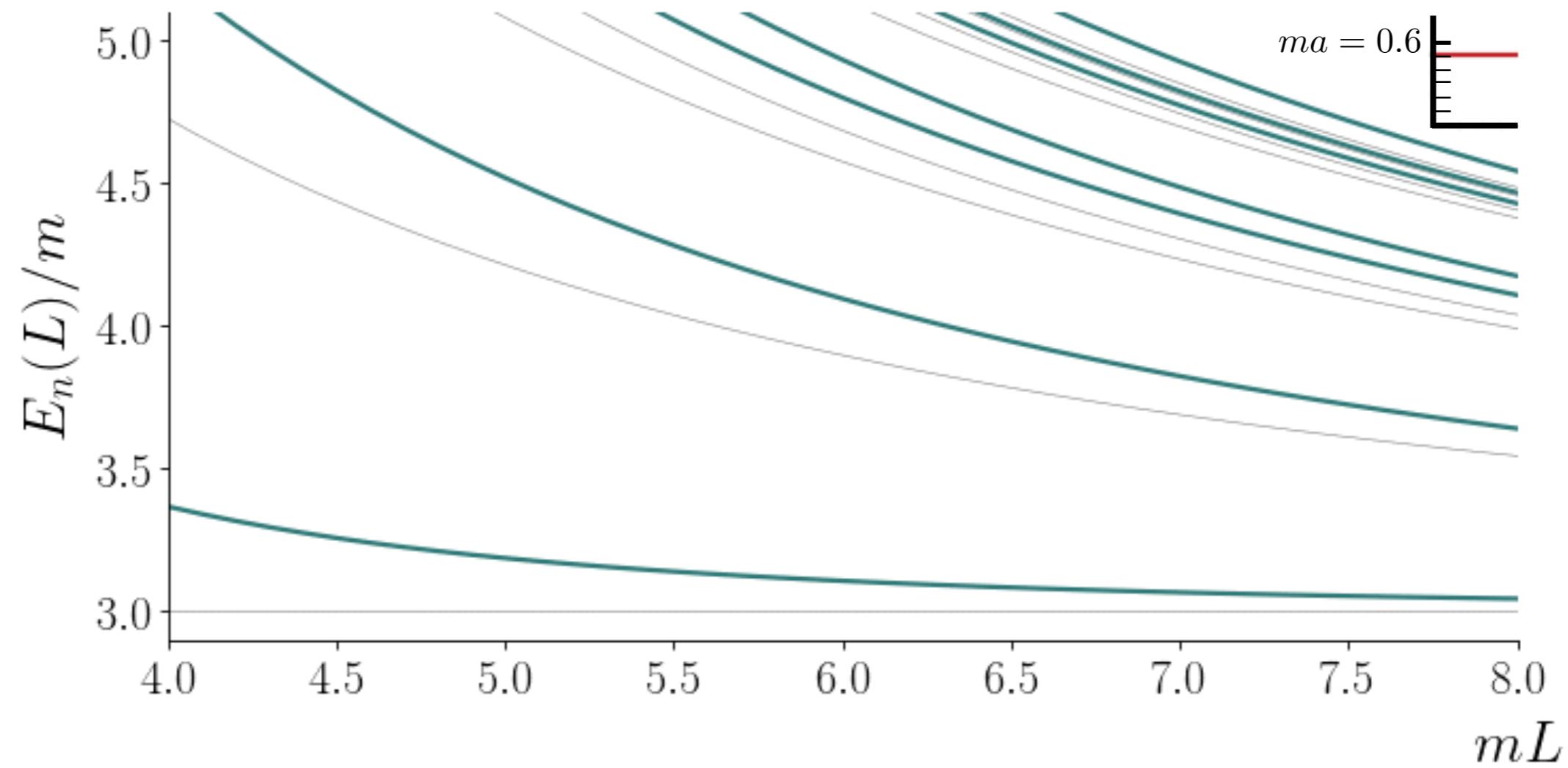
$\mathcal{M}_3$

$$= i\mathcal{M}_2 + i\mathcal{M}_2 + \dots$$

$E_2(L)$

$E_1(L)$

$E_0(L)$

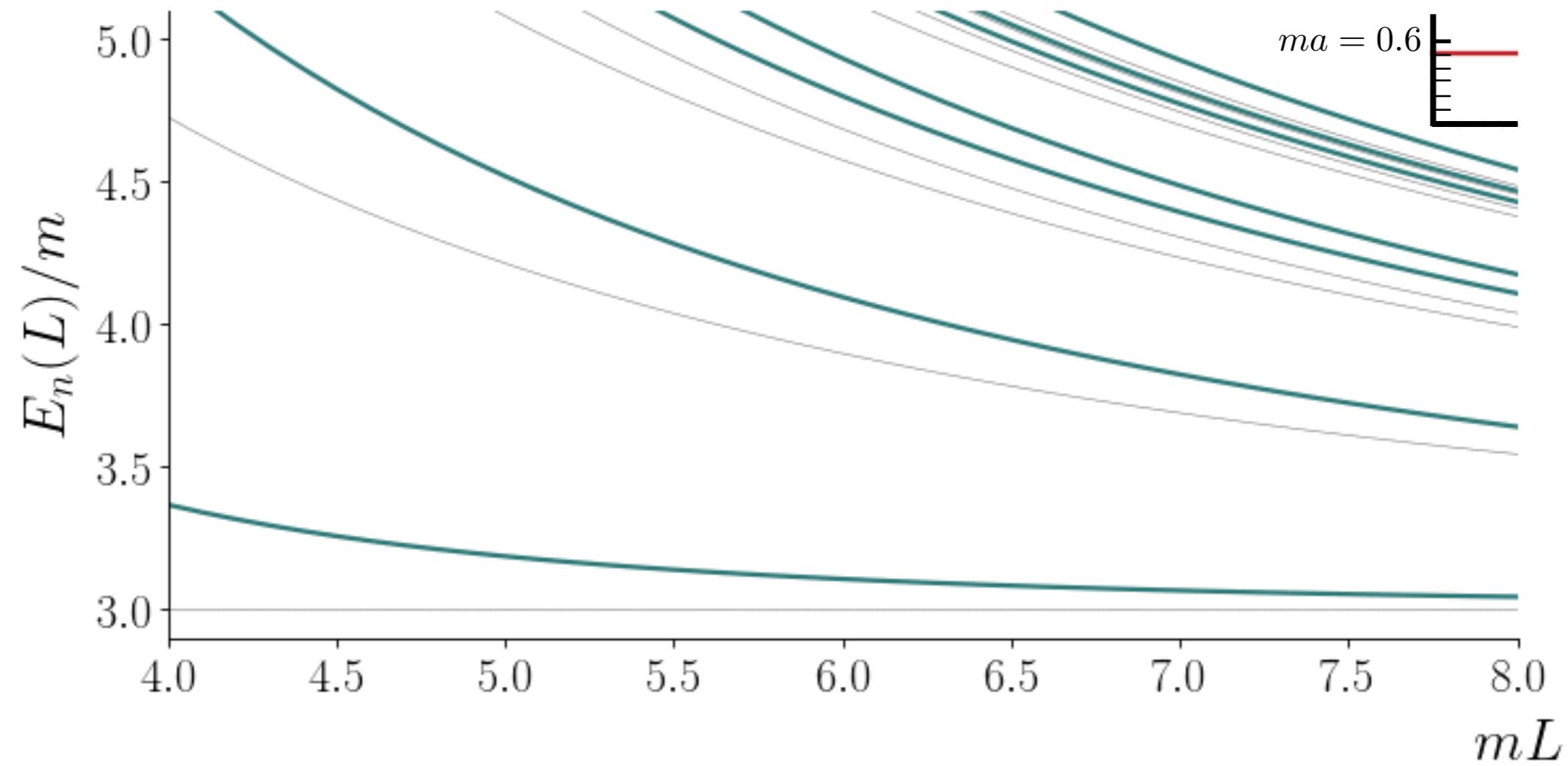


# Two-particle interactions

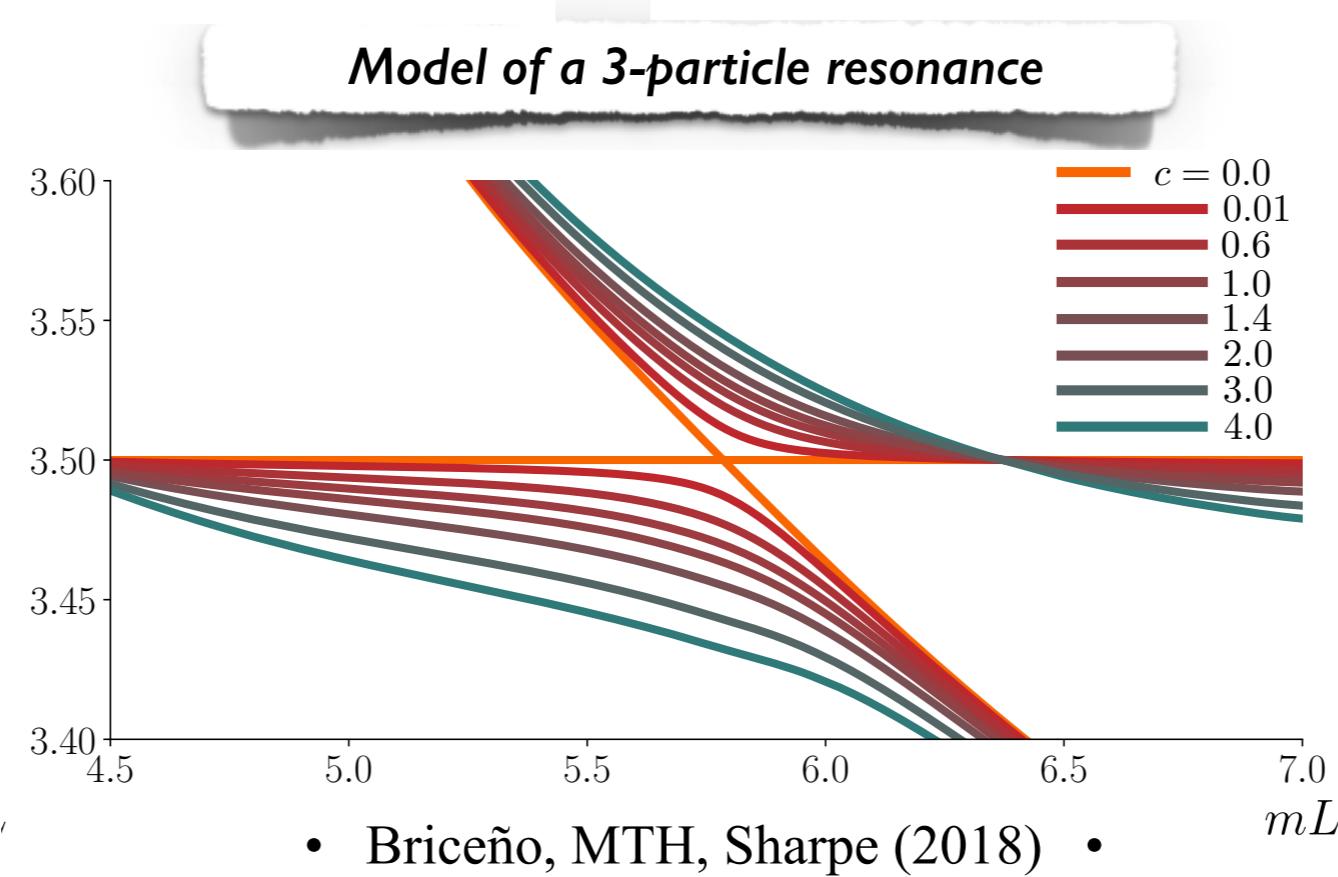
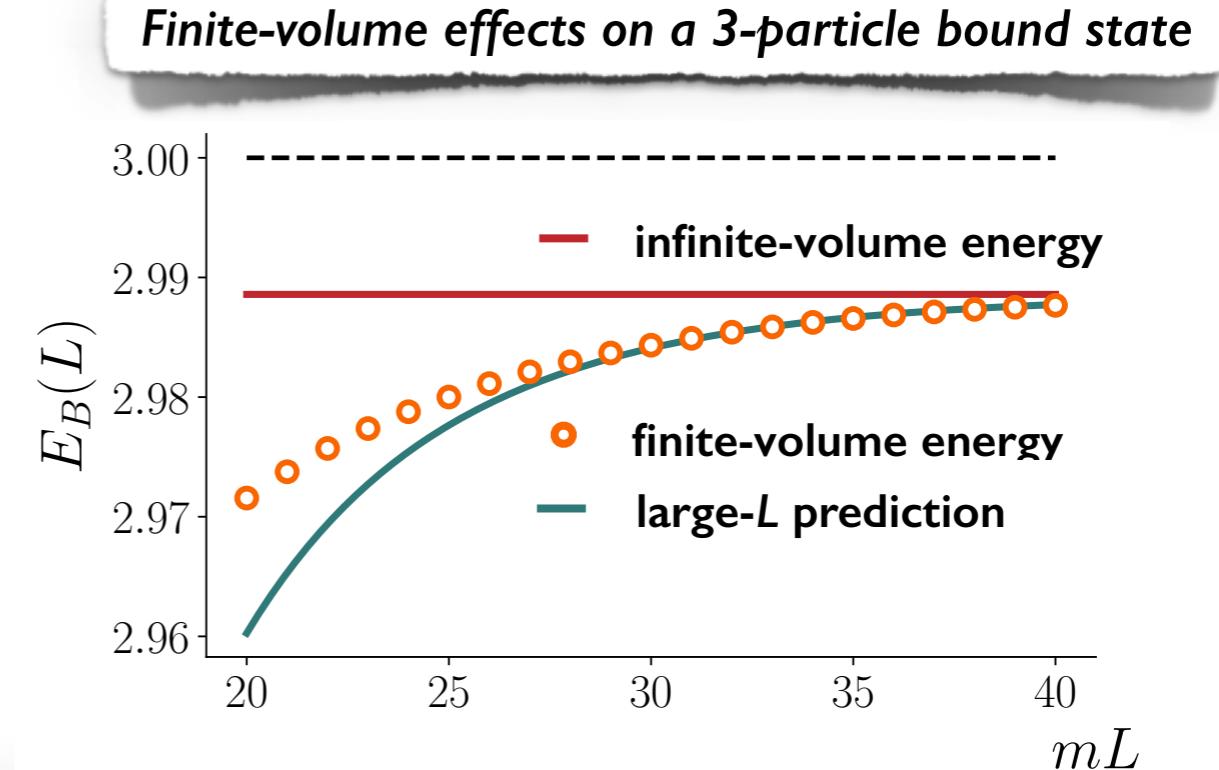
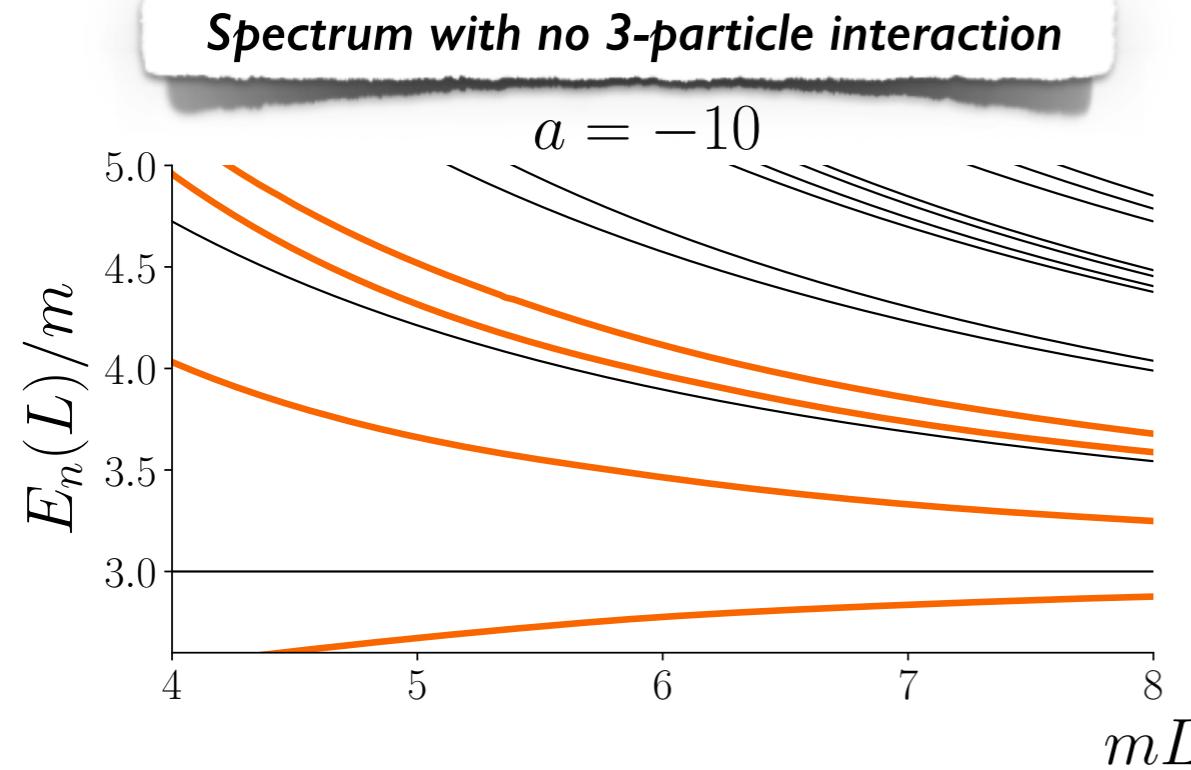
$$\begin{aligned} \text{Diagram with } \mathcal{K}_2 &= -16\pi\sqrt{s} a \\ \text{Diagram with } \mathcal{K}_{\text{df},3} &= 0 \\ \text{Diagram with } \mathcal{M}_2 &= \frac{16\pi\sqrt{s}}{-1/a - ip} \\ \text{Diagram with } \mathcal{M}_3 &= i\mathcal{M}_2 + i\mathcal{M}_2 + \dots \end{aligned}$$

→

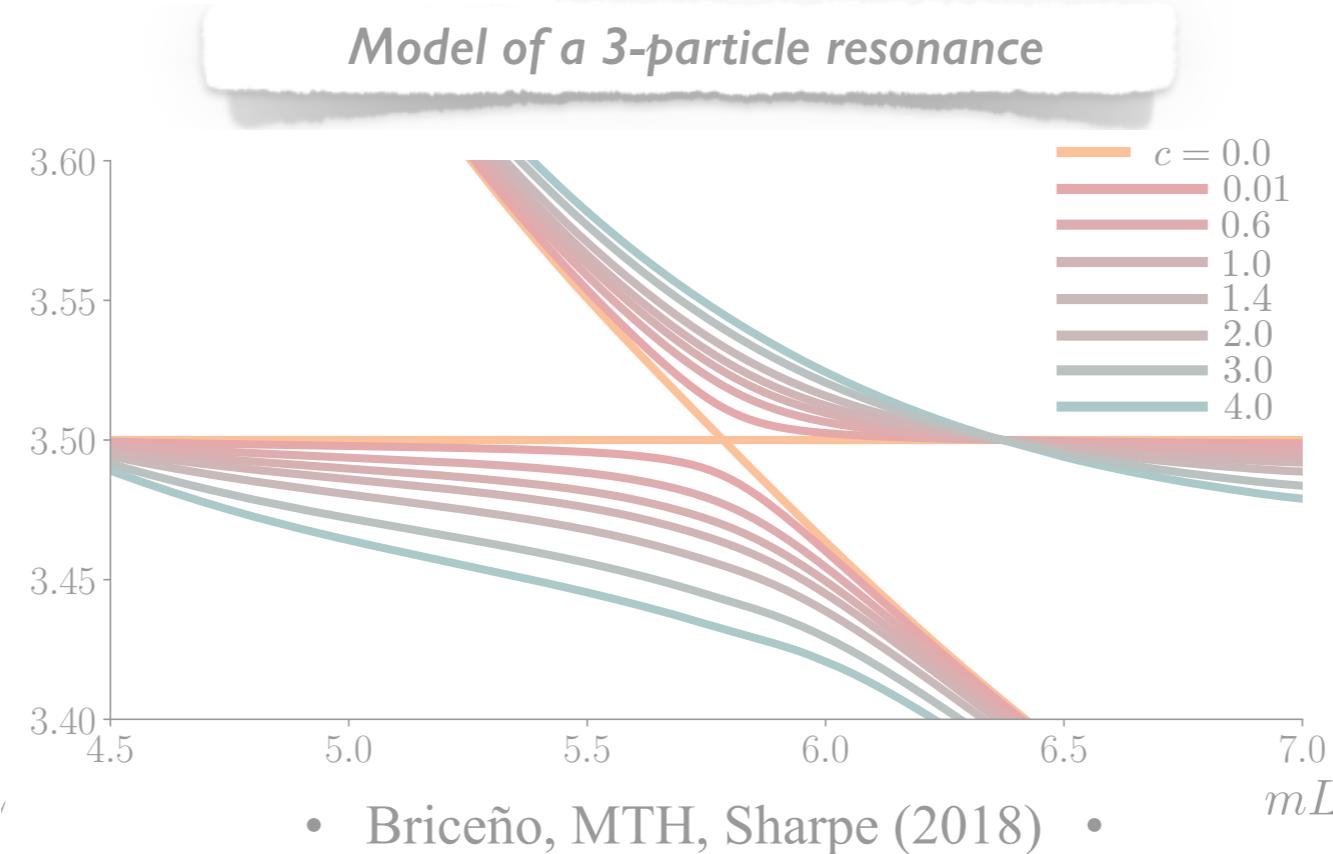
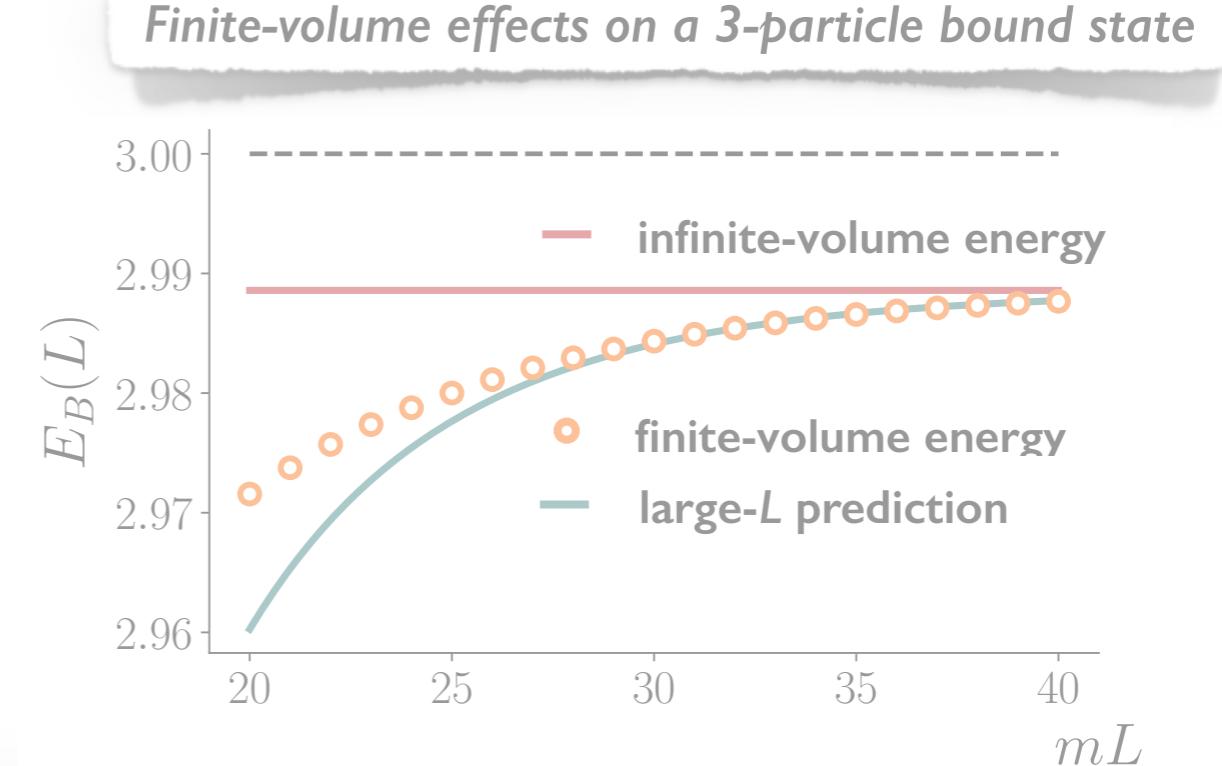
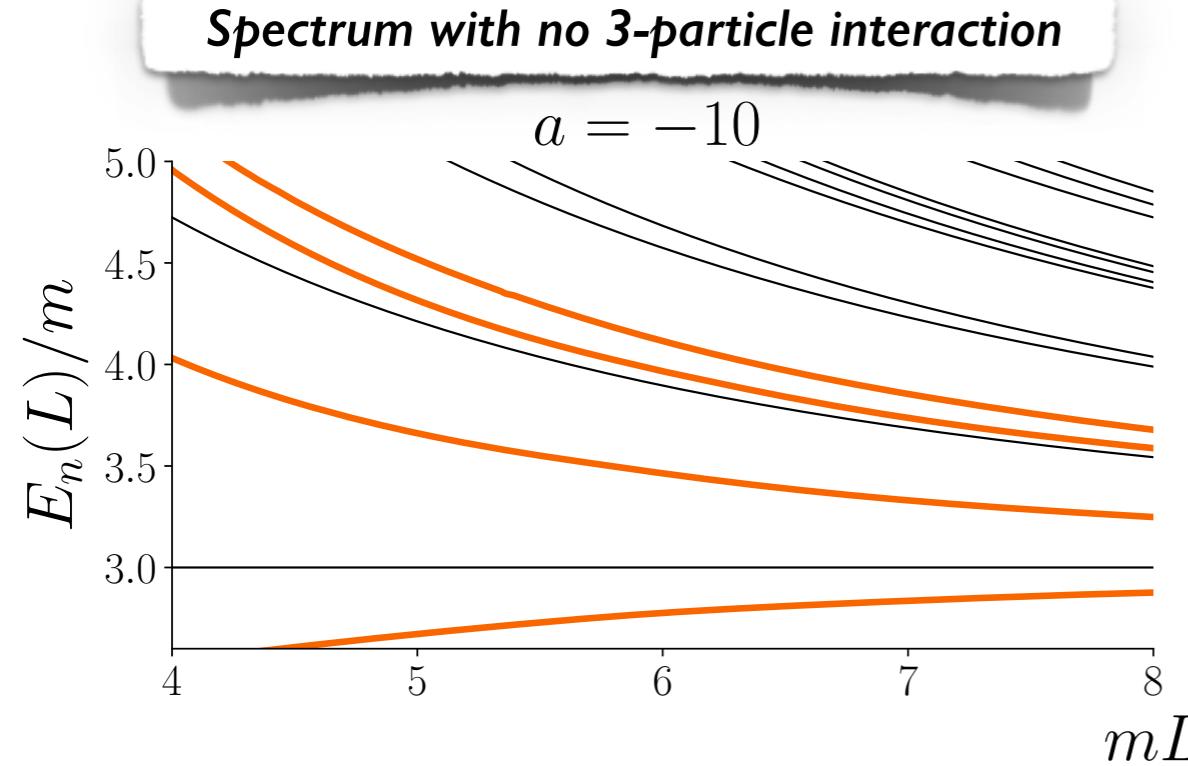
$E_2(L)$   
 $E_1(L)$   
 $E_0(L)$



# Checks and explorations

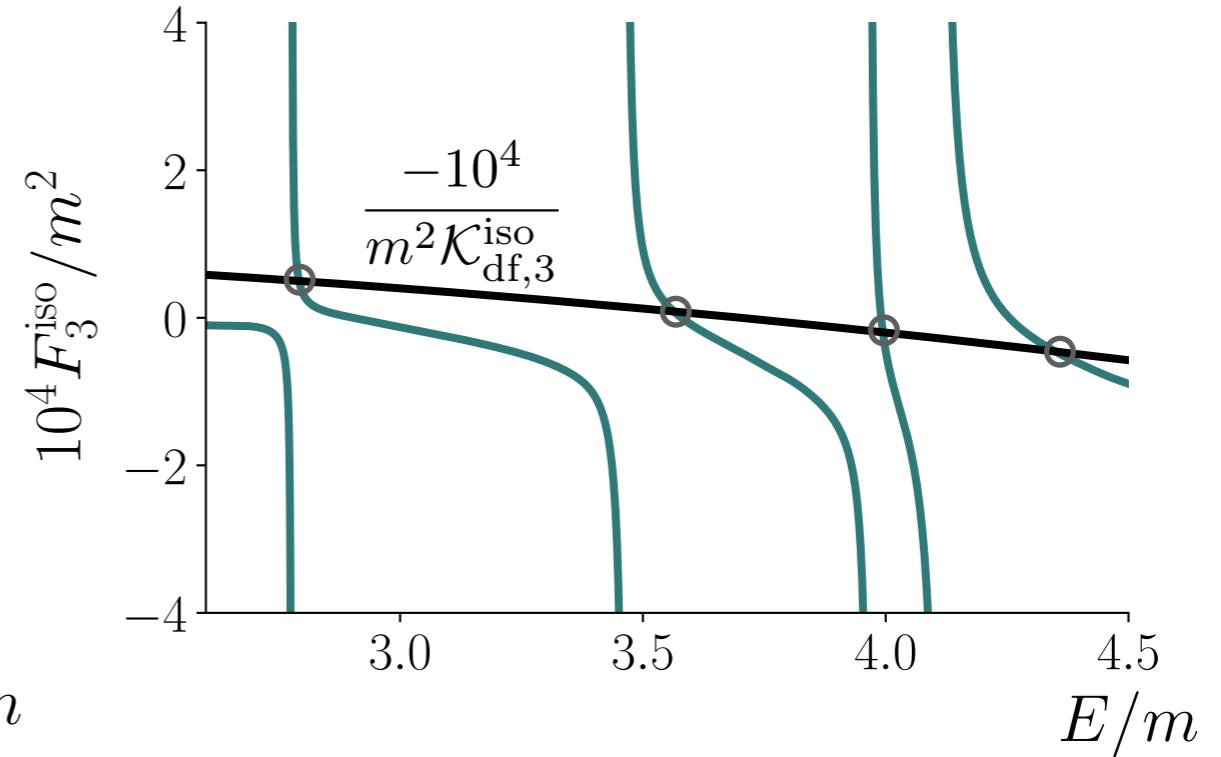
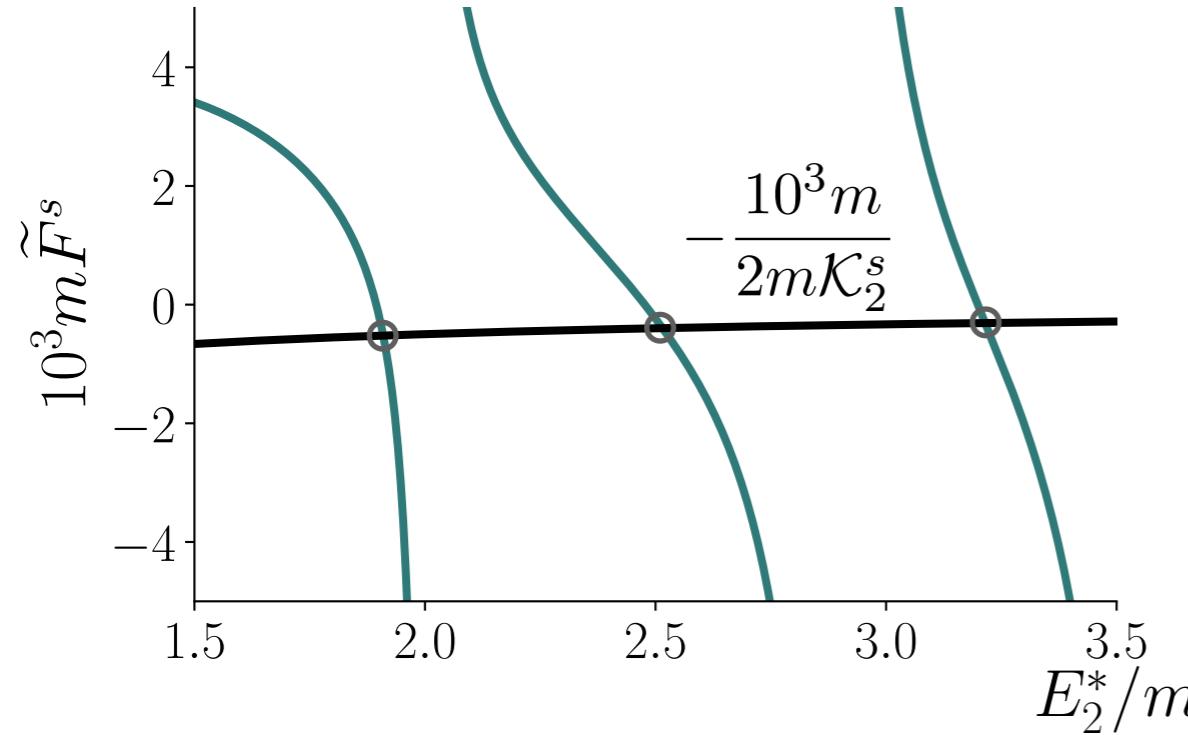


# Checks and explorations

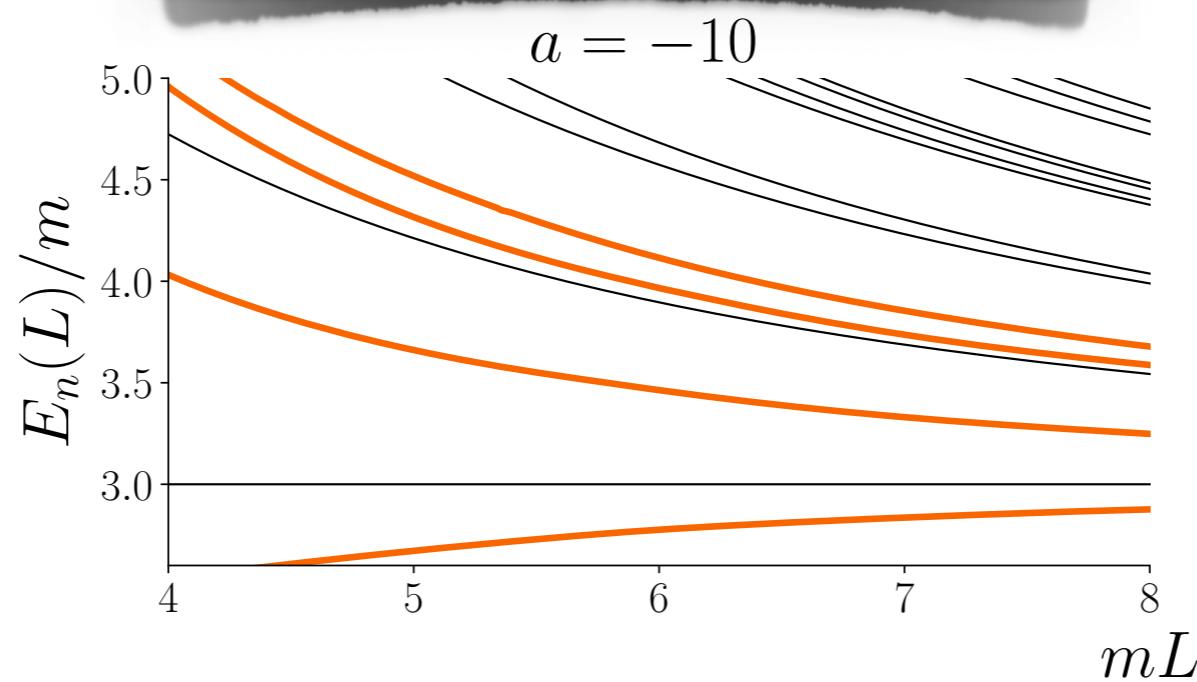


# Attractive scattering

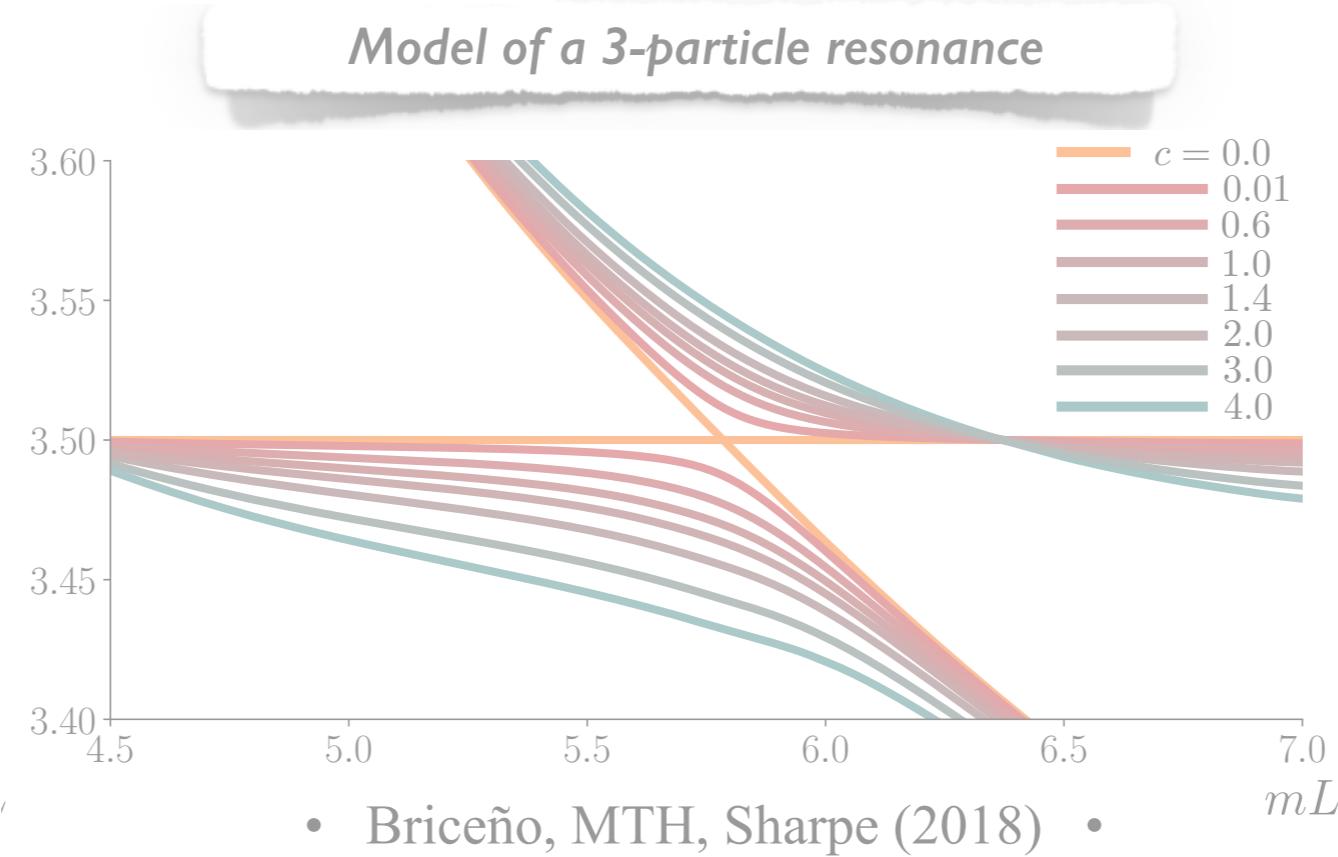
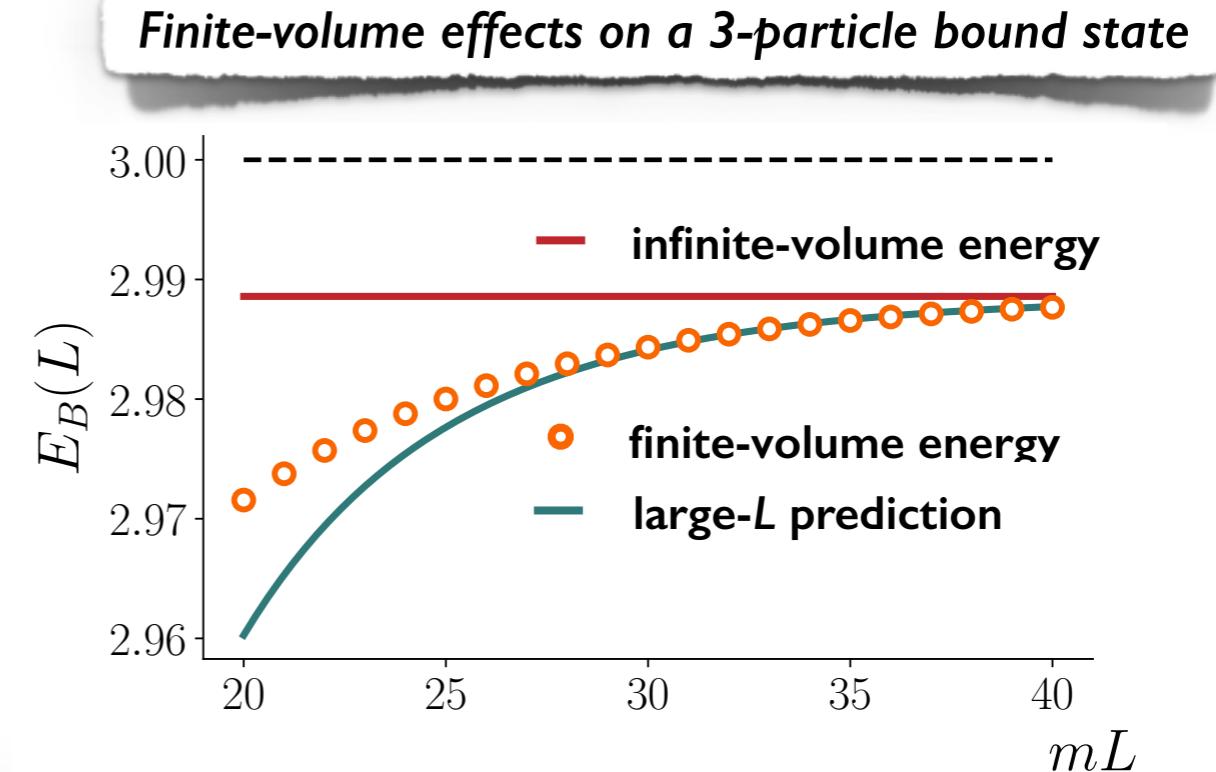
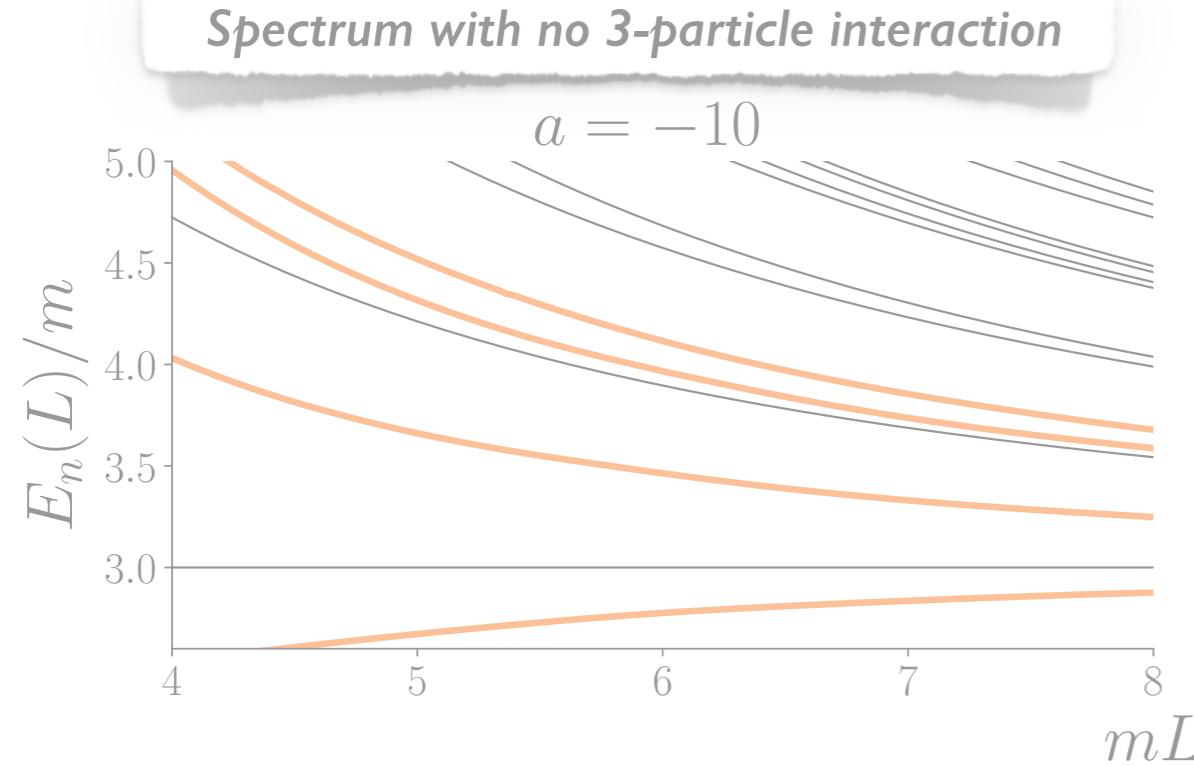
$$ma = -10, \quad mL = 6$$



**Spectrum with no 3-particle interaction**



# Checks and explorations



# Unitary bound state

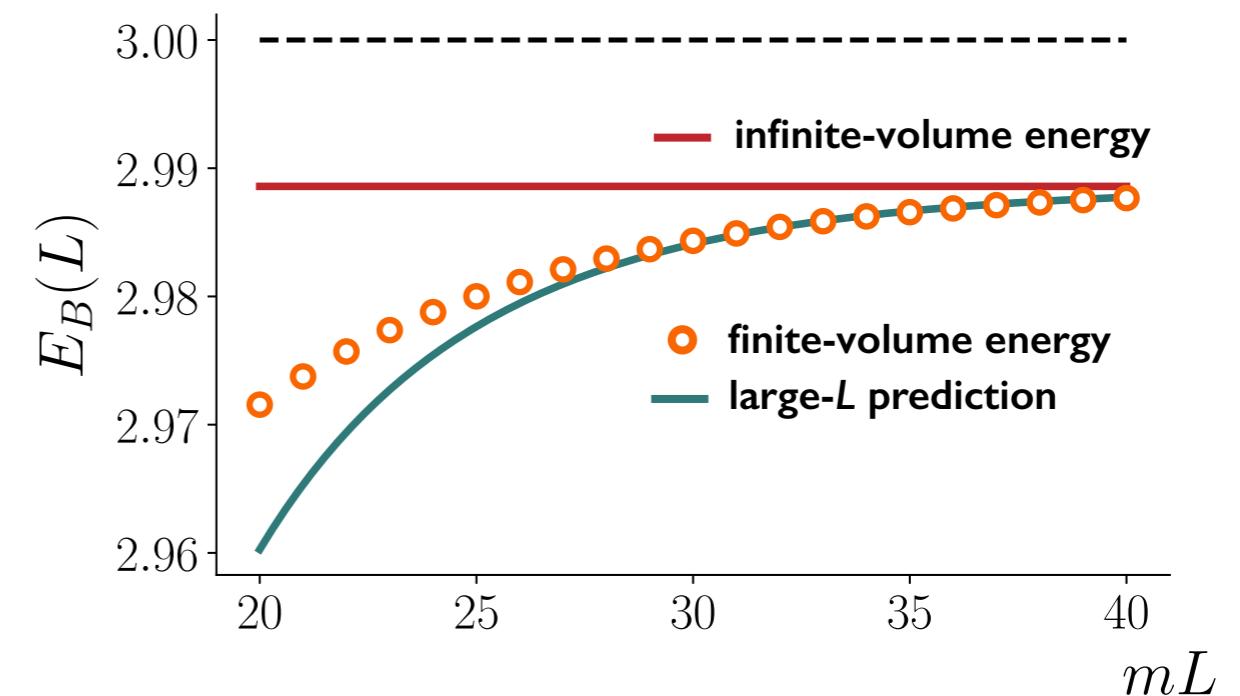
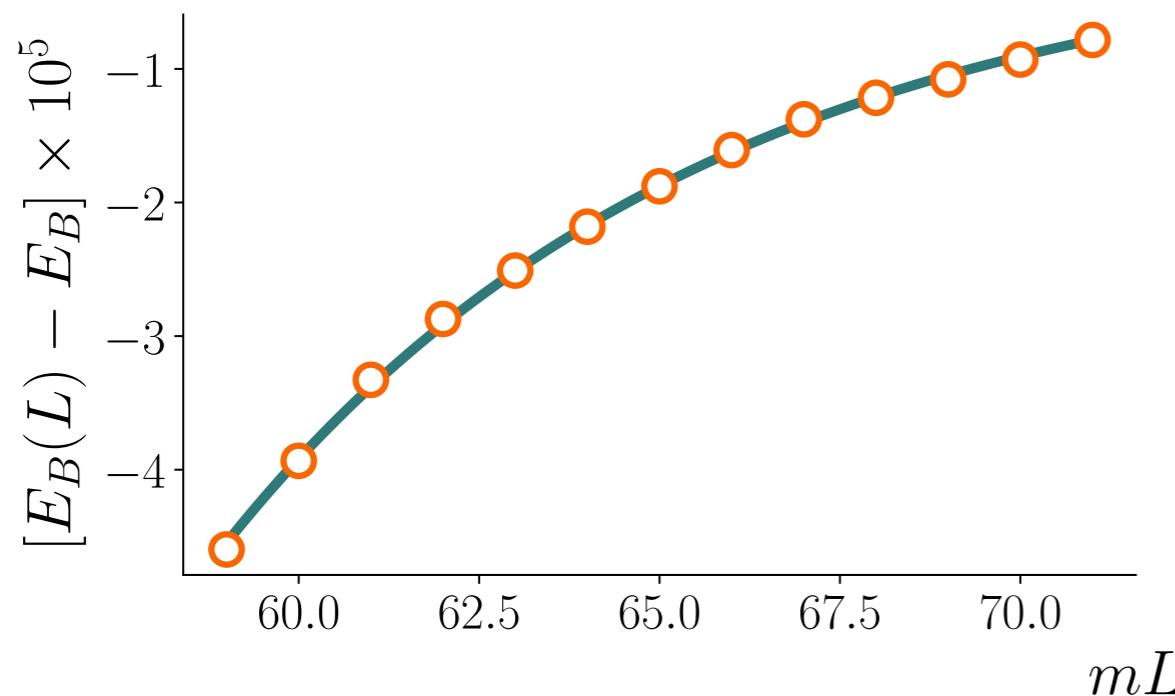
The parameters  $a = -10^4$ ,  $\mathcal{K}_{\text{df},3}^{\text{iso}}(E) = 2500$  lead to a shallow bound state

$$\kappa \approx 0.1m \text{ where } E_B = 3m - \kappa^2/m$$

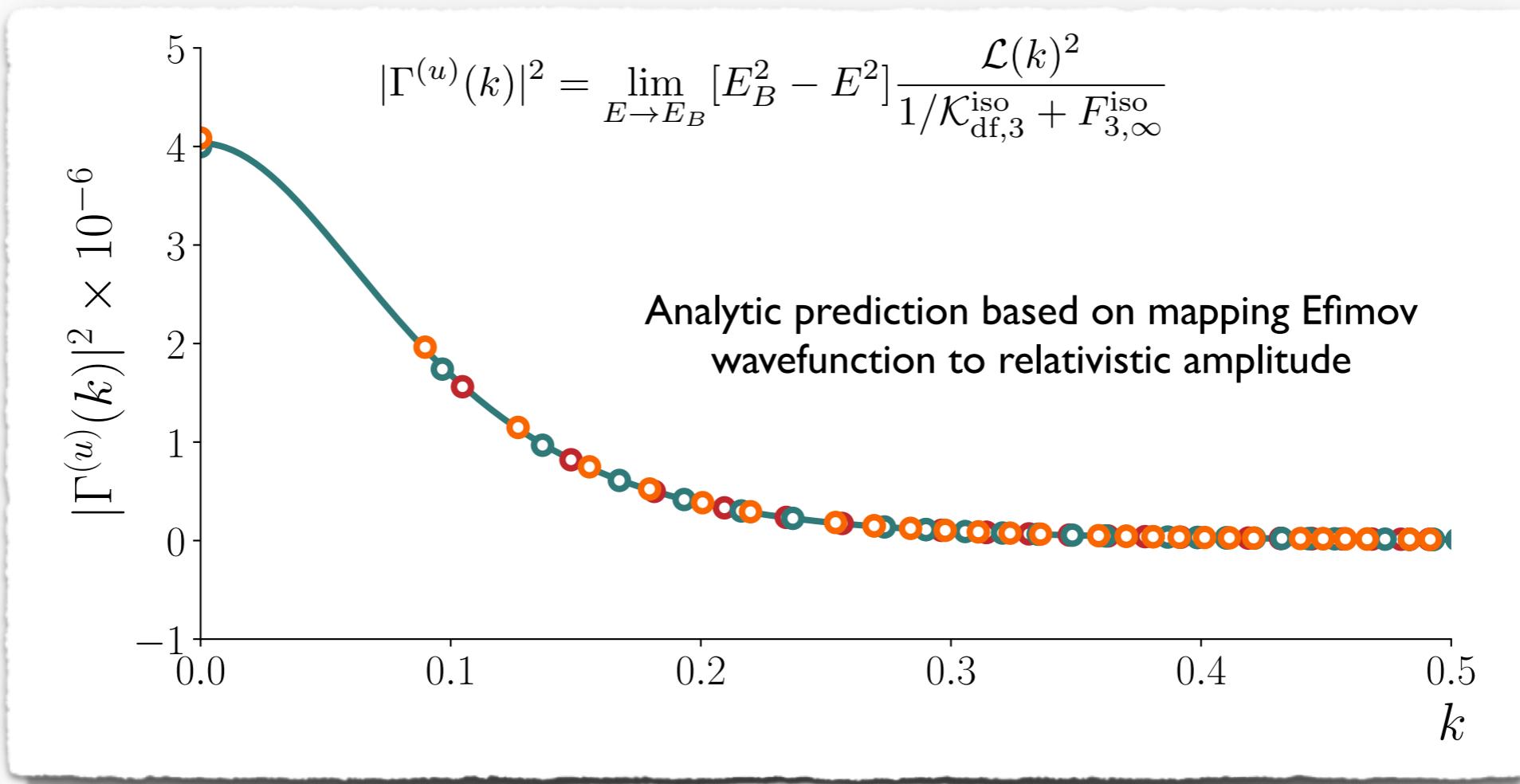
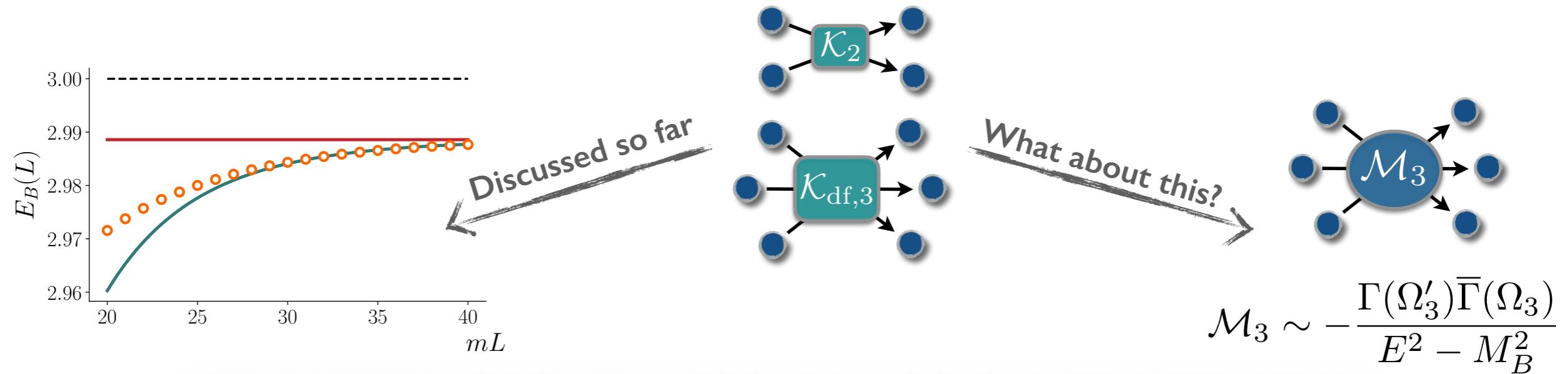
Finite-volume behavior of this state has a known asymptotic form

$$E_B(L) = 3m - \frac{\kappa^2}{m} - (98.35 \dots) |A|^2 \frac{\kappa^2}{m} \frac{e^{-2\kappa L/\sqrt{3}}}{(\kappa L)^{3/2}} \left[ 1 + \mathcal{O}\left(\frac{1}{\kappa L}, \frac{\kappa^2}{m^2}, e^{-\alpha \kappa L}\right) \right]$$

- Meißner, Rios, Rusetsky (2015) •

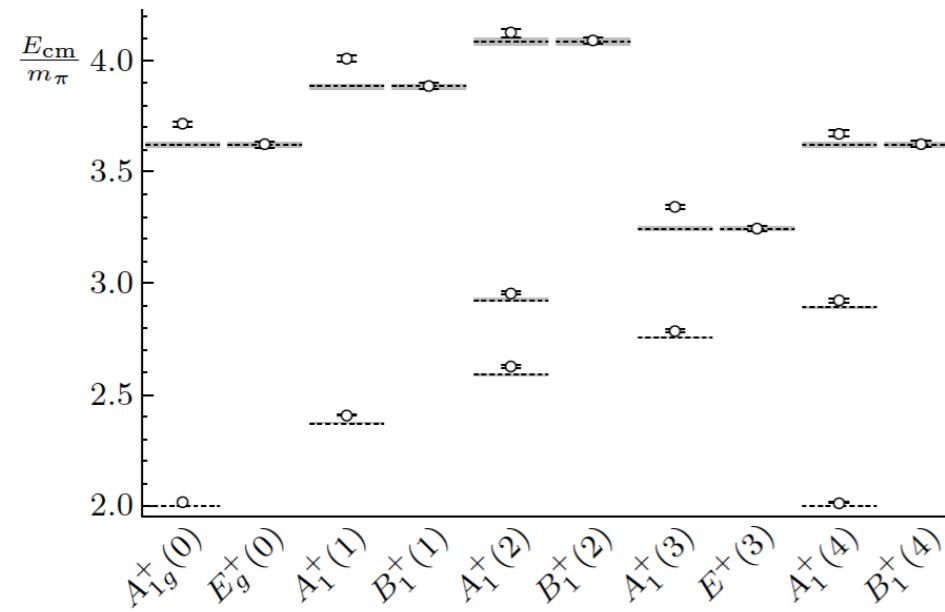


# Bound state scattering amplitude

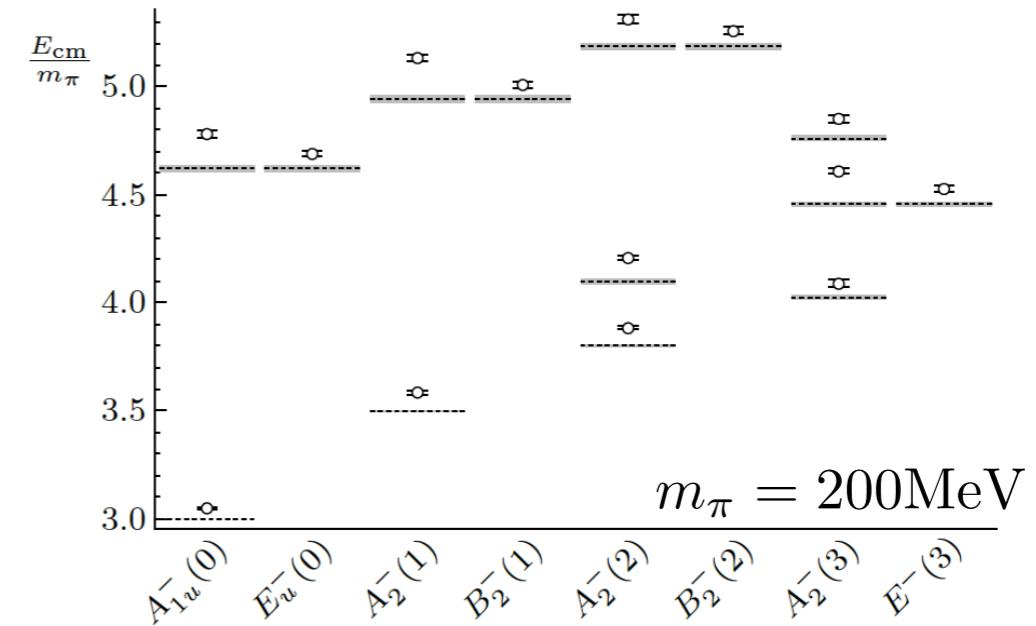


# Lattice QCD (CLS)

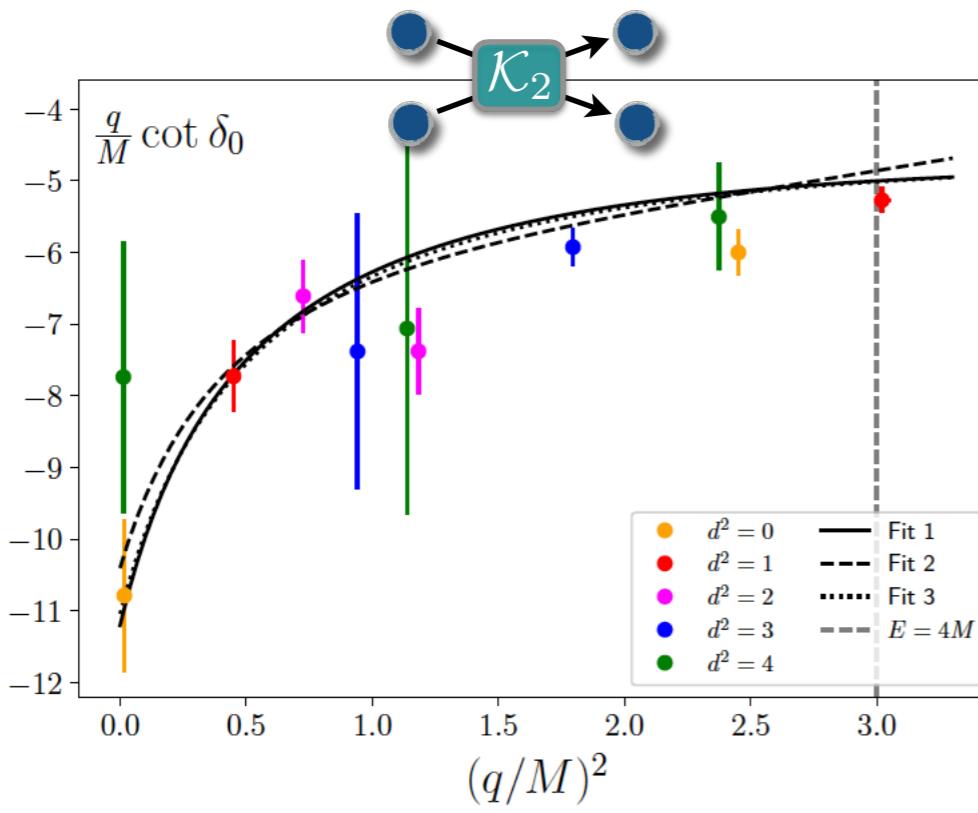
$\pi^+ \pi^+$



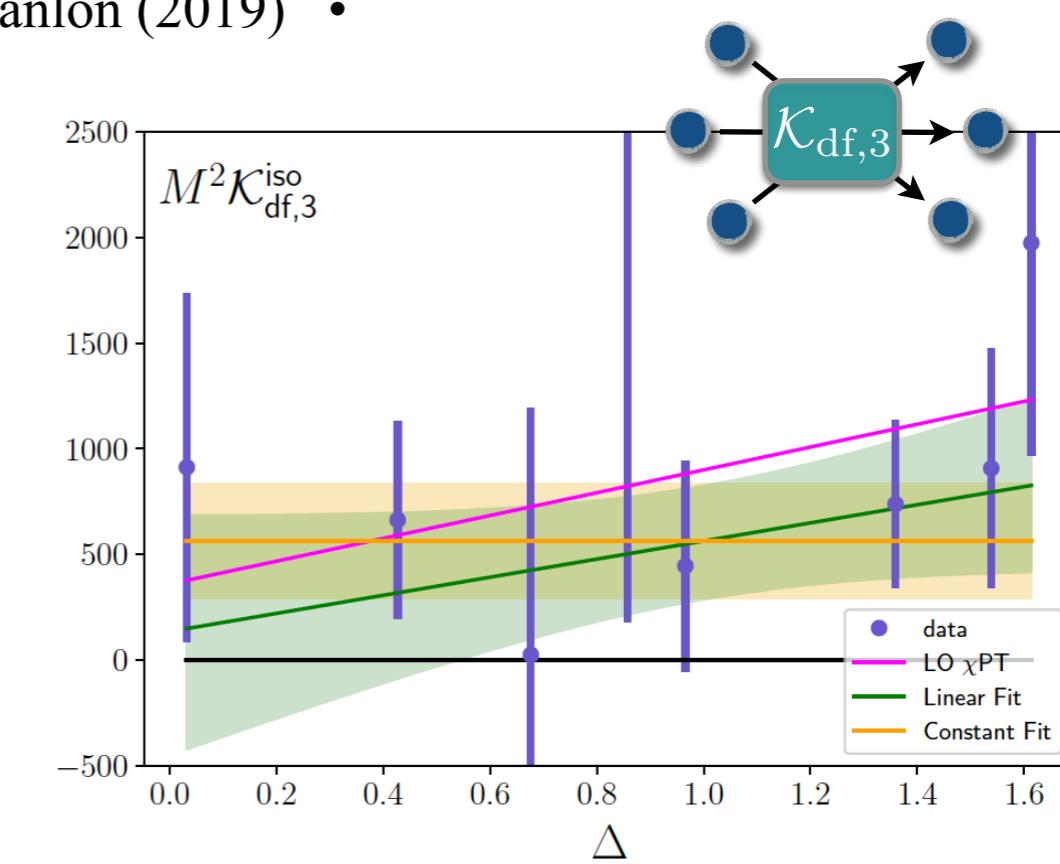
$\pi^+ \pi^+ \pi^+$



• Hörz and Hanlon (2019) •



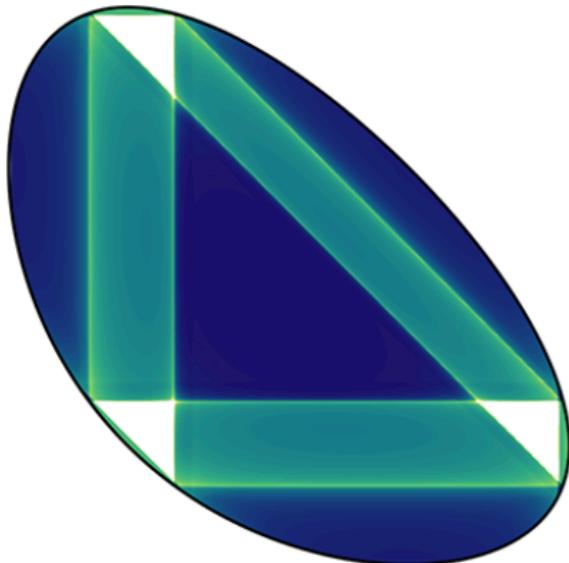
• Blanton, Romero-López, Sharpe (2019) • see also Mai, Döring (2019) •



# Energy-Dependent $\pi^+ \pi^+ \pi^+$ Scattering Amplitude from QCD

Maxwell T. Hansen<sup>ID, 1,2,\*</sup> Raul A. Briceño<sup>, 3,4,†</sup> Robert G. Edwards<sup>ID, 3,‡</sup>  
Christopher E. Thomas<sup>ID, 5,§</sup> and David J. Wilson<sup>ID, 5,||</sup>

(for the Hadron Spectrum Collaboration)



## EDITORS' SUGGESTION

### Energy-Dependent $\pi^+ \pi^+ \pi^+$ Scattering Amplitude from QCD

A three-hadron scattering amplitude is computed using lattice QCD for the first time.

Maxwell T. Hansen *et al.*

[Phys. Rev. Lett. 126, 012001 \(2021\)](#)



$$\pi^+ \pi^+ \pi^+ \rightarrow \pi^+ \pi^+ \pi^+$$

*lattice details*

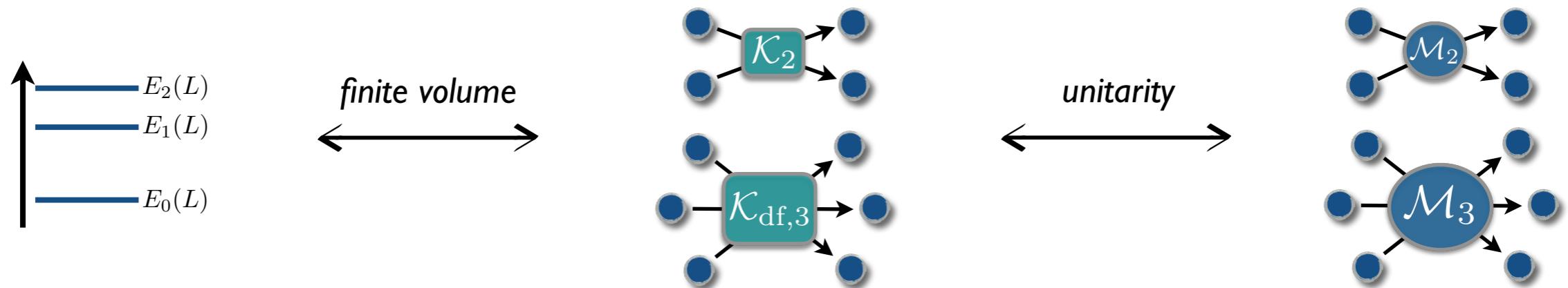
$$N_f = 2 + 1 \quad a_s/a_t = 3.444(6)$$

$$m_\pi \approx 400\text{MeV} \quad a_s \approx 0.12\text{fm}$$

$$L_s/a_s = 20, 24$$

$$\begin{array}{c} \bullet & \bullet & \bullet & \bullet \\ \vdots & \vdots & \vdots & \vdots \\ \bar{a}_t & \bullet & \bullet & \bullet \\ \hline & a_s & & \end{array}$$

## □ Workflow outline



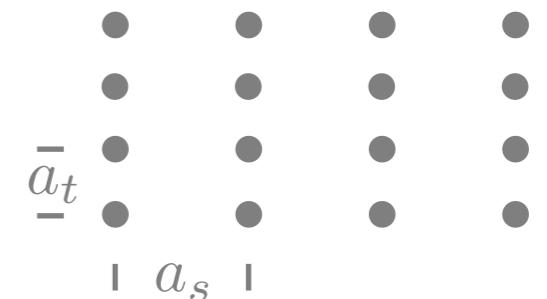
$$\pi^+ \pi^+ \pi^+ \rightarrow \pi^+ \pi^+ \pi^+$$

*lattice details*

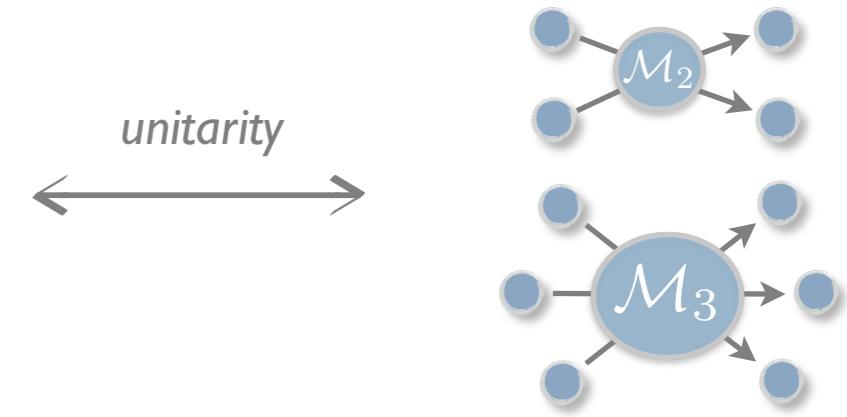
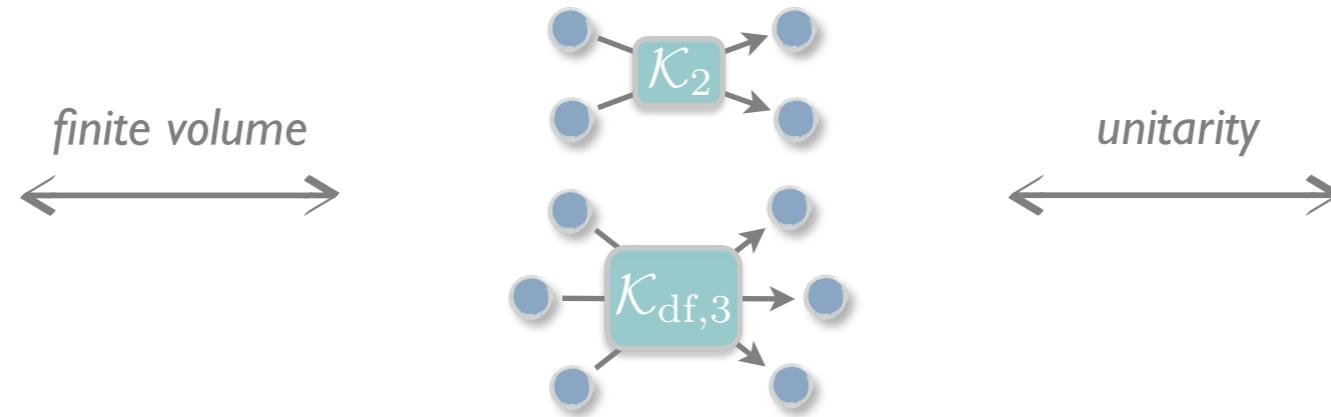
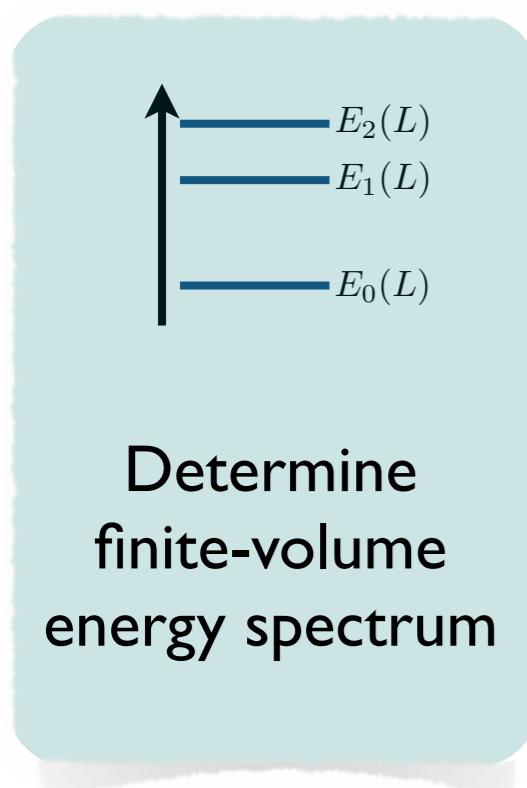
$$N_f = 2 + 1 \quad a_s/a_t = 3.444(6)$$

$$m_\pi \approx 400\text{MeV} \quad a_s \approx 0.12\text{fm}$$

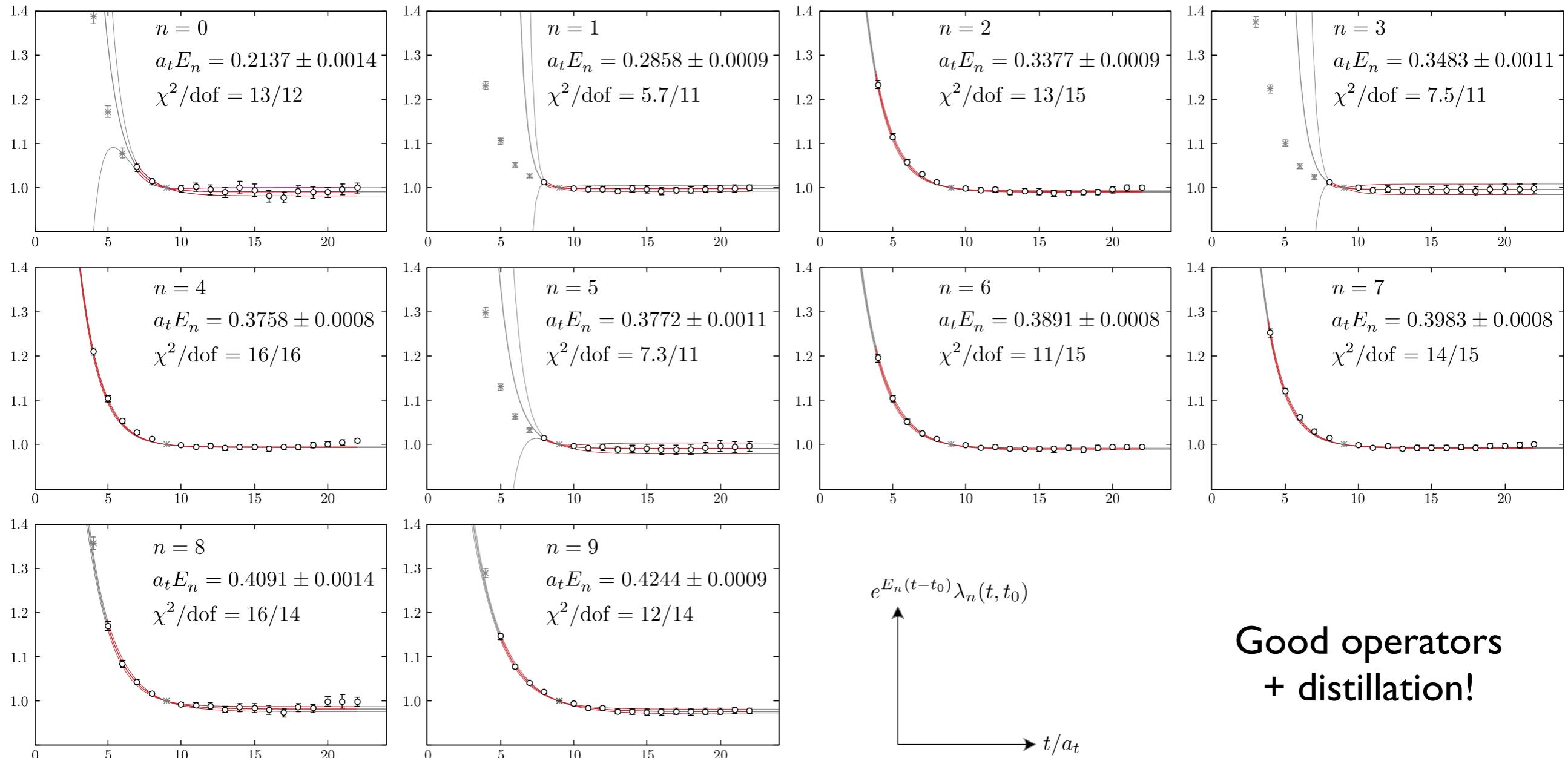
$$L_s/a_s = 20, 24$$



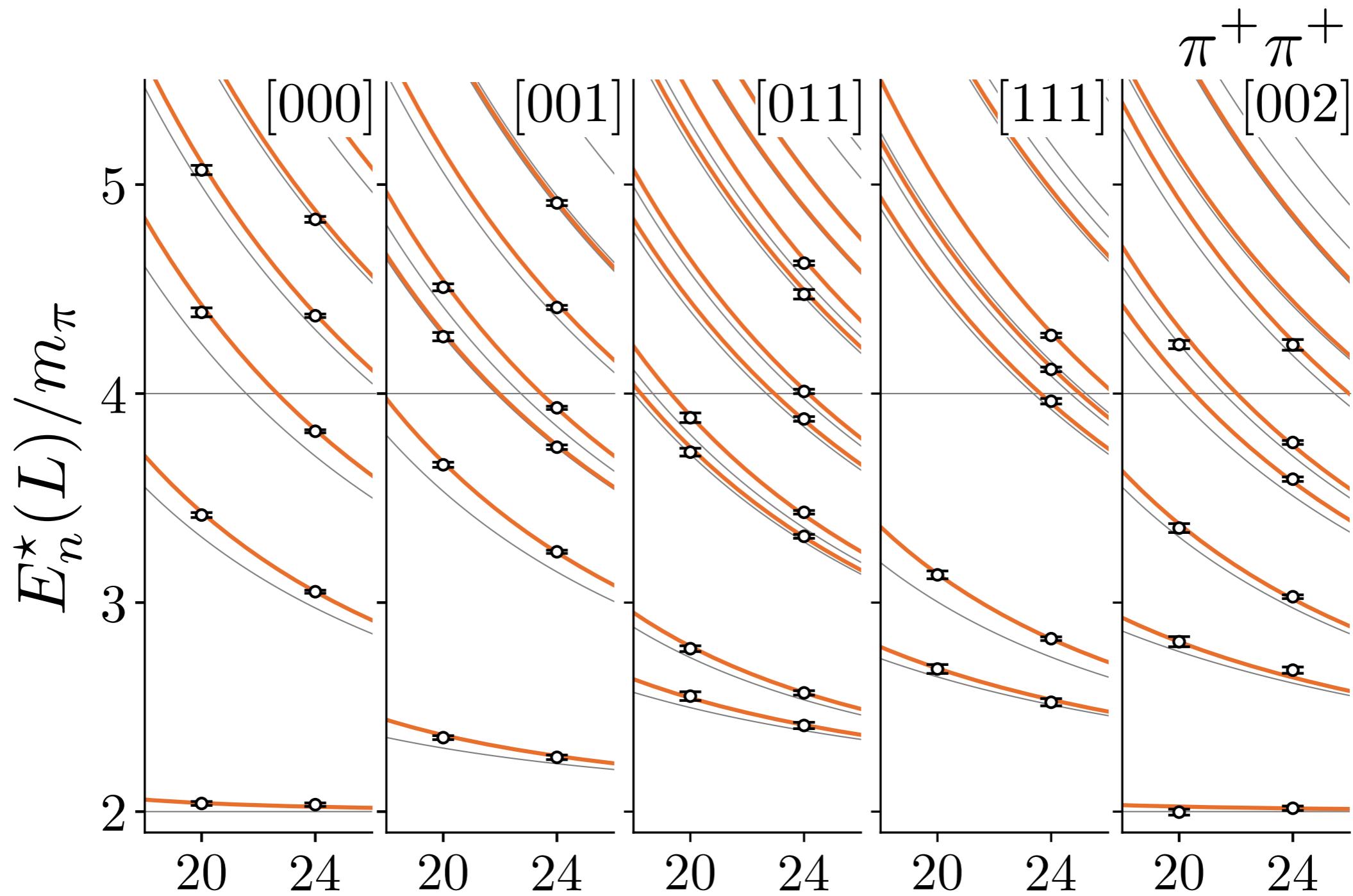
## □ Workflow outline



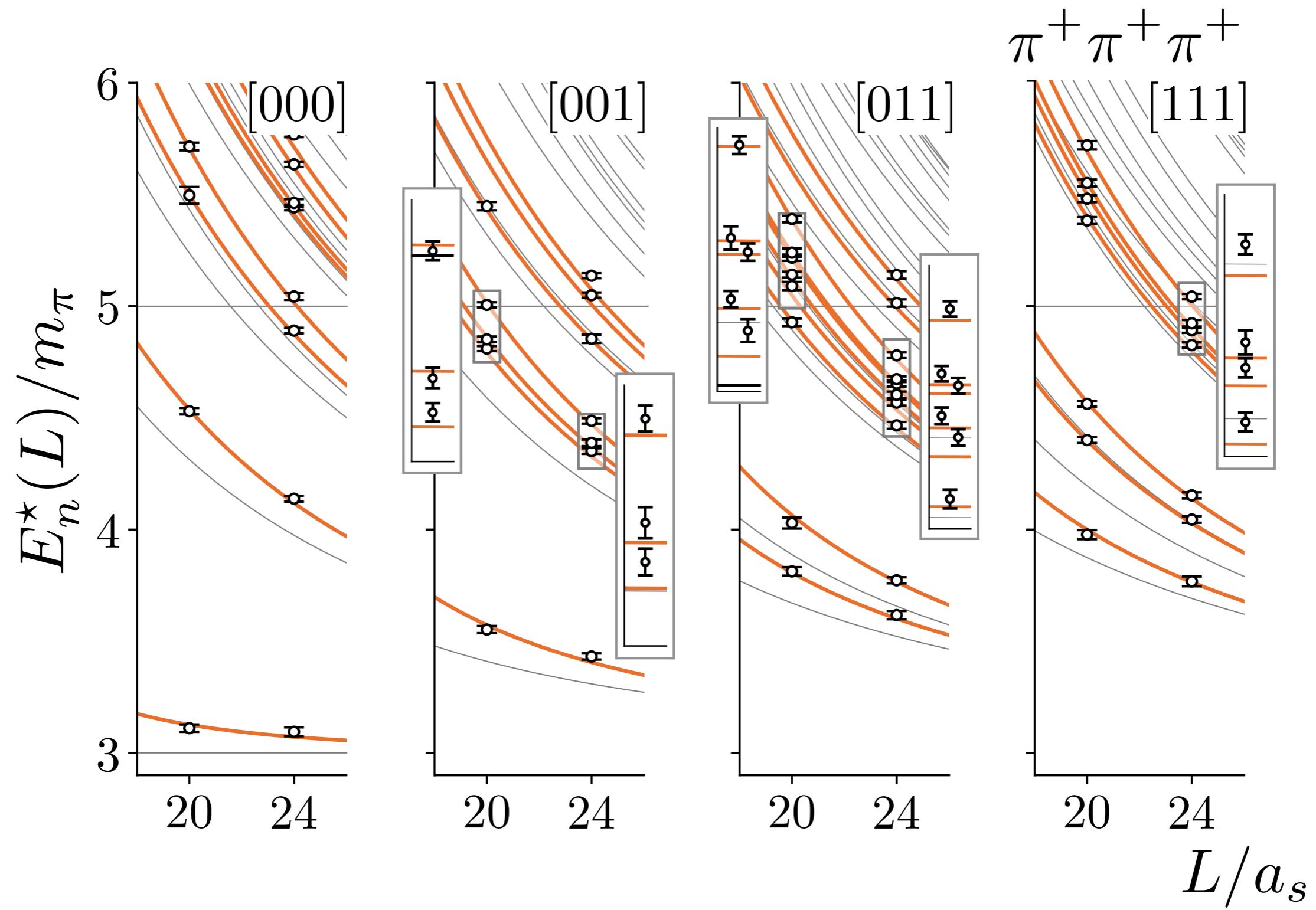
$$I = 3 (\pi^+ \pi^+ \pi^+), \quad P = [000], \quad \Lambda = A_1^-, \quad L/a_s = 24$$



# $\pi^+ \pi^+$ energies



# $\pi^+\pi^+\pi^+$ energies



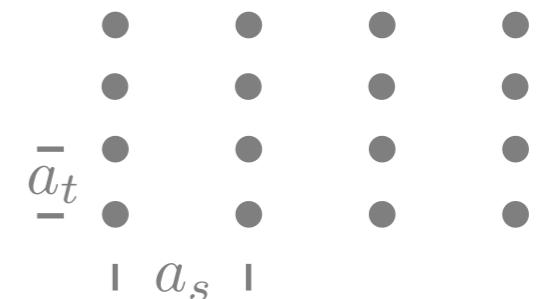
$$\pi^+ \pi^+ \pi^+ \rightarrow \pi^+ \pi^+ \pi^+$$

*lattice details*

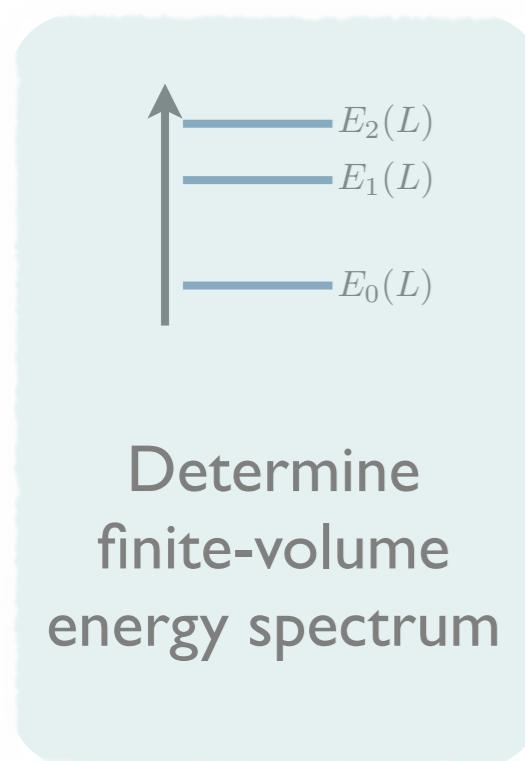
$$N_f = 2 + 1 \quad a_s/a_t = 3.444(6)$$

$$m_\pi \approx 400\text{MeV} \quad a_s \approx 0.12\text{fm}$$

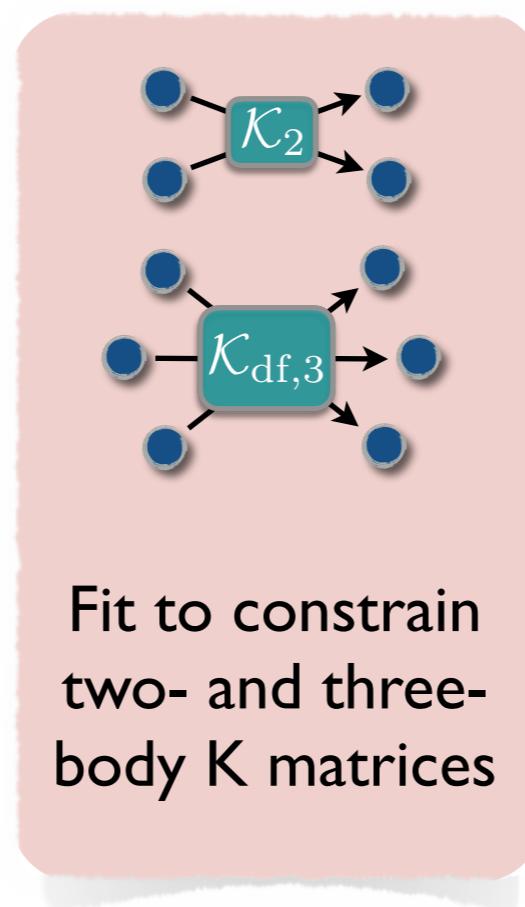
$$L_s/a_s = 20, 24$$



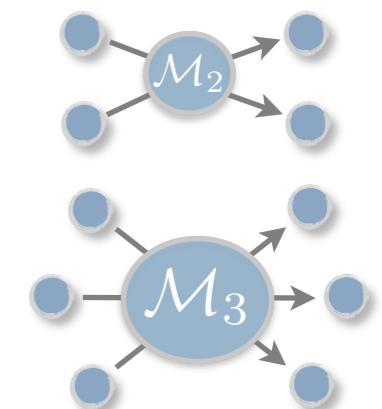
## □ Workflow outline



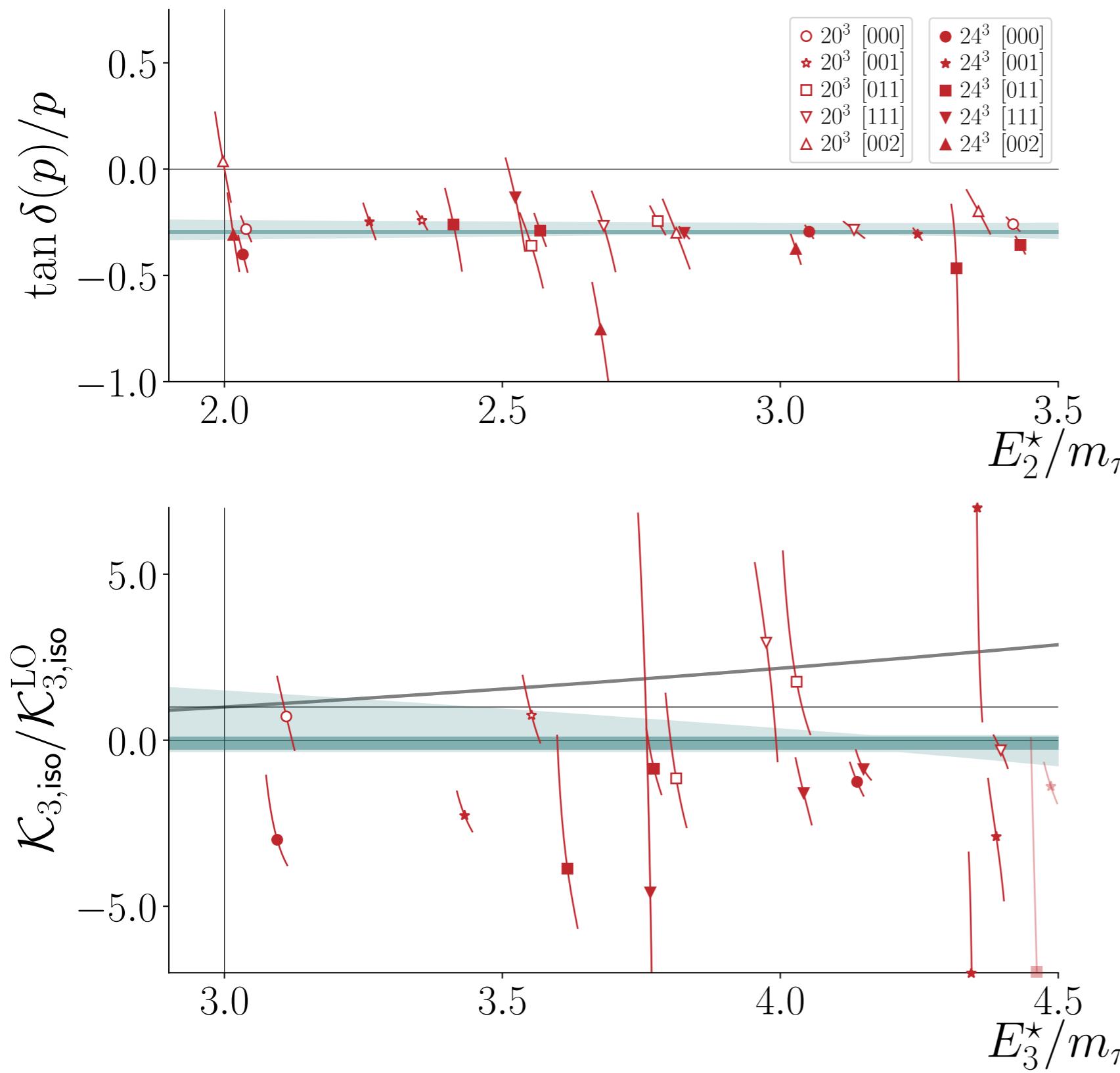
*finite volume*



*unitarity*



# K matrix fits



Finite-volume formalism  
relates energies to K matrices

One-to-one for  $K_{\text{df},3}$   
depending only on  $E_{\text{cm}} = E^\star$

Fit both two and three-body  
K to various polynomials

Cut on the CM  
energy in the fits

$K_{\text{df},3}$  is scheme  
dependent (removed  
upon converting to  $\mathcal{M}_3$ )

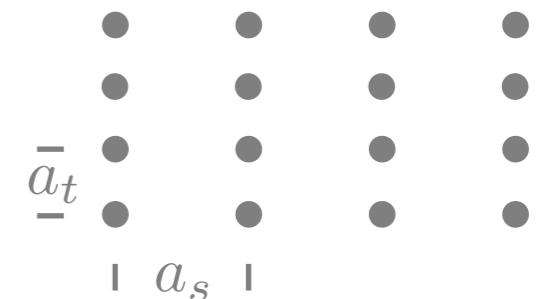
$$\pi^+ \pi^+ \pi^+ \rightarrow \pi^+ \pi^+ \pi^+$$

*lattice details*

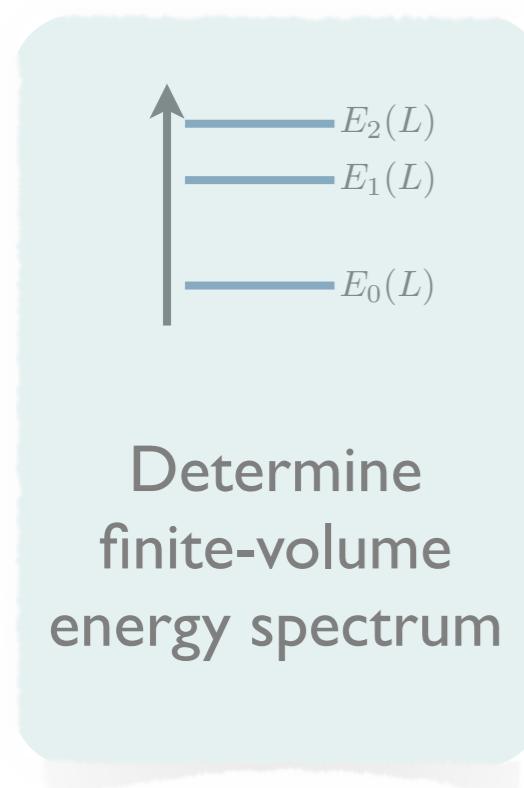
$$N_f = 2 + 1 \quad a_s/a_t = 3.444(6)$$

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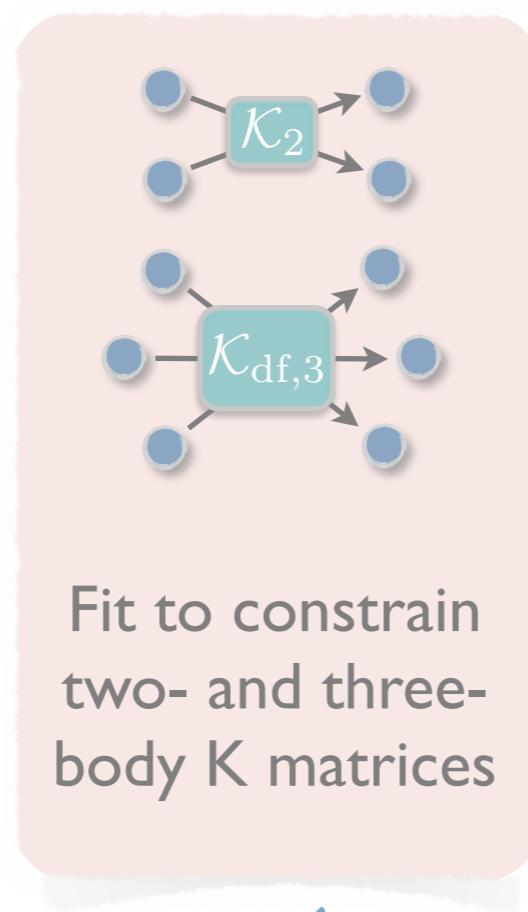
$$L_s/a_s = 20, 24$$



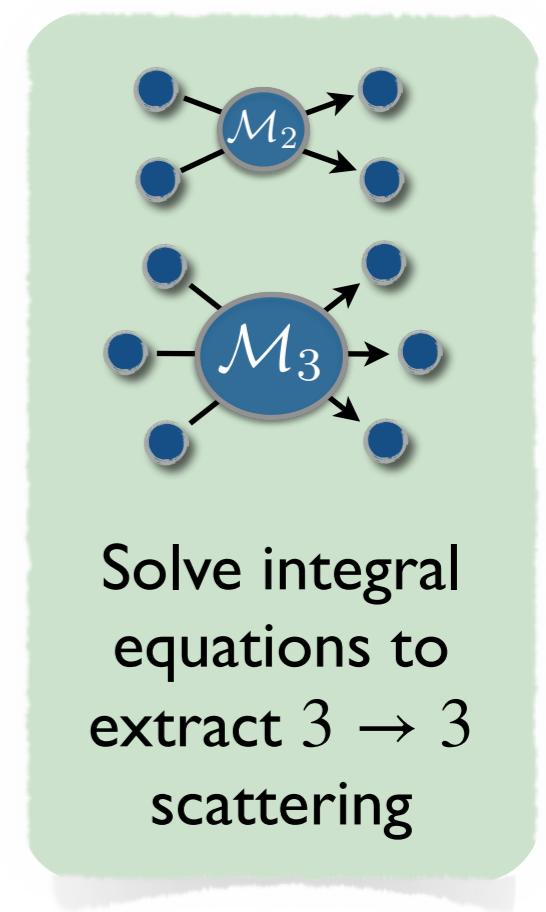
## □ Workflow outline



*finite volume*

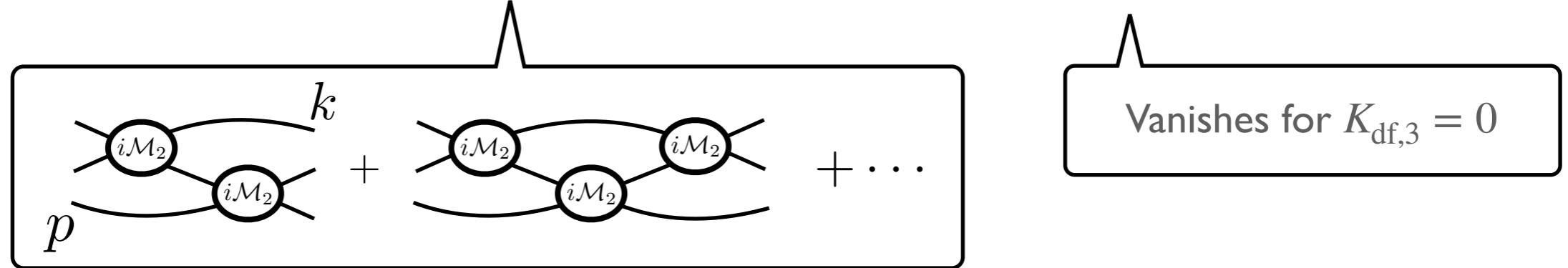


*unitarity*



# Integral equation

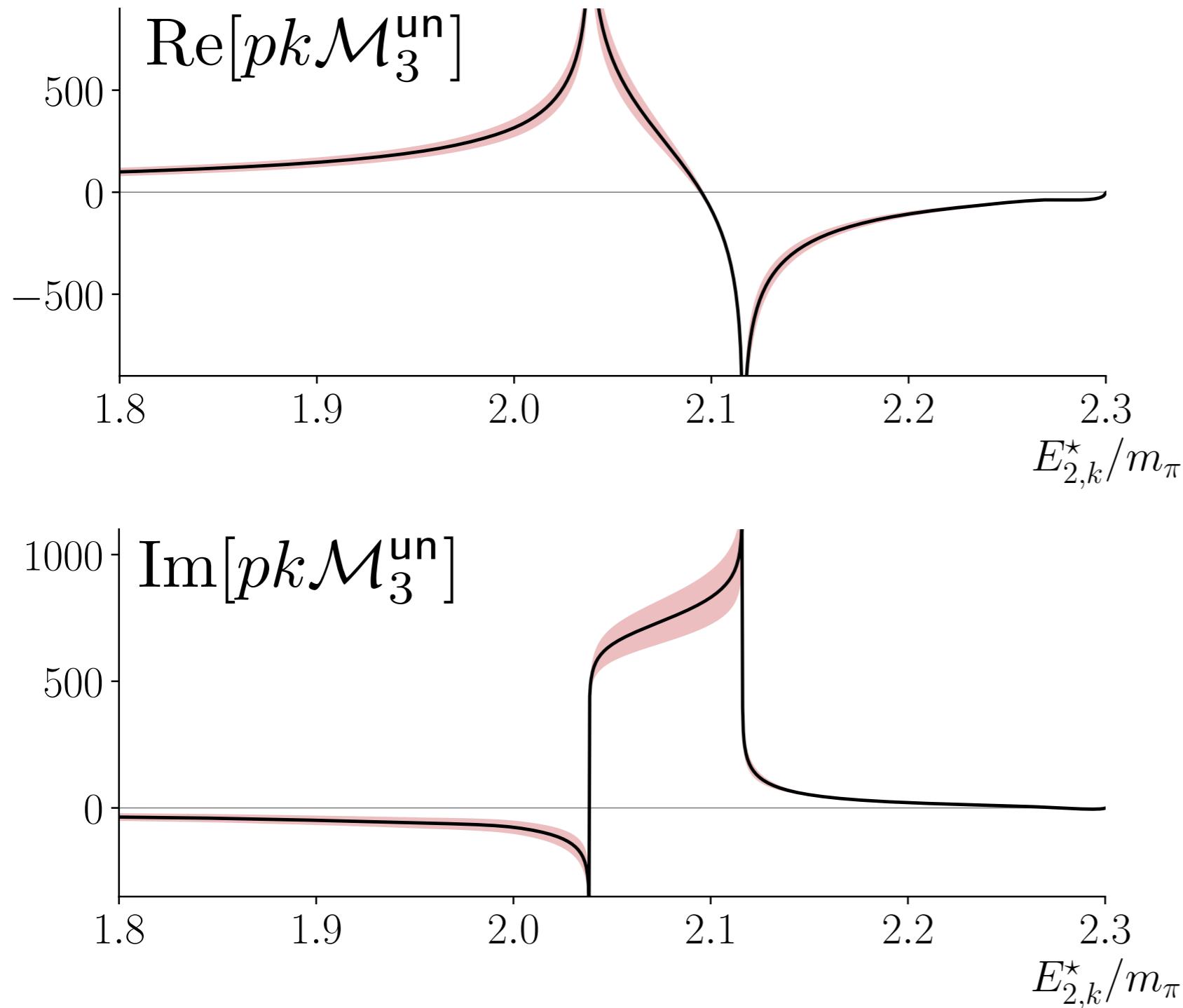
$$\mathcal{M}_3^{\text{un}}(E_3^*, \mathbf{p}, \mathbf{k}) = \mathcal{D}^{\text{un}}(E_3^*, \mathbf{p}, \mathbf{k}) + \mathcal{E}^{\text{un}}(E_3^*, \mathbf{p}) \mathcal{T}(E_3^*) \mathcal{E}^{\text{un}}(E_3^*, \mathbf{k})$$



$$D(N, \epsilon) = -\mathcal{M} \cdot G(\epsilon) \cdot \mathcal{M} - \mathcal{M} \cdot G(\epsilon) \cdot P \cdot D(N, \epsilon)$$

$$\mathcal{D}^{\text{un}}(E_3^*, \mathbf{p}, \mathbf{k}) = \lim_{\epsilon \rightarrow 0} \lim_{N \rightarrow \infty} D_{pk}(N, \epsilon)$$

# Integral equation



Total angular momentum = 0

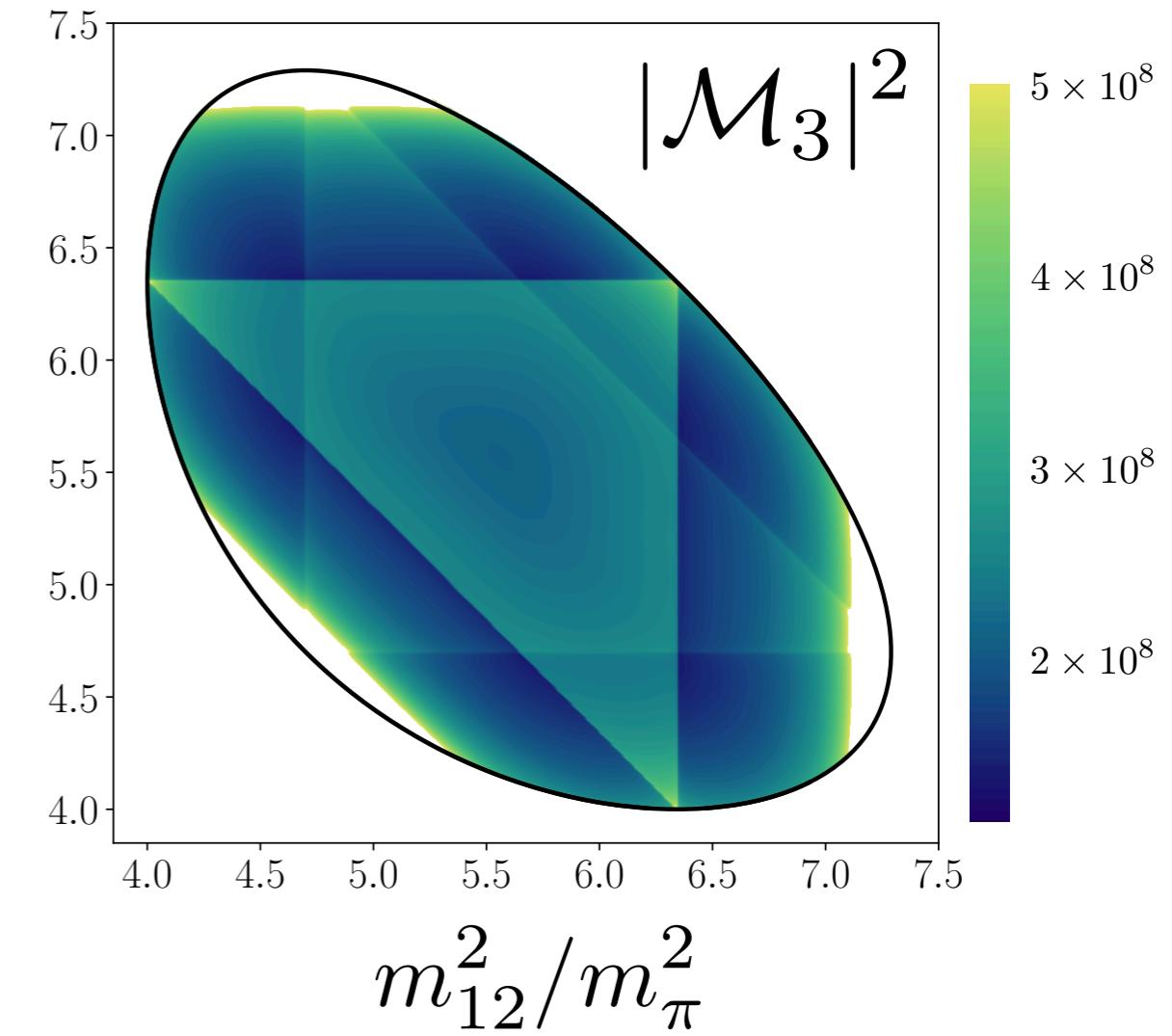
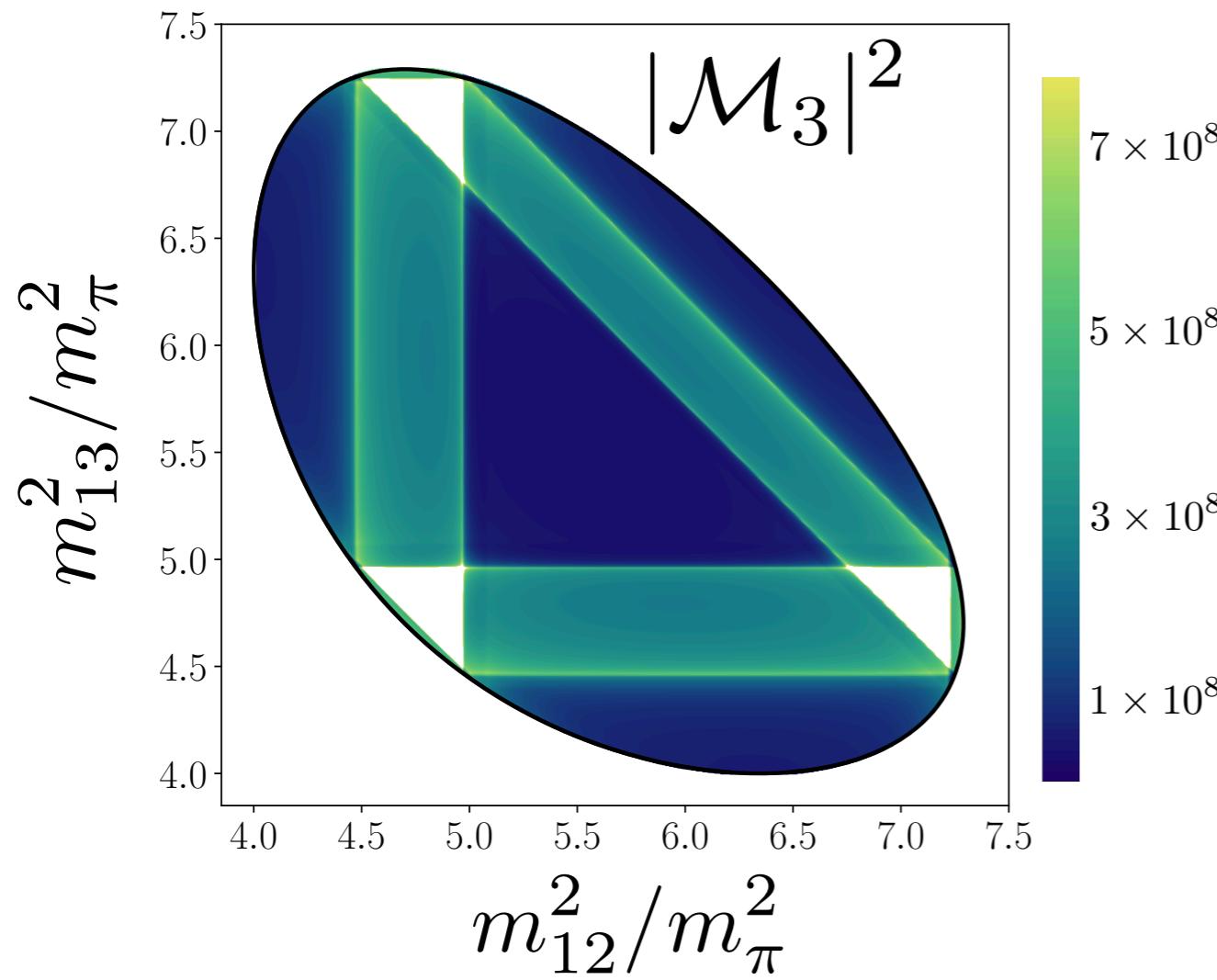
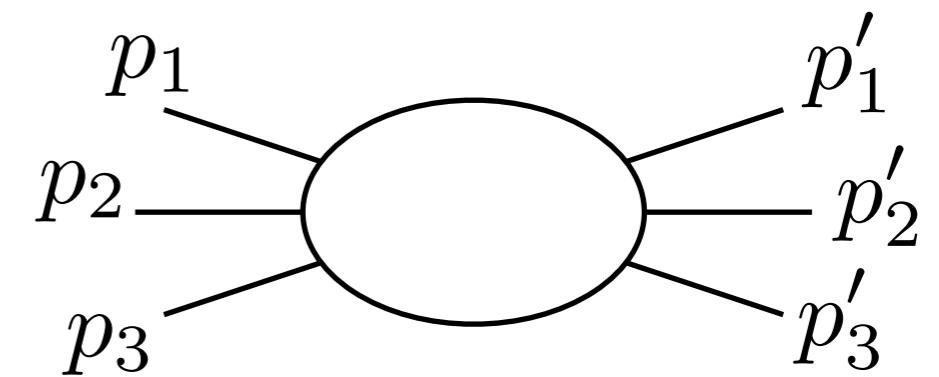
Two-particle sub-system  
angular momentum = 0

Plot at fixed  $E_3^*$  and  $p$

Both two- and three-body  
uncertainties estimated

Still need to symmetrize

$$\mathcal{M}_3 = \sum_{i,j \in \{1,2,3\}} \mathcal{M}_3^{\text{un}}(p'_i, p_j)$$



# Outlook for exotics

## □ Challenging for many reasons

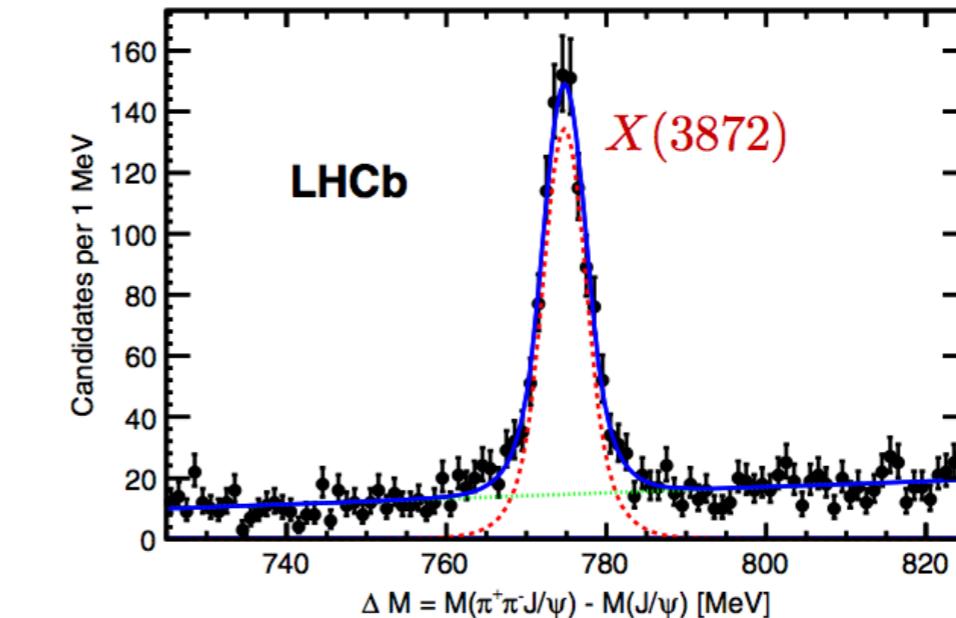
*incredible hierarchy of scales*

$$M_X - M_{D^0 D^{0*}} = 0.01 \pm 0.18 \text{ MeV}$$

$$\Gamma(X) < 1.2 \text{ MeV}$$

- Lebed, Mitchell, Swanson, *PPNP review* (2017)

*many open channels*



- LHCb (PRD92, 2015)

$$X \rightarrow \omega J/\psi, \pi\pi J/\psi, D^{*0}\bar{D}^0, \dots$$

## □ My speculative outlook

*should be possible to derive N-particle volume formalism*

*then relate K-matrices to spectra to identify LQCD strategy*

*hierarchy of scales = not always a problem*

$$\mathcal{K}_2 \propto a \rightarrow \infty$$

# Five-day outline

- Warm-up and definitions
  - Meaning of Euclidean
  - Finite-volume set-up
- $e^{-mL}$  volume effects
  - Mass, matrix element in  $\lambda\phi^4, g\phi^3$
  - Pion mass
  - LO-HVP for  $(g - 2)_\mu$
  - Bethe-Salpeter kernel
- $2 \rightarrow 2$  finite-volume (f.v.) formalism
  - Scattering basics
  - Derivation
  - Generalizations
  - Applications
- $(1+\mathcal{J}) \rightarrow 2$  f.v. formalism
  - Derivation
  - Expansions
  - Applications
- $2 + \mathcal{J} \rightarrow 2$  f.v. formalism
  - Derivation
  - Testing the result
  - Numerical explorations
- Non-local matrix elements
  - Derivation
  - Applications
- $3 \rightarrow 3$  f.v. formalism
  - New complications
  - Derivation ( $E_n(L)$  to  $\mathcal{K}_{df,3}$ )
  - Integral equations ( $\mathcal{K}_{df,3}$  to  $\mathcal{M}_3$ )
  - Testing the result
  - Numerical explorations/calculations
- Finite-volume QED + QCD
- Different formalisms
- Basic examples